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SENSITIVITY ANALYSIS OF TRACK VIBRATIONS DUE TO VERTICAL STIFFNESS VARIATION IN HIGH-SPEED RAILWAYS

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Abstract. *High speed trains, when crossing regions with abrupt changes in vertical stiffness of the track and/or subsoil, may generate excessive ground and track vibrations. There is an urgent need for specific analyses of this problem so as to allow reliable estimates of vibration amplitude. Full understanding of these phenomena will lead to new construction solutions and mitigation of undesirable features. In this paper analytical transient solutions of dynamic response of one-dimensional systems with sudden change of foundation stiffness are derived. Results are expressed in terms of vertical displacement. Sensitivity analysis of the response amplitude is also performed. The analytical expressions presented herein, to the authors' knowledge, have not been published yet. Although related to one-dimensional cases, they can give useful insight into the problem. Nevertheless, in order to obtain realistic response, vehicle-rail interaction cannot be omitted. Results and conclusions are confirmed using general purpose commercial software ANSYS. In conclusion, this work contributes to a better understanding of the additional vibration phenomenon due to vertical stiffness variation, permitting better control of the train velocity and optimization of the track design.*

1 INTRODUCTION

High speed trains, when crossing regions with abrupt changes in vertical stiffness of the track and/or subsoil, may generate excessive ground and track vibrations, ballast projections and increase the wheel, rail and track deterioration with negative impact on the vehicle stability and passengers comfort. Therefore specific methodologies must be established in order to determine these vibrations. Full understanding of these phenomena will lead to new construction solutions and mitigation of undesirable features.

Detailed information about these vibrations can be gathered from 3D finite element models. However, it is known that these analyses require very fine meshes to accurately capture the generated waves, giving rise to time consuming computation impractical for ordinary design. Analyses incorporating stiffness transitions necessitate more careful treatment in the region of the “discontinuity”, which naturally further increases the computing time.

First insight on these phenomena can be gained from simplified one-dimensional models. Steady-state analytical solutions of dynamic response of a beam on elastic or more complex foundation under moving loads are widely available [1, 2]. However, introduction of non-homogeneous properties of the system naturally requires transient response, for which available results in the literature remain scarce.

In this paper, analytical transient solutions of dynamic response of one-dimensional systems with sudden change of foundation stiffness are derived. Results are expressed in terms of vertical displacement. The problem is approached from the simplest level. First, a simply supported beam is considered. Transient solution for moving force is extended to a beam on elastic foundation in Section 2. Dynamic response accounting for localized abrupt stiffness change is derived in Section 3 and sensitivity analysis of displacement at some chosen position with respect to stiffness is performed in Section 4. Next, in Section 5, the cantilever transient solution is extended to account for elastic foundation and two cantilevers are joined together to form a clamped beam by imposing continuity conditions at the point of foundation stiffness change (Section 6). These results permit to study force passage through contiguous locations with different foundation stiffness. All presented results are confirmed using general purpose finite element code ANSYS. Graphs and numerical results are presented. Nevertheless, it must be pointed out that, in order to obtain realistic dynamic response, vehicle spring-mass-damper system interacting with the rail track must also be considered. This point will be object of future research.

2 MOVING LOAD ON SIMPLY SUPPORTED BEAM ON ELASTIC FOUNDATION

The governing equation describing the dynamic response, under a constant moving load, P , of an Euler-Bernoulli’s beam in terms of displacements can be written as [3]:

$$EI \frac{\partial^4 w(x, t)}{\partial x^4} + \mu \frac{\partial^2 w(x, t)}{\partial t^2} + 2\mu\omega_b \frac{\partial w(x, t)}{\partial t} = \delta(x - ct)P. \quad (1)$$

It is assumed that the beam follows linear elastic Hooke’s law, has constant cross-section and constant mass per unit length, μ . As usual, small displacements, Navier’s hypothesis and Saint-Venant’s principle are adopted. E , I and μ_b stand for Young’s modulus, moment of inertia and circular frequency of damping, respectively; w represents the vertical deflection measured from equilibrium position and oriented downwards, x is spatial coordinate measured from left to right end of the beam and t is the time. It is also assumed that the mass of the load is small compared with the mass of the beam and that the load moves with constant

speed, c . δ in equation (1) stands for the Dirac function. Denoting by L the total length of the beam, boundary and initial conditions read as:

$$w(0, t) = 0, \quad w(L, t) = 0, \quad \left. \frac{\partial w^2(x, t)}{\partial x^2} \right|_{x=0} = 0, \quad \left. \frac{\partial w^2(x, t)}{\partial x^2} \right|_{x=L} = 0, \quad (2)$$

$$w(x, 0) = 0, \quad \left. \frac{\partial w(x, t)}{\partial t} \right|_{t=0} = 0 \quad (3)$$

Problem (1-3) can be solved by methods of integral transformations. First, Fourier sine finite integral transformation is implemented and then Laplace-Carson integral transformation is applied, enabling to express the displacement in terms of sine series [3].

In order to account for the effect of elastic foundation, characterized by Winkler's constant k , an additional term must be introduced into equation (1), which takes the following form:

$$EI \frac{\partial w^4(x, t)}{\partial x^4} + \mu \frac{\partial w^2(x, t)}{\partial t^2} + 2\mu\omega_b \frac{\partial w(x, t)}{\partial t} + kw(x, t) = \delta(x - ct)P \quad (4)$$

The additional term changes the image of Laplace-Carson transformation. Nevertheless, basically the same procedure as in [3] can be used to obtain the final expression for the vertical displacements. In order to validate this procedure, an equivalent model was created using ANSYS software [4]. Element BEAM 54 of ANSYS library, with the capacity of introduction of elastic foundation, was used. Rayleigh damping was introduced by means of the coefficient $\alpha = 2\omega_b$. Since ANSYS does not allow moving load implementation, for each time step a new force position had to be considered according to the load speed and the element size. The computational results matched almost exactly the analytical solution, confirming the suitability of the strategy adopted for the analysis and suggesting that it is possible to solve numerically other situations impossible to treat analytically.

Verification graphs are shown for representative cases, without attempt to adjust material and geometrical properties to some real representation of railway track. Parameters ξ , ψ and ζ can be introduced in order to characterize the magnitude of velocity, c , and the damping level, ω_b . Thus, ξ and ζ stand for the ratio of circular frequency of the load, ω , and first circular frequency of the free vibration of simply supported beam without, $\omega_{(1)}$, and with, $\tilde{\omega}_{(1)}$, elastic foundation. Similarly, ψ is the ratio between the damping circular frequency ω_b and $\omega_{(1)}$. One could then write:

$$\xi = \frac{\omega}{\tilde{\omega}_{(1)}}, \quad \psi = \frac{\omega_b}{\omega_{(1)}}, \quad \zeta = \frac{\omega}{\omega_{(1)}}, \quad (5)$$

where: $\omega = \frac{\pi c}{L}$, $\omega_{(1)} = \frac{\pi^2}{L^2} \sqrt{\frac{EI}{\mu}}$, $\tilde{\omega}_{(1)} = \sqrt{\frac{\pi^4}{L^4} \frac{EI}{\mu} + \frac{k}{\mu}}$.

To assess the match between analytical and numerical results, a beam with $L=10\text{m}$, $EI=1000\text{Nm}^2$, $P=1\text{N}$ and $\mu=0,06\text{kg/m}$ was considered. Several combinations of ξ , ψ and ζ were chosen. In Figures 1 and 2, graphs for the cases $\xi=0.476$, $\psi=0.5$, $\zeta=0.5$ and $\xi=2$, $\psi=5$, $\zeta=2.1$ are plotted for ten positions of the load along the beam length, namely at each precise meter. It can be observed that the violet analytical curves completely overwrite the numerical blue ones, which shows the excellent agreement previously mentioned.

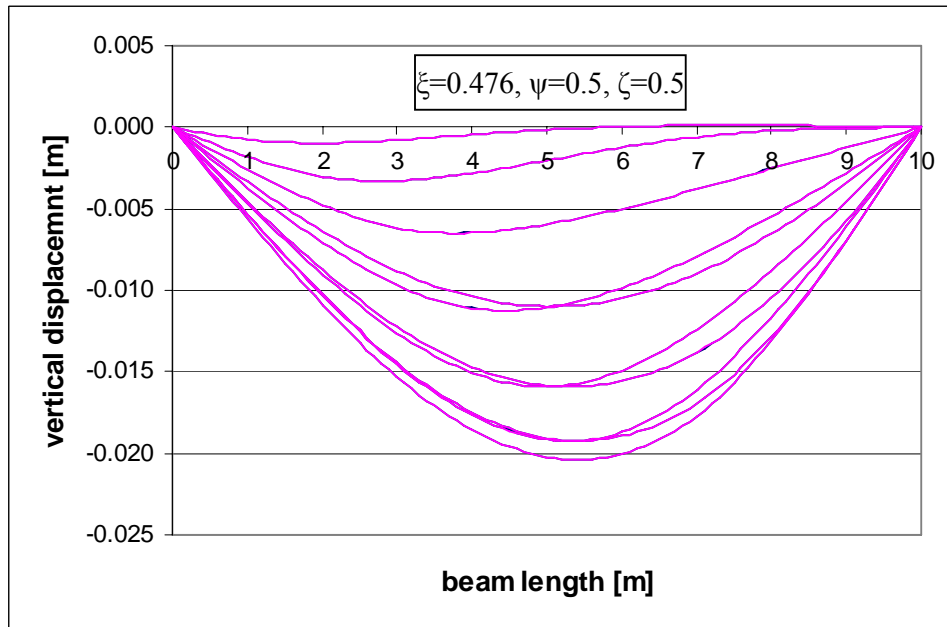


Figure 1: Vertical displacement for $\xi=0.476$, $\psi=0.5$ and $\zeta=0.5$, plotted for 10 positions of the moving load (violet curve: analytical solution; blue curve: numerical results)

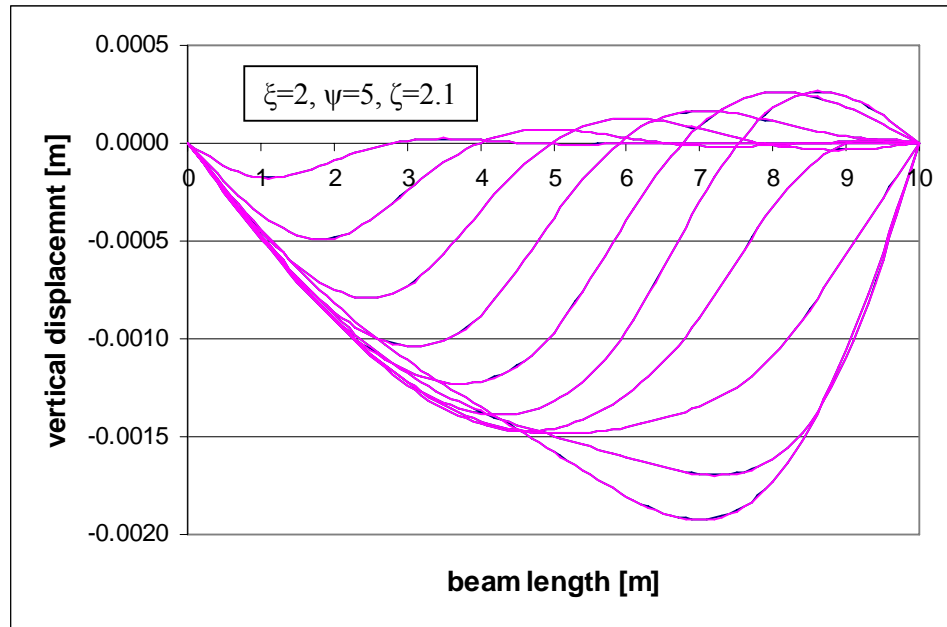


Figure 2: Vertical displacement for $\xi=2$, $\psi=5$ and $\zeta=2.1$, plotted for 10 positions of the moving load (violet curve: analytical solution; blue curve: numerical results)

3 MOVING LOAD ON SIMPLY SUPPORTED BEAM ON ELASTIC FOUNDATION WITH LOCALIZED CHANGE IN VERTICAL STIFFNESS

In this case it is necessary to add an additional term to Equation (4) in the following way:

$$EI \frac{\partial w^4(x, t)}{\partial x^4} + \mu \frac{\partial w^2(x, t)}{\partial t^2} + 2\mu\omega_b \frac{\partial w(x, t)}{\partial t} + kw(x, t) + \delta(x - x_0)k_0 w(x, t) = \delta(x - ct)P \quad (6)$$

where x_0 corresponds to the position of the sudden stiffness change and k_0 stands for the addition of Winkler's constant at the position x_0 . Therefore, this effect corresponds to the introduction of a discrete elastic spring. Although it is impossible to solve Equation (6) exactly, an approximate solution is proposed herein using an approach similar as another described in [3] for a different context. It is assumed that:

$$w(x, t)|_{x=x_0} \cong \frac{2}{L} W(j, t) \sin \frac{j\pi x_0}{L} \quad (7)$$

where $W(j, t)$ stands for the image of $w(x, t)$ in Fourier sine finite integral transformation. After merging Equation (7) into (6), the procedure to obtain the final solution is similar to the one described in previous section. It can then be written:

$$w(x, t) \cong \frac{2}{L} \sum_{j=1}^{\infty} W(j, t) \sin \frac{j\pi x}{L} \quad (8)$$

where

$$W(j, t) = \frac{Pc}{\mu \left((a^2 + b^2 - c^2)^2 + 4a^2c^2 \right)} \cdot \left(\frac{(a^2 + b^2 - c^2)}{c} \sin(ct) - \frac{(b^2 - a^2 - c^2)}{b} e^{-at} \sin(bt) - 2a(\cos(ct) - e^{-at} \cos(bt)) \right), \quad (9)$$

and

$$a = \omega_b, \quad b = \sqrt{\tilde{\omega}_{(j)}^2 - \omega_b^2 + \frac{2k_0}{L\mu} \sin^2 \left(\frac{j\pi x_0}{L} \right)}, \quad c = j\omega. \quad (10)$$

Obviously, when $k_0=0$, Equations (8-10) reduce to the problem (1-3) addressed in the previous section.

Analytical (approximate) results according to (8-10) were compared with numerical ones. Significant increase in elastic foundation from $k = 1\text{N/m}^2$ to localized $k_0 = 1000\text{N/m}$ at position $x_0 = 0.3L$ was considered. To measure this localized effect, parameter κ is introduced:

$$\kappa = \frac{k_0}{kL}. \quad (11)$$

Keeping the definition of parameters ξ , ψ and ζ from previous section, the case for $\xi=2$, $\psi=0.5$, $\zeta=2.1$ and $\kappa=100$ is presented in the graphs depicted in Figures 3, 4 and 5. For the sake of comparison, the analytical solution for the case without k_0 introduction (orange curve) is also included. Three different force positions are shown: at $0.2L$ (before the localized stiffness in-

crease); at $0.3L$ (coincident with the localized stiffness increase) and at $0.7L$ (after the localized stiffness increase).

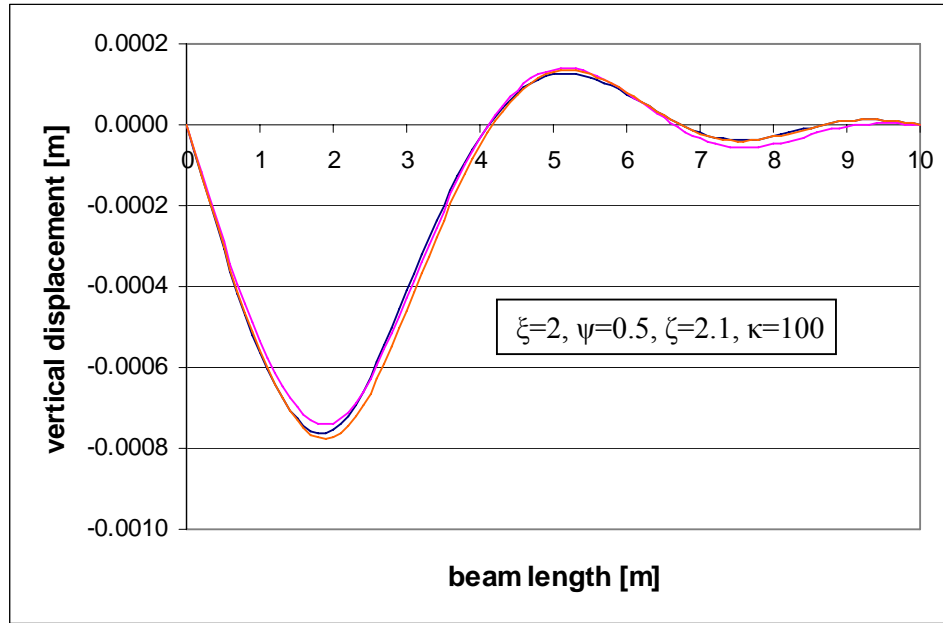


Figure 3: Vertical displacement for $\xi=2$, $\psi=0.5$, $\zeta=2.1$ and $\kappa=100$, plotted at the position $0.2L$ of the moving load (blue curve: numerical results; violet and orange curves: analytical solutions with and without k_0 , respectively)

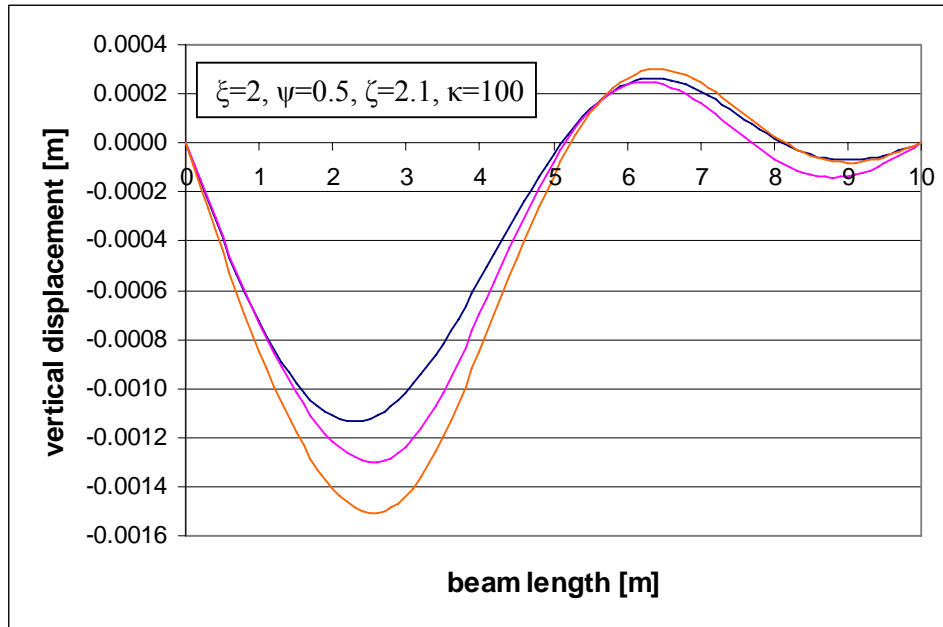


Figure 4: Vertical displacement for $\xi=2$, $\psi=0.5$, $\zeta=2.1$ and $\kappa=100$, plotted at the position $0.3L$ of the moving load (blue curve: numerical results; violet and orange curves: analytical solutions with and without k_0 , respectively)

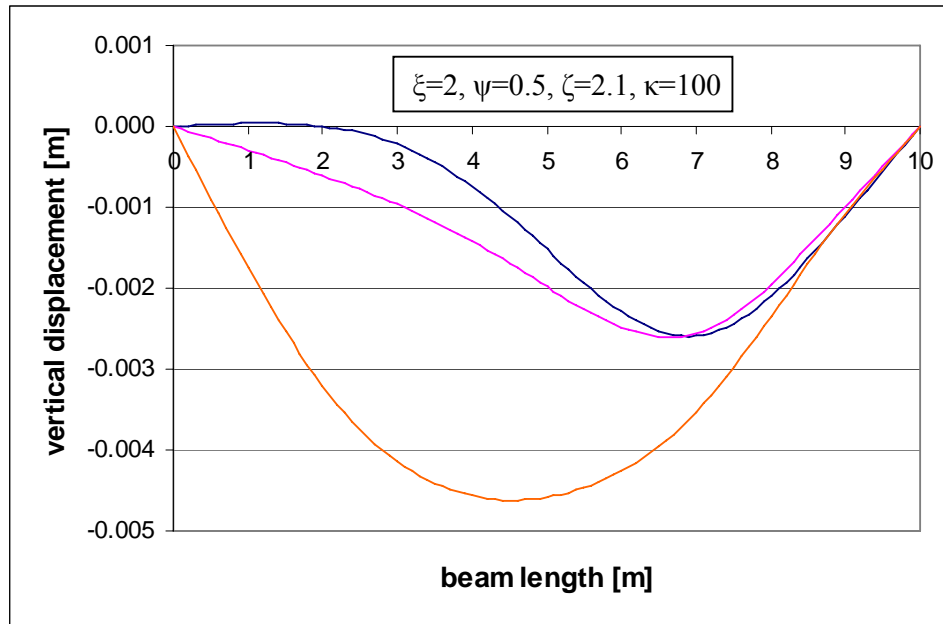


Figure 5: Vertical displacement for $\xi=2$, $\psi=0.5$, $\zeta=2.1$ and $\kappa=100$, plotted at the position $0.7L$ of the moving load (blue curve: numerical results; violet and orange curves: analytical solutions with and without k_0 , respectively)

It is seen that the high value of k_0 does not cause any irregularity in the vertical deflection profile; in fact, maximum deflection is lower than in the original case, due to the fact that k_0 works as an additional flexible beam support, reducing the “beam span”. However, this is not in accordance with what is experienced by real rail vehicles. One of the obvious reasons for this discrepancy is that in such cases interaction of spring-mass-damper system of the vehicle with the beam structure cannot be omitted. Nevertheless, it is observed that the approximate analytical solution agrees reasonably with the numerical one. This allows for direct sensitivity analysis and, in the other hand, making the beam sufficiently long, conclusions can be drawn for real situations, without influence of the support conditions.

In order to evaluate the efficiency of this solution, a case approaching real conditions is analysed. It is assumed that the beam corresponds to the full superior structure of a railway track, including rails, sleepers, ballast and sub-ballast. Dimensions of the beam cross-section equal to $4\text{m}\times 1\text{m}$, material properties $E=200\text{MPa}$, $\rho=1800\text{kg/m}^3$ and a load $P=100\text{kN}$ moving at constant velocity $c=45,3\text{m/s}=163,2\text{km/h}$ are adopted. Beam length is extended to 100m and Winkler’s constant is taken as 200kN/m^2 . Actually this value is quite a low estimate for a real soil. The reason for this is that the numerical solution becomes unreliable for very strong foundation, because then first natural frequencies, $\tilde{\omega}_{(j)}$, are very similar and this can disorder their sequence and consequently attribute different weight to waves forms in transient solution. Analytical solution is more stable numerically, nevertheless, significant number of sine series must be taken into account. Results were calculated by computer package Matlab [5] where 1000 series members were implemented. Analytical and numerical values fit well, as shown in Figure 6. However, the assumption subjacent to Equation (7) for localized stiffness increase is approximate and the coincidence of numerical and analytical values may be occasional. Hence, further research is needed in order to validate and improve this approach.

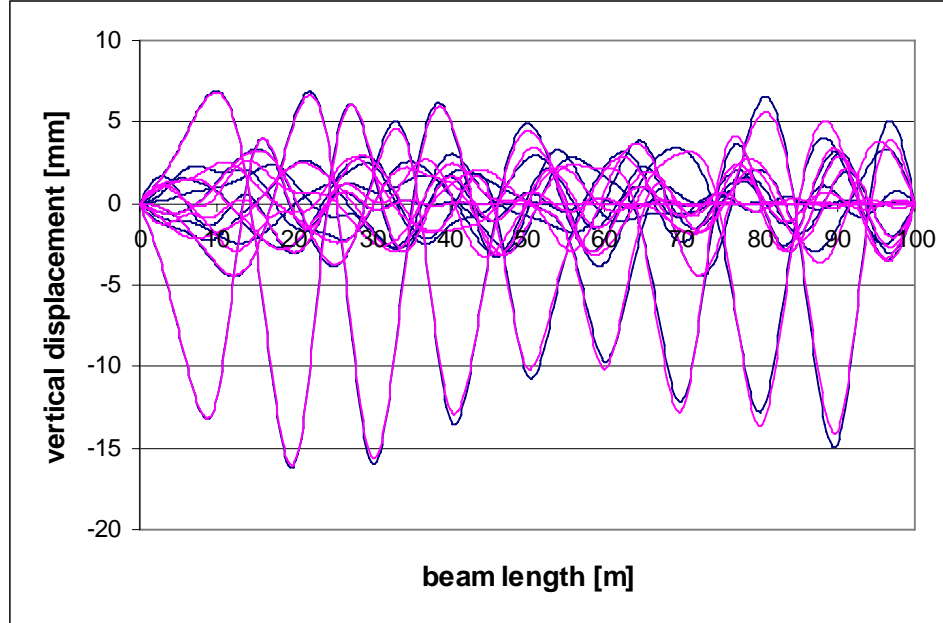


Figure 6: Vertical displacement for $c=45,3\text{m/s}$, no damping and $k=200\text{kN/m}^2$, plotted for 10 positions of the moving load $P=100\text{kN}$ (violet curve: analytical solution; blue curve: numerical results)

At the present, extension of results to consider harmonic loads, multiple loads, other types of damping and simplified vehicle models are under development.

4 SENSITIVITY ANALYSIS OF DYNAMIC RESPONSE OF MOVING LOAD ON SIMPLY SUPPORTED BEAM ON ELASTIC FOUNDATION WITH LOCALIZED ABRUPT CHANGE IN VERTICAL STIFFNESS

All derivatives of solution (8-10) can be done in an analytical way. For instance:

$$\frac{\partial}{\partial k_0} w(x, t) \cong \frac{2}{L} \sum_{j=1}^{\infty} \frac{\partial W(j, t)}{\partial k_0} \sin \frac{j\pi x}{L}, \quad (12)$$

where

$$\begin{aligned} \frac{\partial W(j, t)}{\partial k_0} = & \frac{\partial b}{\partial k_0} \frac{-4Pbc(a^2 + b^2 - c^2)}{\mu \left((a^2 + b^2 - c^2)^2 + 4a^2c^2 \right)^2} \cdot \\ & \left(\frac{(a^2 + b^2 - c^2)}{c} \sin(ct) - \frac{(b^2 - a^2 - c^2)}{b} e^{-at} \sin(bt) - 2a(\cos(ct) - e^{-at} \cos(bt)) \right) + \\ & \frac{\partial b}{\partial k_0} \frac{Pc}{\mu \left((a^2 + b^2 - c^2)^2 + 4a^2c^2 \right)} \cdot \\ & \left(\frac{2b}{c} \sin(ct) - \left(1 + \frac{a^2 + c^2}{b^2} \right) e^{-at} \sin(bt) - t \frac{(b^2 - a^2 - c^2)}{b} e^{-at} \cos(bt) - 2ate^{-at} \sin(bt) \right) \end{aligned} \quad (13)$$

and

$$\frac{\partial b}{\partial k_0} = \frac{\frac{1}{L\mu} \sin^2\left(\frac{j\pi x_0}{L}\right)}{\sqrt{\omega_{(j)}^2 - \omega_b^2 + \frac{k}{\mu} + \frac{2k_0}{L\mu} \sin^2\left(\frac{j\pi x_0}{L}\right)}}. \quad (14)$$

Implementation of equations presented above is straightforward.

5 MOVING LOAD ON CANTILEVER BEAM ON ELASTIC FOUNDATION

In order to study sudden elastic foundation stiffness change, implemented in whole region, cantilever dynamic response must be extended to account for elastic foundation. Then two cantilever solutions, corresponding to beams clamped on left and right hand side, with different value of Winkler constant can be connected together by continuity conditions.

General expression of transient vertical displacement of beam with various boundary conditions subjected to a moving load can be written in the following form, [3]:

$$w(x, t) = \sum_{j=1}^{\infty} \frac{\mu}{W_j} W(j, t) w_{(j)}(x), \quad (15)$$

where

$$w_{(j)}(x) = \sin \frac{\lambda_j x}{L} + A_j \cos \frac{\lambda_j x}{L} + B_j \sinh \frac{\lambda_j x}{L} + C_j \cosh \frac{\lambda_j x}{L} \quad (16)$$

and

$$W_j = \int_0^L \mu w_{(j)}^2(x) dx. \quad (17)$$

Constants from equation (16), λ_j , A_j , B_j , C_j , must be determined numerically in order to satisfy given boundary conditions. Irrespectively of the beam being clamped on the right or on the left hand side, λ_j correspond to roots of the following equation:

$$1 + \cos \lambda_j \cosh \lambda_j = 0 \quad (18)$$

and A_j is calculated from:

$$A_j = -\frac{\sin \lambda_j + \sinh \lambda_j}{\cos \lambda_j + \cosh \lambda_j}. \quad (19)$$

Then for right clamping $C_j = A_j$, $B_j = 1 \forall j$ and for left clamping $C_j = -A_j$, $B_j = -1 \forall j$.

Governing equation is again Equation (4) and the natural frequencies are given by:

$$\omega_{(j)} = \sqrt{\frac{\lambda_j^4 \pi^4}{L^4} \frac{EI}{\mu}}, \quad \tilde{\omega}_{(j)} = \sqrt{\frac{\lambda_j^4 \pi^4}{L^4} \frac{EI}{\mu} + \frac{k}{\mu}} \quad (20)$$

for beams without and with elastic foundation, respectively.

The free end of the cantilever must allow introduction of non-zero vertical force and moment, which will correspond to transversal force and bending moment of the full clamped beam after connection. Then $W(j, t)$ reads as:

$$W(j, t) = \frac{1}{b\mu} \int_0^t (Pw_{(j)}(c\tau) - EI z(0, L, \tau)) e^{-a(t-\tau)} \sin(b(t-\tau)) d\tau, \quad (21)$$

where

$$z(0, L, t) = \frac{V(0, t)}{EI} w_{(j)}(0) - \frac{M(0, t)}{EI} \left. \frac{dw_{(j)}(x)}{dx} \right|_{x=0}, \quad (22)$$

$$z(0, L, t) = -\frac{V(L, t)}{EI} w_{(j)}(L) + \frac{M(L, t)}{EI} \left. \frac{dw_{(j)}(x)}{dx} \right|_{x=L}$$

for right and left clamping, respectively, and

$$a = \omega_b, \quad b = \sqrt{\tilde{\omega}_{(j)}^2 - \omega_b^2}. \quad (23)$$

For the sake of comparison, similar parameters as in Section 2 are introduced according to:

$$\omega = \frac{\lambda_1 c}{L}, \quad \xi = \frac{\omega}{\tilde{\omega}_{(j)}}, \quad \psi = \frac{\omega_b}{\omega_{(j)}}, \quad \zeta = \frac{\omega}{\omega_{(j)}} \quad (24)$$

with Equation (20) implemented. Verification of derived formulas was again performed in ANSYS. Moving load and/or prescribed variation of free end internal forces was tested. In every case analyzed, the match between analytical and numerical solutions was excellent.

6 MOVING LOAD ON CLAMPED BEAM ON ELASTIC FOUNDATION WITH SUDDEN DROP IN VERTICAL STIFFNESS

In order to model the dynamic response of a clamped beam with sudden drop/increase in vertical stiffness, two cantilever solutions, corresponding to beams clamped on left and right hand side, with different value of Winkler constant are connected together by continuity conditions. The point of Winkler constant discontinuity corresponds to the point of beam continuity, therefore equilibrium of internal forces must be preserved and equality of vertical displacement and of its spatial derivative (rotation) must be maintained at that point. The internal forces (the unknowns) can be simply introduced by same values in both clamped beam solutions and solved from the equations imposing continuity of vertical displacement and rotation at the point of Winkler constant discontinuity.

Solution of this problem is not straightforward and must be done numerically, although parameters dependence is preserved. In more detail, the main difficulty lies in Equation (21), where the function z , given by (22) must be integrated over time when the actual time variation of internal forces is unknown. Nevertheless, linear variation can be assumed within each time step. Then, in the way described above, internal forces can be solved at each time step and, obviously, values V_i and M_i for $i=1, \dots, n$ can be used to determine piecewise linear distribution and to calculate V_{n+1} and M_{n+1} . Unfortunately none of the previous integrations can be used in next time step, due to the convolution form of the integral.

Correctness of this procedure was tested in ANSYS. First, clamped beam of $L=20\text{m}$ was solved numerically and internal forces from its middle section were extracted at each time step. Then, deflection curves were completely rebuilt assuming that two cantilever beams form the full structure. Thus, in the first half of the time needed for the load to cross the structure, the left cantilever solution accounted for the moving load and prescribed internal forces variation at the “free” end. The right cantilever was loaded during this period only by prescribed internal forces variation. Then, when the force passed the middle section, solutions

switched their role. Coincidence of results is again very good. In Figures 7 to 9, graphs of deflection curves calculated numerically and recovered analytically are shown. For the sake of simplicity, only the following cases are presented: $\xi=\zeta=1$ and $\psi=0$ for positions of the load at 0.1, 0.2, ..., 1m and results on the right hand side; $\xi=\zeta=1$ and $\psi=0$ for positions of the load at 1, 1.1, 1.2, ..., 2m and results on the left hand side; $\xi=0.268$, $\psi=2$ and $\zeta=1$ for positions of the load at 1, 1.1, 1.2, ..., 2m and results on the right hand side.

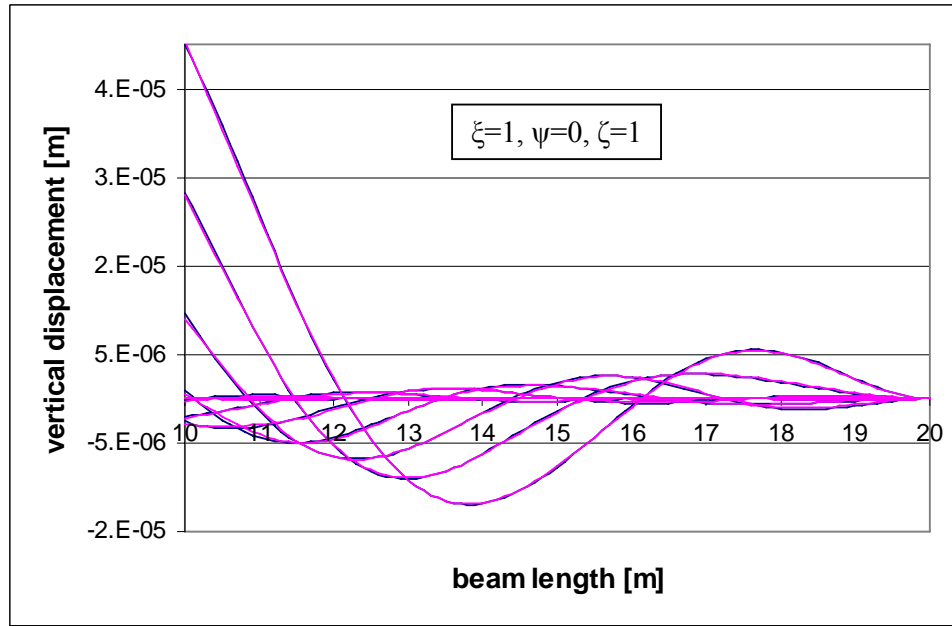


Figure 7: Vertical displacement for $\xi=1$, $\psi=0$ and $\zeta=1$, plotted for 10 positions of the moving load from 0 to 1m (violet curve: analytical solution; blue curve: numerical results)

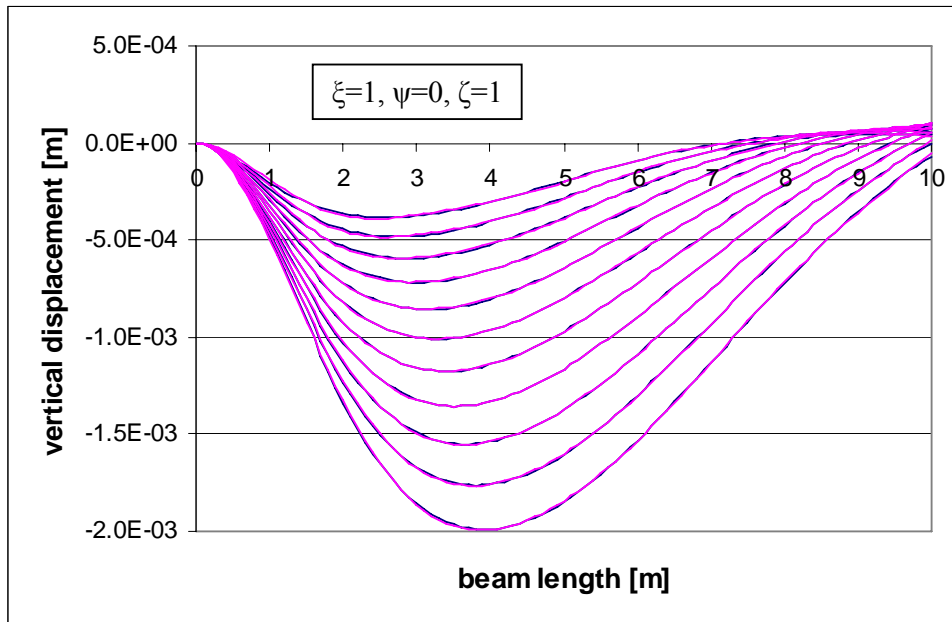


Figure 8: Vertical displacement for $\xi=1$, $\psi=0$ and $\zeta=1$, plotted for 11 positions of the moving load from 1 to 2m (violet curve: analytical solution; blue curve: numerical results)

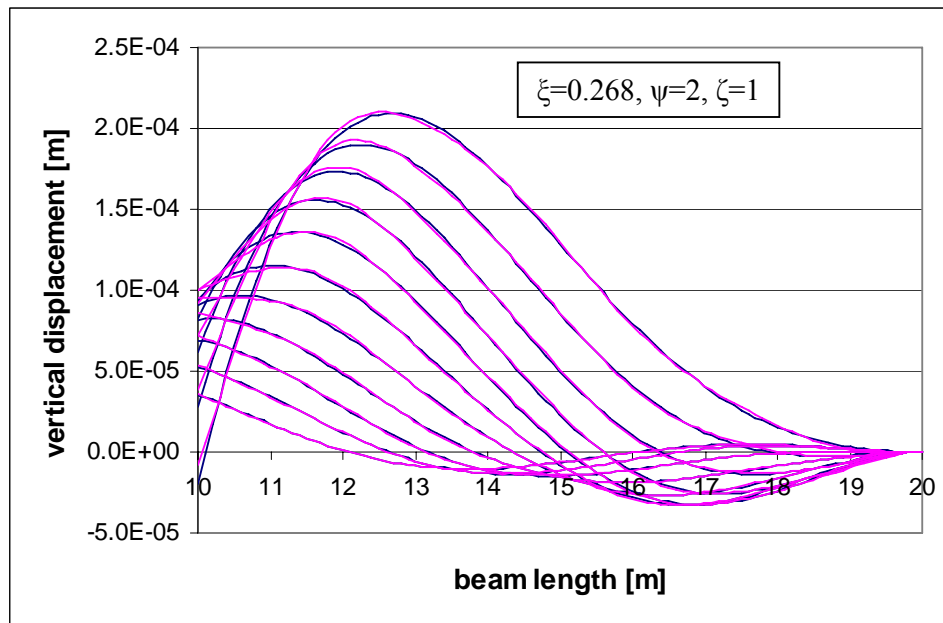


Figure 9: Vertical displacement for $\xi=0.268$, $\psi=2$ and $\zeta=1$, plotted for 11 positions of the moving load from 1 to 2m (violet curve: analytical solution; blue curve: numerical results)

7 CONCLUSIONS

In this paper, analytical transient solutions of dynamic response of one-dimensional systems with sudden change of foundation stiffness are derived. Abrupt localised increase/decrease is solved approximately, sudden drop/increase in foundation stiffness valid in whole region is solved exactly. However, assumptions about time variation of internal forces at the section of discontinuity must be adopted and the analytical solution will include numerical procedure. In both cases, results are expressed in terms of vertical displacement. Sensitivity analysis of the response amplitude is also performed. The analytical formulation for this problem, to the authors' knowledge, has not been published yet. Although related to one-dimensional cases, this study can give first insight into the problem of excessive ground vibrations caused by high speed trains crossing regions with abrupt changes in vertical stiffness of the track and/or subsoil.

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