

# Three-dimensional numerical modeling of tsunami-related internal gravity waves in the Hawaiian atmosphere

Giovanni Occhipinti<sup>1,2</sup>, Pierdavide Coïsson<sup>1</sup>, Jonathan J. Makela<sup>3</sup>, Sébastien Allgeyer<sup>4</sup>, Alam Kherani<sup>5</sup>, Hélène Hébert<sup>4</sup>, and Philippe Lognonné<sup>1</sup>

<sup>1</sup> Institut de Physique du Globe de Paris, Sorbonne Paris Cité, Univ Paris Diderot, UMR 7154 CNRS, F-94100 Saint Maur des Fossés, France

<sup>2</sup> Office National d'Etudes et Recherches Aérospatiales, Palaiseau, France

<sup>3</sup> Department of Electrical and Computer Engineering, University of Illinois at Urbana-Champaign, Urbana, Illinois, USA

<sup>4</sup> CEA-DAM-DIF, F-91297 Arpaçon, France

<sup>5</sup> Instituto Nacional de Pesquisas Espaciais (INPE), Sao Jose dos Campos, Sao Paula, BR-12227010, Brazil

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The tremendous tsunami following the 2011 Tohoku Earthquake produced internal gravity waves (IGWs) in the neutral atmosphere and large disturbances in the overlying ionospheric plasma while propagating through the Pacific ocean. To corroborate the tsunamigenic hypothesis of these perturbations, we use a 3D numerical modeling of the ocean-atmosphere coupling, to reproduce the tsunami signature observed in the airglow by the imager located in Hawaii and clearly showing the shape of the modeled IGW. The agreement between data and synthetics not only supports the interpretation of the tsunami-related-IGW behavior, but strongly shows that atmospheric and ionospheric remote sensing can provide new tools for oceanic monitoring and tsunami detection.

**Key words:** Internal Gravity Waves, Tsunami, Airglow, Numerical Modeling.

## 1. Introduction

After the 2004 Great Sumatra Earthquake and the consequent Indian Ocean Tsunami, the scientific community put its attention to alternative methods in oceanic monitoring to improve the response of tsunami detection and estimation.

Improvement of classic techniques, such as seismic source estimation [e.g., Ammon et al., 2006] and densification of the number of buoys in the oceans [González et al., 2005], was supported by a new effort in remote sensing: nominally the space-altimetry observation of the tsunami in the open sea [Okal et al., 1999, Smith et al., 2005] and the tsunami detection by ionospheric monitoring [Artru et al., 2005, Occhipinti et al., 2006]. The recent devastating tsunamigenic earthquake occurred on 11 March, 2011, off the Pacific coast of Tohoku (Mw=9.3, 05:47:32 UT), indicating once again the importance of moving forward in this direction.

Several theoretical studies in the 1970s, including Hines's pioneering works on internal gravity waves (IGWs), suggested that atmospheric IGWs generated by a tsunami (Fig. 1) may well produce identifiable ionospheric signatures in the plasma [Hines, 1972, Peltier & Hines, 1976]. In essence, during their upward propagation the IGWs are strongly amplified by the effect of the decrease of atmospheric density. The interaction of IGWs with the plasma at ionospheric heights produces strong variations in the plasma velocity and plasma density observable by ionospheric sounding.

The encouraging work of Artru et al. [2005] on the detection of the Peruvian tsunamigenic quake on 23 June, 2001 (Mw=8.4 at 20:33 UT) in the total electron content

(TEC) measured by the Japanese dense GPS network, GEONET, opened the modern debate about the feasibility of tsunami detection by ionospheric sounding. The giant tsunami following the Sumatra-Andaman event (Mw=9.3, 00:58:50 UT, 26 December, 2004, [Lay et al., 2005]), an order of magnitude larger than the Peruvian tsunami, provided worldwide remote sensing observations in the ionosphere, giving the opportunity to deeply explore ionospheric tsunami detection by ionospheric sounding [Liu et al., 2006a, Liu et al., 2006b, Lognonné et al., 2006, DasGupta et al., 2006, Occhipinti et al., 2006, Occhipinti et al., 2008b].

The theoretical and numerical work of Occhipinti et al. (2006) supported by TEC observation from the on-board altimeters Topex/Poseidon and Jason-1 quantitatively confirms the link between the ionospheric perturbation and the sea level displacement induced by the tsunami propagation.

Recent observational works highlight systematic ionospheric tsunami detection following several events with magnitude around 8 (Kuril 2006, Samoa 2009 and Chile 2010) utilizing GPS-derived TEC performed by moderately-sized networks, such as those on Hawaii [Rolland et al., 2010, Galvan et al., 2011].

Subsequent modeling studies estimate the role of the geomagnetic field in the tsunami signature at the E-region and F-region [Occhipinti et al., 2008a] as well as the detection possibility by over-the-horizon radar [Coïsson et al., 2011]

More recently, Hickey et al. [2010] have suggested that appreciable airglow modulations in the 630.0-nm intensity should be expected. However, the first observations of this effect were not made until the 2011 Tohoku earthquake [Makela et al., 2011].

In this work, we model the IGW passing over Hawaii produced by the propagation of the Tohoku tsunami through

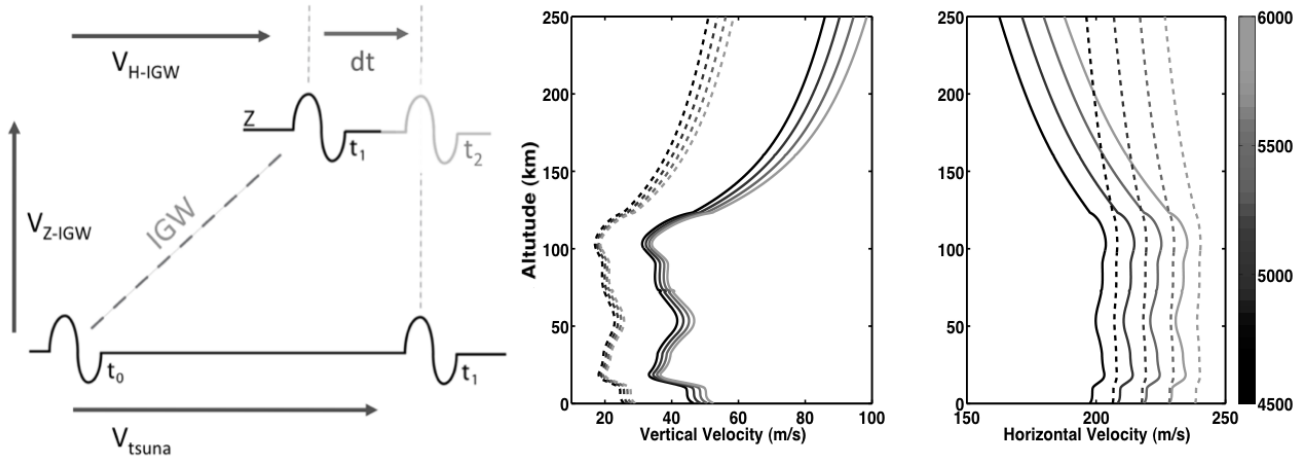


Fig. 1. Left: schematic representation of tsunami-related-IGW propagation with its vertical  $V_{Z-IGW} = \frac{k_z k_h^2 N^2}{\omega D^2}$  and horizontal  $V_{H-IGW} = \frac{k_h N^2 (D - k_h^2)}{\omega D^2}$  group velocities, where  $k_z$  and  $k_h$  are the vertical and horizontal  $k$ -vector,  $g$  the gravity,  $N$  the Brunt-Väisälä frequency and  $D = k_z^2 + k_h^2 + \frac{N^4}{4g^2}$ . Following [Occhipinti *et al.*, 2006, Occhipinti *et al.*, 2008a], we assume that the spatial and temporal dependency of wavefield is expressed in the form  $\exp(k_x x - k_z z - \omega t)$ , where the dispersion equation take the form  $\omega^2 D = k_h^2 N^2$ . As  $V_{H-IGW}$  is lower than the tsunami speed  $V_{tsuna} = \sqrt{gH}$ , the IGW accumulates a delay ( $dt$ ) with respect to the tsunami at the sea level during the upward propagation. The vertical and horizontal group velocity profiles for IGW of period  $T_1 = 14$  minutes (solid lines) and  $T_2 = 26$  minutes (dashed lines) are shown in the middle and right panel respectively; gray color indicates the speed dependence of ocean depth  $H$  (m). At the altitude of 250 km the delay  $dt$  is on order of 6 min and 3 min for  $T_1$  and  $T_2$  respectively. An IGW needs around 80 min and 145 min to reach the altitude of 250 km for  $T_1$  and  $T_2$  respectively. Once the IGW reaches ionospheric altitude, the delay  $dt$  between the IGW and tsunami wavefronts is constant during the oceanic propagation for a constant oceanic depth. The effect of the bathymetry variation can reduce, vanish or invert the delay  $dt$ .

the Pacific Ocean in order to reproduce the airglow observation measured by the ground-based camera located on the Haleakala Volcano on Maui, Hawaii [Makela *et al.*, 2011].

### 1.1 Modelling

Following Occhipinti *et al.* [2006, 2008a] we compute, by three-dimensional numerical modeling, the atmospheric IGWs induced by the propagation of the 11 March 2011 tsunami in the overlying atmosphere. In essence, the IGW modeling is divided into two steps: first, we compute the tsunami propagation using a realistic bathymetry taking into account the Earth sphericity, then we compute the excited IGW in the neutral atmosphere in a cartesian coordinate system.

The ocean surface displacement is calculated using a finite difference scheme that resolves the hydrodynamical equations on a 30" bathymetric grid following Hébert *et al.* (2007). The input earthquake source consists of the USGS finite fault model (<http://earthquake.usgs.gov/earthquakes>).

Tsunami propagation is calculated through the whole Pacific Ocean every minute during the 12 hours following the earthquake.

The vertical velocity produced by the sea level displacement is used as input to the pseudo spectral propagator based on the linearised momentum and continuity equations, for irrotational, inviscid and incompressible flow (see Occhipinti *et al.* [2006, 2008a] for more details).

We propagate the IGW from the ocean surface to 250 km, which is the altitude at which the 630.0-nm airglow emission is assumed to be concentrated [*e.g.*, Kelley *et al.*, 2000; Makela & Kelley, 2003]

Attenuation effects, such as viscosity and thermal conduction, are neglected in our model as we are principally interested to the shape of the tsunami signature in the airglow data. Anyway, we highlight that their effect at 250 km of altitude reduces by a factor 5 the amplitude of the IGWs [Hickey *et al.*, 2009].

We focus the analysis in the Hawaii region, in a geographic grid between  $0^\circ$  and  $40^\circ$  latitude N and  $180^\circ$  and  $220^\circ$  longitude E. According to the modeling, the tsunami was crossing this region between 11 UT and 15 UT with a maximum wave amplitude of about 10-15 cm in the open ocean. Figure 2 shows the IGW model results at 13:20 UT with a well defined structure of gravity wave propagation. The IGW is traveling through this area southeastwards with its wave front being in the SW-NE direction. The retarding effect of the Hawaii archipelago can be appreciated at the center of the horizontal cross-section.

As the bathymetry variation also affects the vertical and horizontal IGW group velocities (Fig. 1 and 4), the low-bathymetry close to Hawaii affects the propagation of the tsunami-related IGW. Additionally, vertical and horizontal IGW group velocities are strongly dependent on the dominant period of the tsunami at sea level. Two main periods have been observed after analyzing the Deep-ocean Assessment and Reporting of Tsunamis (DART) data: 14 min and 26 min [Makela *et al.*, 2011].

As a consequence of the combined effect of the period and the bathymetry on the IGW group velocities, the IGWs generated before the interaction with the Hawaiian archipelagos and the Hawaiian-Emperor seamount chain, go beyond the tsunami wave front during their propagation in the atmosphere.

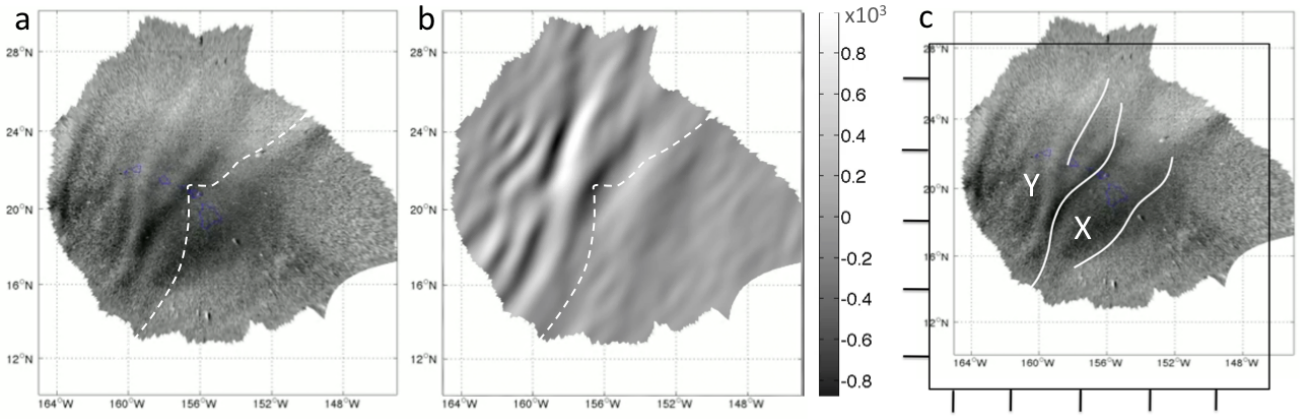


Fig. 3. a) IGW imaged by the 630.0-nm ground-based airglow camera located on the Haleakala Volcano on Maui, Hawaii, at 13:20 UT. b) normalized vertical velocity of the modeled IGW ( $(kg/m^3)^{\frac{1}{2}}m/s$ ). The Y structure as well as the longer wavelength anticipating the Y (X) are present in both, airglow data and AGW synthetics. Those structures are observed between 12:12 to 13:32. The white dotted line in a and b shows the tsunami wavefront line at 13:20 UT. c) graphical estimation of the shift induced by the wind between model and data: the grid and the white lines are estimated on b, then shifted on a to fit with the position of Y and X. The estimated shift of  $2^\circ$  is coherent with previous observations [Occhipinti *et al.*, 2006].

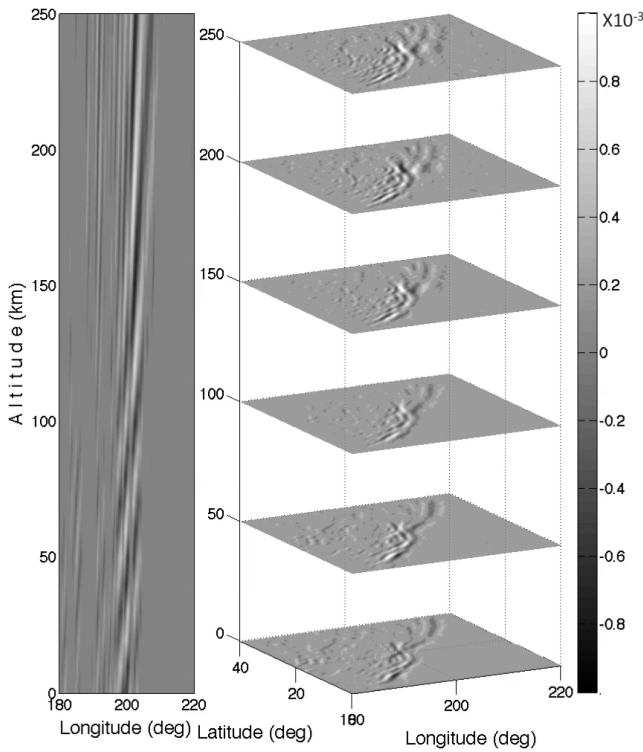


Fig. 2. IGW modeled normalized vertical velocity component of atmospheric wave motion in Hawaii region at 13:20 UT. Left: vertical cross section at  $20^\circ$  latitude N clearly showing the inversion of the velocity component of IGW motion in a relative short distance; Right: horizontal cross sections from the bottom of the atmosphere up to 250 km. The entire movie is available at [www.ipgp.fr/~ninto/EPS/SMmovie.mov](http://www.ipgp.fr/~ninto/EPS/SMmovie.mov).

This effect is visible in the vertical cross-section of the IGW close to Hawaii (Fig. 2) as well as on the IGW propagating at 250 km of altitude, clearly highlighting the differentiation link to the period (Fig. 3a-b): the longer wavelength IGW (indicated by the X in figure 3) is over the tsunami wavefront where the shorter wavelength IGW with the "Y"

shape stay behind the tsunami wavefront. We note that this effect is also visible in the airglow data (Fig. 3a).

## 1.2 Comparison with Hawaiian airglow observation

The tsunami reached Hawaii during the night hours with a predominantly clear sky, optimal conditions to detect the expected IGWs by optical sensors. The highly sensitive, wide-angle camera system located atop the Haleakala Volcano on Maui, Hawaii, was able to record the airglow emission at 630.0 nm with a temporal resolution of approximately 5 minutes during the time when the tsunami was crossing that area. The spatial coverage of the observation is approximately  $10^6 km^2$  in the ionosphere at the assumed emission altitude of 250 km. The 630.0-nm airglow emission is produced by the dissociative recombination of  $O_2^+$  [Link & Cogger, 1988] and is sensitive to both changes in the electron and neutral densities.

Comparison between the modeled IGW and the airglow observation in Figure 3 shows several interesting likenesses explaining the nature of the airglow observation:

The Y shape presents in our IGW modeling is also recognizable in the data, this is clear evidence of the tsunami nature of the airglow observation. Additionally, we highlight a longer wavelength signal (X) in both airglow data and IGW modeling, arriving before the Y curved feature and also before the tsunami arrival time (white dotted line). This effect has been detailed in the previous section and is the consequence of the low bathymetry around Hawaii.

The Y curved feature also allows us to quantify the shift of  $2^\circ$  between airglow data and the synthetic IGW (Fig. 3c). A similar shift was observed by Occhipinti *et al.* [2006] in comparison between altimeters-TEC data and synthetics, and is most likely due to the zonal and meridional wind neglected in our modeling.

## 1.3 Conclusions

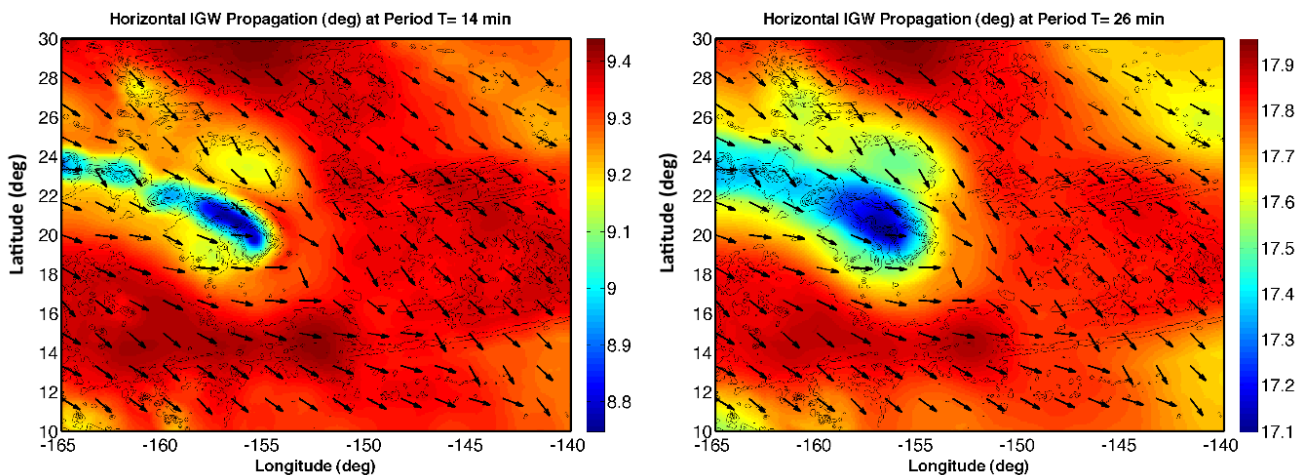


Fig. 4. Horizontal propagation distance (colorscale in degrees) and direction (arrows) covered by the IGW from the generation at the sea level to the altitude of 250 km at the main period of 14 min (left) and 26 min (right). In essence, the IGW generated at given position at the sea level will be observed at 250 km of altitude after covering a distance shown by the color and in the direction indicated by the arrow. The low bathymetry close to Hawaii reduces the horizontal velocity of the IGW and tsunami (Fig. 1), consequently IGW generated before Hawaii go beyond the IGW generated in the shallow bathymetry region as well as beyond the tsunami. This effect affects more the period at 26 min than 14 min.

The tremendous tsunami following the 11 March 2011 Tohoku mega-quake generated a huge amount of data all around the Pacific Ocean. This event particularly highlight, for the first time, the possibility to detect tsunami relate IGW with the ground-based airglow camera located in Hawaii [Makela *et al.*, 2011].

The observed airglow signal presents several questions and uncertainties that are explored here with the support of the three-dimensional modeling of the tsunami related IGW propagating in the atmosphere overlying the Hawaiian archipelagos.

The comparison with the modeling, first, highlights the tsunamic nature of the airglow data by the characteristic "Y" shape appearing in both synthetics and data; second, it explains the anticipated signal arriving before the tsunami front-wave, and finally, it allows quantification of the shift of  $2^\circ$  mainly due to the wind effect.

This preliminary result motivate the development of a full modeling of the airglow observation taking into account, not only the wind, but above all the different ionized species taking part in the formation of the airglow signal, as well as the observational geometry of the ground-based camera.

Waiting for a full airglow modeling, the effort of this preliminary work highlights that remote sensing of tsunamis via the atmospheric/ionospheric monitoring by ground-based or on-board camera is a mature technique for oceanic monitoring and can take a role in the future of tsunami detection.

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