Simplified Model for the Non-Linear Behaviour Representation of Reinforced Concrete Columns Under Biaxial Bending

H. Rodrigues
Civil Engineering Department, University of Aveiro, Portugal
Faculty of Natural Sciences, Engineering and Technology - Oporto Lusophone University

A. Andrade-Campos
Department of Mechanical Engineering, University of Aveiro, Portugal

X. Romão, A. Arêde
Department of Civil Engineering, Faculty of Engineering, University of Oporto, Portugal

H. Varum, A. Costa
Department of Civil Engineering, University of Aveiro, Portugal

SUMMARY:
In the present paper a simplified model is proposed for the force-deformation behaviour of reinforced concrete members under biaxial loading combined with axial force. The starting point for the model development was an existing fixed-length plastic hinge element model that accounts for the non-linear hysteretic behaviour at the element end-sections, characterized by trilinear moment-curvature laws. To take into account the section biaxial behaviour, the existing model was adopted for both orthogonal lateral directions and an interaction function was introduced to couple the hysteretic response of both directions.
To calibrate the interaction function it were used numerical results, obtained from fibre models, and experimental results. For the parameters identification, non-linear optimization approaches were adopted, namely: the gradient based methods followed by the genetic, evolutionary and nature-inspired algorithms.
Finally, the simplified non-linear model proposed is validated through the analytical simulation of biaxial test results carried out in full-scale reinforced concrete columns.

Keywords: RC columns, non-linear biaxial behaviour, simplified non-linear model, optimization techniques

1. INTRODUCTION

The non-linear models should reproduce the influence of the structural system's geometry and also of the distribution of mass and stiffness, in particular the effect of irregularities in terms of stiffness and/or mass, in plan and in elevation. A vast number of models have been developed to represent the material's non-linear behaviour and are typically divided into different categories, namely global models, microscopic models and discrete finite element (member) models (H Rodrigues *et al.*, 2010; Scott *et al.*, 2008; Taucer *et al.*, 1991).
The importance of study of the tridimensional response of RC buildings is recognized, associated with tridimensional earthquake actions or to building irregularities, which induces a biaxial bending demand combined with axial load in the columns. Different modelling strategies have been proposed for the simulation of the biaxial cyclic behaviour of RC elements with axial force. However, it is recognized that the available biaxial models are not mature enough to be used in practice, nor to be incorporated into codes/standards as has occurred with uniaxial simplified models. Detailed reviews of the models available can be found in the CEB Report Nº220, "RC Frames under Earthquake Loading – State of the Art Report" (CEB, 1996), and in Fardis (Fardis, 1991). Besides the theory of fibre models (Petrangeli *et al.*, 1999; Spacone *et al.*, 1992; Taucer *et al.*, 1991), existing analytical models follow the concepts of classical plasticity (Pecknold, 1974), Mroz's theory of multi-surface plasticity (ElMandooh Galal & Ghobarah, 2003; Powell & Chen, 1986; Takizawa *et al.*, 1976), Bouc-Wen (Wen, 1976), hysteresis modelling (Casciati, 1989; Kunnath & Reinhorn, 1990; Romão *et al.*, 2004; C. H. Wang & Wen, 2000), bounding surface plasticity (Bousias *et al.*, 2002; M.G. Sfakianakis & M.N. Fardis, 1991; M.G. Sfakianakis & M.N. Fardis, 1991) or lumped damage models (M.E. Marante &

Flórez-López, 2002; Maria Eugenia Marante & Flórez-López, 2003; Mazza & Mazza, 2008).

The present paper proposes an upgraded simplified model for the representation of the biaxial bending response of columns with axial force, which is based in an existing uniaxial hysteretic Costa-Costa model (Costa & Costa, 1987).

The general formulation of the proposed model is established in the analogy and comparison with the biaxial formulation of the Bouc-Wen smooth hysteretic model [9]. Although retaining most of the physical meaning embodied in the Bouc-Wen model, the structural modelling strategy adopted retains the simplicity and versatility of the original piecewise linear (PWL) numerical tool.

2. COSTA AND COSTA UNIAXIAL HYSTERETIC MODEL

For the development and validation of the simplified biaxial model, the Costa-Costa uniaxial hysteretic model (CEB, 1996; Costa & Costa, 1987) was adopted, which is briefly described in the next paragraph.

This model represents a generalisation of the original Takeda model (CEB, 1996; Takeda *et al.*, 1970) with a trilinear skeleton curve for monotonic loading, defined by the cracking point $(d_c;F_c)$ and the yield point $(d_y;F_y)$, and includes pinching, stiffness degradation and strength degradation effects.

Unloading-reloading loops prior to yielding in either direction are bilinear, with slopes equal to those of the pre-cracking and post-cracking branches in the virgin loading branch. After yielding, pinching is modelled by a bilinear reloading curve. The first branch of the reloading stage has an inferior slope (see Figure 1 - branch 8), while the second branch heads to the most extreme point of any previous post-yield excursion in the direction of the reloading. For this the stiffness K_s is multiplied by the factor $(d_y/d_m)^\beta$, where d_y represents the yield displacement, d_m the maximum response displacement and β is a positive constant. Thus:

$$K_s = \frac{F_m}{d_m - d_0}\left(d_y/d_m\right)^\beta \qquad (2.1)$$

where F_m is the force at the previous maximum response point and d_0 is the deformation at the load reversal point.
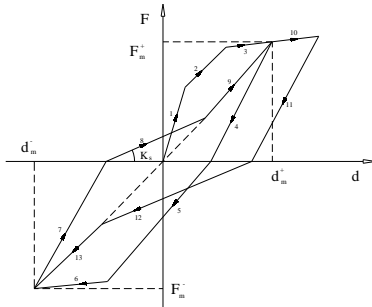

Figure 1. Pinching effect in the Costa-Costa (Costa & Costa, 1987) model
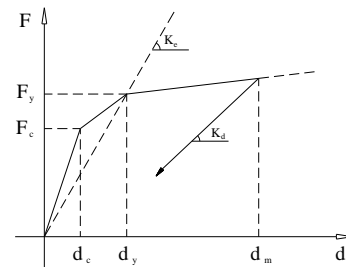
Figure 2. Unloading stiffness in the Costa-Costa (Costa & Costa, 1987) model

The unloading stiffness K_d, after yielding, is reduced relative to the elastic stiffness K_e by the factor $(d_y/d_m)^\alpha$, where α is a positive constant (see Figure 2).

Post-yield strength and stiffness degradation with cycling loading is modelled by directing the reloading branch, after modification for pinching, towards a point at a displacement equal to d_m and at a force $F'_m=(1-\lambda)\cdot F_m$, where λ is the Wang and Shah damage index (M. L. Wang & Shah, 1987). After reaching this end point of the reloading branch, further loading takes place parallel to the post-yielding stiffness of the virgin loading curve.

3. THE BIAXIAL BOUC-WEN MODEL

The original formulation of the Bouc-Wen model, cast within the endochronic theory framework was presented by Bouc (Bouc, 1971), for uniaxial behaviour representation (in terms of Force-Displacement, *F-u*), and it was later generalized by Wen (Wen, 1976). The generalized model expresses the restoring force as a combination of an elastic force and a plastic force:

$$\text{F} = \alpha \cdot K \cdot u + (1-\alpha) \cdot F_y \cdot Z \tag{3.1}$$

where K is the initial stiffness, α the post-yielding stiffness ratio, F_y the yielding force, and Z is a hysteretic parameter.

Later, this uniaxial formulation was extended by Park, Wen and Ang (Park *et al.*, 1986; Wen, 1976) to define a biaxial force-deformation model with coupled differential equations. This model was then used and modified by Kunnath and Reinhorn (Kunnath & Reinhorn, 1990) to model the behaviour of RC columns under biaxial loads. Later on, the model was generalised by Casciati (Casciati, 1989) and also by Wang and Wen (C. H. Wang & Wen, 2000), which resulted in two different formulations of the initial biaxial model. Since the Wang and Wen (C. H. Wang & Wen, 2000) formulation is simpler, it was selected for to implementing the biaxial model proposed in this work. Nevertheless, the same mathematical reasoning can be applied to the Casciati (Casciati, 1989) form.

The biaxial construction of the Bouc-Wen model in the Wang and Wen (C. H. Wang & Wen, 2000) form follows the same general idea as for the uniaxial case. The restoring forces for both directions are defined by:

$$\begin{cases} F_x = \alpha_x \cdot K_x \cdot u_x + (1-\alpha_x) \cdot F_x^y \cdot Z_x \\ F_y = \alpha_y \cdot K_y \cdot u_y + (1-\alpha_y) \cdot F_y^y \cdot Z_y \end{cases} \tag{3.2}$$

in which the involved parameters have the same meaning as for the uniaxial case, but are now referred to the two orthogonal directions *X* and *Y*, by the subscripts x and y respectively. The hysteretic parameters Z_x and Z_y are then defined by the following coupled differential equations, where all the parameters involved have also the same meaning as for the uniaxial case and *sign()* refers to the mathematical *signum* function.

$$\begin{cases} \dot{Z}_x = \dfrac{A \cdot \dot{u}_x - |\dot{u}_x| \cdot Z_x \cdot |Z_x|^{n-1} \cdot [\beta + \gamma \cdot sign(\dot{u}_x \cdot Z_x)] - |\dot{u}_y| \cdot Z_x \cdot |Z_y|^{n-1} \cdot [\beta + \gamma \cdot sign(\dot{u}_y \cdot Z_y)]}{u_x^y} \\[4mm] \dot{Z}_y = \dfrac{A \cdot \dot{u}_y - |\dot{u}_y| \cdot Z_y \cdot |Z_y|^{n-1} \cdot [\beta + \gamma \cdot sign(\dot{u}_y \cdot Z_y)] - |\dot{u}_x| \cdot Z_y \cdot |Z_x|^{n-1} \cdot [\beta + \gamma \cdot sign(\dot{u}_x \cdot Z_x)]}{u_y^y} \end{cases} \tag{3.3}$$

As for the case of the uniaxial model, Equation 3.2 can also be reformulated into an incremental form:

$$\begin{cases} \dot{F}_x = \alpha_x \cdot K_x \cdot \dot{u}_x + (1-\alpha_x) \cdot F_x^y \cdot \dot{Z}_x \\ \dot{F}_y = \alpha_y \cdot K_y \cdot \dot{u}_y + (1-\alpha_y) \cdot F_y^y \cdot \dot{Z}_y \end{cases} \tag{3.4}$$

considering that the global restoring forces F_{x_i} and F_{y_i} result from:

$$\begin{cases} F_{x_i} = F_{x_{i-1}} + \dot{F}_{x_i} \\ F_{y_i} = F_{y_{i-1}} + \dot{F}_{y_i} \end{cases} \tag{3.5}$$

Considering now the definition of the incremental orthogonal forces $\dot{F}_{xi}$ and $\dot{F}_{yi}$ given by, the first part of this system can be written, by a simple mathematical transformation, as:

$$\begin{cases} \dot{F}_{x_i} = \alpha_x \cdot K_x \cdot \dot{u}_{x_i} + (1-\alpha_x) \cdot K_x \cdot \left(A \cdot \dot{u}_{x_i} - |\dot{u}_{x_i}| \cdot Z_{x_i} \cdot |Z_{x_i}|^{n-1} \cdot [\beta + \gamma \cdot sign(\dot{u}_{x_i} \cdot Z_{x_i})] \right) \\[4mm] \dot{F}_{y_i} = \alpha_y \cdot K_y \cdot \dot{u}_{y_i} + (1-\alpha_y) \cdot K_y \cdot \left(A \cdot \dot{u}_{y_i} - |\dot{u}_{y_i}| \cdot Z_{y_i} \cdot |Z_{y_i}|^{n-1} \cdot [\beta + \gamma \cdot sign(\dot{u}_{y_i} \cdot Z_{y_i})] \right) \end{cases} \tag{3.6}$$

which matches to the uniaxial incremental restoring forces calculated for each direction without biaxial interaction. In order to simplify the notation, these incremental forces will be denoted as $\dot{F}_{uni-xi}$ and $\dot{F}_{uni-y_i}$, respectively. The remaining part of the system corresponds to the correction factors C_{fx_i} and C_{fy_i} accounting for the interaction between the two loading directions:

$$\begin{cases} C_{fx_i} = -(1-\alpha_x)\cdot K_x \cdot |\dot{u}_{y_i}| \cdot Z_{x_i} \cdot |Z_{y_i}|^{n-1} \cdot [\beta + \gamma \cdot sign(\dot{u}_{y_i} \cdot Z_{y_i})] \\ C_{fy_i} = -(1-\alpha_y)\cdot K_y \cdot |\dot{u}_{x_i}| \cdot Z_{y_i} \cdot |Z_{x_i}|^{n-1} \cdot [\beta + \gamma \cdot sign(\dot{u}_{x_i} \cdot Z_{x_i})] \end{cases} \tag{3.7}$$

Therefore, based on Equations 3.6 and 3.7, a condensed form of Equation 4 can be written as:

$$\begin{cases} \dot{F}_{x_i} = \dot{F}_{uni-x_i} + C_{fxi} \\ \dot{F}_{y_i} = \dot{F}_{uni-y_i} + C_{fyi} \end{cases} \tag{3.8}$$

Considering that the incremental forces $\dot{F}_{uni-xi}$ and $\dot{F}_{uni-y_i}$ can be obtained with any uniaxial hysteretic model (particularly with well-established PWL models), the presented framework introduces a simple and flexible form to represent biaxial bending in columns. This formulation requires the same information needed for the corresponding uniaxial PWL model, only introducing an additional correction term which couples the two loading directions.

Since the uniaxial hysteretic model that may be considered to obtain the incremental forces $\dot{F}_{uni-xi}$ and $\dot{F}_{uni-y_i}$ could be very different from the original uniaxial Bouc-Wen formulation, the type and level of biaxial interaction between the two orthogonal directions can be also different. Romão et al. (Romão et al., 2004) proposed an additional parameter δ, which was included in Equation 3.8 in order to scale the level of interaction between the two loading directions. The final formulation of the proposed method is then defined as written in Equation 3.9, which states that for each incremental displacements vector $(\dot{u}_{xi}; \dot{u}_{y_i})$, the incremental forces $\dot{F}_{uni-xi}$ and $\dot{F}_{uni-y_i}$ can be separately calculated, as:

$$\begin{cases} \dot{F}_{x_i} = \dot{F}_{uni-x_i} + \delta \cdot C_{fx_i} \\ \dot{F}_{y_i} = \dot{F}_{uni-y_i} + \delta \cdot C_{fy_i} \end{cases} \tag{3.9}$$

The values of the scaling interaction factor δ will be defined as for the best-fitting of the numerical results to the experimental results obtained with biaxial tests, additional information about the presented formulation can be found in the literature (Romão et al., 2004).

4. PARAMETER IDENTIFICATION FOR THE SCALING INTERACTION FACTOR

4.1 Optimization method

Numerical non-linear simplified models may play an important role in the design of new structures and in the assessment of existing ones. More complex techniques and models have been developed to simulate, with increasing accuracy, the behaviour of different materials and structures. However, many of these simulation models require the determination and calibration of a large number of parameters adjusted to the specific material and structural problem.

The identification of parameters for the mathematical models adopted to describe the behaviour of physical systems is a common problem in engineering. The complexity of the models, as well as the number of parameters associated, normally increases with the complexity of the physical system. The determination of these parameters should be based on the comparison of the mathematical model results and experimental results. However, when the required number of experimental tests and

parameters increases, it may become impractical to identify the accurate parameters (Andrade-Campos *et al.*, 2007; Bruhns & Anding, 1999). In these cases, it may be used an inverse formulation for the identification of parameters. This approach often leads to the resolution of a non-linear optimization problem.

In this work, aiming at calibrating the parameters for the hysteretic biaxial model based on uniaxial models with an interaction function, a gradient based-method (Levenberg-Marquardt (LM) method) was adopted, sequentially associated with an evolutionary algorithm method (real search space EA), grouped. With these global cascade algorithms, it is intended to aggregate the advantages of both algorithms and to minimise their disadvantages (Andrade-Campos *et al.*, 2009).

For the application of these models, the SDL/SiDoLo optimisation lab computer program was used (Andrade-Campos, 2011). The program was designed for specific engineering inverse problems, such as parameter identification and initial shape optimization problems. It inherits the wealth of experience gained in such problems by the previous SIDOLO code, and adds the latest developments in direct search optimization algorithms (Andrade-Campos, 2011; Andrade-Campos *et al.*, 2007).

4.2 Prediction and optimization of the scaling factor equation

The calibration of the interaction function, using the optimization strategies presented in the previous section, was performed in two phases. In the first phase, it was intended to select the analytical expression-type most suitable for the interaction scaling factor (δ) function. After the expression type selection, the second phase, consists in the calibration of the parameters of the interaction function as well as of the interaction scaling factor. The parameters' calibration corresponds to the best-fit of the results for a group of numerical analyses.

In order to define the interaction scaling factor (δ) which characterizes the level of interaction, as a function dependent of the section properties and loading direction, a set of numerical analyses was performed. To this aim, twenty seven rectangular RC columns were defined, varying in cross-section dimensions (30x30, 30x50 and 30x60), reinforcing steel ratio (1%, 1.5% and 2%) and axial load ratio (0.1, 0.2 and 0.3). For the column, a cantilever 1.5m high was considered. The response of each column was obtained with pushover analysis for different directions Columns were modelled with a force-based element formulation and considering a fibre discretization at the section level. Two uniaxial (0° and 90°) and five biaxial (at 11.25°, 22.5 °, 45 °, 67.5 °, 78.75 ° angles) pushover analyses were performed using the computer program SeismoStruct (SeismoSoft, 2004).

An interaction scaling factor was determined for each biaxial response of each column, based on the Costa-Costa uniaxial model coupled with the proposed interaction function previously presented in Equation 3.9. For these analyses, a gradient-based optimization algorithm was used (Andrade-Campos *et al.*, 2007). In order to reduce the number of variables, for the shape factors (γ, β and n) necessary to calculate C_{fx_i} and C_{fy_i} were assumed equal to the values suggested by Kunnath and Reinhorn were adopted ($\beta = \gamma = 0.5$ and $n = 2$).

Aiming at evaluating the goodness-of-fit of the given interaction scaling factor, for each column and for each direction calculation are made for, the difference between the simulated response with the simplified (*sim*) model (with the interaction function parameters) and with the refined numerical model (reference values, *ref*). These differences are evaluated for each direction (X or Y) by the Relative Global Error (RGEdirection), as given in Equation 4.1. The combination of the error for the two directions (RGEtotal) is calculated as presented in Equation 4.2.

$$\text{RGE}_{\text{direction}}[\%] = 100 \sum_{i=1}^{c} \left(1 - \frac{A_i^{sim}}{A_i^{ref}} \right) \tag{4.1}$$

$$\text{RGE}_{\text{total}} = \sqrt{\text{RGE}_\text{X}{}^2 + \text{RGE}_\text{Y}{}^2} \tag{4.2}$$

where A_i^{sim} and A_i^{ref} are, respectively, the simulated and reference values of the potential energy associated with each pushover response.

After obtained the optimal values of the scaling factor (δ) for each analysis (each column and each direction), an expression type was selected (see Equation 4.3), which depends on the column properties, namely the cross-section dimensions (h and b), the axial load ratio (ν); and loading direction (α).

$$\delta = C_1 \cdot \left(\frac{h}{b}\right)^{C_2} \cdot e^{C_3 \cdot \nu} \cdot (\tan \alpha)^{C_4} \qquad (4.3)$$

The four constants (C1, C2, C3, and C4) are to be obtained by optimization for all pushover curves, as will be shown in the following Section.

With the adopted equation for the scaling factor, the interaction function parameters were optimized for all biaxial pushover curves, using a cascade optimization strategy.
At this stage, the Levenberg-Marquardt (LM) gradient-based method and an evolutionary method (real search space EA) were grouped in sequential/cascade strategies. Thus, as mentioned before, in order to combine the advantages of both algorithms and minimise their disadvantages, the following sequence *LM+EA+LM* was used.

An important aspect in a cascade algorithm is the choice of the criteria to switch from one optimizer to another. In the present case, a heuristic approach was adopted based on numerical experiments. The criteria, as suggested in (Andrade-Campos *et al.*, 2009), were: i) Switching from LM to EA: if, from one iteration to another, the relative decrease in the quadratic objective function is less than $1\times10\text{-}15$, or the maximum admissible iteration number (predetermined value) is reached. ii) Switching from EA to LM: if stagnation of more than 500 generations is observed or the relative decrease in the quadratic objective function is less than $1\times10\text{-}15$ or the maximum admissible iteration number (predetermined value) is reached. The obtained results are summarized in Table 1 and the convergence evolution of the cascade optimization strategies in the parameter identification are presented in Figure 3.

Table 1. Parameters achieved by the cascade optimization strategy

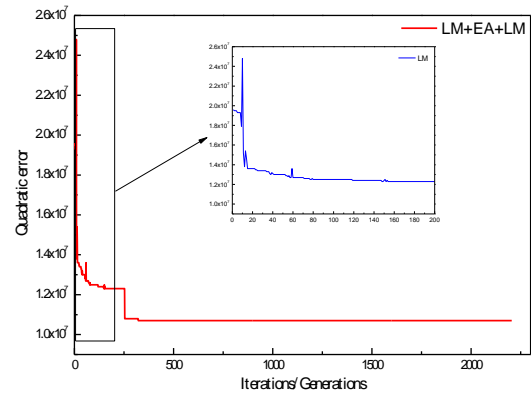| Parameter	Value
β	0.37
γ	0.90
n	2.00
C_1	1.00
C_2	0.23
C_3	0.45
C_4	-0.52

Figure 3. Convergence evolution of the cascade optimization strategies in the parameter identification

Figure 4 shows the plot of the Relative Global Error, calculated for each direction (using Equation 4.1), for two situations, in order to compare the reference simulation with the simplified model results, with and without considering the interaction function, represented in the figure with filled and unfilled marks respectively). By comparing the Relative Global Errors for the two situations the error reduction in each direction is clear when the interaction function is considered.

Figure 5 includes a selected group of examples of pushover curves for different columns and different pushover loading angles. In each plot, for both directions, the obtained pushover curves are represented: i) by the refined numerical fibre model (reference curves – blue lines with square marks); ii) by the simplified model without the biaxial bending interaction function (red lines with trianglular marks); and iii) by the simplified model but with the interaction function (the optimised solution – green lines with diamond marks). Also, the examples in Figure 5 confirm the error reductions, in both column directions, obtained by adopting the interaction function combined with the optimized scaling factor.
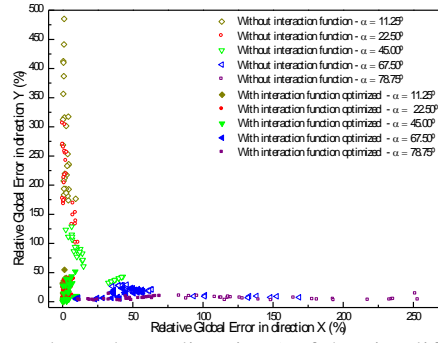
Figure 4. Relative Global Error (X and Y column directions) of the simplified model results, with (filled marks) and without (unfilled marks) interaction function, compared with the refined numerical model results
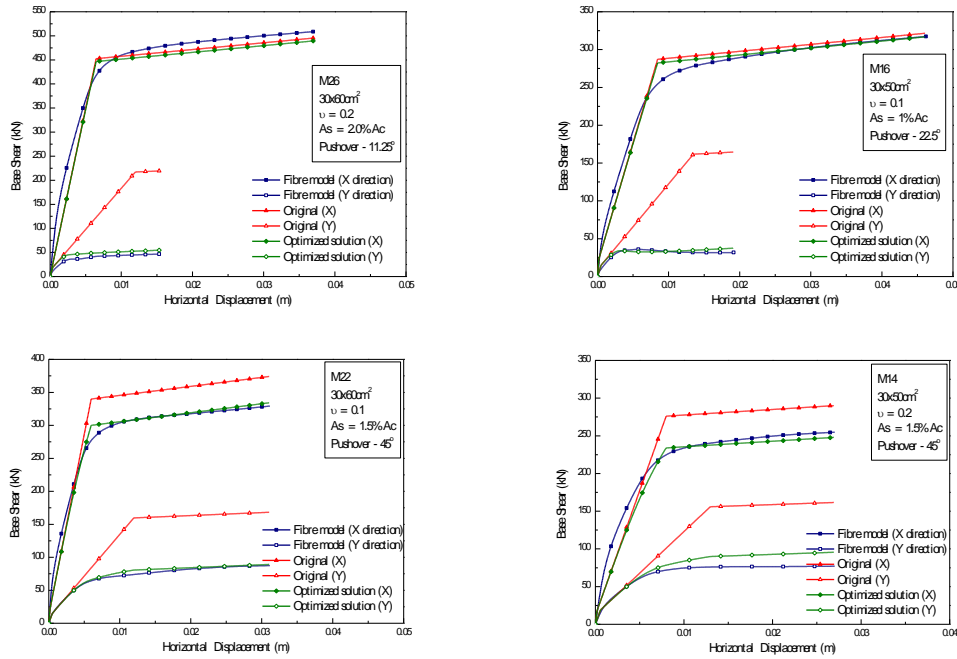

Figure 5. Examples of pushover curves for different columns and different pushover loading angles for the refined numerical model, the simplified model without the biaxial bending interaction function and the simplified model with the interaction function

5. VALIDATION OF THE MODEL WITH RESULTS FROM CYCLIC TESTS

5.1 Introduction

For the validation of the proposed simplified model with interaction functions, the experimental results of cyclic tests on 8 RC columns were used of the test campaign (H. Rodrigues *et al.*, 2012; H. Rodrigues, 2012). First, the uniaxial tests were modelled in order to obtain the primary skeleton curves for each independent direction. Then, using the obtained curves, the biaxial tests were simulated and the results of the numerical model with the interaction function are compared with the test results.

5.2 Analysis of the results

For the validation of the proposed simplified interaction model, a trilinear envelope curve was considered for the primary curve of each column direction, by best-fit adjustment to the uniaxial experimental results.

For all uniaxially tested columns, the experimental results were well reproduced with the adjusted

uniaxial trilinear curves and with the hysteretic rules of the original model, as represented in Figure 6. In the figures the following are plotted for each column: the numerical calculations with the interaction model (blue); the experimental results (red); and the trilinear primary curve (green).

As observed in these figures, the strength degradation is difficult to represent, particularly for the last cycles of the experimental response. However, significant differences are only observed for demands corresponding to drifts greater than 2.5%. The energy dissipation evolution is also well represented. Again, only for the last cycles (associated with the differences in terms of strength degradation), an overestimation of the dissipated energy is obtained with the numerical model.

In order to validate the rules and parameters of the interaction model for the simulation of the biaxial response of the tested RC columns, the trilinear curve adjusted to the uniaxial test results was considered for the primary curve in each direction.

The prediction of the experimental biaxial response obtained with the simplified interaction model for the tested RC columns is represented in Figures 7, 8 and 9. The same line legend is adopted as for the uniaxial cases.

As can be observed, the maximum strength was properly obtained with the simplified biaxial model. For the columns strong directions, a underestimation of the maximum strength of 15% was observed, while overestimation of 25% is reached in the weak direction. The unloading stiffness and the pinching effect were reasonably reproduced in most cases.
The strength degradation was also reasonably approximately in the examples under analysis. Only in the latter stages of the columns' response, close to the columns' failure, considerable differences were detected in terms of strength degradation.

The evolution of the accumulated energy dissipation is reasonable well simulated until noticeable strength degradation is observed. In the stronger column direction an overestimation of around 20% was reached and in the weak direction the underestimation is around 25%.

In general, a good agreement between the predicted numerical results and the experimental hysteretic response was observed, indicating that the proposed strategy may be suitable to simulate the response of columns to biaxial loading based on uniaxial behaviour curves associated with properly calibrated coupling interaction functions.
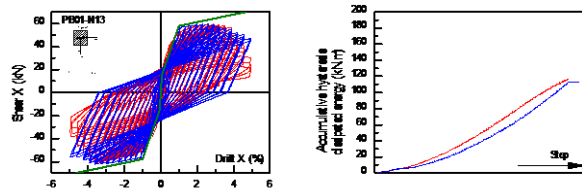

Figure 6 – Base-shear versus drift of column N13 – Uniaxial test
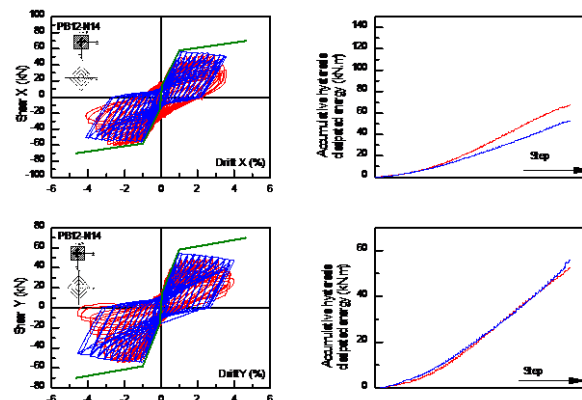

Figure 7. Base-shear versus drift of column N14 – Biaxial test, rhombus displacement pattern
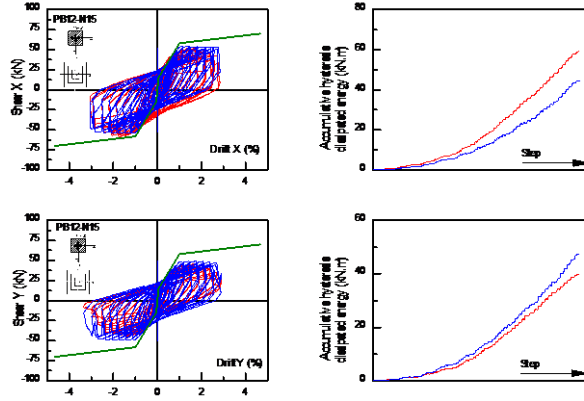
Figure 8. Base-shear versus drift of column N15 – Biaxial test, quadrangular displacement pattern
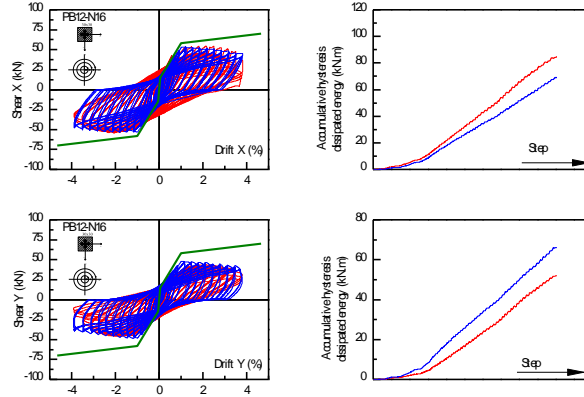

Figure 9. Base-shear versus drift of column N16 – Biaxial test, circular displacement pattern

6. CONCLUSIONS AND FINAL COMMENTS

There are still a number of unsolved problems associated with modelling of RC elements under biaxial loading. Simplified biaxial models may be adopted if they can adequately reproduce the main characteristics of the element's response (such as the strength and stiffness degradation, ductility, and energy dissipation capacity) relative to the columns uniaxial response.

In the present paper, a simplified interaction model for the response of RC columns to biaxial loading is described, based on existing uniaxial models. The proposed model corresponds to an upgrade of the existing Costa-Costa uniaxial hysteretic model, and adopts an interaction function based on the Bouc-Wen biaxial hysteretic model, coupling the two loading directions. The model parameters were calibrated using optimization techniques, based on the results of a parametric study on the tridimensional response of RC columns models with a refined model. The validity of the proposed model was demonstrated through the analytical simulation of biaxial tests on RC columns. The obtained numerical results were adequate and proved the efficiency of the model, which is a simple tool capable of reproducing the response of RC elements considering the biaxial interaction.

The proposed simplified model can be a useful tool in the design and assessment of RC structures where the response is dependent on biaxial bending of the elements. This non-linear model accounts for mechanical features such as hysteretic behaviour rules, strength and stiffness degradation, and the pinching effect. However, additional research is still necessary to objectively define the interaction function parameters, which establish the coupling of the response in the two loading directions. Moreover, the model application to experimental results obtained by other authors should be made. The implementation of the proposed model in a structural analysis program to obtain the response of RC multi-storey buildings (which strongly depends on the biaxial response of the columns) is another important task that should be archived.

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