

Letters to the Editors

Off-shell K matrix for the Morse function*

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In a recent paper (Talukdar *et al*, 1977) (hereafter cited as paper I) we have derived a wave function method for computing the off-shell K matrix elements for potentials regular at the origin. The method has been found to have certain computational advantages. More specifically, we could relate the K matrix elements for the exponential potential to Bessel functions. The object of the present communication is to demonstrate explicitly that basic results of paper I can also be used to obtain the K matrix for the Morse function in terms of tabulated transcendental functions. The Morse function represents a static soft core potential which can account for the behaviour of s -wave nucleon-nucleon phase shifts at high energies, in particular that these phase shifts become negative (Darewych & Green, 1967). The Morse function has been used to approximate the velocity dependent nucleon-nucleon potential (Green *et al*, 1967; Lodhi, 1969, 1970).

The s -wave part of the off-shell K matrix is given by (Talukdar *et al*, 1977)

$$\langle p | K | q \rangle = -\frac{2}{\pi pq} \int_0^{\infty} \sin pr V(r) \phi(k, q, r) dr, \quad \dots (1)$$

where $\phi(k, q, r)$ is an off-shell wave function regular at the origin. It satisfies a standing wave boundary condition and can be written as

$$\begin{aligned} \phi(k, q, r) = & -\frac{1}{4} \pi q \langle k | K | q \rangle [f(k, r) + f(-k, r)] \\ & + \frac{1}{2i} [f(k, q, r) - f(k, -q, r)] \quad \dots (2) \end{aligned}$$

The functions $f(k, r)$ and $f(k, q, r)$ represent the on and off-shell Jost solutions (Ruda & Whiting 1973). The object $\langle k | K | q \rangle$ stands for the half-off-shell K matrix given by

$$\langle k | K | q \rangle = \frac{2 \operatorname{Im} f(k, q)}{\pi q |f(k)| \cos \delta(k)} \quad \dots (3)$$

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In eq. (3) $f(k)$ and $f(k, q)$ represent the on and off-shell Jost functions and $\delta(k)$, the phase of the Jost function, which is obviously the phase shift. The on and off-shell Jost solutions for the Morse potential

$$V(r) = V_0(e^{-2(r-d)/b} - 2e^{-(r-d)/b}) \quad \dots \quad (4)$$

can be written in the forms (Tulukdar *et al.*, 1975).

$$f(k, r) = \alpha^{kb} e^{ikr} z^{-ikb} e^{-4z} {}_1F_1(a, c, z) \quad \dots \quad (5)$$

and

$$\begin{aligned} J(k, q, r) = & \sum_{n'=0}^{\infty} G_{n'}(V_0, b, d, q, \rho) \frac{\sigma(\sigma+c-1)}{(\sigma+n')(\sigma+n'+c-1)} \\ & \times e^{-2z} Z^{n'+2-ikb} {}_2F_2 \left(\begin{matrix} 1 & \sigma+n'+a \\ \sigma+n'+1, & \sigma+n'+c \end{matrix} \middle| Z \right), \end{aligned} \quad \dots \quad (6)$$

where $Z = \xi e^{-r/b}$

and $G_{n'}(V_0, b, d, q, \rho) = (2V_0 b)^{nq} e^{dq} \frac{\rho^{n'}}{n'!}$

with

$$\rho = \frac{1}{2}, \quad \xi = \alpha \nu, \quad \alpha = 2V_0^{1/2} b, \quad \nu = e^{d/b},$$

$$\sigma = ikb - iq b,$$

$$a = \frac{1}{2} - ikb - V_0^{1/2} b,$$

and $c = 1 - 2ibk$.

Combining eq (1), (2), (5) and (6) we get the off-shell K matrix in the form

$$\begin{aligned} \langle p | K | q \rangle = & \frac{2V_0^q}{\pi p q} \sum_{m, n \geq 0} \frac{(-1)^m \xi^{m+n}}{2^m m!} \left[B(k, q) A(k) \frac{(a)_n}{(c)_n n!} \right. \\ & \left. \langle \xi^{-kb} (\nu I(2, -k, 0) - 2I(1, -k, 0) + \xi^{kb} (\nu I(2, k, 0) - 2I(1, k, 0))) \right. \\ & \left. + \frac{1}{2i} \sum_{n'=0}^{\infty} \left[D_{n'}(q, \sigma) \frac{(\sigma+n'+a)_n}{(\sigma+n'+1)_n (\sigma+n'+c)_n} \xi^{(n-iqb)} \{ \nu I(2, -q, n') - 2I(1, -q, n') \} \right. \right. \\ & \left. \left. - D_{n'}(-q, \sigma') \frac{(\sigma'+n'+a)_n}{(\sigma'+n'+1)_n (\sigma'+n'+c)_n} \xi^{n'+iqb} \{ \nu I(2, q, n') - 2I(1, q, n') \} \right] \right] \quad (7) \end{aligned}$$

where

$$B(k, q) = -\frac{1}{4} \pi q \langle k | K | q \rangle,$$

$$A(k) = \xi^{ikb} = \alpha^{ikb} e^{ikd},$$

$$(\alpha)_n = \frac{\Gamma(\alpha+n)}{\Gamma(\alpha)},$$

$$Dn'(\pm q, x) = G_n'(V_0, b, d, \pm q, \rho) \frac{x(x+c-1)}{(x+n')(x+n'+c-1)}$$

and

$$I(j, g, n') = \frac{pb^2}{p^2b^2 + (m+n+j+n'+igb)^2}$$

The Pochhammer symbol $(\alpha)_n$ as well as the integral $I(j, g, n')$ characterising the off-shell K matrix in eq. (7) converge uniformly (Luke, 1969) as n and $n' \rightarrow \infty$. It is therefore possible to sum the series on a computer and use it to check on programmes which attempt to compute the K matrix by using any iterative procedure.

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