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Optical and electronic properties of orthorhombic and trigonal AXO₃ (A=Cd, Zn; X=Sn, Ge): First principle calculation

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ABSTRACT

Electronic structure and optical properties of the CdXO₃ and ZnXO₃ (X=Ge, Sn) compounds have been investigated based on density functional theory. According to the predictive results, reveal that the CdXO₃ and ZnXO₃ would be candidates for a high performance lead free optical crystal, which will avoid the environmental toxicity problem of the lead-based materials.

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1. Introduction

CdXO₃ and ZnXO₃ (X=Sn, Ge) are multifunctional materials, that have structure of perovskite oxides [1–3]. There are two points of view on the structures of ZnXO₃ and CdXO₃, perovskite-type oxide [4] and ilmenite structures [5]. It was found that ilmenite-type ZnXO₃ contains only cations with the electronic configuration of $(n-1)d^{10}ns^0$. However, until now little attention has been paid to oxides containing only main group cations with such electronic configuration, so these compounds gave us a new strategy to search more polar crystals [6]. To search for new high performance ferroelectrics based on the above strategy we have investigated CdXO₃ and ZnXO₃(X=Ge, Sn), because the investigation of perovskiteand ilmenite-based oxide systems with ferroelectric properties has gained large interest within the scientific community in recent years [7, 8]. It is well known that LiNbO₃ compounds is so important optical materials with a wide applications in optoelectronics and other nonlinear systems. Therefore, it seems to us noncentrosymmetric (NCS) compounds CdXO₃ and ZnXO₃ with the same ilmenite structure will be used in different application areas, like optoelectronics, too

In present work, we investigated the electronic and optical properties of the $CdXO_3$ and $ZnXO_3(X=Ge, Sn)$ compounds. The method of calculation is given in section 2; the results are discussed in section 3. Finally, the summary and conclusion are given in section 4.

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[206]/74 👄 H. OZISIK ET AL.

2. Method of calculation

In the present paper, all of our calculations that are performed using the ab-initio totalenergy and molecular-dynamics program VASP (Vienna ab-initio simulation program) [9-12] developed within the density functional theory (DFT) [13], the exchange-correlation energy function is treated within the GGA (generalized gradient approximation) by the density functional of Perdew et al. [14]. The potentials used for the GGA calculations take into account the $3d^{10}4s^24p^2$ valence electrons of each Ge-, and $4d^{10}5s^25p^2$ valence electrons of each Sn. The valence electronic configuration were 4d¹⁰5s² for each Cd-, $3d^{10}4s^2$ for each Zn-, and $2s^22p^4$ for each O- atom. When including a plane-wave basis up to a kinetic-energy cutoff equal to 19 Ha (the electronic energy convergence is 10^{-8} eV), the properties investigated in this work are well converged. The Brillouin-zone integration was performed using special k points sampled within the Monkhorst-Pack scheme [15]. We found that a mesh of $9 \times 9 \times 9$ k-points was required to describe well of these the density of states and electronic structures. In structure relaxations, the force convergence for ionic is set to 10^{-3} eV/A. This k-point mesh guarantees a violation of charge neutrality less than 0.008eV. Such a low value is a good indicator for an adequate convergence of the calculations.

The primitive cell structures (ZnXO₃ and CdXO₃) contains 2 molecules and 10 atoms. When we started the calculations, we have optimized the structural properties first. The lattice parameters obtained as a result of this optimization are shown in Table 1 with the experimental and theoretical results. The structural parameters obtained are in a good agreement with the experimental and theoretical values [1, 16–26]. We have used these structural properties in all our subsequent calculations. Following geometry optimization, the Kohn-Sham electronic band structures and the partial densities of states per atom and per orbital were calculated.

3. Results and discussion

In the first step of our calculations, we have carried out the equilibrium lattice constants of $CdXO_3$ and $ZnXO_3(X=Ge, Sn)$ compounds. The optimization lattice parameters obtained after the geometry optimization are presented in Table 1. Experimental value are shown as well, for comparison. The lattice parameters very close to experimental data when we look at the GGA result. We have also calculated formation energies for these compounds. The results are also presented in Table 1. The calculated negative formation energies mean that these phases for $CdXO_3$ and $ZnXO_3$ compounds are thermodynamically stable at ground state. Generally, ilmenite type rhomboedric structures exhibit a higher structural stability than that of perovskite type orthorhombic structures since the ilmenite type rhomboedric structures possesses the lower formation energies, and $ZnGO_3$ compound in ilmenite type rhomboedric structures may be the most stable among them.

The Kohn-Sham electronic band structure gives a picture of electronic eigenenergies as a function of a set of quantum numbers which form the components of a wavevector \mathbf{k} in the first Brillouin zone (BZ). For perovskite and ilmenite CdXO₃ and ZnXO₃, the path in the BZ used for the DFT computations are formed straight segments connecting a set of high symmetry points. Figure 1 present the electronic band structures obtained for the CdXO₃ and ZnXO₃ together with the respective partial and total DOS in a manner

| Lattice | Material | а | b | с | E ₀ (eV/f.u.) | V ₀ (Å ³ /f.u.) | $E_g (eV)^*$ | $\Delta { m H_f}$ (eV/p.a) | Refs. |
|---------------------------|--------------------|--------|--------|--------|--------------------------|---------------------------------------|--------------------|----------------------------|---------------|
| Ilmenite type | ZnGeO₃ | 5.046 | | 14.077 | -28.1930 | 51.73 | 2.106 (I) | -1.6806 | Present |
| Rhomboedric R-3(148) | | 4.9568 | | 13.860 | | 49.15 | | | Exp.[16] |
| A: 6c (0, 0, z) | ZnSnO₃ | 5.378 | | 14.29 | -27.3057 | 59.66 | 1.302 (I) | -1.5802 | Present |
| B : 6c (0, 0, z) | | 5.419 | | 14.348 | | | | | US-PBE [17] |
| O : 18f (x, y, z) | | 5.284 | | 14.091 | | | | | Exp. [18] |
| | | 5.2118 | | 13.90 | | 56.29 | | | LDA-CA [19] |
| | | 5.419 | | 14.348 | | 60.82 | | | GGA-PBE[20] |
| | | 5.284 | | 14.091 | | 56.78 | | | Exp. [21] |
| | $CdGeO_3$ | 5.193 | | 15.257 | -27.2324 | 59.38 | 1.398 (I) | -1.5602 | Present |
| | | 5.098 | | 14.883 | | 55.83 | | | Exp. [22] |
| | CdSnO ₃ | 5.556 | | 15.272 | -26.5540 | 68.05 | 0.981 (I) | -1.5016 | Present |
| Perovskite type | ZnGeO ₃ | 5.074 | 5.14 | 7.451 | -27.7240 | 48.59 | 1.588 (D) | -1.5868 | Present |
| Orthorhombic Pnma (62) | ZnSnO₃ | 5.383 | 5.375 | 7.922 | -27.1404 | 57.31 | 1.069 (D) | -1.5471 | Present |
| A :4c (x, 0.25, z) | | 5.428 | 5.422 | 7.994 | | | | | US-PBE [17] |
| B :4b (0, 0, 0.50) | $CdGeO_3$ | 5.31 | 5.362 | 7.582 | -26.7519 | 53.97 | 0.817 (D) | -1.4641 | Present |
| O :8d (x, y, z) | | 5.209 | 5.253 | 7.434 | | 50.86 | | | Exp. [22] |
| 4c (x, 0.25, z) | | 5.2114 | 5.2608 | 7.4263 | | 50.90 | | | Exp [23] |
| | | 5.3165 | 5.3750 | 7.5903 | | 54.23 | | | GGA-PW91 [24] |
| | | 5.136 | 5.369 | 7.579 | | 49.10 | 0.82 | | US-PBE [25] |
| | | 5.152 | 5.197 | 7.334 | | 54.08 | 1.67 | | US-LDA [25] |
| | | 5.211 | 5.261 | 7.443 | | 51.01 | | | Exp. [25] |
| | CdSnO ₃ | 5.556 | 5.676 | 8.041 | -26.4046 | 63.40 | 0.650 (D) | -1.4717 | Present |
| | | 5.4989 | 5.6068 | 7.9488 | | 61.27 | 0.94 | | PAW-PBE [26] |
| | | 5.1927 | 5.2856 | 7.4501 | | 51.12 | 1.7 | | PAW-LDA [26] |
| | | 5.4588 | 5.5752 | 7.8711 | | 59.89 | 3.0 | | Exp. [1] |
| | | | | | | 0.42 | | | US-PBE [26] |
| | | | | | | | 1.14 | | US-LDA [26] |

Table 1. The calculated equilibrium lattice parameters (a, b, and c in Å) and electronic bandgaps (E_g in eV) for AXO₃ (A = Zn, Cd and X = Ge, Sn) in orthorhombic perovskite and rhomboedric ilmenite structures.

* (I) and (D) refers indirect and direct bandgap, separately.

similar to our previous work [27]. Each ilmenite unit cell has two units (Z = 2), which leads to 68 valence electrons per cell, while the corresponding values for perovskite are Z = 4, and 136 valence electrons.

The band structure of CdXO₃ and ZnXO₃ along the principal symmetry directions have been calculated by using the equilibrium lattice constants as shown in Table 1 in perovskite orthorhombic and ilmenite phases. As a result of our calculations the band structure of the ZnGeO₃ and ZnSnO₃ have a direct gap (Γ -high symmetry point), which are 1.588 eV and 1.069 eV in the perovskite phase, respectively. For CdGeO₃ and CdSnO₃, we received the same results in perovskite phase. They have a direct band at Γ -high symmetry point with the forbidden gap 0.817 eV and 0.650 eV, respectively. In the ilmenite phase, the gaps are 2.106 eV and 1.302 eV for ZnGeO₃ (indirect band transition between B- Γ high symmetry points) and ZnSnO₃ (X- Γ), respectively. For CdGeO₃ (B- Γ) and CdSnO₃ (Z- Γ) in ilmenite phase, the gaps are 1.398 eV and 0.981 eV, respectively. The general features of the energy bands such as band gaps, DOS and orbital hybridization are similar for all investigated compounds. Furthermore our results are in agreement with the results obtained in previous calculations [1, 25, 26]. We have summarized the band gap energies for ZnXO₃ and CdXO₃ in Table 1 with the available theoretical and experimental results. Our results in perovskite type orthorhombic structures can roughly compare with band structure for SrSnO₃ and CaGeO₃ compound in same structure [8, 25]. The band structures of CdXO₃ and ZnXO₃ is qualitatively similar to that of SrSnO₃ and CaGeO₃. The band gap for SrSnO₃ (2.9 eV) [8] is higher than that for CdXO₃ and ZnXO₃ compounds.



Figure 1. The calculated electronic band structure of AXO₃ (A = Zn, Cd and X = Ge, Sn) compounds. (a) Ilmenite-ZnGeO₃, (b) Perovskite-ZnGeO₃, (c) Ilmenite-ZnSnO₃, (d) Perovskite-ZnSnO₃, (e) Ilmenite-CdGeO₃, (f) Perovskite-CdGeO₃, (g) Ilmenite-CdSnO₃, (h) Perovskite-CdSnO₃.

The real $\varepsilon_1(\omega)$ and imaginary $\varepsilon_2(\omega)$ parts of $\varepsilon(\omega) = \varepsilon_1(\omega) - i\varepsilon_2(\omega)$ were calculated by using Kramers-Kroning transformation [28]. Figure 2 shows the real and imaginary parts of $\varepsilon(\omega)$ for both ilmenite and perovskite ZnXO₃ and CdXO₃ compounds in a manner similar to our



Figure 2. The calculated real and imaginary parts of linear dielectric function of AXO_3 (A = Zn, Cd and X = Ge, Sn) compounds. (a) ZnGeO₃ (Ilmenite-type), (b) ZnGeO₃ (Perovskite-type), (c) ZnSnO₃ (Ilmenite-type), (d) ZnSnO₃ (Perovskite-type), (e) CdGeO₃ (Ilmenite-type), (f) CdGeO₃ (Perovskite-type), (g) CdSnO₃ (Ilmenite-type), (h) CdSnO₃ (Perovskite-type).

recent work [29]. The calculated $\varepsilon_2(\omega)$ show main peaks in the range of 2.0 to 20 eV for both phases. These peaks related for $\varepsilon_2(\omega)$ are related to the interband transition from the valence to conduction bands (interband transition from O2s and O2p to Ge and Sn s- and p- states).

Conclusion

This work present ab initio studies of electronic structure and optical properties of $CdXO_3$ and $ZnXO_3(X=Ge, Sn)$ compounds. By a search of the total energy minimum the equilibrium volume are found, which are in good agreement with experimentally determined one. The DFT calculations of the electronic structure and optical properties of the investigated compounds have been performed using the calculated equilibrium lattice parameters. It is shown that all compounds considered are semiconductor nature. The lowest CB is formed of the valence orbitals of the Cd(Zn) and X (Sn,Te) atoms and major contribution comes from d-orbitals of Cd(Zn) and p-orbitals of X(Ge, Sn) atoms. Topmost valence band is formed of the valence orbitals of the O-atoms and major contribution comes from p-orbitals of O. Finally, optical properties were investigated. The relations of the optical properties to the interband transitions were also discussed.

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