Elevated temperature lasing from injection microdisk lasers on silicon

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Abstract

A combination of high operation temperatures and small sizes of diode lasers directly grown on silicon substrates is essential for their application in future photonic integrated circuits. In this Letter, results are presented on quantum dot III-V-on-Si microdisk diode lasers tested at elevated temperatures. To the best of our knowledge, the first uncooled microlasers with diameter of 30 μ m capable of operating in continuous wave regime at 60°C are demonstrated. In lasing regime, emission spectra contain one most intense line with the full width at half maximum of 30 pm; side mode suppression ratio reaches 18 dB. Because of self-heating, an actual temperature of the active region is close to 100°C. Under pulsed excitation, the maximal lasing temperature is 110°C.

Keywords: Microcavity devices, quantum dots, semiconductor lasers

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1. Introduction

Compact size and capability of operation at elevated temperatures facilitate integration of light emitters with silicon-based electronic and photonic elements. Furthermore, a laser heterostructure is preferably to be directly grown on silicon [1]. Microdisk (MD)/microring (MR) lasers are prospective for such integration, since high quality factors can be achieved even in resonators of a few micrometers in diameter [2, 3]. Suppressed lateral migration of charge carriers along the quantum dot (QD) plane makes QD lasers less sensitive to deep etching [4] and promotes use of ODs in microlasers intended for operation at elevated temperatures [5, 6]. Strong carrier localization within individual QDs is also favorable for their resistance against non-radiative recombination in III-V-on-Si structures. A significant progress has recently been demonstrated in QD structures monolithically grown on Si including realization of edge-emitting lasers [7-11] and optically pumped microlasers [12]. Recently, room temperature (RT) lasing has been achieved in electrically pumped MDs [13] and MRs [14] on Si with the smallest diameter of 14 and 10 μ m, respectively. For larger MR devices (100 µm in diameter) [14], the maximum operation temperature (pulsed and continuous wave (CW)) of 100°C was demonstrated. The lasing spectra of those MRs contain a number of closely spaced modes of comparable intensity within a broad spectral interval from 1.28 to 1.3 µm.

Devices of smaller size would be helpful for singlemode lasing. However, there are several effects which impede their high temperature operation. First, relative contribution of inevitable non-radiative recombination at sidewalls increases proportionally to the perimeter-to-area ratio. In this respect, a MD laser can be advantageous over a MR one owing to absence of an inner hole. For example, a CW RT threshold current density (J_{th}) of 594 A/cm² achieved in MDs with diameter of 30 µm is only 1.4 times higher than J_{th} measured in the broad-area laser made of the same QD-on-Si epitaxial structure [13]. Second, self-heating is expected to be more significant in smaller lasers because of hindered heat dissipation. In comparison with a MR laser of the same diameter, a MD laser has larger area of its contact with substrate, which is beneficial for lower thermal resistance. Specific thermal resistance of 4 °C×cm²/mW was found in QD-on-Si MD lasers with diameters from 14 to 30 µm [13].

In this article, we report on uncooled operation of QD-on-Si MD lasers at elevated temperatures. Lasing was achieved under CW and pulsed excitation up to 60 and 110° C, respectively, in MDs with diameter of 30 µm. These are the smallest microlasers on silicon operating at these elevated temperatures. Nearly single mode emission is demonstrated with a dominant mode linewidth as narrow as 30 pm. Up to 90°C lasing proceeds on the ground state (GS) optical transition of QDs with wavelength about 1.35 µm. At higher temperatures, lasing wavelength jumps to 1.3 µm that corresponds to the excited state (ES) transition.

2. Methods

A III-V epitaxial structure was grown by molecular beam epitaxy on a silicon (001) substrate with 4° offcut to the [011] plane. Upon oxide desorption, a 30-nm AlAs nucleation layer, a 570-nm GaAs buffer layer, an Al_{0.4}Ga_{0.6}As/GaAs superlattice (SL) of 200 nm, and a 100 nm GaAs buffer layer were deposited followed by *n*-type doped InGaAs/GaAs dislocation filter layers (DFL) [15] of total thickness of 1.7 μ m. An AlGaAs/GaAs laser heterostructure grown on the top of the DFL consists of a 1.4 μ m *n*-Al_{0.4}Ga_{0.6}As lower cladding layer, an undoped GaAs waveguiding layer, a 1.4 μ m *p*-Al_{0.4}Ga_{0.6}As/GaAs SLs inserted at interfaces. An active region represents 7 layers of InAs/InGaAs dot-in-well structures [16] separated by 38.5 nm GaAs spacers. Total thickness of the structure cross-section (the left panel of figure 1) reveals good interface planarity and absence of microcracks.

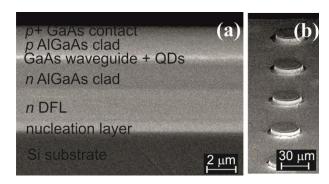


Figure 1. SEM images of cross section view of the epitaxial structure (a) and axonometric view of five 30-µm microdisks (b).

The structure was processed into an array of 4.5- μ m tall MDs [6] with diameter of 30 μ m (the right panel of figure 1) using photolithography and inductively coupled BCl₃/Ar plasma dry etching process. No sidewall coating or passivation was applied. Individual circular-shaped AgMn/NiAu *p*-ohmic contacts were placed on top of the MDs and a common AuGe/Ni/Au *n*-ohmic contact was deposited onto the etched surface between the MDs. MD lasers were put onto a copper holder, which temperature was varied by a heater from RT (about 25°C) to 110°C.

Microlaser characteristics were measured in CW and pulsed regimes. In the latter case, we used 0.5-µs-long injection pulses with 10 kHz repetition rate. The MDs were electrically pumped using needle probes. DC current was varied from 0 to 25 mA. In pulsed regime, injection current was limited at 300 mA. A turn-on voltage of 1.6 V and a series resistance of 40 Ohm were derived from the current-voltage curves (5-20 mA) taken at RT. An adjustable infrared objective with x10 magnification was used to collect light emitted from an individual microlaser. Emission spectra were acquired with spectral resolution of ~30 pm using a Horiba FHR 1000 monochromator and an InGaAs CCD array.

3. Results and discussion

An emission spectrum of a MD contains a series of nearly equally spaced sharp lines which are attributed to whispering gallery modes of different azimuthal orders. A representative CW spectrum measured at 17 mA current and the holder temperature of 60°C is depicted in figure 2. As argued below, this injection current is beyond the lasing threshold estimated as 15 mA for these operation conditions. The most intense (dominant) mode is located at 1362.2 nm. It is much brighter than the background electroluminescence in the whole interval of spontaneous emission. The neighboring longer-wavelength (side) mode ($\lambda = 1367.5$ nm) is more than one order of magnitude less intense in comparison with the dominant one. Other modes are hardly seen since they are one more order of magnitude less intense. High purity of emission spectrum distinguishes the MD lasers under study from recently demonstrated MR lasers on silicon,¹⁴ where a plurality of modes of comparable intensity is observed.

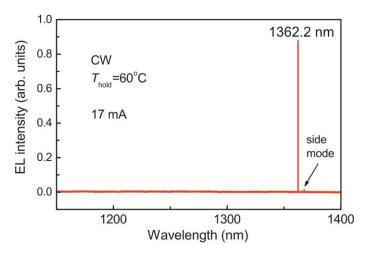


Figure 2. Overall emission spectrum obtained at 17 mA.

Figure 3 depicts in detail those parts of the spectrum where the dominant and side modes are located. The data are presented for injection current varied around the threshold, from 14.5 to 16 mA. As current increases in this interval, intensity of the side mode remains nearly unchanged whereas the dominant mode demonstrates a rapid growth. The side mode suppression ratio (SMSR), which is defined as peak-to-peak ratio of the dominant mode to the side one, reaches its maximal value of 17.8 dB at 17 mA. The cavity resonances are shifted to longer wavelengths as injection current increases. In CW regime, the mode redshift is caused by self-heating of the uncooled device. Spectral positions of the dominant and side modes against injection current are presented in the inset of figure 3. For both modes, the slope $(d\lambda/dI)$ of 0.144 nm/mA was derived. This corresponds to a current induced temperature increment (dT/dI) of 1.77 °C/mA taking into account a temperature coefficient $(d\lambda/dT)$ of 81.5 pm/°C [6]. An actual temperature of the MD laser was then calculated by adding the temperature can reach 100°C at the maximal current used (25 mA). At threshold, the device is additionally heated by approximately 26°C.

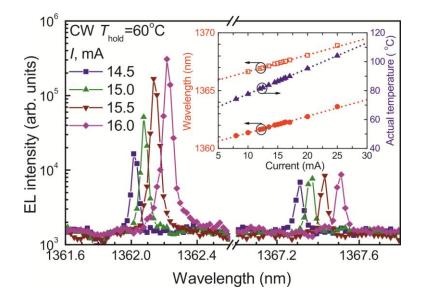


Figure 3. Close-up spectra (logarithmic scale) taken at 60° C and different injection currents. Inset: Current-induced variation of spectral positions of the dominant (circles) and side (squares) modes measured at holder temperature of 60° C, CW regime. Triangles: estimated actual temperature of MD. Dashed curves: linear fits.

Realization of CW lasing at 60°C is confirmed by dependences of the dominant mode integrated intensity, linewidth and SMSR on injection current (figure 4). The light-current curve reveals a pronounced knee. The linewidth simultaneously drops down to the spectral resolution limit, and a rapid growth of SMSR takes place. We estimated the threshold current I_{th} as 15 mA, which corresponds to J_{th} of 2.1 kA/cm². As injection current increases beyond 17 mA, SMSR starts decreasing and even changes its sign. This is accompanied by a decay of the mode intensity and some broadening of the line. All these observations reflect the fact that the longer-wavelength mode ($\lambda = 1367.5$ nm) becomes the dominant one. Similar behavior is also observed at lower temperatures. The lower is temperature the larger is interval of injection current in which a certain lasing mode persists.

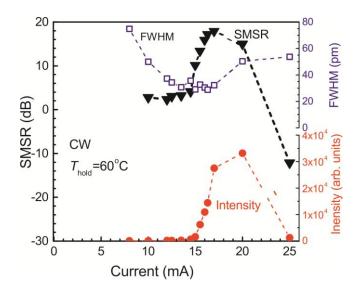


Figure 4. SMSR (triangles), FWHM (squares), and integrated line intensity (circles) of the dominant mode (λ =1362.2 nm) as functions of injection current in CW regime at 60°C.

The figure 5 shows log-log plot of the dominant mode intensity as a function of injection current at RT, 45°C and 60°C. All the dependences show a characteristic S-shape indicating the transition from spontaneous to lasing regime of emission [17]. As holder temperature increases, I_{th} shifts to higher values. Meanwhile, at 75°C the dominant mode does not anymore demonstrate a rapid growth, and its intensity remains low with increase in injection current. The maximal temperature, at which CW lasing is achieved in these QD microdisk lasers on silicon, is 60°C. To the best of our knowledge, this is the highest CW operation temperature of any III-V microlasers directly grown on silicon and having a microlaser lateral size of few tens of micrometers.

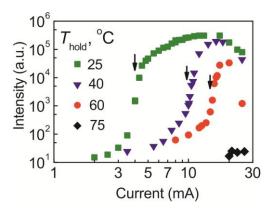


Figure 5. Integrated intensity of the dominant modes at different holder temperatures (squares 25°C, triangles 45°C, circles 60°C, rhombs 75°C). Arrows indicate threshold current values where applicable.

Figure 6 summarizes the minimal J_{th} measured at different holder temperatures in CW and pulsed regimes. At RT, the threshold current is 4.2 mA ($J_{th} = 0.59 \text{ kA/cm}^2$) in CW and 3.8 mA ($J_{th} = 0.54 \text{ kA/cm}^2$) in pulsed mode. The maximal temperature, at which lasing was realized, is limited by our experimental setup, 110°C, with J_{th} of 13.45 kA/cm². Spectral positions of the dominant modes taken at the thresholds are shown in figure 6 together with wavelengths of the GS and ES optical transitions of QDs. It is seen that up to 90°C lasing proceeds within the GS transition; the wavelength is around 1.35 µm. The spectral positions of the lasing mode measured at 100 and 110°C are significantly shifted to a shorter wavelength (~1.3 µm) that corresponds to the ES transition. This indicates that the maximal gain achievable at the GS transition is insufficient to balance optical loss at sufficiently high temperatures. We suggest that additional optical loss at high temperatures is caused by enhanced free carrier absorption in the laser waveguiding layer.

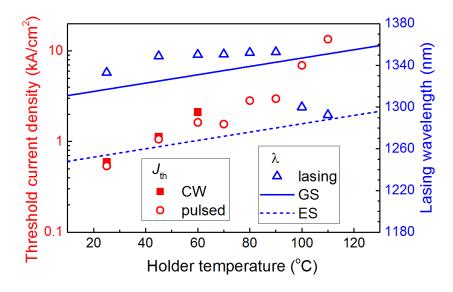


Figure 6. Temperature dependences of threshold current densities in CW (squares) and pulsed (circles) regimes, spectral positions of lasing wavelength (triangles), ground-state (solid line) and excited-state (dashed line) QD emission maxima.

4. Conclusion

In summary, we examine uncooled QD-based microdisk lasers on silicon at elevated temperatures. We demonstrate, for the first time, lasing operation of 30- μ m in diameter microlasers at 60°C in CW and 110°C in pulsed mode. Under these operation conditions (excepting 100-110 °C), lasing wavelength is well beyond 1.3 μ m being within the GS optical transition. As holder temperature reaches 100°C, lasing wavelength switches to the ES transition (~1.3 μ m). In CW lasing regime, the $\lambda/\Delta\lambda$ ratio exceeds 45 000. At the maximal temperature of CW operation lasing spectrum remains nearly singlemode with the dominant mode to the side mode intensity ratio about 18 dB.

Acknowledgments

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References

- [1] Cornet C, Léger Y and Robert C 2016 Integrated Lasers on Silicon (ISTE Press Elsevier). https://www.elsevier.com/books/integrated-lasers-on-silicon/cornet/978-1-78548-062-1
- [2] Srinivasan K, Borselli M, Painter O, Stintz A and Krishna S 2006 Opt. Express 14 1094.
- [3] Kryzhanovskaya N V, Zhukov A E, Maximov M V, Moiseev E ., Shostak I I, Nadtochiy A M, Kudashova Yu V, Lipovskii A A, Kulagina M M, and Troshkov S I 2015 *IEEE J. Selected Topics Quantum Electron.* 21 1900905.
- [4] Ouyang D, Ledentsov N N, Bimberg D, Kovsh A R, Zhukov A E, Mikhrin S S, and Ustinov V M 2003 Semicond. Sci. Technol. 18 L53.
- [5] Munsch M, Claudon J, Malik N S, Gilbert K, Grosse P, Gérard J-M, Albert F, Langer F, Schlereth T, Pieczarka M M, Höfling S, Kamp M, Forchel A, and Reitzenstein S 2012 Appl. Phys. Lett. 100 031111.
- [6] Kryzhanovskaya N V, Moiseev E I, Kudashova Yu V, Zubov F I, Lipovskii A A, Kulagina M M, Troshkov S I, Zadiranov Yu M, Livshits D A, Maximov M V, and Zhukov A E 2015 *Electron. Lett.* **51** 1354.
- [7] Wang T, Liu H, Lee A, Pozzi F, and Seeds A 2011 Opt. Express 19 11381.
- [8] Lee A, Jiang Q, Tang M, Seeds A, and Liu U 2012 Opt. Express. 20, 22181.
- [9] Lee A D, Jiang Q, Tang M, Zhang Y, Seeds A J, and Liu H 2013 *IEEE J. Selected Topics Quantum Electron.* **19** 1901107.
- [10] Liu A Y, Zhang C, Norman J, Snyder A, Lubyshev D, Fastenau J M, Liu A W K, Gossard A C, and Bowers J E 2014 *Appl. Phys. Lett.* **104** 041104.
- [11] Chen S, Li W, Wu J, Jiang Q, Tang M, Shutts S, Elliott S N, Sobiesierski A, Seeds A J, Ross I, Smowton P M, and Liu H 2016 *Nat. Photon.* 10 307.
- [12] Wan Y, Li Q, Liu A Y, Gossard A C, Bowers J E, Hu E, and Lau K M 2016 Opt. Lett. 41 1664.
- [13] Kryzhanovskaya N V, Moiseev E I, Polubavkina Yu S, Maximov M V, Kulagina M M, Troshkov S I, Zadiranov Yu M, Guseva Yu A, Lipovskii A A, Tang M, Liao M, Wu J, Chen S, Liu H, and Zhukov A E 2017 *Optics Letters*, **42**, 3319.

- [14] Wan Y, Norman J, Li Q, Kennedy M J, Liang D, Zhang C, Huang D, Liu A Y, Torres A, Jung D, Gossard A C, Hu E L, Lau K M, and Bowers J E 2017 *CLEO: Applications and Technology* (San Jose, CA, USA, 14–19 May 2017).
- [15] Tang M, Chen S, Wu J, Jiang Q, Kennedy K, Jurczak P, Liao M, Beanland R, Seeds A J, and Liu H 2016 *IEEE J. Selected Topics Quantum Electron.* 22, 1900207.
- [16] Liu H Y, Sellers I R, Badcock T J, Mowbray D J, Skolnick M S, Groom K M, Gutiérrez M, Hopkinson M, Ng J S, David J P R, and Beanland R 2004 *Appl. Phys. Lett.* **85**, 704.
- [17] Jaffrennou P, Claudon J, Bazin M, Malik N S, Reitzenstein S, Worschech L, Kamp M, Forchel A, and Gérard J-M 2010 *Appl. Phys. Lett.* **96**, 071103.