# An experimental study of Primary Cosmic Rays at the knee energy region by observation of Extensive Air Showers (E AS) 

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Summary. - Simultaneous measurements have been made of the radial (lateral) electron density distribution and the radial muon density distribution at various measured muon energies in the range $2.5-100 \mathrm{GeV}$ in vertically incident EAS in the size range $3.15 \times 10^{4}-1.79 \times 10^{6}$ (primary energy range $2.4 \times 10^{14}-8.3 \times 10^{15} \mathrm{eV}$ ) particles detected near sea level. The characteristics of these radial distributions in terms of the measured shower parameters have been determined and used to draw conclusions about the average nuclear mass of the primaries of these EAS.

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## 1. - Introduction

It is well established that the Primary Cosmic $R$ ay energy spectrum steepens in the energy range $10^{14}-10^{16} \mathrm{eV}$ known as the knee energy region. The spectrum steepening is thought to be due to mechanisms related to acceleration and propagation of Primary Cosmic Rays (PCR). An Extensive Air Shower experiment at primary energy in the knee energy region in which the energy spectrum and lateral distribution of muons can be measured accurately over a wide range can indicate the trend of average behaviour of PCR mass at the knee region. In an EAS of size $\left(\mathrm{N}_{\mathrm{e}}\right) 10^{4}-10^{6}$ particles at sea level the total number of muons and muon lateral distribution, if the primary is a heavy nucleus, will be different from those in a primary proton initiated shower and the difference can be detected from the data of a proper set-up experiment. In the present paper the results of an air shower experiment covering the knee energy region will be described. This experiment with the provision for differential measurement of both low- and high-energy muons simultaneously by two shielded magnet spectrographs has yielded both muon energy spectrum and radial density distributions at various muon energies. The properties of these distributions have been analysed to infer the trend of average PCR mass composition at the knee region.

## 2. - North Bengal University ( $N B U$ ) air shower experiment

The North Bengal University EAS array for observation of air showers has been developed in stages since 1980 (B asak et al. [1]). The set-up has been designed to detect small- and medium-size air showers with a close-packed array (of spacing 8 m ) using an array of 35 scintillation detectors each of size $50 \mathrm{~cm} \times 50 \mathrm{~cm}$, two shielded muon magnet spectrographs (with a spacing of 4 m , maximum detectable momentum (MDM) $500 \mathrm{GeV} \mathrm{c}^{-1}$, each of area $1 \mathrm{~m} \times 1 \mathrm{~m}$ and cut off at an energy of 2.5 GeV ) and a neon flash tube (NFT) chamber as a low-energy muon detector. Two spectrographs each with a lever arm of 6.3 m were set up pointing to the zenith to collect muons above 2.5 GeV in incident vertical air showers. With such a close-packed array, the determination of shower size and other shower parameters has been more precise.
21. EAS detector array characteristics. - The array of 35 close-packed plasticscintillation detectors (fig. 1) has been operating at a site of atmospheric depth $\sim 1000 \mathrm{gcm}^{-2}$ at the NBU campus. The response measured in terms of the relative light output of plastic scintillator (manufactured by Bhabha Atomic Research Centre, India) for incident EAS electrons is $\sim 100 \%$ and the relative efficiency in terms of single-particle pulse height is nearly uniform from centre to edge of each plasticscintillation detector. The pulses from all the 35 scintillation detectors are digitized by

| $\begin{gathered} \square \\ (0,28) \end{gathered}$ | $\begin{array}{cc}  & 0 \\ & (14,28) \\ & \\ 0 & \\ (7,24) & \end{array}$ |  | $\stackrel{\square}{(21,24)}$ | - -- Scintillation counter O--Magnetic spectrograph - |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | $\stackrel{\square}{(35,24)}$ |  |
| $\begin{gathered} \square \\ (0,20) \end{gathered}$ |  | $\stackrel{\square}{(14,20)}$ |  |  |  |  | $\stackrel{\square}{(42,20)}$ |
|  | $\stackrel{\square}{(7,16)}$ |  | $\underset{(21,16)}{M_{1}}$ | $\begin{gathered} M_{2} \\ 0 \end{gathered}$ | $\stackrel{\square}{(35,16)}$ |  |
| $\stackrel{\square}{(0,12)}$ |  | $\stackrel{\square}{(14,12)}$ |  |  |  | $\underset{(42,12)}{\square}$ |
|  | $\begin{gathered} \square \\ (7,8) \end{gathered}$ |  | $\begin{gathered} \square \\ (21,8) \end{gathered}$ |  | $\stackrel{\square}{(35,8)}$ |  |
| $\begin{gathered} \square \\ (0,4) \end{gathered}$ |  | $\stackrel{\square}{\square}$ |  | $\begin{gathered} \square \\ (28,4) \end{gathered}$ |  | $(42,4)$ |
| $Y$ (2, ${ }^{\text {r }}$ |  |  |  |  |  |  |
| $\underset{\substack{(0,0) \\ \mathrm{x} \text { in } \mathrm{m}}}{\substack{\mathrm{in}}}$ | $\stackrel{\square}{(7,0)}$ |  | $\begin{gathered} \square \\ (21,0) \end{gathered}$ |  | $\begin{gathered} \square \\ (35,0) \end{gathered}$ |  |

Fig. 1. - A diagram of $N B U$ air shower array set-up.


Fig. 2. - The efficiency of detection of the array as a function of distance from the centre of the array for $\mathrm{s}=1.2$ and $\mathrm{N}_{\mathrm{e}}=8 \times 10^{4}$.


Fig. 3. - Same as in fig. 2 but for $\mathrm{N}_{\mathrm{e}}=10^{5}$.


Fig. 4. - Same as in fig. 2 but for $N_{e}=5 \times 10^{5}$.


Fig. 5. - Same as in fig. 2 but for $\mathrm{N}_{\mathrm{e}}=10^{6}$.


Fig. 6. - The average triggering probability for the array of detectors as a function of shower size for core distances at $20 \mathrm{~m}, 30 \mathrm{~m}$ and 40 m .
an analog-to-digital converter one after the other and connected to the memory unit for storing the digital information and subsequent transfer to the printer for printing on a paper tape. The printed outputs give the information about the particle densities on each detector in an individual shower.

The efficiency of detecting showers of different $\mathrm{N}_{\mathrm{e}}$ and age (s) within the sensitive detecting area of $\sim 1200 \mathrm{~m}^{2}$ and the average triggering probability of the array for showers falling within $15^{\circ}$ of the zenith are shown in figs. 2, 3, 4, 5 and 6 .

Two magnetic spectrographs and the NFT chamber were operated under an EAS trigger which was also used for the photographic recording of muon trajectories within a track location uncertainty of $\pm 0.14 \mathrm{~cm}$.

The selection criteria for shower detection consists of the following steps:

1) Shower is recorded by the detecting system if the registered electron density in any four adjacent detectors of the 8 central triggering detectors is greater than 4 particles $/ \mathrm{m}^{2}$.
2) The electron densities at 19 points are registered simultaneously with the photographic recording of the trajectories of muons by the two spectrographs at four points in individual shower.
3) F or each recorded shower, the core location, the shower size and the photon electron cascade age (s) of the EAS are determined by fitting the registered electron


Fig. 7. - Plots of lateral electron density distribution at shower size $1.09 \times 10^{5}-1.79 \times 10^{6}$ with shower age $(\mathrm{s})=1.2$.
densities with the following formula for the electron density $\Delta(r)$ :

$$
\begin{equation*}
\Delta(r)=N_{e} / r_{1}^{2}\left[f\left(r / r_{1}, s\right)\right], \tag{1}
\end{equation*}
$$

where $f\left(r / r_{1}, s\right)=c(s)\left(r / r_{1}\right)^{a_{1}+a_{2}(s-1)}\left(1+r / r_{1}\right)^{b_{1}+b_{2}(s-1)}$.
$\mathrm{c}(\mathrm{s})$ is the normalisation constant, $\mathrm{r}_{1}=24 \mathrm{~m}$ (M olière unit of displacement at sea level) and $a_{1}, a_{2}, b_{1}, b_{2}$ are constants. Some examples of aver age radial electron density distributions at $N_{e}=1.09 \times 10^{5}-1.79 \times 10^{6}$ particles together with the graphs of fitting function are shown in fig. 7. The shower size recorded by the array is $3.15 \times$ $10^{4}-1.79 \times 10^{6}$ particles.

## 3. - Analysis and error estimation

A standard $\chi^{2}$-minimization procedure based on the method of steepest descent has been used to determine the air shower parameters and to simulate the errors in the air shower parameters.

The errors in the determination of EAS parameters have been evaluated through the standard procedure of artificial shower analysis. A shower of known parameters is


Fig. 8. - E rror distribution in core location, shower size and age parameter.
allowed to be incident at any point on the array selected at random and particle density in each detector is calculated according to the chosen lateral distribution function. To reproduce the experimental conditions, the statistical fluctuations in the number of particles in each detector and the systematic errors in the conversion of pulse height into particle density are superposed on each density. F or a set of densities for each shower, $\chi^{2}$-minimization procedure is applied to estimate the shower parameters. The estimated shower parameters give the deviations from the actual ones used for an artificial shower and lead to the following error estimates:
a) $\Delta X=2.40 \mathrm{~m}, ~ \Delta \mathrm{Y}=2.77 \mathrm{~m}$,
b) $\Delta N_{e} / N_{e}=9.61 \%$,
c) $\Delta \mathrm{S}=0.13$.

Some histograms for the deviations of the parameters are shown in fig. 8.

## 4. - Results on "near vertical showers"

4.1. Determination of shower age(s) and energy of the shower. - Using all the registered electron densities in a shower event, the best-fitted value of the shower age (s) was determined by the method of least square. F or each shower of fixed size and age (s), the mean energy of the primary particle (proton) was obtained with a maximum error of $10 \%$ (which includes fluctuation in E AS development and the error in shower size measurement) from the energy scale established on the basis of hybrid M onte Carlo model (Trzupek et al. [2]) for EAS at sea level as given by

$$
\begin{equation*}
\mathrm{E}_{0}(\mathrm{eV})=3.03 \times 10^{10} \times \mathrm{N}_{\mathrm{e}}^{0.87} \tag{2}
\end{equation*}
$$

42. M easurement of muon density. - The average muon density in a shower as a function of the radial distance from the shower core was estimated for each of the


Fig. 9. - Variation of muon density with muon energy at various radial distances.


Fig. 10. - L ateral muon density distribution at the threshold energy of 2.5 GeV for showers of size range $5.97 \times 10^{4}-9.02 \times 10^{5}$ of $\mathrm{s}=1.2$.
various shower groups in the shower size range $3.15 \times 10^{4}-1.79 \times 10^{6}$ particles using the average density defined as

$$
\begin{equation*}
\varrho_{\mu}\left(\geqslant \mathrm{E}_{\mu}, \mathrm{N}_{\mathrm{e}}(\mathrm{~s}), \mathrm{r}\right)=\mathrm{N}_{\mu}\left(\geqslant \mathrm{E}_{\mu}, \mathrm{N}_{\mathrm{e}}(\mathrm{~s}), \mathrm{r}\right) / \mathrm{N}_{\mathrm{t}}\left(\mathrm{~N}_{\mathrm{e}}(\mathrm{~s}), \mathrm{r}\right) \mathrm{A}, \tag{3}
\end{equation*}
$$

where $\mathrm{N} \mu\left(\geqslant \mathrm{E}_{\mu}, \mathrm{N}_{\mathrm{e}}(\mathrm{s}), \mathrm{r}\right)$ is the total number of muons recorded in a particular distance r for a particular shower size $\mathrm{N}_{\mathrm{e}}$ in a certain period of time above a threshold energy $\left(\geqslant E_{\mu}\right),\left(\mathrm{N}_{\mathrm{t}}\left(\mathrm{N}_{\mathrm{e}}(\mathrm{s}), r\right)\right.$ represents the total number of showers of size $\mathrm{N}_{\mathrm{e}}(\mathrm{s})$ at the same distance interval recorded at the same time and $A$ is the effective area of the muon detectors.

In a particular shower size group the muons are divided into groups on the basis of specified threshold energies (indicated as $\mathrm{E}_{\mu}$ ) and then each group is distributed into a number of bins in respect of radial distances from the shower core to determine the average muon density as a function of radial distance and threshold energy.

Some results on muon energy spectra and lateral distributions are given in figs. 9 and 10 .
4.3. The dependence of muon density on shower size and radial distance. - The variation of muon density at fixed radial distances has been studied as a function of shower size at various muon threshold energies by assuming a dependence of the form


Fig. 11. - Variation of muon density at radial distance ranges $8-12 \mathrm{~m}$ and $30-40 \mathrm{~m}$ with shower size at muon threshold ( $\geqslant \mathrm{E}_{\mu}$ ) energy 2.5 GeV .
given by

$$
\begin{equation*}
\varrho_{\mu}\left(\geqslant \mathrm{E}_{\mu}, \mathrm{r}\right) \sim \mathrm{N}_{\mathrm{e}} \beta\left(\mathrm{E}_{\mu}, \mathrm{r}\right) . \tag{4}
\end{equation*}
$$

The fit to the observed data has yielded the results shown in figs. 11,12 and 13 . It is seen that at a fixed muon threshold energy $\beta$ decreases slowly with radial distance for all shower sizes in the range $3.15 \times 10^{4}-1.79 \times 10^{6}$ particles and the trend of variation of $\beta$ with radial distance is similar for all muon threshold energies. All these results together seem to indicate that the muon distribution function does not change with primary energy. The values of the exponent $(\beta)$ obtained by fitting the observed data are given in table l.

The measured lateral distribution of muons in showers of various sizes are fitted to a relation of the form

$$
\begin{equation*}
\varrho_{\mu}\left(\geqslant \mathrm{E}_{\mu}, \mathrm{r}, \mathrm{~N}_{\mathrm{e}}\right) \sim \mathrm{r}^{-\alpha}\left(\geqslant \mathrm{E}_{\mu}\right) . \tag{5}
\end{equation*}
$$

A plot of this function for $N_{e}=4.48 \times 10^{5}$ and $E_{\mu} \geqslant 2.5 \mathrm{GeV}$ is shown in fig. 14 as an example. The value of $\alpha$ derived from this plot is 0.647 . The least-square fitted line obtained from the plot of $\log \varrho_{\mu}$ vs. $\log r$ at fixed $N_{e}$ and $E_{\mu}$ gives the value of $\alpha$. To determine such values of $\alpha$ the shower cores except for the last radial bin (as for the


Fig. 12. - Same as in fig. 11 but for muon threshold $\left(\geqslant \mathrm{E}_{\mu}\right)$ energy 10 GeV .
example $\bar{r}=34.1 \mathrm{~m}$ shown in fig. 14) were taken inside the edges of the array where the efficiency of the array is large. F or the last radial bin, which is also close to the array boundary, obviously the error of $r$ and $N_{e}$ are comparatively large but the overall effect of these errors on determination of $\alpha$ is small. The mean values of $\alpha$ with r.m.s errors for $\mathrm{E}_{\mu} \geqslant 2.5 \mathrm{GeV}$ are given in table II.

Values of $\alpha$ from the fit of the data for $\mathrm{E}_{\mu} \geqslant 2.5 \mathrm{GeV}$ are shown as a function of $\mathrm{N}_{\mathrm{e}}$ in fig. 15. The trend of $\alpha$ vs. $E_{0}$ curve shows that $\alpha$ is a function of energy up to $4.6 \times 10^{15} \mathrm{eV}$ and becomes constant at higher energies. The lateral muon density distribution in a shower initiated by a heavy primary is flatter compared to the corresponding one initiated by a lighter primary at the same energy. In other words, for a lighter primary composition it is expected to have steeper muon lateral distribution and hence a large value of $\alpha$. Therefore, the variation of $\alpha$ with $\mathrm{E}_{0}$ as shown in fig. 15 possibly indicates that the effective primary mass is decreasing with energy increasing from $4.3 \times 10^{14} \mathrm{eV}$ to around $4.6 \times 10^{15} \mathrm{eV}$.

## 5. - Discussion

The method used in the present analysis consists of examining the accurately measured muon densities as a function of shower size and the muon lateral distribution


Fig. 13. - Same as in fig. 11 but for muon threshold ( $\geqslant \mathrm{E}_{\mu}$ ) energy 50 Gev .
(for various muon threshold energies) as a function of shower size. The analysis also includes the results of an estimation of errors to the EAS parameters. The primary energy of an EAS event is not possible to be determined precisely from the observed particle density distribution due to the presence of fluctuations in an EAS development. H owever, the mean primary energy range ( $2.4 \times 10^{14}-8.3 \times 10^{15} \mathrm{eV}$ ) concerned in the present experiment was determined by comparing the observed vertical shower size with the results of hybrid $M$ onte Carlo model for proton primary at sea-level covering this primary energy range. Some representative examples of measured muon energy spectra and radial distributions are presented in figs. 9 and 10 to show that these provide a firm experimental base for comparing with different models of EAS. An analysis of these results given in subsect. 43 leads to the following conclusion.

Table I. - Values of the exponent $\beta$ at different distance ranges for various muon energies.

| M uon energy (GeV) | Value of $\beta$ (for $\mathrm{r}=8-12 \mathrm{~m}$ ) | Value of $\beta$ (for $\mathrm{r}=30-40 \mathrm{~m}$ ) |
| :--- | :--- | :--- |
| 2.5 | $0.692 \pm 0.009$ | $0.653 \pm 0.012$ |
| 10 | $0.699 \pm 0.010$ | $0.655 \pm 0.013$ |
| 50 | $0.661 \pm 0.011$ | $0.587 \pm 0.021$ |



Fig. 14. - Variation of muon density with radial distance for $\mathrm{N}_{\mathrm{e}}=4.48 \times 10^{5}$.

The radial muon density distribution at low and high muon energies does not change with the primary energy.

For low-energy muons, radial muon density distribution steepens in the primary energy range $4.3 \times 10^{14}-4.6 \times 10^{15} \mathrm{eV}$ with radial distance ( $r$ ) indicating that the effective primary mass decreasing between $4.3 \times 10^{14}-4.6 \times 10^{15} \mathrm{eV}$.
Y. K awamura et al. [3] and M. Ichimura et al. [4] have derived similar conclusions about primary mass in the knee energy region from direct observations from their new

Table II. - Values of the exponent $\alpha$ for various shower sizes at muon threshold energy ( $\geqslant \mathrm{E}_{\mu}$ ) 2.5 GeV .

| Shower size | $\alpha \pm \mathrm{d} \alpha$ |
| :--- | :--- |
| $5.97 \times 10^{4}$ | $0.553 \pm 0.009$ |
| $1.09 \times 10^{5}$ | $0.578 \pm 0.010$ |
| $2.21 \times 10^{5}$ | $0.606 \pm 0.012$ |
| $4.48 \times 10^{5}$ | $0.647 \pm 0.014$ |
| $9.02 \times 10^{5}$ | $0.661 \pm 0.016$ |
| $1.79 \times 10^{6}$ | $0.662 \pm 0.016$ |



Fig. 15. - Variation of $\alpha$ with shower size $\left(\mathrm{N}_{\mathrm{e}}\right)$ at the muon threshold ( $\geqslant \mathrm{E}_{\mu}$ ) energy 2.5 GeV .
emulsion chamber experiments. From an analysis of low-energy EAS muons Blake et al. [5] reported similar trend for average primary mass in the primary energy region $6.0 \times 10^{14}-5 \times 10^{15} \mathrm{eV}$.

This work is being continued to improve statistics on this aspect of study using muon lateral distributions at different energies of muons in EAS. The existing array is being expanded to operate at larger shower sizes so that the present method of analysis can be applied with better accuracies on the determination of EAS parameters.

## REFERENCES

[1] Basak D. K., Chakrabortı N., Ghosh B., Goswamı G.C. and Chaudhurı N., Nucl. Instrum. M ethods, 227 (1984) 167.
[2] Trzupek A., Lu Y. and Poirier J., XXIII ICRC, Calgary, 4 (1993) 359.
[3] K awamura Y. et al., Phy. Rev. D, 40 (1989) 729.
[4] Ichimura M. et al., ICRR Report, 287 (1992) 92-95.
[5] Blake P. R. and Tummey S. P., XXIII ICR C, Calgary, 4 (1993) 363.

