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# Angular Distribution of Electrons in Extensive Air Showers

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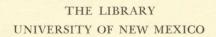
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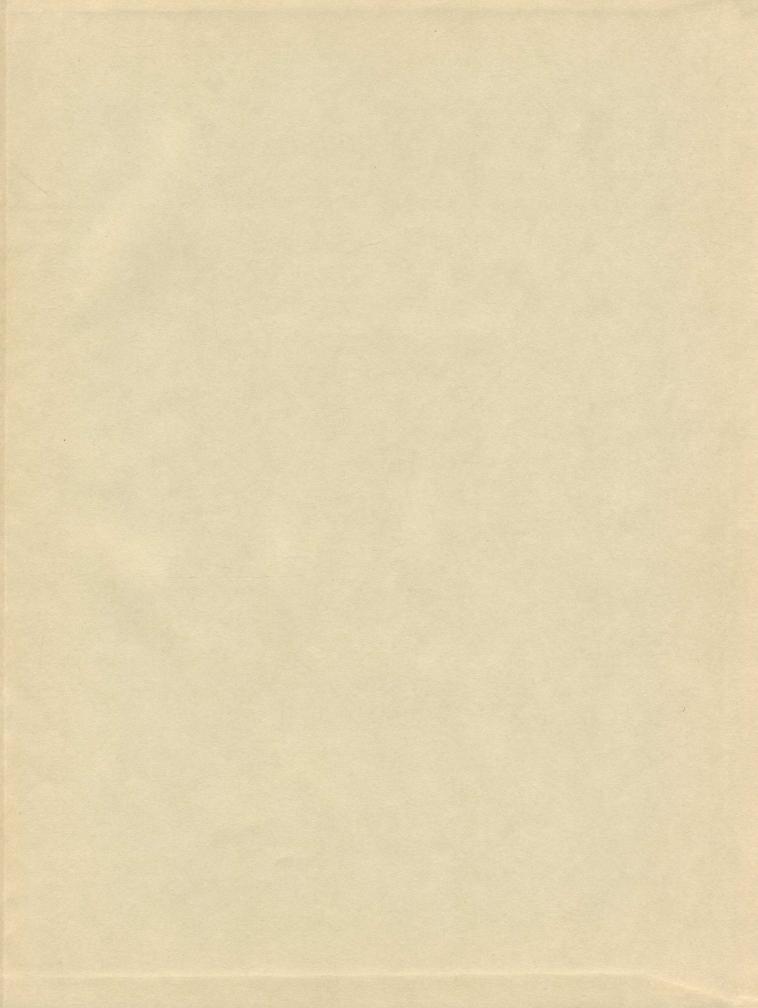


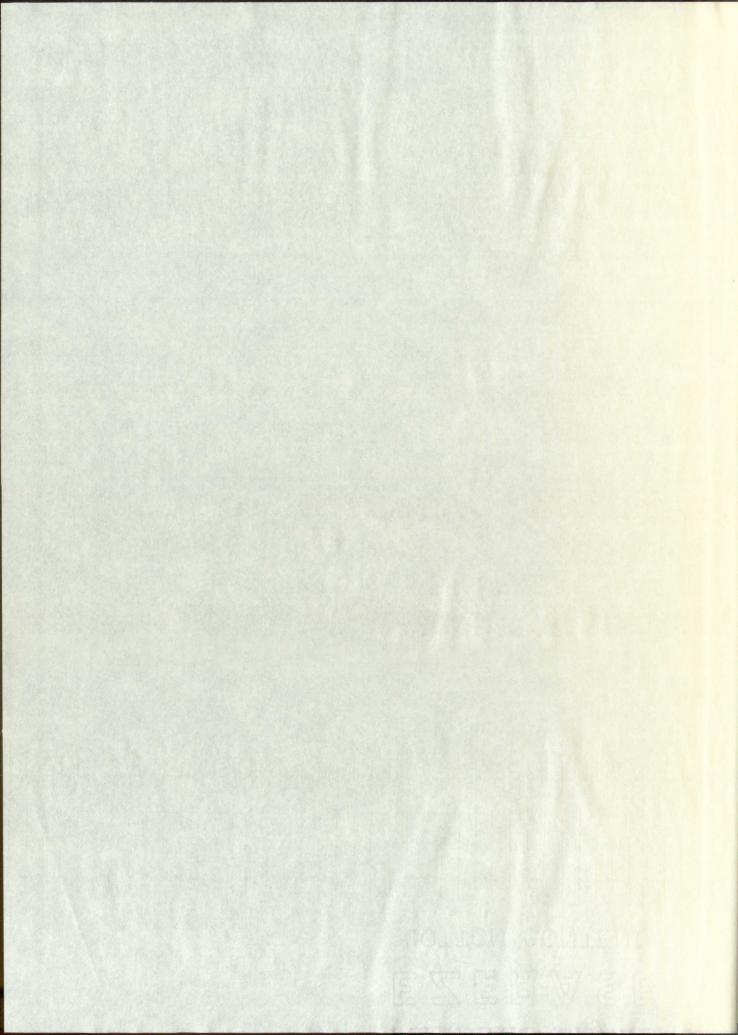
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# ANGULAR DISTRIBUTION OF ELECTRONS IN EXTENSIVE AIR SHOWERS

By
Don D. Laniewski

A Thesis

Submitted in Partial Fulfillment of the Requirements for the Degree of

Master of Science in Physics

The University of New Mexico

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This thesis, directed and approved by the candidate's committee, has been accepted by the Graduate Committee of the University of New Mexico in partial fulfillment of the requirements for the degree of

MASTER OF SCIENCE

Eslastetter

may 30, 1961

# ANGULAR DISTRIBUTION OF ELECTRONS IN EXTENSIVE AIR SHOWERS

by

Don D. Laniewski

Thesis committee

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#### I. INTRODUCTION

The earth's atmosphere is constantly being bombarded by high energy particles from outer space known as "primary cosmic rays". As they penetrate the atmosphere, these particles, which consist of protons, alpha particles and heavier nuclei, lose energy through collisions and give rise to secondary particles consisting of protons, neutrons, and both neutral and charged pi-mesons. The neutral pi-mesons almost immediately disintegrate into photons which then multiply into electromagnetic cascades. Charged pi-mesons disintegrate into mu-mesons and neutrinos. Mu-mesons interact with nuclei weakly; therefore they lose energy by ionization until they decay into electrons and neutrinos. Thus in the atmosphere one may find protons, neutrons, pi-mesons, mu-mesons, electrons, and photons. Electrons are the main products of an extensive air shower and comprise approximately 90% of the shower. Since these electrons are the most numerous particles near the axis of a large shower and are easily detected, extensive air showers are most easily studied by means of the electronic component. Thus, whatever concern there may be about the nucleonic cascade of which the electron is an outgrowth, it is necessary first to understand the development of the electronic component and its lateral and angular distribution. Also, for the correct analysis of any experimental data, it is of particular importance to know something about the

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angular distribution of the particles at a prescribed distance from the axis of the shower. These effects are quite often concealed because of the zenith-angle distribution.

These problems are quite complicated mathematically, 2,3 and it was the purpose of this experiment to gain some information on these problems.

Since the actual tracks left by the electronic component can be photographed and analyzed, a cloud chamber was used at the principal piece of equipment in this experiment. The results of this investigation furnished the basis of this paper. Details of the experiment, analysis of data, interpretation of results, and comparison of these results with those of other experiments and theories are treated in subsequent parts of the paper.

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#### II. APPARATUS

#### 1). Description of Chamber

The cloud chamber, Plate I, designed and built by John R. Green, was operated at the rear of the Physics Department of the University of New Mexico during the experiment. It is constructed of brass with the outside dimensions being 12 3/4 inches high, 12 inches wide, and 16 inches deep. The cloud chamber is a constant-pressure change type, utilizing the release of excess pressure to cool the chamber gas sufficiently to form the saturated vapor necessary to form droplets.

The chamber was operated at an expanded pressure of approximately 1.15 atmospheres of argon with sufficient alcohol to give tracks with an expansion ratio of 1.13.

The front and both sides were covered with \( \frac{1}{6} \) inch glass for viewing and observation. The center of the chamber contained \( \frac{1}{6} \) inch of lead plate.

## 2). Operation

The cloud chamber was used in conjunction with other experiments being performed in connection with extensive air showers. Triggering of the cloud chamber will be described in more detail later in the paper. Simply, a pulse was used to trigger a thyratron, the discharge of which was used to energize an electromagnet which in turn released an arm that normally maintained an air-tight condition between the rear

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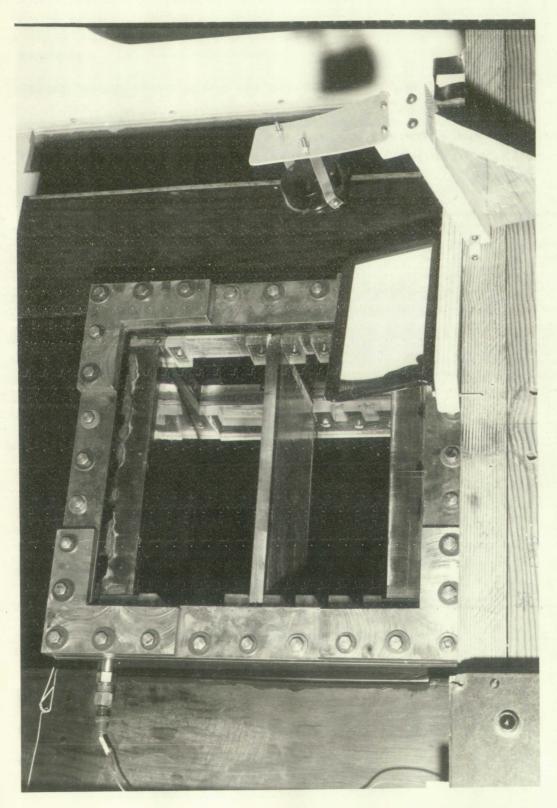


Plate 1.



trap door and a rubber gasket. Greater detail can be obtained in Reference 4. The chamber was then allowed to expand rapidly. After the expansion had taken place, and a photograph taken, the chamber was reset automatically for the next expansion. This resetting was operated by a motor which was started by the expansion. The operations that followed were simply:

- 1). The arm of the magnet was brought back into position.
- 2). The chamber was expanded twice at a very slow rate. Such slow expansions were utilized for clearing the chamber of small droplets that did not evaporate readily.
- 3). The film of the photographing camers was advanced. During the recycling of the chamber, the chamber was insensitive to rapid expansions. This period lasted approximately two minutes.

## 3). Photography

The chamber was photographed by means of a stereoscopic camera placed four feet from the front of the chamber. To eliminate the need for fast shutter synchronization, the camera had no shutter and was open constantly; therefore it was necessary to keep the camera and chamber in a dark enclosure by means of a "cloth-tent".

The film used was Kodak Linograph Ortho 35 mm. film

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The film ower was hours of the colline to be come the

and was processed normally. Illumination was provided by photoflash tubes at both sides of the chamber.

#### 4). Location

The cloud chamber was operated at the rear of the Physics Department at the University of New Mexico, Albuquerque, New Mexico (elevation 1575 meters). It was situated approximately 8 feet from the main scintillator tank and housed in a building attached to the scintillator structure. See Plate II. The roof of the housing was made of 3/4 inch exterior plywood. The housing was kept at a constant temperature (approximately 23±2°C.) by sufficient insulation and the use of heaters and a refrigeration unit.

## 5). Geometry of the Experiment

The experimental set-up consisted of five liquid scintillation detectors. Two of the scintillators that are 10 feet in diameter are placed to form a vertical telescope.

The same structure that housed these scintillators also housed the cloud chamber as discussed earlier. The three other scintillators are three feet in diameter and are arranged in the form of an isosceles triangle with the larger scintillators. Figure 1 shows the arrangement.

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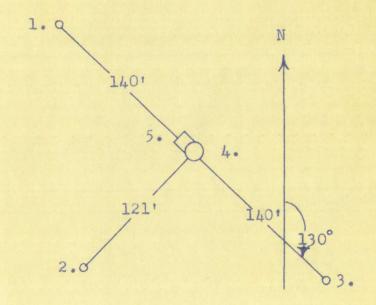
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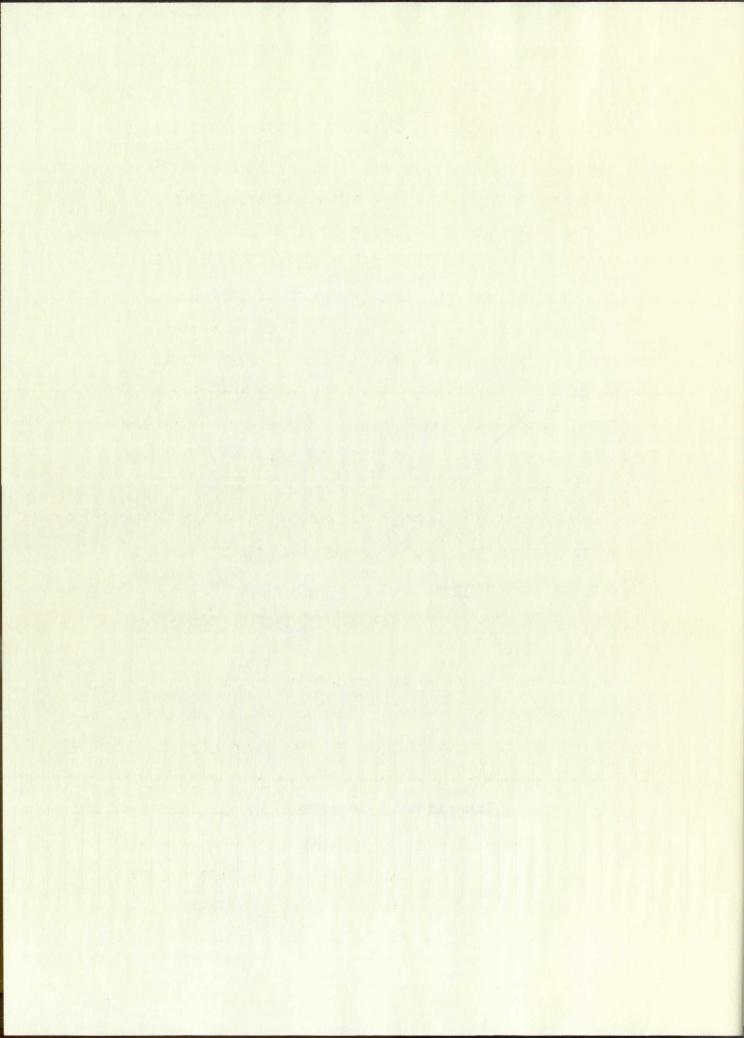
Plate 2.



1,2, and 3. 3-ft. Scintillators
4. Vertical Telescope
5. Cloud Chamber



Experimental Geometry



#### III. EXPERIMENTAL DATA

### 1). General

means of pulses from the upper scintillator of the vertical telescope. The output pulse was discriminated so that a triggering pulse was supplied to the cloud chamber only when the particle density was greater than approximately 100 particles per square meter at the location of the large scintillator. This corresponds to a minimum shower size of approximately 4x10 particles. The minimum size for which the entire array, including the three smaller scintillators, was sensitive was  $6x10^5$  particles. Therefore, the events selected by the single scintillator are representative, on the average, of a smaller size shower than that of the main array.

Each time the chamber was triggered and expanded, a photograph of the event was taken.

At the same time, as an independent experiment by another student, the density of the shower, as observed in each of the scintillating tanks, was automatically recorded by an I.B.M. machine. These data furnished the basis for the determination of the core location and size of the shower. The author of this paper utilized these core locations in part of the analysis of his own experiment.

Of approximately four hundred pictures taken, only sixty were found to be useful in the part of the experiment

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having to do with core locations. A large percentage of the unused pictures consisted of showers for which it was not possible to determine the location of the core. This was brought about by the fact that the tanks that determined the location of the core were not run in four-fold coincidence with the scintillator that controlled the cloud chamber.

Most of the unusable pictures were due to periods of operation of an improperly expanded chamber.

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#### IV. DATA AND ANALYSIS

#### Part I

### 1). General

After the film was processed, those events for which the entire array was sensitive were sorted out. Enlargements were made on 8x10 inch projection paper and then analyzed. The following data were extracted from the enlargements:

- 1). The angles measured on the enlargements between the zenith and the direction of the incident particle in the upper half of the chamber.

  Measurements were not made in the lower portion of the chamber because of the scattering in the lead.
- 2). The same angles measured as in 1)., but only for those incident particles for which a minimum multiplicity of three could be observed below the ½ inch of lead.
- 3). The general direction of the particles, that is, whether they appeared to be incident from the left or right side of the chamber.
- 2). Angular Distribution as a Function of Distance from the Core.

The angles measured from the vertical for all the incident particles were summed and averaged for each event.

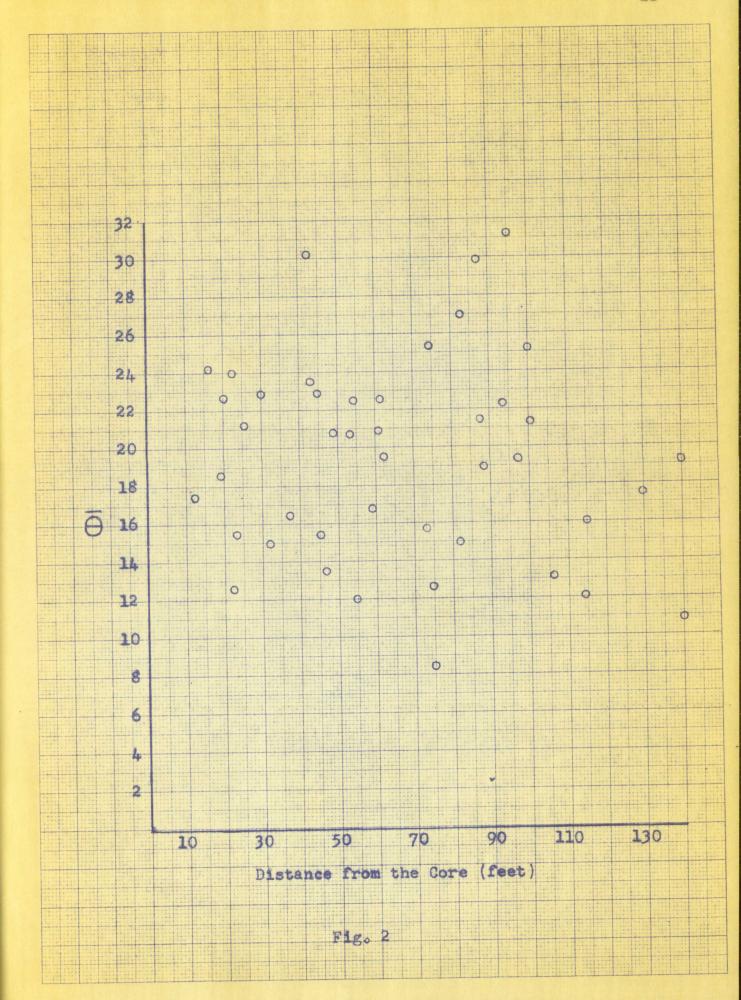
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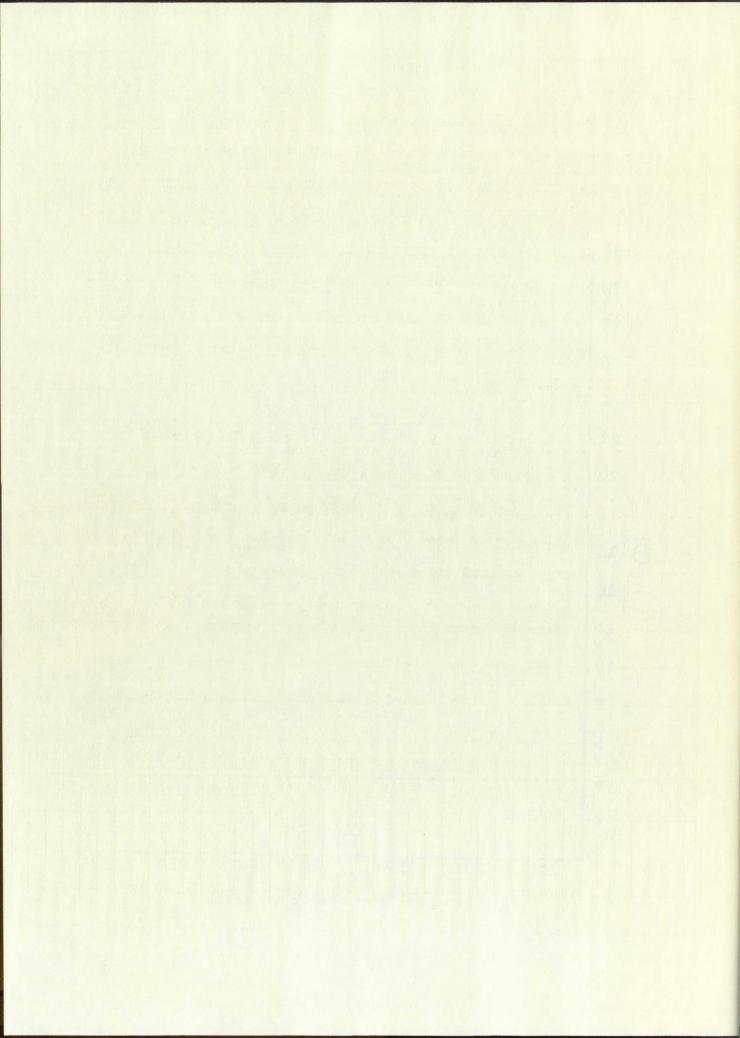
These averaged angles were plotted as a function of distance from the point of impact of the shower. The results are shown in the following table and graphically in Figure 2.

R (feet)	0 (degrees)	R (feet)	e (degrees)
12	17.4	61.	20.9
16	24.2	62	19.5
19	18.6	73	15.8
20	22.7	74	25.3
22	12.5	75	8.4
23	15.4	75	12.7
25	21.1	82	15.2
30	22.9	82	27.0
32	15.0	87	21.5
37	16.5	88	18.9
42	30.1	93	27.3
43	23.6	95	31.2
45	15.4	97	19.4
45	22.8	100	25.2
47	13.4	100	13.1
49	20.8	106	15.8
53	20.8	115	12.0
54	22.5	130	17.5
55	12.0	140	11.1
59 61	22.6	140	19.2

Calculations by Galbraith show that the scattering angle of electrons is inversely proportional to the total energy of the electrons; that is, the highest energy particles are least scattered and are found nearest the core. As seen from the graph in Figure 2, however, the observed average projected angle is not a function of the distance from the core. In fact, there appears to be no correlation between the distance and the angles measured.

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According to calculations of Roberg and Nordheim, a low energy electrons diffuse in random directions. Their spectrum show about 1/6 of all electrons in the shower are composed of these low energy electrons. An example of such low energy electrons can be seen in Plate III. Since such low energy particles were included in obtaining the results shown in Figure 2, the distribution seems to indicate that low energy electrons are of considerable importance at all radial distances.

# 3). Angular Distribution of High Energy Electrons as a Function of Distance From the Core.

electrons, the photographs were re-analyzed, and only those tracks that resulted in a minimum multiplicity of three particles upon interaction with a 1/2 inch lead plate in the center of the chamber were considered. This was done to set some criterion for the energy of the incident particles for which the angular distributions would be determined. From the calculations of Ivanenko, such a multiplicity would require the initiating particle to have a minimum energy of approximately 165 Mev. Typical of these is the event shown in Plate IV. As was done previously, the projected angles were measured and averaged, and were plotted as a function of the distance of impact of the shower.

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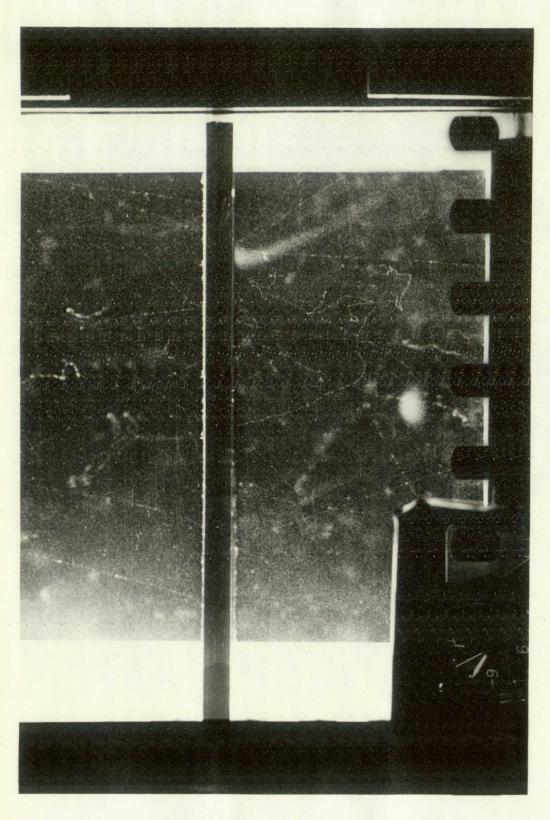


Plate 3.



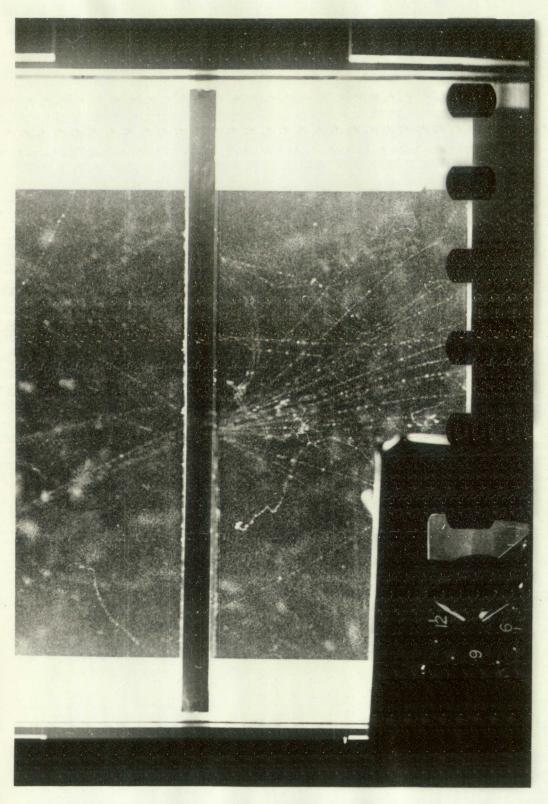
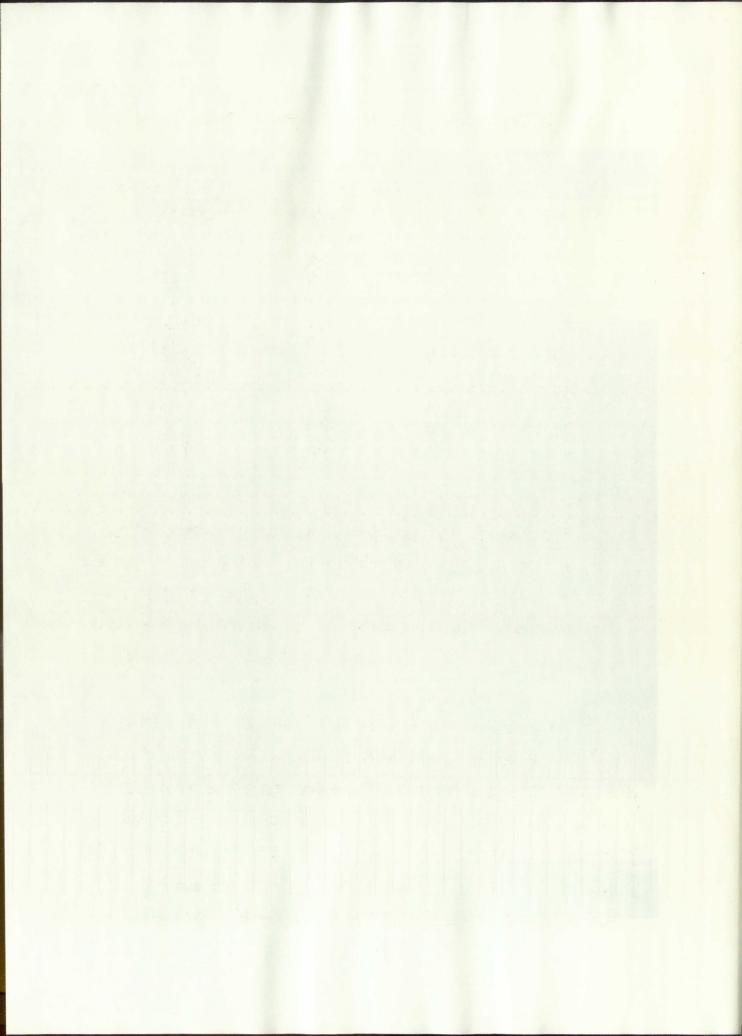


Plate 4.



The distribution is shown in the following table and graphically in Figure 3.

R (feet)	⊕ (degrees)	R (feet)	e (degrees)
12 16 22 25 30 32 42 45 45 47 49 53	14.3 24.0 9.3 10.0 12.7 5.0 11.0 26.0 3.0 20.6 10.3 11.2 11.6	61 62 75 82 82 88 93 97 100 100 106 115 140	15.0 16.0 2.0 6.0 25.0 22.5 3.0 12.5 4.5 4.0 14.0 3.5 3.6

Even though the low energy electrons (which would have given large angle scattering) were eliminated, the mean angle distribution is still found to be poorly correlated with the distance from the shower. As is evident from the graph, however, the average mean angle has dropped considerably.

Ivanenko showed a definite increase in the mean scattering angle as a function of distance from the core. Comparison with his results, however, would require the mean projected angle from the core direction as a function of the energy of the particles at a given distance from the core of the shower. The core direction could not be included in the analysis by the author since this determination

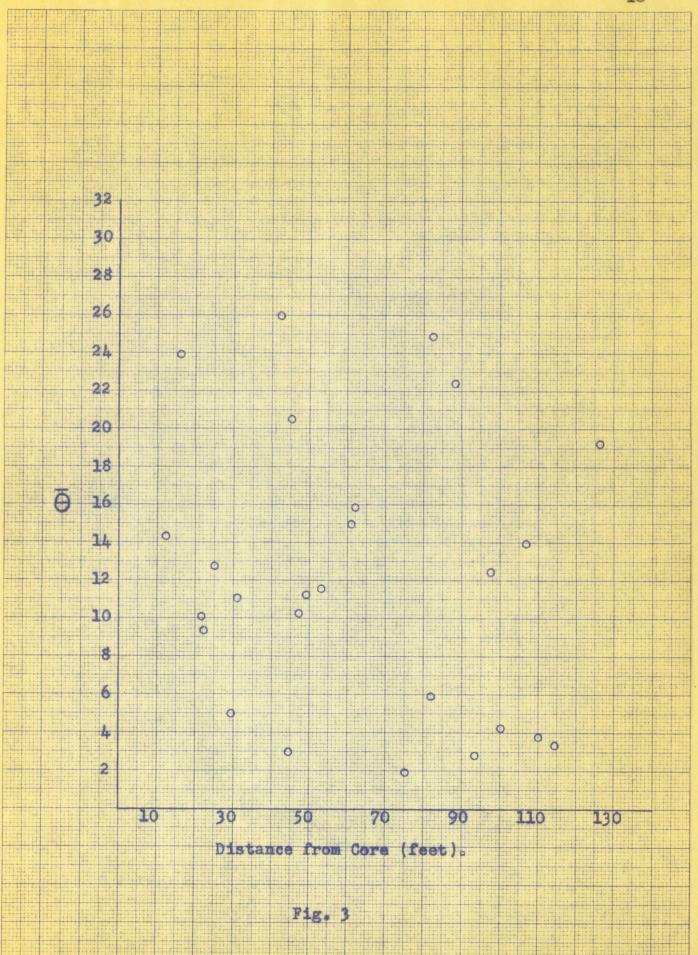
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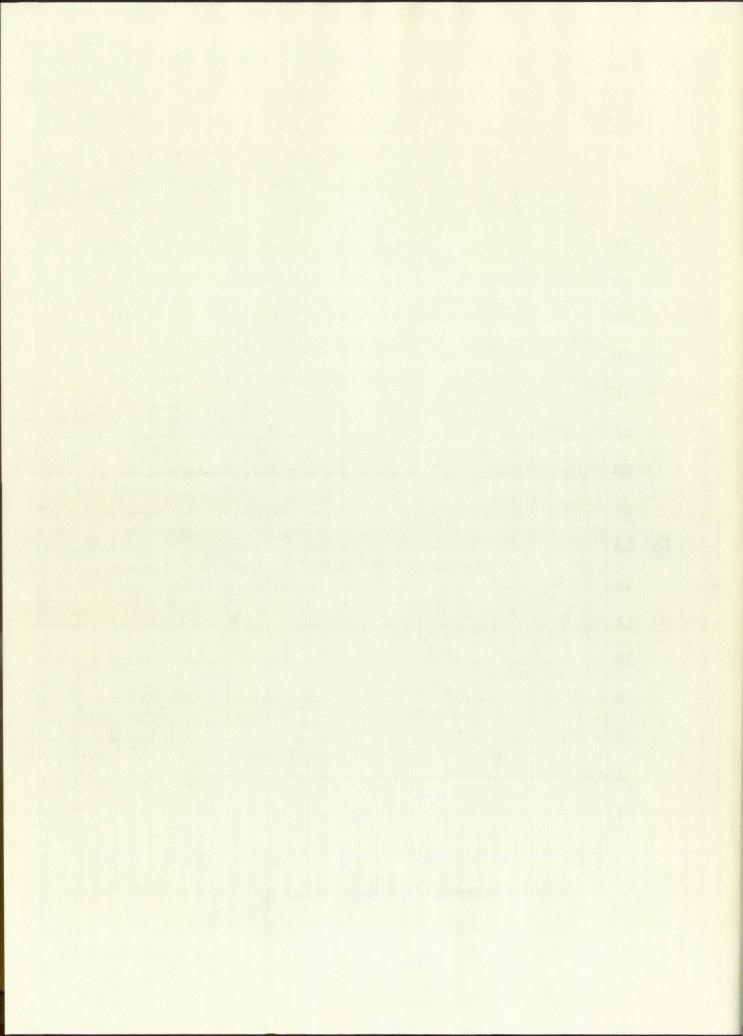
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was not possible with the experimental set-up. Since the results in Figure 3 do not show the distribution as predicted by Ivanenko, it was deemed necessary to investigate the situation further.

A distribution was made utilizing the same values of  $\bar{\theta}$  and the core distance but also utilizing the direction of the core as observed from the photographs. The photographs of the incident particles were classified as either positive or negative, depending upon whether the particles appeared to arrive from the northern or southern directions respectively. Because of the location of the camera and two-dimensional photography, it was impossible to determine which showers arrived from the east-west directions. Reference was made to the azimuth angle of the point of impact of the shower to investigate the possible influence of the showers arriving with large angles to the zenith. The results are shown in the following table.

with Azimuth Angles 0 to 1800	No. of Particles with Azimuth Angles 180 to 360°
3	5
9	9
	with Azimuth Angles 0 to 1800

If all the showers arrive from the zenith, then those photographs classified as positive should correspond to cores

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whose locations were characterized by the azimuth angles 0 to 180°. Similarly, all photographs classified negative would have core locations characterized by azimuth angles 180 to 360°. Obviously, one can see from the table that this is not the case. The fraction of showers arriving with appreciable zenith angles is certainly not negligible. The large number of events in the negative direction can be attributed to the geometry of the experiment which favored showers whose cores fell on that side. It is apparent from the distribution that the observed angular distributions are strongly influenced by the zenith angle distribution.

The results of this analysis make it apparent that one cannot determine the angular distribution of the particles with regard to the direction of the shower axis as a function of the distance from the axis of the shower without a simultaneous knowledge of the zenith angle.

Because of the large influence of the zenith angle distribution on the observed results, it was determined to analyze the data to determine an approximate zenith angle distribution for the shower sxis.

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### 1). General

The change of shower counting rate with varying altitude, varying barometric pressure at a fixed elevation, and varying zenith angle are closely related since they all result primarily from the growth and decay of the electronic cascades as they traverse increasing amounts of matter. Interest in these phenomena persists because the development of extensive air showers is influenced not only by the properties of pure electronic cascades, but also by the shape of the primary cosmic ray energy spectrum and the nature of the process by which the primary energy is transferred to the soft component, both at the origin of the shower and throughout the associated nuclear cascade. For this paper we will be concerned with the zenith angle distribution.

There are several methods of measuring the zenith angle distribution. One method of measuring the distribution is to compare the coincidence rates obtained with counters whose axes are vertical and horizontal. One analy developed method is to measure the difference in time of response of counters that are separated from each other in a horizontal plane. This method can yield the space angle distribution as well as the distribution of projected angles. An analagous method is to measure the angles of the tracks of all shower particles seen in a cloud chamber. These experiments are undesirable since they yield the angular distribution,

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not of the shower axis, but of the shower particles, many of which have been badly scattered. The method employed for this paper was to consider only those cloud chamber pictures that showed high energy electrons of nearly parallel tracks of which the average direction could be well determined. Since information in coincidence from all five tanks is no longer required, all photographs obtained were now analyzed. It is known that the response of a single scintillator, such as triggered the chamber, favors showers whose cores strike close to the apparatus, 13 the average distance being approximately fifteen to twenty meters. At this distance from the core the average energy of the particles is still several hundred Mev. and the angular deviation from the core direction only a few degrees. Thus, this method should result in a fair approximation to the zenith angle distribution and should not be strongly influenced by the angular deviation of the particles from the direction of the axis of the shower. A typical event used is shown in Plate V.

### 2). Zenith-Angle Distribution

Let us assume that the zenith-angle distribution takes the form 14)

where A is the normalizing constant equal to 
$$\frac{n+1}{2\pi}$$
, so that
$$\int_{A}^{2\pi} \frac{\pi}{2\pi} d\theta d\theta d\theta = 1 \qquad (2)$$

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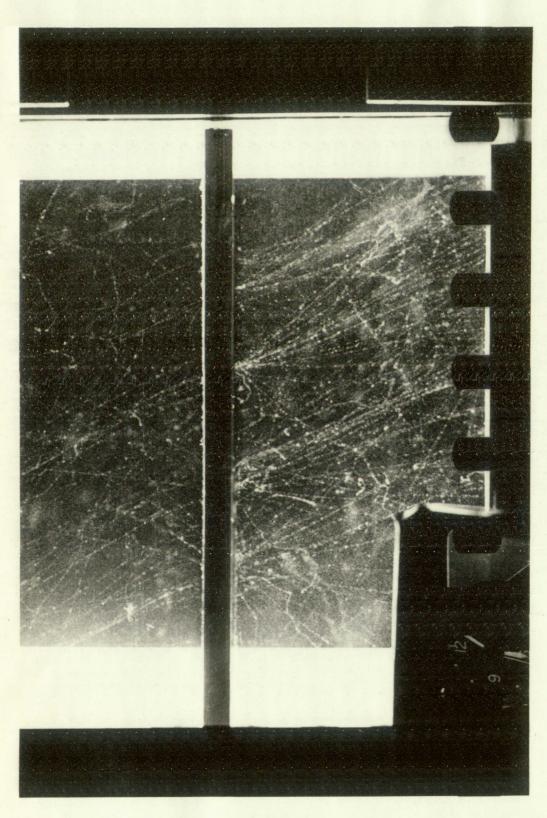
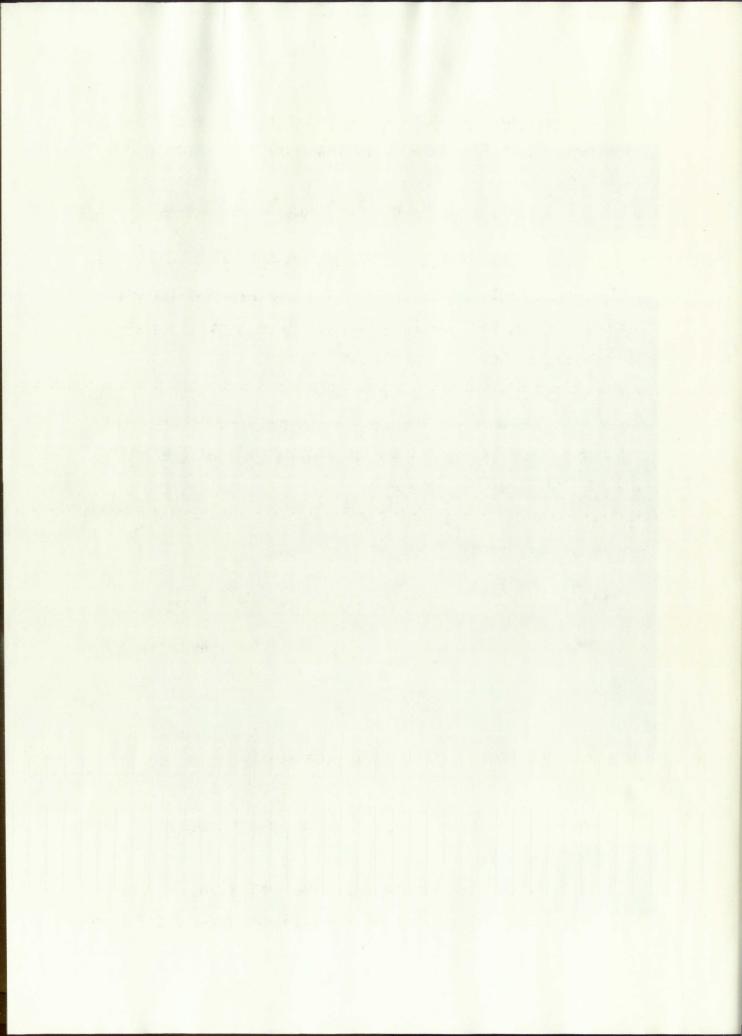


Plate 5.



This zenith angle distribution gives the probability per unit solid angle for showers arriving at the zenith angle  $\theta$ . Therefore

$$P(\theta, \phi)$$
 sine de  $d\phi = \frac{n+1}{2\pi}$  cos e sine de  $d\phi$  (3)

Since these distributions are computed for space-angle distributions, a relationship must be developed for the distribution of projected angles which are used in this paper. From Figure 4 the relation is seen to be

where  $\theta_p$  is the projected angle on the plane of observation and  $\theta$  and  $\phi$  are the space zenith and azimuth angles respectively. Squaring, we obtain

Inserting this in the integral, we obtain

$$\frac{2}{\tan \theta_p} = \iint_{0}^{\frac{\pi}{2}} \frac{n+1}{2\pi} \cos \theta \sin \theta \tan \theta \cos \phi d\theta d\phi$$

The integral can be evaluated in a straightforward manner to give

$$\tan \theta_{p} = \frac{1}{n-1} \tag{5}$$

From the above function, we can find the value of n.

# 3). Experimental Determination of the Zenith Angle Distribution.

To determine the zenith angle distribution, it was necessary to review all of the photographs taken during the

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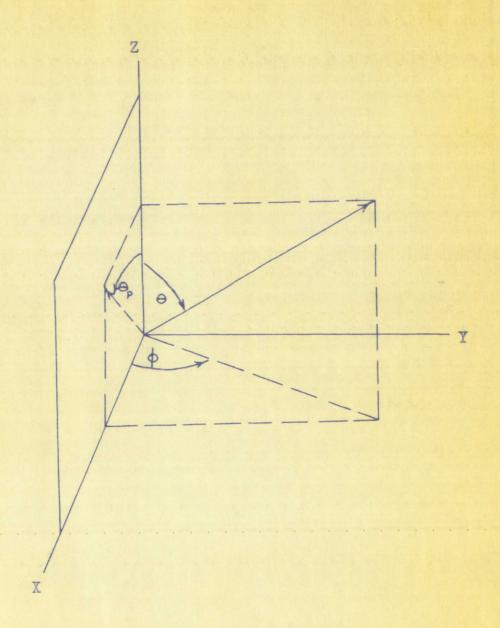
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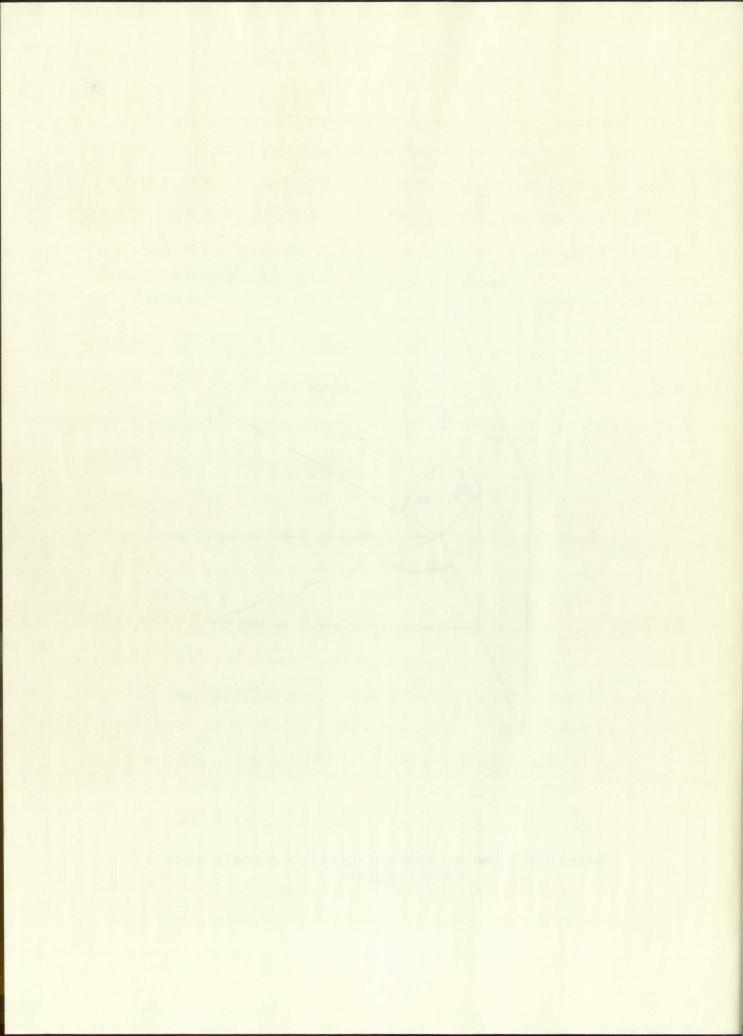
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Relation Between Observed Projected Angle And Space Angles



experiment. This was done by projecting each negative individually and selecting only those photographs that showed high energy electrons or a minimum of three nearly parallel tracks. The angles were measured from the vertical and averaged for each event. The tangent of each average angle was then squared and added. The mean square tangent was found to be

 $\tan \theta_{p} = 0.186$ 

or = 23.33°

Employing equation (1), we obtain n = 6.37.

## 4). Comparison with other Experiments.

The value of n obtained agrees favorably with other experiments performed to determine zenith angle distributions. Kraybill, ousing the cos e distribution, determined the mean zenith angle for mountain elevations to be approximately 20°. Other work by Daudin, williams, brown and McKay, suring the same distribution, determined the mean zenith angle to be 30°, 25°, and 30° respectively, at various mountain altitudes.

One large deviation came from the work of Guseva,
Zatsepsin, and Khristiansen in the U.S.S.R. at approximately 3860 meters. Using a cloud chamber and a hodoscope
arrangement, they found the zenith angle to be very steep

experiment. This and aslessing reagances has a trace to the second showed in the second second second showed in the second secon

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4). Comparison Fire Ward Tracket . (4

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with a value of n ranging from 7.5 to 13.5 with the most probable value being 11.

Experiments performed at Cornell University by K. Greisen 15 show that the angular distribution there more nearly follows the distribution

 $\frac{x}{e^{\frac{x}{\lambda}}}$  (sec6-1)

where  $\theta$  is the zenith angle and x the vertical depth of the atmosphere (at Cornell, 1006 gm./cm.). This function gives a mean zenith angle of  $27^{\circ}$ .

The agreement between the various experiments is satisfactory in all cases except the one in the U.S.S.R.. In
this case the discrepancy is strong and has not been explained;
but the error quoted is large and the distribution somewhat
uncertain. If this discrepancy can be ignored, the remaining
agreement with the other values gives support of the method
used in this experiment as a satisfactory means to study
angular distributions of extensive air showers.

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Dr. Jon B. Croom who at the send the account of the deal thin THE PROPERTY OF THE PARTY OF TH Dr. James F. Product, while the Taylor of the action THE WITH LABOUR PARTY AND

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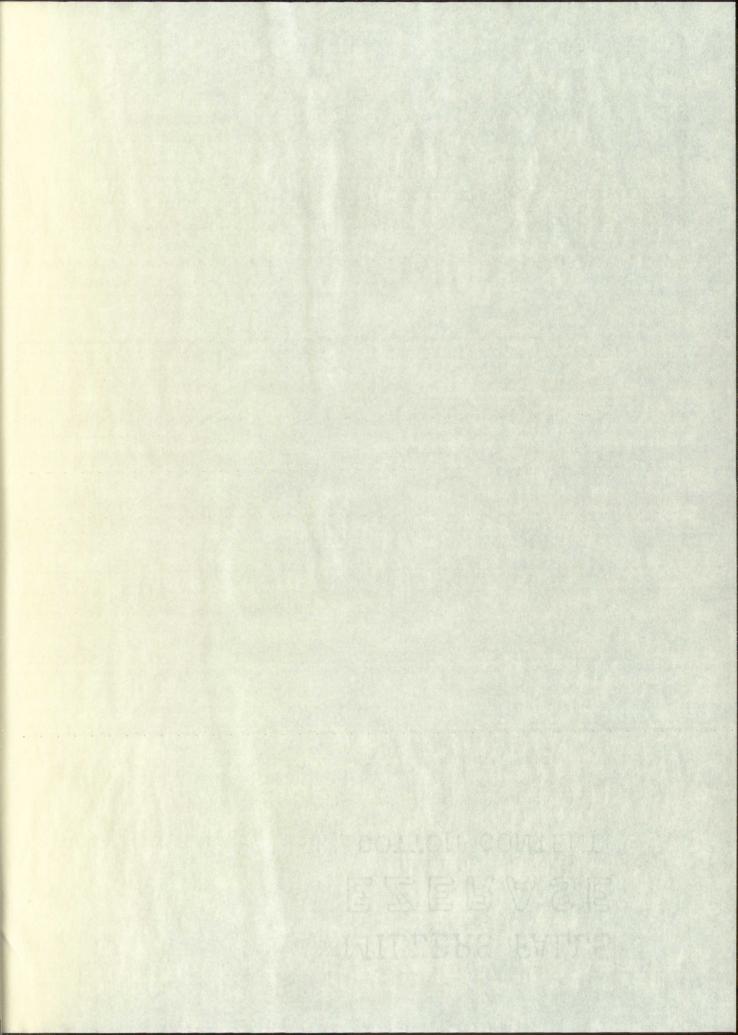
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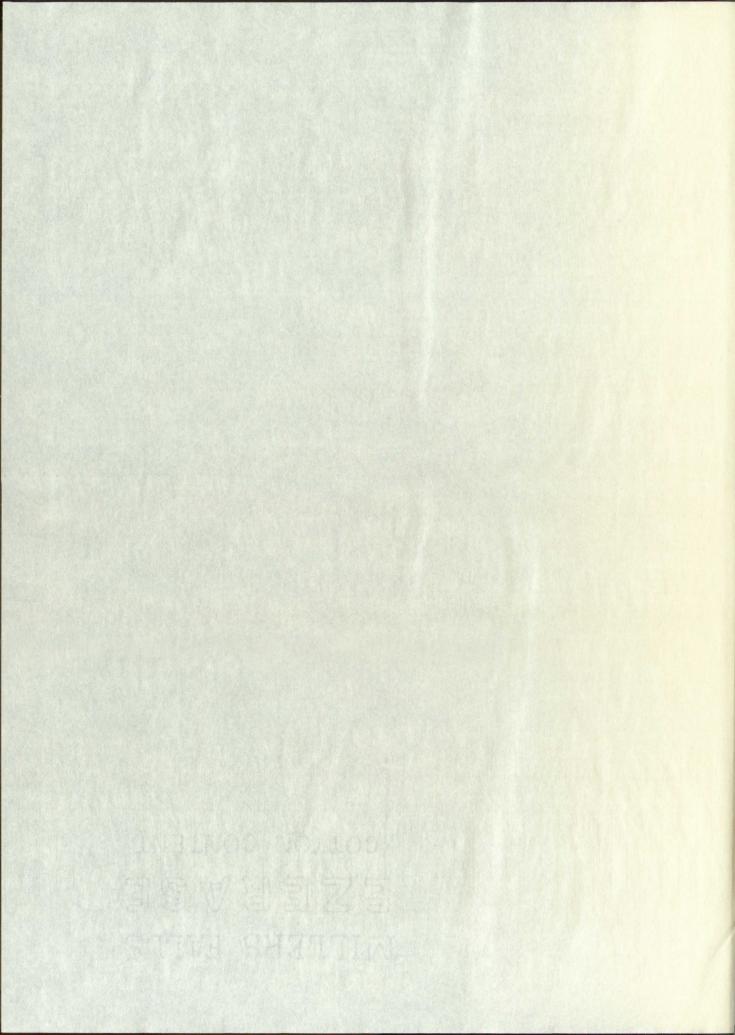
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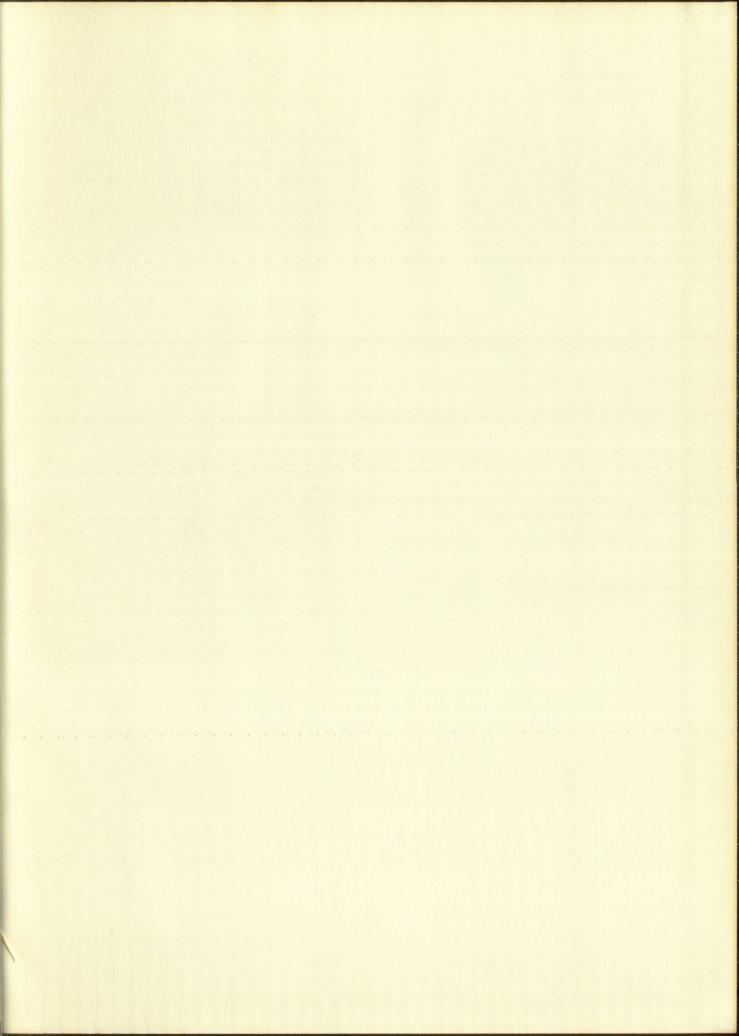
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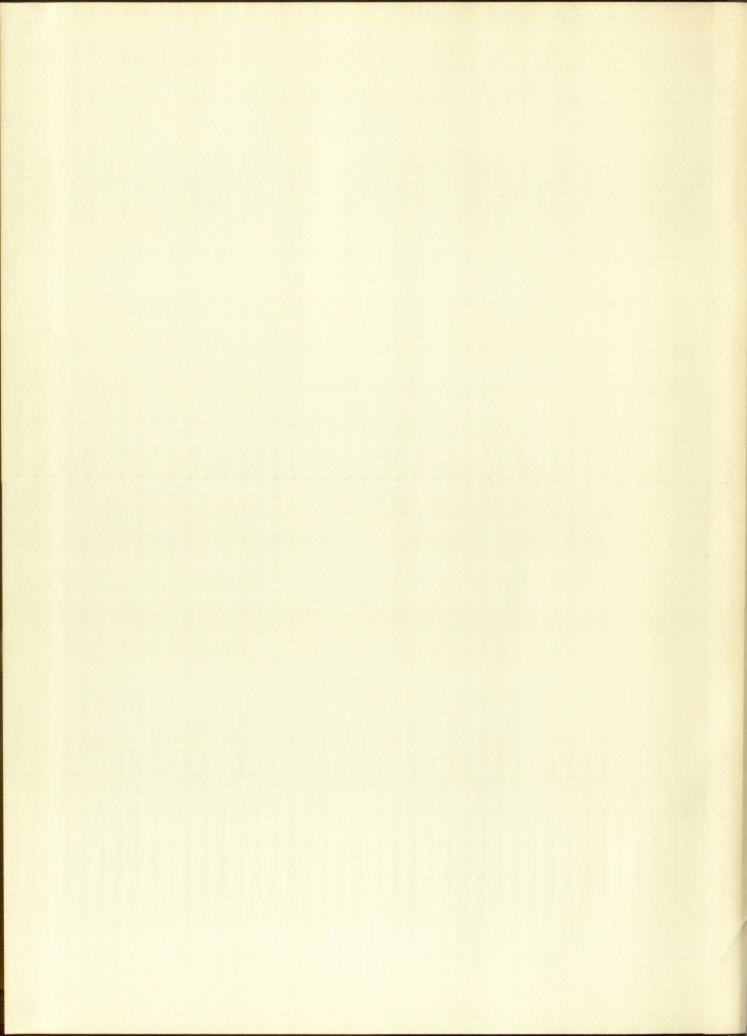
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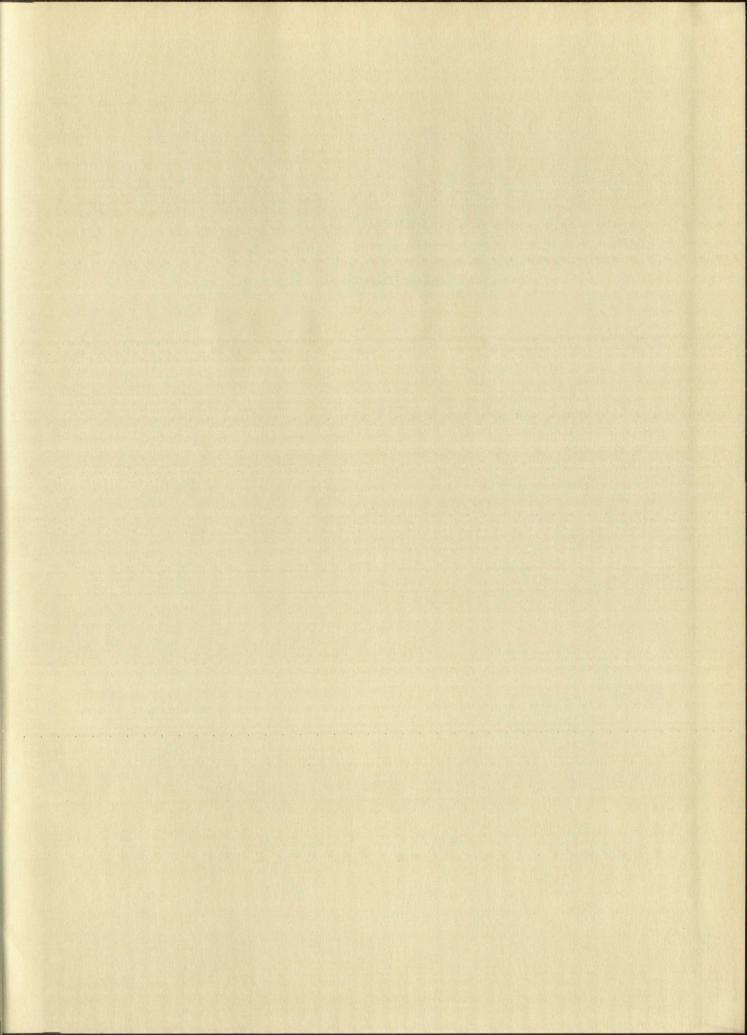
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