# A SEARCH FOR AN INTERMEDIATE-MASS BLACK HOLE IN THE CORE OF THE GLOBULAR CLUSTER NGC 6266 

Bernard J. McNamara ${ }^{1}$, Thomas E. Harrison ${ }^{1}$, Holger Baumgardt ${ }^{2}$, and Pouria Khalaj ${ }^{3}$<br>${ }^{1}$ Department of Astronomy, New Mexico State University, Box 30001, MSC 4500, Las Cruces, NM 88003-8001, USA; bmcnamar@nmsu.edu, tharriso@nmsu.edu<br>${ }^{2}$ School of Mathematics and Physics, University of Queensland, Brisbane, QLD 4068, Australia; h.baumgardt@uq.edu.au<br>${ }^{3}$ Department of Physics, Sharif University of Technology, Tehran, Iran; pooria_khalaj@physics.sharif.edu<br>Received 2011 July 26; accepted 2011 December 15; published 2012 January 16


#### Abstract

It has long been thought that intermediate-mass black holes (IMBHs) might be located in the cores of globular clusters. However, studies attempting to confirm this possibility have been inconclusive. To refine the search for these objects, Baumgardt et al. completed a series of $N$-body simulations to determine the observational properties that a host globular cluster should possess. Keys to revealing the presence of an IMBH were found to be the shape of the cluster's core proper motion dispersion profile and its surface density profile. Among the possible host clusters identified by Baumgardt et al., NGC 6266 was found to be the most suitable object to search. Hubble Space Telescope Wide Field Planetary Camera 2 images with an epoch difference of eight years were, therefore, used to measure this cluster's internal proper motion dispersion profile from $0.8 \operatorname{arcsec}$ to $17 \operatorname{arcsec}$ from the cluster center. This profile and the surface density profiles obtained by Noyola \& Gebhardt and Trager et al. were then compared to those produced by $N$-body simulations of NGC 6266 with and without an IMBH. We find that a centrally located IMBH is not required to match these profiles, but that an IMBH with a $1 \sigma$ upper limit mass of less than a few thousand $M_{\odot}$ cannot be excluded. To establish the existence of this object, the exact location of the density center and more precise velocity measurements within the inner $1 \operatorname{arcsec}$ of this center are required. Our best-fitting model of NGC 6266 without an IMBH yields a cluster mass of $M=8.22 \pm 0.17 \times 10^{5} M_{\odot}$, leading to a mass-to-light ratio of $M / L_{V}=2.05 \pm 0.04$.


Key word: globular clusters: individual (NGC 6266)

## 1. INTRODUCTION

The cores of globular clusters have long been thought to be likely sites for the formation of an intermediate-mass black hole (IMBH). Over 20 years ago Peterson et al. (1989) suggested that the globular cluster M15 possessed a centrally located $1000 M_{\odot}$ IMBH. This claim was supported by the study of Gebhardt et al. (1994) that used surface brightness profiles and radial velocity measurements of 144 stars located within 10 arcsec of the cluster center. Gerssen et al. $(2002,2003)$ subsequently probed the inner core region of M15 using Hubble Space Telescope (HST) by measuring the radial velocities of 64 stars located within $2 \operatorname{arcsec}$ of the cluster center. Although the presence of an IMBH was not required to explain their measurements, they could not rule out such an object if it had a mass of $1.7-3.2 \times 10^{3} M_{\odot}$. Finally, using radial velocity measurements within 1 arcsec of the M15 core, Kiselev et al. (2008) claim that an IMBH is present with a mass between 1 and $9 \times 10^{3} M_{\odot}$.

Other globular clusters have also been suggested to possess IMBHs. Indeed, measurements of the ratios of globular cluster core sizes to their half-mass radii have been used to support the view that a significant fraction of these clusters possess an IMBH (Trenti 2006). Gebhardt et al. $(2002,2005)$ used radial velocity data to argue that a $20,000 M_{\odot}$ IMBH is located in the center of G1, the largest globular cluster in M31. Several other investigators have suggested that a $40,000 M_{\odot}$ IMBH is located in the center of $\omega$ Centauri (Noyola et al. 2008, 2010; Jalali et al. 2012). Recently, Lützgendorf et al. (2011) also found kinematic evidence for a $17 \pm 9 \times 10^{3} M_{\odot}$ IMBH in the center of NGC 6388.

Theoretical models provide several different pathways for the formation of an IMBH within a globular cluster. These
objects could be created by the runaway mergers of massive stars (Portegies Zwart \& McMillan 2002; Portegies Zwart et al. 2004; Taniguchi et al. 2000; O'Leary et al. 2006), or the accretion of significant amounts of interstellar material onto a stellar mass black hole (Kawakatu \& Umemura 2004). Indirectly, support for the existence of IMBHs is also provided by an extrapolation of the $M_{\mathrm{BH}}-M_{\text {Bulge }}$ relation for supermassive black holes in galactic nuclei (Magorrian et al. 1998; Ebisuzaki et al. 2001). This extrapolation leads to the prediction of a typical black hole mass of a few thousand solar masses for globular clusters.

Despite this body of evidence, the existence of IMBHs in globular clusters is still unsettled. Illingworth \& King (1977) showed that the surface brightness profile of M15 and its central velocity dispersion could be explained by a modest population of centrally located neutron stars. Dubath et al. (1994) found a much lower value of $14 \mathrm{~km} \mathrm{~s}^{-1}$ for the M15 central velocity dispersion than the $25 \mathrm{~km} \mathrm{~s}^{-1}$ reported by Peterson et al. (1989), reducing the need for an IMBH. Additional radial velocity data obtained by Dull et al. (1997) further reduced the core velocity dispersion by about $1.4 \mathrm{~km} \mathrm{~s}^{-1}$. Using multimass Fokker-Planck models they found that post-collapse models only required the presence of a central population of neutron stars containing about $3 \%$ of the total cluster mass. Subsequently, Baumgardt et al. (2003a) showed that numerical N -body simulations could also explain the M15 brightness profile without an IMBH. Finally, McNamara et al. (2003) measured the M15 internal proper motion dispersion profile from $0^{\prime} .7$ to $15^{\prime \prime} .6$ of the cluster center and found that it did not require the presence of an IMBH. Baumgardt et al. (2003b) also challenged the assertion that G1 contains an IMBH. Based on an N -body simulation, they found it was not necessary to invoke the presence of an IMBH. The mass of the IMBH suspected to be

Table 1
HST WFPC2 First and Second Epoch Images

| Data Set | Date | Exp <br> $(\mathrm{s})$ | Data Set | Date | Exp <br> $(\mathrm{s})$ | Data Set | Date <br> (s) |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| U67E020DR | 2000 Aug 21 | 1 | UA2V010EM | 2008 Apr 6 | 1 | UA2V010GM | 2008 Apr 6 | 30 |
| U67E020ER | 2000 Aug 21 | 1 | UA2V010CM | 2008 Apr 6 | 1 | UA2V0107M | 2008 Apr 6 | 140 |
| U5AQ1701R | 2000 May 4 | 3 | UA2V010DM | 2008 Apr 6 | 1 | UA2V0108M | 2008 Apr 6 | 140 |
| U5AQ1702R | 2000 May 4 | 30 | UA2V0101M | 2008 Apr 6 | 3 | UA2V0109M | 2008 Apr 6 | 140 |
| U5AQ1703R | 2000 May 4 | 30 | UA2V0104M | 2008 Apr 6 | 3 | UA2V010AM | 2008 Apr 6 | 140 |
| U67E0209R | 2000 Aug 21 | 140 | UA2V0102M | 2008 Apr 6 | 3 | UA2V010HM | 2008 Apr 6 | 140 |
| U67E020AR | 2000 Aug 21 | 140 | UA2V0103M | 2008 Apr 6 | 3 | UA2V010IM | 2008 Apr 6 | 140 |
| U67E020BR | 2000 Aug 21 | 140 | UA2V0105M | 2008 Apr 6 | 30 | UA2V010JM | 2008 Apr 6 | 140 |
| U67E020CR | 2000 Aug 21 | 140 | UA2V0106M | 2008 Apr 6 | 30 | UA2V010KM | 2008 Apr 6 | 140 |
| UA2V010BM | 2008 Apr 6 | 1 | UA2V010FM | 2008 Apr 6 | 30 |  |  |  |

present on the core of $\omega$ Centauri has also been significantly revised downward from $40,000 M_{\odot}$ to a $1 \sigma$ upper limit of 12,000 $M_{\odot}$ (Anderson \& van der Marel 2010; van der Marel \& Anderson 2010). Other methods for detecting IMBHs in clusters have also been suggested, like observations of the mass segregation profile of stars in clusters (e.g., Gill et al. 2008), searches for radio and X-ray emission coming from accreting IMBHs (e.g., Maccarone \& Servillat 2008), or IMBH detection through tidal disruption of passing white dwarfs (Rosswog et al. 2008). However, none of these methods has so far provided conclusive evidence for the existence of IMBHs in star clusters.
In this study we use the internal proper motion dispersion profile of NGC 6266 to investigate whether this globular cluster contains an IMBH. We find that it is not necessary to invoke an IMBH to explain its proper motion dispersion and surface density profiles, but that the presence of an IMBH with a mass of less than about $2000 M_{\odot}$ cannot be ruled out.

In the following text, Section 2 presents the rationale for the search for an IMBH in NGC 6266; the observational data and procedure used to measure its proper motion dispersion profile are presented in Sections 3 and 4. Sections 5 and 6 describe the method used to compute the internal proper motion dispersion profile and the $N$-body model simulations. Our assessment of the need for an IMBH in NGC 6266 is discussed in Section 7.

## 2. TARGET SELECTION: THE CHOICE OF NGC 6266

Baumgardt et al. (2005) used $N$-body simulations to determine the observational parameters a globular cluster should have if it contained a centrally located IMBH. They found that candidate clusters would have a central surface brightness profile with a power-law slope between -0.1 and -0.3 and a King concentration factor of about $c \approx 1.5$. Noyola \& Gebhardt (2006) measured these properties for 37 galactic globular clusters using the Wide Field Planetary Camera 2 (WFPC2) of the HST. Based on these values, Baumgardt et al. (2005) identified nine clusters as likely hosts for an IMBH: NGC 5286, 5694, 5824, 6093, 6266, 6388, 6397, 6541, and 6715.

Among these clusters, NGC 5824, NGC 6397, and NGC 6541 have relatively high central concentrations and were, therefore, judged not to be the best objects to search. NGC 6715 has a long half-mass relaxation time and was, therefore, given lower weight. NGC 5286, 5694, 6093, and 6388 have distances of $>10 \mathrm{kpc}$ making it difficult to accurately measure their core proper motion dispersion profiles. In contrast, NGC 6266 is a nearby, massive cluster whose properties fall within the range predicted for globular clusters that might possess an IMBH. It was, therefore, selected as our target object.

## 3. THE HST WFPC2 IMAGES

All of the images used in this study were obtained with the WFPC2 on board the HST. This instrument was in service from 1993 December to 2009 May when it was replaced by the Wide Field Camera 3. The WFPC2 consisted of four $800 \times 800$ pixel cameras located near the $H S T$ focal plane. The field of view of three of these cameras formed an "L" shaped pattern. Their scale was $0.1 \operatorname{arcsec}$ pixel $^{-1}$. The fourth camera was located at the junction of the base and upper portion of the "L," had a scale of $0.046 \operatorname{arcsec}$ pixel $^{-1}$, and a field of view of $34^{\prime \prime} \times 34^{\prime \prime}$. It was referred to as the Planetary Camera or PC. The images used in this study were obtained with the PC. First epoch images, which have NGC 6266 roughly centered in the PC, had already been obtained in 2000 May and August. They used the F555W filter and had exposure times of 1,30 , and 140 s . The approximate visual magnitude range of the stars measured in this study is $\Delta V=14-19$. Second epoch NGC 6266 images were obtained using the PC, with the same filter and exposure times, in 2008 April. On the first and second short exposure images, the cluster location differed by 48 and 31 pixels in $x$ and $y$, respectively. For the deeper exposures, the cluster center positions differed by 102 and 151 pixels in $x$ and $y$. The images used to measure the proper motions are listed in Table 1. In addition to the above images, two blue F439W images (U5AQ1704R and U5AQ1705R) were downloaded from the HST archive. They were used along with F555W images to construct a color-magnitude diagram of our program stars.

## 4. DETERMINATION OF THE PROPER MOTIONS

The NGC 6266 proper motions were determined following the procedure employed by McNamara et al. (2003) in their study of M15 which also utilized WFPC2 images. Stellar positions were measured on the HST images in the manner described by Anderson \& King (1999, 2003). This procedure analyzes images with small offsets to measure the undersampled $H S T /$ PC point-spread function. The accuracy of the derived positions was assessed by comparing positions from images obtained at the same epoch. The internal agreement between these $x$ and $y$ values was about $\pm 0.02$ pixels for unblended, well-exposed stars, but this precision declines for bright and faint stars.

The goal of this study was to measure the cluster's internal proper motion dispersion profile. Therefore, target stars were chosen based on their potential to yield high-quality proper motions. To be included in this study a star had to meet the following criteria: (1) the counts in its central pixel had to exceed the local background level by at least a factor of 10 ,


Figure 1. Plot of the relative proper motions of all 886 stars measured for this study. The $x$ - and $y$-axes have units of arcseconds per century. The cluster concentration at zero motion is evident. Filled circles correspond to stars whose motions were measured on at least three image sets, crosses correspond to motions measured on one or two image sets. To avoid the complication this latter group poses to the internal motion analysis all stars were required to be measured on at least three image sets.
(2) its four brightest pixels had to exceed a preset minimum value, (3) its total count had to exceed a minimum value, (4) none of its pixels could be saturated, and (5) it could not be located within 4 pixels of an image defect or a neighbor that had a pixel count greater than $33 \%$ of the target star's central pixel count. This last criterion was especially restrictive in the crowded field of NGC 6266. Changing this criterion by 1 pixel altered the sample size by about $40 \%$. An additional important factor that affected the sample size was the different location of the cluster on the first and second epoch images. However, these offsets did not affect the inner core profile where the influence of the IMBH would be the strongest. Taken together, the above criteria greatly limit the number of target stars. However, since the eight-year internal motion within NGC 6266 is expected to amount to only 0.06 pixels, only high-quality motions were desired. The dense stellar background within this cluster makes it difficult to precisely measure the positions of faint stars.

Image distortions were removed using higher order terms in the proper motion solutions rather than a distortion map. The usual terms for zero point, scale, and orientation were included as well as quadratic terms in $x$ and $y$, an $x y$ cross term, and a radially dependent magnitude term. Cluster stars served as reference objects. The orientation of the reference frame is that of HST image UA2V0108M. This choice has no bearing on the conclusions of this study. The same general least-squares routine as employed by McNamara et al. (2003) was employed in this investigation.

Proper motions were derived for 886 stars. Figure 1 shows their vector point diagram. A few field stars or poorly measured objects are evident, but the vast majority of stars are cluster members. A color-magnitude diagram of these stars based on HST archive F555W and F439W images of NGC 6266 is shown in Figure 2. To be used in the internal motion analysis


Figure 2. Instrumental color-magnitude diagram of NGC 6266 based on archive HST F814W and F555W images. Solid circles show all 466 NGC 6266 stars used to compute the internal proper motion profile. Stars that were measured on fewer than three image sets are shown as crosses and were not used to compute the profile. The sharp boundary between these two groups at faint magnitudes reflects our desire to target high signal-to-noise stars.
a star's motion was required to have been measured on at least three image sets. Four hundred and sixty-six stars met this criterion, 45 of which were located within about $3 \operatorname{arcsec}$ of the cluster center. Stars meeting these criteria are shown as filled circles in Figure 2. Having at least three independent measurements helped ensure that the proper motion had been accurately measured. Histograms of these proper motions are shown in Figure 3. Gaussian distributions provide satisfactory fits. The measured dispersions are $\sigma_{x}=0.051 \pm 0.001 \operatorname{arcsec}$ per century and $\sigma_{y}=0.048 \pm 0.001$ arcsec per century. Table 2 provides a list of these proper motions. From left to right the columns provide: a sequential target number, the number of image sets on which the star's motion was measured, the position of the star on HST image UA2V0108M, an instrumental visual magnitude and $B-V$ color, and the proper motion and its error in units of arcseconds per century in $x$ and $y$, respectively. The position of the cluster on these images and their different limiting magnitude did not allow the colors of all the program stars to be measured.

## 5. COMPUTATION OF THE INTERNAL PROPER MOTION DISPERSION PROFILE

The internal proper motion dispersion profile was computed using the procedure described by Jones (1970), Meylan \& Mayor (1991), and McNamara et al. (2003). For each radial bin the observed proper motion dispersion in one coordinate, $\sigma_{o}$ is

$$
\begin{equation*}
\sigma_{o}^{2}=\frac{1}{n-1} \sum_{k=1}^{n} \mu_{k}^{2} \tag{1}
\end{equation*}
$$

where $\mu_{k}$ is the proper motion of star $k$ and $n$ is the number of stars in the sample. The mean proper motion of the sample is assumed to be zero. The internal proper motion dispersion is

Table 2
NGC 6266 Proper Motion Data

| Number | Image Sets | $x$ <br> (pixels) | $y$ <br> (pixels) | Inst $V$ | Inst $B-V$ | $\mu_{x}$ | Error | $\mu_{y}$ | Error |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | ---: | :---: |
| 1 | 3 | 63.73 | 493.08 | -11.62 | 4.28 | 0.014 | 0.007 | 0.085 | 0.002 |
| 2 | 5 | 70.65 | 406.58 | -11.04 | 3.23 | 0.002 | 0.035 | -0.012 | 0.040 |
| 3 | 6 | 72.76 | 252.16 | -9.23 | 4.07 | 0.001 | 0.010 | 0.053 | 0.010 |
| 4 | 6 | 73.96 | 275.53 | -8.88 | 4.07 | -0.004 | 0.011 | -0.008 | 0.023 |
| 5 | 3 | 74.76 | 147.16 | -10.90 | 4.17 | -0.095 | 0.007 | 0.017 | 0.005 |
| 6 | 3 | 78.24 | 123.32 | -11.30 | 4.23 | 0.048 | 0.009 | -0.001 | 0.006 |
| 7 | 3 | 81.74 | 665.63 | -11.06 | 4.29 | -0.005 | 0.009 | 0.033 | 0.008 |
| 8 | 3 | 88.37 | 648.60 | -11.09 | 4.28 | -0.055 | 0.012 | -0.003 | 0.011 |
| 9 | 6 | 95.39 | 379.64 | -9.78 | 4.15 | 0.043 | 0.013 | 0.032 | 0.006 |
| 10 | 3 | 96.79 | 188.33 | -12.01 | 4.32 | 0.005 | 0.003 | -0.016 | 0.004 |



Figure 3. Proper motion histograms of the 466 stars used in this study and Gaussian fits to their distributions. The horizontal axis units are arcseconds per century. The dispersions in $x$ and $y$, respectively, were 0.051 and $0.048 \operatorname{arcsec}$ per century. The average motions of the sample in $x$ and $y$ were -0.002 and 0.000 arcsec per century, respectively.
then given by the equation

$$
\begin{equation*}
\sigma_{i}^{2}=\sigma_{o}^{2}-\frac{1}{n} \sum_{k=1}^{n} \zeta_{k}^{2} \tag{2}
\end{equation*}
$$

where $\zeta_{k}$ is the standard deviation of the $k$ th star's proper motion. The error in $\sigma_{i}$ was computed as $\sigma_{i}(1 / 2 N)^{1 / 2}$ (Pryor \& Meylan 1993; Meylan \& Mayor 1991; Trumpler \& Weaver 1953). The proper motion dispersion profile was computed using equally spaced radial bins except for the last bin which extended out to the limit of the HST image. The adopted spacing was largely dictated by the number of stars in the inner bin sample. Our desire was to have the first bin contain $\geqslant 8$ stars so that the error in the rms dispersion would be small enough to provide meaningful information. A bin size of 1.5 arcsec was found to satisfy this need. To test the sensitivity of the measured profile to the adopted cluster center the profile was computed using the center positions determined by Djorgovski \& Meylan (1993), Harris (1996), Beccari et al. (2006), and Noyola \& Gebhardt


Figure 4. From top to bottom the internal proper motion dispersion profile for NGC 6266 based on the cluster center found in Djorgovski \& Meylan (1993), Harris (1996), Beccari et al. (2006), and Noyola \& Gebhardt (2006). The bottom profile is based on an average of the above center locations. The $x$-axis measures distance from the assumed cluster center in arcseconds. The $y$-axis provides the internal proper motion dispersion in arcseconds per century, offset from each other for clarity. The proper motion dispersion profile is essentially flat between $0.7 \operatorname{arcsec}$ and $18 \operatorname{arcsec}$ from the cluster center.
(2006). These locations were within about 40 pixels of each other on a PC image. A fifth profile was computed using the average of these four locations. The resultant profiles are shown in Figure 4. In general, they have the same flat shape indicating that the computed profile is relatively insensitive to the assumed core location. The profile based on the Harris cluster center shows an upturn in the inner core region, but its $1 \sigma$ error bar is consistent with a flat profile. In our subsequent analysis, the profile based on the cluster center determined by Noyola \& Gebhardt (2006) is used. This profile, computed using a slightly smaller bin size of 1.2 arcsec, is provided in Table 3. Profiles were also computed using stellar proper motions determined from at least two and four image sets. They were very similar to those shown in Figure 4 even though the number of stars in a bin could differ by more than a factor of two.

Table 3
NGC 6266 Dispersion Profile Using the Noyola \& Gebhardt (2006) Cluster Center

| Mean $r$ | Number | $\sigma_{x}$ | Error $x$ | $\sigma_{y}$ | Error $y$ | $\sigma$ | Error |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.77 | 8 | 0.060 | 0.015 | 0.048 | 0.012 | 0.053 | 0.009 |
| 1.88 | 18 | 0.050 | 0.008 | 0.045 | 0.008 | 0.047 | 0.006 |
| 3.09 | 19 | 0.035 | 0.006 | 0.054 | 0.009 | 0.041 | 0.005 |
| 4.12 | 32 | 0.046 | 0.006 | 0.044 | 0.006 | 0.045 | 0.004 |
| 5.48 | 42 | 0.054 | 0.006 | 0.048 | 0.005 | 0.050 | 0.004 |
| 6.76 | 32 | 0.044 | 0.006 | 0.054 | 0.007 | 0.048 | 0.004 |
| 7.83 | 47 | 0.047 | 0.005 | 0.047 | 0.005 | 0.047 | 0.003 |
| 8.97 | 52 | 0.049 | 0.005 | 0.040 | 0.004 | 0.044 | 0.003 |
| 10.13 | 34 | 0.047 | 0.006 | 0.046 | 0.006 | 0.046 | 0.004 |
| 11.41 | 38 | 0.054 | 0.006 | 0.049 | 0.006 | 0.051 | 0.004 |
| 12.59 | 36 | 0.055 | 0.006 | 0.048 | 0.006 | 0.051 | 0.004 |
| 13.83 | 30 | 0.049 | 0.006 | 0.046 | 0.006 | 0.047 | 0.004 |
| 16.91 | 78 | 0.046 | 0.004 | 0.041 | 0.003 | 0.043 | 0.002 |

## 6. $N$-BODY MODELING

To test whether NGC 6266 contains an IMBH $N$-body simulations were constructed. All simulations were made with the Graphical Processing Unit (GPU)-enabled version of the collisional $N$-body code NBODY6 (Aarseth 1999) on the GPU cluster of the University of Queensland. The simulations started from King (1962) models with varying initial central concentrations $c$, half-mass radii $r_{h}$, and IMBH masses. Stellar evolution was modeled according to the stellar evolution routines of Hurley et al. (2000) and using a metallicity of $[\mathrm{Fe} / \mathrm{H}]=-1.18$ as given by Harris (1996). All simulations started with $N=50,000$ stars distributed according to a Kroupa (2001) mass function in the mass range $0.1<m / M_{\odot}<100$. The time period covered by the simulations was $T=12$ Gyr. Ten snapshots spaced by 50 Myr intervals between 11.0 and 11.5 Gyr were overlaid and compared with NGC 6266.

The surface density profile was determined from the luminosities and positions of all stars in the combined snapshot file, excluding only the brightest giants. The velocity dispersion was determined from all stars with apparent $V$ magnitudes in the range $14<V<18.5$ at the distance of NGC 6266, similar to the magnitude range of stars for which proper motions were measured. To decrease the statistical noise of the $N$-body data, the infinite projection method of Mashchenko \& Sills (2005) was used when calculating the surface brightness and velocity dispersion profiles.

Since direct $N$-body simulations with the number of stars typical for globular clusters (roughly $2 \times 10^{6}$ in the case of NGC 6266) are currently not possible, smaller clusters were simulated and then scaled to NGC 6266. To perform the scaling, we determined the projected half-light radius for the combined data set of the different snapshots and scaled the model cluster down in radius such that its projected half-light radius was equal to the half-light radius of NGC 6266 as given by Harris (1996), $r_{\mathrm{hp}}=0.92$. This corresponds to 1.82 pc at the distance of NGC 6266. Simultaneously, the mass of the simulated cluster was increased such that the relaxation time remains unchanged and the velocities of the stars were increased correspondingly. A similar method has been used by Baumgardt et al. (2003b) to model G1 and by Jalali et al. (2012) to model $\omega$ Cen and we refer the reader to these papers for a more complete discussion of scaling. To determine the best-fitting model for NGC 6266, a grid of simulations was constructed by varying the concentration parameter $c$ of the King model, the initial half-mass radius $R_{\mathrm{Hi}}$ of the cluster, and the mass of the central IMBH. Five different
values for $c=0.5,1.0,1.5,2.0,2.5$ and five different values for $R_{\mathrm{Hi}}=2.0,3.0,5.0,7.0,10.0 \mathrm{pc}$ were used with either no IMBH or an IMBH with either $1 \%$ or $2 \%$ of the cluster mass at the end of the simulations. For each of the 75 grid points we made three simulations starting from different random number seeds and averaged the results. In order to increase the number of models which we can compare to NGC 6266, we also interpolated linearly between the grid points.

We determined the velocity dispersion and surface brightness profile of each cluster after scaling together with their uncertainties using the infinite projection method of Mashchenko \& Sills (2005). The surface brightness profile was determined from all stars still bound to the cluster. Stars were weighted by their luminosities excluding only very bright stars with $L>10 L_{\odot}$. The velocity dispersion was determined from those stars which fall into the same luminosity range as the observed stars. We then used a $\chi^{2}$ test to find the best-fitting cluster model, taking into account observational errors as well as model uncertainties. For the $\chi^{2}$ test we used the surface brightness profile as determined by Noyola \& Gebhardt (2006) and the velocity dispersion profile for the Noyola \& Gebhardt cluster center as determined in this work. Since the surface brightness profile from Noyola \& Gebhardt only extends out to $100^{\prime \prime}$, the profile given by Trager et al. (1995) was used to increase the radial coverage to $300^{\prime \prime}$.

Figure 5 shows the $\chi^{2}$ maps for the best-fitting models without an IMBH. Shown are $\chi^{2}$ maps for the surface density, the velocity dispersion, and a $\chi^{2}$ map for the combined fit. Good-fitting models for the velocity dispersion alone lie in a relatively small region going from $R_{\mathrm{Hi}}=6 \mathrm{pc}$ for clusters starting from low-concentration models to $R_{\mathrm{Hi}}=2 \mathrm{pc}$ for highconcentration ones. This is due to the dynamical evolution of the clusters and the scaling in half-mass radius which we apply to them at the end of the runs. In order to match the velocity dispersion of NGC 6266, the unscaled half-mass radii of our model clusters have to be around 12 pc at $T=12 \mathrm{Gyr}$. Since core-collapse times decrease with the initial concentration of a cluster, and since clusters expand after core collapse due to heating by binary stars, extended low-concentration clusters and dense high-concentration clusters lead to the best fit of the velocity dispersion. Best-fitting clusters for the surface density follow a different distribution and come mainly from models starting with $R_{\mathrm{Hi}}=4.5 \mathrm{pc}$, so that the best-fitting models for the combined data set come from a relatively small region in half-mass radius versus concentration space.

Our best-fitting model for the combined fit of velocity dispersion and surface density profile has an initial King concentration $c=0.76$ and an initial half-mass radius of $R_{\mathrm{Hi}}=4.56 \mathrm{pc}$ with a $\chi^{2}$ value of 41.2 . The reduced $\chi^{2}$ value of this model is 0.70 , indicating a good fit to the surface density and velocity dispersion profile.

Figure 6 depicts the fit to the observed velocity and surface density profile of the best-fitting no-IMBH model. The surface density profile of NGC 6266 is well fitted by the no-IMBH model except for radii around $10^{\prime \prime}$ where the model is slightly below the observations. The no-IMBH model also fits the outermost regions quite well despite being isolated, indicating that NGC 6266 is not strongly effected by the Galactic tidal field. Figure 6(b) shows that the observed velocity dispersion is essentially flat and that the cluster model without an IMBH provides a good fit to this profile. From our best-fitting model of NGC 6266 without an IMBH we derive a cluster mass of $M=8.22 \pm 0.17 \times 10^{5} M_{\odot}$, leading to a mass-to-light ratio of $M / L_{V}=2.05 \pm 0.04$. The derived cluster mass is about $25 \%$


Figure 5. $\chi^{2}$ maps for (a) surface density, (b) velocity dispersion, and (c) the combined map for clusters without IMBHs. The asterisk marks the best-fitting model in each of the panels and the contour lines indicate lines of constant $\Delta \chi^{2}$ corresponding to $1 \sigma$ to $5 \sigma$. The best-fitting model for the combined data set starts from an initial half-mass radius of $R_{\mathrm{Hi}}=4.56 \mathrm{pc}$ with an initial central concentration of $c=0.76$.


Figure 6. Fits of the best-fitting models without an IMBH and with IMBHs containing $1 \%$ and $2 \%$ of the cluster mass to the velocity dispersion and surface density profile of NGC 6266. The best-fitting model with IMBHs has density cusps in the center and fits the surface density data less well than the no-IMBH model. They also provide worse fits to the velocity dispersion profile due to their central rises in velocity dispersion as a result of the IMBH.
higher than what a Jeans analysis of surface density and velocity dispersion profile would give (approximately $6.5 \times 10^{5} M_{\odot}$ ). This is due to mass segregation, which causes massive mainsequence stars to be centrally concentrated, giving them a smaller than average velocity dispersion. Comparison of our $M / L_{V}$ value with the theoretical values shown in Figure 7 of Mieske et al. (2008) shows that the $M / L_{V}$ value derived here is in good agreement with that expected for $10-12 \mathrm{Gyr}$ old stellar populations following standard, Kroupa (2001) or Chabrier (2003) initial mass functions (IMFs). Our runs without an IMBH finally indicate that NGC 6266 is on its way to core collapse, but has not yet reached this state.

If we make a simultaneous search in a three-dimensional (3D) grid with King concentration $c$, half-mass radius $R_{\mathrm{Hi}}$, and IMBH mass on the axes, we find that the best fit for a model with an IMBH is the same as for the no-IMBH case, i.e., the model starting with $c=0.76$ and an initial half-mass radius of $R_{\mathrm{Hi}}=4.56 \mathrm{pc}$ and no IMBH fits better than any model with an IMBH. The reason for the worse fit of the IMBH models is shown in Figure 6. Models with IMBHs are significantly below the observed surface density profile at radii of a few arcsec. They also provide a less good fit to the velocity dispersion profile since the measured velocity dispersion profile is essentially flat in the center while the IMBH models predict a central rise. This is especially a problem for models with more massive IMBHs.

We note that the best-fitting IMBH models have weak cusps in the central surface density. The central slope measured by

Noyola \& Gebhardt (2006) for NGC 6266 is $-0.11 \pm 0.08$. This value is near the lower limit of the slope that $N$-body simulations predict would be produced by the presence of an IMBH (Noyola \& Baumgardt 2011; Vesperini \& Trenti 2010). It is also close to the slope for $\omega$ Centauri where a similar search was only able to place a $1 \sigma$ upper limit on a possible IMBH of $12,000 M_{\odot}($ van der Marel \& Anderson 2010). Both of these results are consistent with the analysis of Noyola \& Baumgardt (2011).

The data for NGC 6266, therefore, do not require an IMBH, although we cannot rule out the possibility that a small-mass IMBH might be present in the core of NGC 6266 either. The formal $1 \sigma$ upper limit to this mass is about $400 M_{\odot}$ if the Noyola \& Gebhardt center is used. This corresponds to $M_{\mathrm{IMBH}} / M_{C}<$ $0.048 \%$. To test the sensitivity of this result to the assumed cluster center, models were fitted to all of the profiles shown in Figure 4. The total cluster mass and, therefore, the $M / L$ ratio were found to be essentially unchanged: the $M / L$ value changed by only 0.03 , an amount similar to the error found using the Noyola \& Gebhardt based profile. The reason for this is that most of the information for the cluster mass comes from large radii where the profile does not change significantly if the center is shifted by a small amount. However, the profile computed using the center determined by Beccari et al. (2006) has an upturn in the central region. If this profile is used, the best-fitting model has an IMBH mass of $4000 M_{\odot}$. This result highlights an important point. Limits on the IMBH mass are
very sensitive to the shape of the dispersion profile near the cluster center and this shape can vary depending on the adopted location of the center. In our analysis we have adopted the cluster center determined by Noyola \& Gebardt, but this position may not be correct. The presence of an IMBH with a few thousand solar masses in the center of NGC 6266 can therefore not be excluded with the current data.

## 7. SUMMARY

We have measured the internal proper motion profile of the core region of NGC 6266 and compared it to state-of-the-art $N$-body simulations of this cluster with and without a centrally located IMBH. From our best-fitting model of NGC 6266 without an IMBH we derive a cluster mass of $M_{C}=8.22 \pm 0.17 \times 10^{5} M_{\odot}$. The implied mass-to-light ratio of NGC 6266 is $M / L_{V}=2.05 \pm 0.04$, which is compatible with predictions of stellar evolution models using standard (e.g., Kroupa 2001) mass functions. We find that it is not necessary to invoke the presence of an IMBH to explain the measured proper motion dispersion profile. When coupled with the cluster's surface brightness profiles measured by Noyola \& Gebhardt (2006) and Trager et al. (1995), the $1 \sigma$ limit on the IMBH mass is about $400 M_{\odot}$. Most other density centers also have no-IMBH models as best-fitting models; however the velocity dispersion profile obtained when using the Beccari et al. (2006) density center leads to a best-fitting model with an IMBH mass of $4000 M_{\odot}$. IMBHs of a few thousand solar masses are therefore compatible with the present data. The IMBH mass limit found in this study, $M_{\text {IMBH }} / M_{C}<0.5 \%$, is consistent with that estimated from studies of other globular clusters. For instance, at the $1 \sigma$ level, van der Marel \& Anderson (2010) and Miocchi (2010) found that the allowed mass of an IMBH in $\omega$ Centauri was $M_{\text {IMBH }} / M_{C}=0.48 \%$ and $0.58 \%$, respectively. To better constrain the presence of an IMBH in NGC 6266 a more precise central position along with a proper motion dispersion profile that extends to within a fraction of an arcsecond of this location are needed. It is not clear that this latter goal will be possible within the foreseeable future since multiple epoch images with a resolution greater than that of WFPC2 will be required.

This work was supported by NASA grant HST-GO-11111.01. H.B. also acknowledges support from the Australian Research Council through Future Fellowship grant FT0991052.

## REFERENCES

Aarseth, S. J. 1999, PASP, 111, 1333
Anderson, J., \& King, I. R. 1999, PASP, 111, 1333

Anderson, J., \& King, I. R. 2003, PASP, 115, 113
Anderson, J., \& van der Marel, R. P. 2010, ApJ, 710, 1032
Baumgardt, H., Hut, P., Makino, J., McMillan, S., \& Zwart, S. P. 2003a, ApJ, 582, L21
Baumgardt, H., Makino, J., \& Hut, P. 2005, ApJ, 620, 238
Baumgardt, H., Makino, J., Hut, P., McMillan, S., \& Zwart, S. P. 2003b, ApJ, 589, L25
Beccari, G., Ferraro, F. R., Possenti, A., et al. 2006, AJ, 131, 2551
Chabrier, G. 2003, PASP, 115, 763
Djorgovski, S., \& Meylan, G. 1993, in ASP Conf. Ser. 50, Structure and Dynamics of Globular Clusters, ed. S. Djorgovski \& G. Meylan (San Francisco, CA: ASP), 325
Dubath, P., Meylan, G., \& Mayor, M. 1994, ApJ, 426, 192
Dull, J., Cohn, H., Lugger, P., et al. 1997, ApJ, 481, 267
Ebisuzaki, T., Makino, J., Tsuru, T. G., et al. 2001, ApJ, 562, L19
Gebhardt, K., Pryor, C., Williams, T. B., \& Hesser, J. E. 1994, AJ, 107, 2067
Gebhardt, K., Rich, R. M., \& Ho, L. C. 2002, ApJ, 578, L41
Gebhardt, K., Rich, R. M., \& Ho, L. C. 2005, ApJ, 634, 1093
Gerssen, J., van der Marel, R., Gebhardt, K., et al. 2002, AJ, 124, 3270
Gerssen, J., van der Marel, R., Gebhardt, K., et al. 2003, AJ, 125, 376
Gill, M., Trenti, M., Miller, M. C., et al. 2008, ApJ, 686, 303
Harris, W. E. 1996, AJ, 112, 1487
Hurley, J. R., Pols, O. R., \& Tout, C. A. 2000, MNRAS, 315, 543
Illingworth, G., \& King, I. V. 1977, ApJ, 218, L109
Jalali, B., Baumgardt, H., Kissler Patig , M., et al. 2012, A\&A, submitted
Jones, B. 1970, AJ, 75, 563
Kawakatu, N., \& Umemura, M. 2004, ApJ, 601, L21
King, I. 1962, AJ, 67, 471
Kiselev, A., Gnedin, Yu., Shakht, N., et al. 2008, Astron. Lett., 34, 529
Kroupa, P. 2001, MNRAS, 322, 231
Lützgendorf, N., Kissler-Patig, M., Noyola, E., et al. 2011, A\&A, 533, A36
Maccarone, T. J., \& Servillat, M. 2008, MNRAS, 389, 379
Magorrian, J., Tremaine, S., Richstone, D., et al. 1998, AJ, 115, 2285
Mashchenko, S., \& Sills, A. 2005, ApJ, 619, 243
McNamara, B. J., Harrison, T. E., \& Anderson, J. 2003, ApJ, 595, 187
Meylan, G., \& Mayor, M. 1991, A\&A, 250, 113
Mieske, S., Hilker, M., Jordán, A., et al. 2008, A\&A, 487, 921
Miocchi, P. 2010, A\&A, 514, 52
Noyola, E., \& Baumgardt, H. 2011, ApJ, 743, 52
Noyola, E., \& Gebhardt, K. 2006, AJ, 132, 447
Noyola, E., Gebhardt, K., \& Bergmann, M. 2008, ApJ, 676, 1008
Noyola, E., Gebhardt, K., Kissler Patig, M., et al. 2010, ApJ, 719, L60
O’Leary, R. M., Rasio, F. A., Fregeau, J. M., Ivanova, N., \& O'Shaugnessy, R. 2006, ApJ, 637, 937
Peterson, R. C., Seitzer, P., \& Cudworth, K. M. 1989, ApJ, 347, 251
Portegies Zwart, S. F., Baumgardt, H., Hut, P., Makino, J., \& McMillan, S. L. W. 2004, Nature, 428, 724

Portegies Zwart, S. F., \& McMillan, S. L. W. 2002, ApJ, 576, 899
Pryor, C., \& Meylan, G. 1993, in ASP Conf. Ser. 50, Structure and Dynamics of Globular Clusters, ed. S. Djorgovski \& G. Meylan (San Francisco, CA: ASP), 357
Rosswog, S., Ramirez-Ruiz, E., \& Hix, W. R. 2008, ApJ, 679, 1385
Taniguchi, Y., Shioya, Y., Tsuru, T., \& Ikeuchi, S. 2000, PASJ, 52, 533
Trager, S. C., King, I. R., \& Djorgovski, S. 1995, AJ, 109, 1912
Trenti, M. 2006, arXiv:astro-ph/0612040v1
Trumpler, R. J., \& Weaver, H. F. 1953, Statistical Astronomy (Berkeley, CA: Univ. California Press)
van der Marel, R. P., \& Anderson, J. 2010, ApJ, 710, 1063
Vesperini, E., \& Trenti, M. 2010, ApJ, 720, 179

