

1 **Radiative and chemical response to interactive stratospheric sulfate**  
2 **aerosols in fully coupled CESM1(WACCM)**

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17 **Key Points:**

- 18 • WACCM accurately calculates radiative and chemical responses to stratospheric sulfate,  
19 validating its use for geoengineering studies
- 20 • Interactive OH chemistry is key to the study of aerosol formation from large stratospheric  
21 SO<sub>2</sub> perturbations
- 22 • OH depletion extended the calculated average initial e-folding time for oxidation of SO<sub>2</sub>  
23 from the 1991 Pinatubo eruption by >50%

## 24 **Abstract**

25 We present new insights into the evolution and interactions of stratospheric aerosol using an  
26 updated version of the Whole Atmosphere Community Climate Model (WACCM). Improved  
27 horizontal resolution, dynamics, and chemistry now produce an internally generated quasi-  
28 biennial oscillation, and significant improvements to stratospheric temperatures and ozone  
29 compared to observations. We present a validation of WACCM column ozone and climate  
30 calculations against observations. The prognostic treatment of stratospheric sulfate aerosols  
31 accurately represents the evolution of stratospheric aerosol optical depth and perturbations to  
32 solar and longwave radiation following the June 1991 eruption of Mt. Pinatubo. We confirm the  
33 inclusion of interactive OH chemistry as an important factor in the formation and initial  
34 distribution of aerosol following large inputs of sulfur dioxide (SO<sub>2</sub>) to the stratosphere. We  
35 calculate that depletion of OH levels within the dense SO<sub>2</sub> cloud in the first weeks following the  
36 Pinatubo eruption significantly prolonged the average initial e-folding decay time for SO<sub>2</sub>  
37 oxidation to 47 days. Previous observational and model studies showing a 30-day decay time  
38 have not accounted for the large (30-55%) losses of SO<sub>2</sub> on ash and ice within 7-9 days post-  
39 eruption, and have not correctly accounted for OH depletion. We examine the variability of  
40 aerosol evolution in free-running climate simulations due to meteorology, with comparison to  
41 simulations nudged with specified dynamics. We assess calculated impacts of volcanic aerosols  
42 on ozone loss with comparisons to observations. The completeness of the chemistry, dynamics,  
43 and aerosol microphysics in WACCM qualify it for studies of stratospheric sulfate aerosol  
44 geoengineering.

## 45 **1 Introduction**

46 In this study, we describe a new version of an earth system model capable of representing the  
47 formation and interactions of stratospheric sulfate aerosol from source gases, and use it to study  
48 the evolution and radiative and chemical impacts of SO<sub>2</sub> inputs from large volcanic eruptions.

49 We use this model in a series of companion papers [*Kravitz et al.*, 2017 submitted; *MacMartin et*  
50 *al.*, 2017 submitted; *Richter et al.*, 2017 submitted; *Tilmes et al.*, 2017 submitted] to study the  
51 effects of different stratospheric sulfate geoengineering strategies. The detailed comparisons to  
52 observations presented here establish confidence in this model, and provide new insights into the  
53 role of interactive chemistry in the evolution of dense SO<sub>2</sub> clouds in the stratosphere.

54 Geoengineering, also known as climate engineering, describes a set of technologies designed to  
55 offset some of the effects of anthropogenic greenhouse gas emissions [*McNutt et al.*, 2015].

56 There are many proposed methods of offsetting anthropogenic climate change, and one method  
57 that has arguably received the most attention is stratospheric sulfate aerosol geoengineering  
58 [*Budyko and Budyko*, 1977; *Crutzen*, 2006]. This method involves injecting large amounts of  
59 sulfur-bearing precursor gases, often sulfur dioxide (SO<sub>2</sub>) into the stratosphere. These gases then  
60 photochemically convert to highly reflective sulfate aerosols, which scatter sunlight back to  
61 space, cooling the Earth's surface and lower atmosphere.

62 The idea of stratospheric sulfate aerosol geoengineering has gained the most traction of all  
63 proposed methods because of its natural analogue of large volcanic eruptions. Such volcanic  
64 eruptions similarly enhance the stratospheric sulfate aerosol layer, resulting in a cooling of  
65 Earth's climate that can last several years [e.g., *Robock*, 2000]. The 1815 eruption of Mt.

66 Tambora in what is now Indonesia was followed by the “year without a summer” in 1816 in New

67 England and Europe – which extended to several years in China – as well as severe disruptions  
68 to the Indian monsoon and to other global climate patterns [*Wood, 2014; Raible et al., 2016*].  
69 The 1991 eruption of Mt. Pinatubo (15.1°N, 120.3°E) produced a rapid global-averaged cooling  
70 at the Earth’s surface of several tenths of a degree Celsius over the following year, despite the  
71 significant warming effects of a coincident El Niño event [*Hansen et al., 1992; Soden et al.,*  
72 *2002; Bender et al., 2010*].

73 Accurately simulating the climate effects of large volcanic eruptions, and in turn stratospheric  
74 sulfate aerosol geoengineering, in a climate model requires the model to simulate processes that  
75 represent all the components of sulfate aerosol formation and microphysical growth; interaction  
76 of aerosols with radiation, dynamics, and chemistry; and sedimentation of the aerosols. Only  
77 recently have climate models included these processes, to allow for the interactive simulation of  
78 stratospheric sulfate aerosol evolution based on emissions of sulfur-bearing precursor gases.  
79 Inclusion of these processes has been shown to greatly improve the treatment of volcanic aerosol  
80 properties and their effects on stratospheric chemistry compared to observations [*Timmreck et*  
81 *al., 1999a; English et al., 2013; Mills et al., 2016; Solomon et al., 2016; Ivy et al., 2017*]. In  
82 addition, simulation of the interactions between stratospheric processes and surface climate  
83 requires coupling to an ocean and sea-ice model, which is often lacking in models with  
84 prognostic aerosol capabilities. These processes are essential for studying the atmospheric and  
85 surface climate impacts of stratospheric sulfate aerosol geoengineering.

86 We describe an updated version of the earth system model described in *Mills et al.* [2016] with  
87 the above processes and interactions included. Updates include increased horizontal resolution  
88 and a self-generating quasi-biennial oscillation (QBO). We use this model here to study the  
89 chemical, microphysical, and radiative effects of historical volcanic eruptions that have occurred

90 during the satellite era (1979-present), with comparison to observations. We demonstrate the  
91 importance of interactive calculations of the abundance of oxidants, such as the hydroxyl radical  
92 (OH), to understanding the observations of SO<sub>2</sub> evolution following large volcanic eruptions.

## 93 **2 Materials and Methods**

### 94 **2.1 WACCM**

95 The Community Earth System Model, version 1 [CESM, *Hurrell et al.*, 2013], is a state-of-the-  
96 art global climate model that includes interactive atmosphere, ocean, land, and sea-ice  
97 components. The atmosphere component of CESM1 is the Community Atmosphere Model  
98 (CAM), which includes a high-top version known as the Whole Atmosphere Community  
99 Climate Model [WACCM, *Marsh et al.*, 2013]. *Mills et al.* [2016] describe the development of a  
100 prognostic treatment of stratospheric sulfate aerosol in CESM1(WACCM) with the more  
101 realistic formulations of radiation, planetary boundary layer turbulence, cloud microphysics, and  
102 aerosols that were introduced in version 5 of CAM [*Neale et al.*, 2010]. *Mills et al.* [2016]  
103 presented and validated volcanic aerosol properties derived from SO<sub>2</sub> emissions over the period  
104 1990-2014, but did not examine radiative forcing. In this paper, we validate radiative forcing  
105 from volcanic aerosol following the 1991 Pinatubo eruption calculated with WACCM. Such  
106 validation is critical for the use of this model in studies of the radiative impacts of stratospheric  
107 sulfate aerosol derived from SO<sub>2</sub> emissions. The horizontal resolution of the atmosphere  
108 component in this model, which we call WACCM hereafter, is 0.95° latitude x 1.25° longitude,  
109 which is double the resolution in each horizontal dimension of previous versions of  
110 CESM1(WACCM) [*Marsh et al.*, 2013; *Mills et al.*, 2016].

111 WACCM extends from the Earth's surface to 140 km in altitude. In our configuration, WACCM  
112 includes comprehensive, fully interactive middle-atmosphere chemistry with 95 solution species,  
113 2 invariant species, 91 photolysis reactions, and 207 other reactions. The chemical scheme  
114 includes gas-phase chemical species in the O<sub>x</sub>, NO<sub>x</sub>, HO<sub>x</sub>, ClO<sub>x</sub>, and BrO<sub>x</sub> chemical families,  
115 along with CH<sub>4</sub> and its degradation products, and the sulfur-bearing gases dimethyl sulfide  
116 (DMS), OCS, SO<sub>2</sub>, SO, S, SO<sub>3</sub>, and H<sub>2</sub>SO<sub>4</sub>. Gas-phase and heterogeneous reactions important in  
117 the stratosphere are included, allowing simulation of the impacts of sulfate aerosols on the  
118 chemical composition of the atmosphere, such as the seasonal ozone hole over Antarctica in  
119 austral spring [Mills *et al.*, 2016]. Our model's middle-atmosphere chemistry is a subset of the  
120 chemistry used in Mills *et al.* [2016], excluding species and reactions that are significant only in  
121 the troposphere. The reduced chemistry produces up to 45% more OH in the troposphere than  
122 Mills *et al.* [2016], resulting in slightly reduced tropospheric lifetimes for species such as CH<sub>4</sub>  
123 and SO<sub>2</sub>. Climate forcings in WACCM include aerosols (tropospheric and stratospheric,  
124 anthropogenic and natural), solar variability, and time-varying mixing ratios of greenhouse gases  
125 (determined by lower boundary conditions and interactive chemistry).

126 WACCM includes a modal treatment of aerosols that is coupled to cloud microphysics [Liu *et*  
127 *al.*, 2012], and which has been extended to include stratospheric sulfate [Mills *et al.*, 2016]. To  
128 simulate the formation and evolution of sulfate aerosol prognostically, our chemical mechanism  
129 includes precursor sulfur-bearing gases and oxidation pathways producing H<sub>2</sub>SO<sub>4</sub>. Source gases  
130 include OCS, which is an important source of background stratospheric aerosol, as well as SO<sub>2</sub>  
131 from anthropogenic sources. The H<sub>2</sub>SO<sub>4</sub> resulting from this oxidation creates new sulfate  
132 aerosols by the microphysical processes of nucleation and condensation. The processes of  
133 coagulation, evaporation, and sedimentation are included in the aerosol microphysical evolution.

134 The model that we use includes a number of improvements to CAM5 physics beyond what was  
135 used in *Mills et al.* [2016]. We use a new surface topography for the CAM finite-volume  
136 dynamical core based on *Lauritzen et al.* [2015]. We include an improved representation of  
137 atmospheric dust, including refined physical parameterizations of dust and improved soil  
138 erodibility, size distributions, wet deposition and optics [*Albani et al.*, 2014]. The cloud  
139 microphysical scheme has been updated to Morrison-Gottelman version 2 (MG2), which  
140 includes prognostic precipitation [*Gottelman and Morrison*, 2015]. An error in the energy  
141 formulation has been corrected [*Williamson et al.*, 2015]. The vertical remapping scheme has  
142 been updated to improve energy conservation. In the original implementation, temperature was  
143 retrieved from total energy remapping (minus kinetic energy), which was shown to produce  
144 significant temperature perturbations at high altitude. In the new implementation, temperature is  
145 remapped over a log-pressure coordinate, which preserves the geopotential at the model lid  
146 during remapping.

147 Ice nucleation has been updated to include effects of pre-existing ice crystals, and to consider in-  
148 cloud variability in ice saturation ratio [*Shi et al.*, 2015]. The ice nucleation scheme was  
149 developed for the troposphere, and contains several assumptions that may adversely affect ice  
150 nucleation in the upper troposphere and lower stratosphere in our model. The heterogeneous ice  
151 nucleation code assumes that only dust aerosols nucleate ice, and only those in the larger of the  
152 two dust aerosol modes. In deriving the fraction of dust in this coarse mode, the code considers  
153 only the ratio of dust to sea salt, neglecting the presence of sulfates in the coarse mode. Thus, in  
154 the upper troposphere and stratosphere, where sulfate dominates aerosol composition, the very  
155 small dust fraction is greatly overestimated because the sea salt mass there is small compared to  
156 the dust mass. Hence the code overestimates heterogeneous freezing in the upper troposphere

157 and lower stratosphere, but only in model gridpoints where dust is present, and hence is not an  
158 issue in most of the stratosphere. In addition, coarse aerosols that nucleate ice are not moved to  
159 the in-cloud population, and are available to nucleate ice each additional time step, leading to  
160 further overestimates of heterogeneous freezing. Homogeneous freezing of aerosols is  
161 considered only for sulfates in the Aitken mode. The neglect of sulfates in the larger  
162 accumulation and coarse modes likely underestimates ice production by homogeneous freezing,  
163 particularly under geoengineered conditions. Because the impacts on ice nucleation are  
164 compensating, the sign of model biases introduced is unclear. Heterogeneous reactions on  
165 stratospheric ice account for a small (~1%) proportion of Antarctic ozone loss. These issues may  
166 have more significant impacts on the interaction of aerosols with ice clouds, which can absorb  
167 outgoing longwave radiation. The erroneous treatment of sulfates as heterogeneous ice nuclei  
168 where dust is present may produce unrealistic increases in cirrus clouds in the upper troposphere  
169 under geoengineering conditions, and the resulting longwave absorption would reduce the  
170 cooling efficiency of geoengineering unrealistically. These issues will be addressed in future  
171 versions of CESM.

172 Our WACCM configuration includes the same 70 vertical layers as described in *Mills et al.*  
173 [2016]. WACCM uses the *Lindzen* [1981] gravity wave propagation scheme, using gravity wave  
174 source specifications for orographic, frontal, and convectively generated gravity waves following  
175 *Richter et al.* [2010]. In this version of WACCM, we have increased the efficiency of  
176 convectively generated gravity waves generated by convection to 0.40, from 0.10 used in *Mills et*  
177 *al.* [2016]. This change, together with increased horizontal resolution, allows for an internal  
178 generation of the QBO. This new development allows for the examination of the effects of SO<sub>2</sub>  
179 injections on the QBO, which are presented in *Richter et al.* [2017 submitted], motivated in part



180 by previous work that has suggested that stratospheric sulfate geoengineering could severely  
181 alter the QBO [*Aquila et al.*, 2014]. Additional tuning of to the gravity wave parameterization in  
182 WACCM has significantly reduced the bias of the temperatures in the Antarctic polar vortex,  
183 which is critical to calculating ozone loss [*Garcia et al.*, 2017].

184 WACCM is fully coupled to the Community Land Model version 4.0 [CLM4.0, *Lawrence et al.*,  
185 2011]. The land model includes interactive carbon and nitrogen cycles, as in CESM1(WACCM)  
186 [*Marsh et al.*, 2013] . In addition, biogenic surface emissions into the atmosphere are calculated  
187 in CLM4.0 using the Model of Emissions of Gases and Aerosols from Nature, version 2.1  
188 [MEGAN2.1, *Guenther et al.*, 2012]. WACCM is also coupled to ocean and sea ice components  
189 that may be interactive, or constrained by data representative of observations. The interactive  
190 components are the Parallel Ocean Program, version 2 [POP2, *Danabasoglu et al.*, 2012] and the  
191 Los Alamos National Laboratory sea ice model, version 4 [CICE4, *Holland et al.*, 2012].

## 192 2.2 Model simulations

193 Interactive stratospheric aerosol is a new development common to an increasing number of  
194 climate models participating in the forthcoming Coupled Model Intercomparison Project, phase 6  
195 [CMIP6, *Eyring et al.*, 2016]. Some of these models also include interactive chemistry. Some  
196 include an interactive QBO. Interactive aerosols from volcanic eruptions will disperse and  
197 evolve differently in different ensemble runs, depending on meteorology [*Jones et al.*, 2016] and  
198 the phase of the QBO [*Trepte and Hitchman*, 1992]. We discuss the impacts of these issues on  
199 chemistry and climate variability by comparing fully coupled free-running (FR) simulations to  
200 those constrained by nudging to specified dynamics (SD). To quantify the importance of  
201 interactive chemistry, we conducted 2 simulations with non-interactive specified chemistry.

202 Table 1 presents a list of simulations performed for this work, each of which is described below.  
203 Simulations were carried out on the Yellowstone high-performance computer platform  
204 [*Computational and Information Systems Laboratory*, 2012].

205 In order to demonstrate radiative balance in the updated model between incoming solar and  
206 outgoing longwave radiation prior to the rapid introduction of anthropogenic greenhouse gases in  
207 the industrial era, we first conducted a 50-year pre-industrial fully coupled free-running  
208 simulation, FRPI, using constant year 1850 climate forcing conditions. The land, ocean, and sea  
209 ice components were initialized with the pre-industrial equilibrium conditions (1 January, year  
210 402, of the 1850 fully coupled control) used to initialize the CESM1 Large Ensemble simulations  
211 [*Kay et al.*, 2015]. Initial conditions for the atmosphere are consistent with pre-industrial  
212 conditions.

213 Our FRVOLC experiment was designed to examine the model's representation of climate  
214 following historical conditions, including volcanic eruptions, from 1975-2016, in an ensemble of  
215 four fully coupled free-running simulations individually named FRVOLC1, FRVOLC2,  
216 FRVOLC3, and FRVOLC4. The land, ocean, and sea ice components are interactive in these  
217 simulations, and were initialized with conditions from January 1, 1975 of four independent  
218 CESM1 transient simulations used for the Large Ensemble simulation, which were picked to  
219 sample contrasting initial ocean states. The atmosphere was initialized from a simulation  
220 conducted for the Chemistry Climate Modeling Initiative for the atmosphere component  
221 [*Solomon et al.*, 2015], regridded to our model's higher horizontal resolution, with the addition of  
222 spun-up initial conditions for aerosols and sulfur gases from a previous run of our model.

223 Coupled free-running experiments allow self-consistent representations of stratospheric aerosols,  
224 with interactions between atmospheric chemistry and dynamics, and ocean, sea ice and land.  
225 However, the number of unconstrained variable climate states in such simulations pose  
226 difficulties for comparisons to observations. Meteorology at the time of volcanic eruption can  
227 play an important role in the latitudinal distribution of aerosol, as can the state of the QBO,  
228 which can affect the transport of stratospheric aerosol from the tropics to higher latitudes [*Trepte*  
229 *and Hitchman*, 1992]. Ocean states, including the El Niño-Southern Oscillation, strongly affect  
230 observations of the Earth's radiation budget, complicating comparisons to coupled free-running  
231 simulations. We therefore rely on SD simulations, with prescribed historical sea surface  
232 temperatures, and atmospheric winds and temperatures nudged to historical meteorology, to  
233 constrain climate variability in WACCM, allowing more detailed comparisons to observations of  
234 chemistry and climate responses to stratospheric aerosol.

235 We performed two SD experiments of the years 1990-2015 using initial conditions from the  
236 FRVOLC1 simulation. The SD experiments use meteorological fields from the NASA Global  
237 Modeling and Assimilation Office Modern-Era Retrospective Analysis for Research and  
238 Applications (MERRA) [*Rienecker et al.*, 2011]. Horizontal winds and temperatures are nudged  
239 toward the MERRA reanalysis fields between the surface and 50 km, with a relaxation time of  
240 50 hours. SDVOLC includes SO<sub>2</sub> emissions from explosive volcanic eruptions, and SDVC is a  
241 “volcanically clean” run without SO<sub>2</sub> emissions from explosive volcanic eruptions.

242 Using WACCM, we performed two specified chemistry experiments [*Smith et al.*, 2014], for  
243 which chemical oxidants (OH, HO<sub>2</sub>, O<sub>3</sub>, and NO<sub>3</sub>) are prescribed. SCVOLC included SO<sub>2</sub> from  
244 volcanic eruptions, and simulates years 1979-1999, with an initial condition for January 1, 1979  
245 from FRVOLC1. SCVC does not include SO<sub>2</sub> from eruptions, and simulates years 1990-1999,

246 with an initial condition for January 1, 1990 from SCVOLC. These specified chemistry runs used  
247 free-running atmospheric dynamics, and prescribed sea surface temperatures and sea ice.

248 Our FRVOLC, SDVOLC, and SCVOLC experiments use a database of SO<sub>2</sub> emissions from  
249 volcanic eruptions based on version 2 of the Volcanic Emissions for Earth System Models [*Neely*  
250 *and Schmidt*, 2016]. The database includes 222 days of eruption for the years 1975-2016, the  
251 dates, spatial coordinates, and SO<sub>2</sub> mass of which are described in Table S1 of the supporting  
252 information. As in *Mills et al.* [2016], eruptive emissions occur over a 6-hour period from 1200  
253 to 1800UT on the date of the eruption. The climatic phase of the 1991 Pinatubo eruption  
254 coincided with the closest pass of Typhoon Yunya (50 km north), which likely affected the initial  
255 transport of SO<sub>2</sub> from the eruption, and which also prevented the retrieval of atmospheric wind  
256 profiles during the eruption [*Holasek et al.*, 1996; *Guo et al.*, 2004b]. To account for the  
257 observed initial transport of SO<sub>2</sub> from the Pinatubo eruption southward, we spread the emissions  
258 from that eruption evenly between 15.13°N and the equator at 120.3°E, as in previous studies  
259 [*Timmreck et al.*, 1999a; 1999b; *Dhomse et al.*, 2014; *Sheng et al.*, 2015; *Mills et al.*, 2016]. As  
260 discussed in *Mills et al.* [2016], we emit Pinatubo SO<sub>2</sub> evenly between 18 and 20 km, which  
261 allows for self-lofting, giving best agreement with MLS observations of the SO<sub>2</sub> cloud [*Read et*  
262 *al.*, 1992].

## 263 **3 Results**

### 264 **3.1 WACCM climate**

265 As we are presenting a new version of WACCM, we begin by validating the radiative balance of  
266 the model in pre-industrial conditions, and the climate and chemistry in present-day conditions

267 with respect to observations, before examining volcanic responses. We ignore the first 24 years  
268 of the FRPI simulation, to allow the components to equilibrate. The difference between the  
269 absorbed solar radiation (ASR) and outgoing longwave radiation (OLR) at the model's top gives  
270 a net radiative flux of  $-0.027 \pm 0.442 \text{ W m}^{-2}$  ( $1-\sigma$  confidence) over the last 26 years, indicating  
271 that the pre-industrial atmosphere is in radiative equilibrium. The Diagnosing Earth's Energy  
272 Pathways in the Climate project, version 2 [Allan *et al.*, 2014; Liu *et al.*, 2015] combines 60°S-  
273 60°N Earth Radiation Budget Satellite (ERBS) broadband non-scanner measurements during the  
274 Earth Radiation Budget Experiment [ERBE, Minnis *et al.*, 1993] with additional data to provide  
275 continuous global monthly observations of ASR and OLR from 1985 to present. Our historical  
276 FRVOLC ensemble calculates a net radiative flux for years 1985-1999 of  $0.56 \pm 0.63 \text{ W/m}^2$ ,  
277 which is in general agreement with  $0.35 \pm 0.66 \text{ W/m}^2$  from the merged ERBS data. Figure 1  
278 compares the global annual surface temperature anomaly for 1979-2015 from the FRVOLC  
279 ensemble to reconstructions from Hadley Centre–Climatic Research Unit Version 4  
280 (HadCRUT4) infilled with kriging [Cowtan and Way, 2014] and GISS Surface Temperature  
281 Analysis [Hansen *et al.*, 2010; GISTEMP Team, 2017]. Anomalies are calculated with respect to  
282 the 1979-2015 average for each data set. Shading shows the range of the annual global mean  
283 values over the 4 FRVOLC ensemble members, and red asterisks show the mean of the  
284 FRVOLC ensemble. Lines show 5-year running averages of the annual anomalies for  
285 simulations and observations. WACCM shows similar decadal variability to the observations,  
286 including significant cooling after the major eruptions of El Chichón (1982, 17.4°N, 93.2°W)  
287 and Pinatubo (1991). Observations lie largely within the range of the FRVOLC ensemble  
288 variability. This gives confidence that the climate response of the model to long-term changes in  
289 greenhouse gases agrees with observations. Least squares linear fit trends for 1975-2016 are 2.45

290  $\pm 0.11$  K/century for the FRVOLC ensemble, which is slightly larger than those for HadCRUT4  
291 ( $2.16 \pm 0.16$ ) and GISTEMP ( $1.92 \pm 0.12$ ).

292 WACCM has a very good representation of the mean temperature in the troposphere and middle  
293 atmosphere. As shown in Figure 2, throughout the entire troposphere the temperature bias  
294 relative to ERA-Interim reanalysis [ERA-Interim, *Dee et al.*, 2011] is less than 2 K (Figure 1), with only  
295 the tropical mid-troposphere carrying a bias greater than 1 K throughout the year (not shown).  
296 WACCM has a cold bias near the tropopause of -2 to -4 K in both the extratropics and the tropics  
297 throughout most of the year. This represents an improvement over CESM1(WACCM) [*Marsh et*  
298 *al.*, 2013], which ran at  $1.9^\circ$  latitude x  $2.5^\circ$  longitude and exhibited extratropical biases ranging  
299 from -6 to -8 K. We attribute this improvement to the improved horizontal resolution. *Charlton-*  
300 *Perez et al.* [2013] showed that the CMIP5 multi-model average also carries a -4 K bias in the  
301 extratropical tropopause temperatures.

302 Lower stratospheric (below 10 hPa) mean temperatures in WACCM are in excellent overall  
303 agreement with observations, as shown in Figure 2. In the upper stratosphere (1-10 hPa),  
304 WACCM has a warm bias ( $< 10$  K) between  $-60^\circ\text{S}$  and  $60^\circ\text{N}$ , and a cold bias of up to 16K at the  
305 winter pole. WACCM also has a cold bias ( $< 12$  K) in the south polar stratosphere in autumn  
306 (MAM) and spring (SON). These biases are of the same magnitude or smaller than those in  
307 CESM1(WACCM) [not shown, *Marsh et al.*, 2013].

308 The summer mesopause in WACCM is near 87 km (log-pressure height) with temperature in  
309 January of 131 K (at  $80^\circ\text{S}$ ) and 130 K in July (at  $80^\circ\text{N}$ ). This is in reasonable agreement with  
310 SABER observations [*Xu et al.*, 2007] which show mesopause temperatures of 134 K in January  
311 (at  $80^\circ\text{S}$ ) and 127 K in July (at  $80^\circ\text{N}$ ). Zonal mean winds averaged over DJF and JJA for

312 WACCM are shown in Figure 3 along with winds from the Upper Atmosphere Research Satellite  
313 Reference Atmosphere Project (URAP) climatology [Swinbank and Ortland, 2003]. The overall  
314 DJF stratospheric and mesospheric wind structure is in good agreement with URAP. The model  
315 climatology is improved over CESM1(WACCM), with a few remaining biases. The stratospheric  
316 NH jet in DJF is  $\sim 10 \text{ m s}^{-1}$  stronger than observed (associated with slightly colder than observed  
317 temperatures) and the summer stratospheric jet is too weak between  $60^{\circ}\text{S}$  and  $90^{\circ}\text{S}$  above 10  
318 hPa. In JJA, the NH summer jet has very good agreement to URAP and is much improved  
319 between  $0$  and  $30^{\circ}\text{N}$  compared to WACCM3 and CESM1(WACCM), which carried a  $30 \text{ m s}^{-1}$   
320 bias in this region [Richter et al., 2010; Marsh et al., 2013]. The SH stratospheric westerly jet is  
321 too strong in WACCM, however unlike in WACCM3 and CESM1(WACCM) it is tilting in the  
322 correct direction (towards the equator as height increases).

### 323 3.2 QBO and stratospheric chemistry

324 WACCM has an internally generated QBO as shown in Figure 4. The period of the QBO in the  
325 FRVOLC ensemble varies between 19 and 36 months, with a mean periods for each ensemble  
326 member varying from 23 to 27 months. In observations, the QBO period ranges between 20 and  
327 34 months, with a mean of 28 months. The amplitude of the westerly QBO phase is between 15  
328 and  $20 \text{ m s}^{-1}$ , exactly as in observations. The easterly QBO phase amplitude ranges between 20  
329 and  $25 \text{ m s}^{-1}$ , and is hence weaker than observed by  $10 \text{ m s}^{-1}$ . Further improvements to the  
330 representation of the QBO in WACCM require a substantial increase in the vertical resolution  
331 [Richter et al., 2014].

332 The stratospheric water vapor “tape recorder” [Mote et al., 1996] is well represented in  
333 WACCM, consistent with a good representation of tropical tropopause temperatures. A good

334 representation of water vapor in the stratosphere is important for climate because of the role of  
335 water vapor as a greenhouse gas. Water vapor also strongly impacts stratospheric chemical  
336 cycles affecting ozone, which is an important radiatively active gas. In comparison to Aura MLS  
337 satellite observations [*Livesey et al.*, 2016] between 2004 and 2014, the magnitude of both dry  
338 and wet phase of the tape recorder follow exactly the observed range (Figure 5). The slope of the  
339 tape recorder in the model is in good agreement with the observations, with a slightly stronger  
340 tropical upwelling in the lower stratosphere.

341 WACCM zonal mean stratospheric ozone column shows very good agreement with observations,  
342 and excellent agreement in high latitudes (Figure 6). The representation of ozone in WACCM is  
343 improved over previous versions of WACCM. Figure 6 compares the zonal average stratospheric  
344 column ozone with a 10°x10° horizontally gridded product, based on MLS and NASA Ozone  
345 Monitoring Instrument (OMI) observations, averaged between 2005 and 2010 [*Ziemke et al.*,  
346 2011]. This figure shows agreement within 8% in high latitudes for four different seasons of the  
347 mean value of the FRVOLC ensemble with the observations. The observations generally lie  
348 within one standard deviation of ensemble variability at high latitudes, with the exception of  
349 60°S in April. We attribute this improved performance to the improved horizontal resolution and  
350 dynamical improvements associated with modifications to the gravity wave parameterization  
351 [*Garcia et al.*, 2017]. The model slightly underestimates column ozone in the Tropics, which  
352 may be due to overly rapid transport.

### 353 3.3 Volcanic aerosol evolution

354 In order to validate that our model produces a reasonable response to stratospheric SO<sub>2</sub>  
355 perturbations, we compare the period from January 1, 1990 to January 1, 2000 in our simulations



356 to observations. This period includes the eruption of Mt. Pinatubo on June 15, 1991, which  
357 produced the best-observed large (>10 Tg) injection of SO<sub>2</sub> into the stratosphere to date.  
358 Comparison of simulated surface climate response to the volcanic eruption based on observations  
359 is complicated by climate variability, including a coincident El Niño event that tended to  
360 counteract the reduction in global average temperatures following the eruption, as well as other  
361 underlying climate oscillations [e.g., *Canty et al.*, 2013]. We therefore constrain our calculations  
362 by using SD-WACCM, which incorporates a data ocean model with observed sea surface  
363 temperatures, as well as nudged atmospheric temperatures and winds.

364 The first step in the production of sulfate aerosol from stratospheric SO<sub>2</sub> input is chemical  
365 oxidation by the OH radical, which, via intermediate steps, produces H<sub>2</sub>SO<sub>4</sub> gas. Figure 7 shows  
366 the time evolution of the total burden of volcanic (SDVOLC minus SDVC) SO<sub>2</sub> in WACCM,  
367 compared to observations of the stratospheric burden following the eruption. *Guo et al.* [2004a]  
368 present and evaluate the SO<sub>2</sub> observations from the Total Ozone Mapping Spectrometer (TOMS)  
369 and the Television Infrared Observation Satellite Optical Vertical Sounder (TOVS) in the first 15  
370 days after the Pinatubo eruption. That work suggested that much of the initial 14 to 23 Tg of SO<sub>2</sub>  
371 (7 to 11.5 Tg of sulfur) from Pinatubo was rapidly catalyzed on ash and ice, fast processes that  
372 are not currently included in WACCM. As in *Mills et al.* [2016], we input 10.0 Tg of SO<sub>2</sub> (5.0  
373 Tg of sulfur) from Pinatubo in WACCM on the day of the eruption, matching the burden from  
374 TOMS and TOVS observations 7-9 days after the beginning of the eruption, when more than  
375 99% of the ash and ice particles had been removed [*Guo et al.*, 2004b]. This "climatically  
376 relevant" sulfur input from Pinatubo is consistent with the 3.7 to 6.7 Tg peak sulfur content of  
377 stratospheric aerosol mass following the eruption derived from satellite observations from the  
378 High-resolution Infrared Radiation Sounder [*HIRS, Baran and Foot, 1994*] and the Improved

379 Stratospheric and Mesospheric Sounder [*Lambert et al.*, 1993], as discussed in *Dhomse et al.*  
380 [2014]. The evolution of aerosol mass burden calculated in WACCM following Pinatubo is  
381 consistent with HIRS observations, as shown in Figure 1 of Mills et al. [2016].

382 Additional SO<sub>2</sub> observations shown in Figure 7, from the Solar Backscatter Ultraviolet  
383 Radiometer-2 (SBUV/2), the Microwave Limb Sounder (MLS) aboard the Upper Atmosphere  
384 Research Satellite, and a high-resolution infrared spectrometer aboard an aircraft [*Mankin et al.*,  
385 1992] are presented as shown in *Read et al.* [1993]. Our calculations show general agreement  
386 with these observations within their limitations. As indicated in Figure 7, the eruption of Cerro  
387 Hudson (45.9°S, 73.0°W) emitted 1.5 Tg SO<sub>2</sub> roughly 2 months after Pinatubo. This additional  
388 input was not observed by MLS, which integrated SO<sub>2</sub> above 21 km, well above the height of the  
389 Cerro Hudson plume. The dashed line shows the SO<sub>2</sub> burden in WACCM above 50 hPa, for  
390 comparison to the MLS observations.

391 *Read et al.* [1993] used these observations to derive a 33-day e-folding decay time with an  
392 extrapolated initial SO<sub>2</sub> injection of 17 Tg. Our calculations point to an interpretation of this  
393 apparent steady exponential decay as the superposition of two more variable processes: loss on  
394 ash and ice, and oxidation by OH. We note that the slope of the semi-logarithmic plot of SO<sub>2</sub>  
395 burden versus time shown for SDVOLC minus SDVC in Figure 7 indicates a much longer initial  
396 lifetime, decreasing to a constant slope by 30 days after the eruption. The reason for this is the  
397 rapid consumption of OH by SO<sub>2</sub> oxidation within the initial dense SO<sub>2</sub> cloud, which limits the  
398 availability of OH, and hence the SO<sub>2</sub> oxidation rate. As Figure 8 shows, OH is reduced by more  
399 than 95% within the cloud as it is transported in the first weeks. Figure 9 shows the daily e-  
400 folding decay time of volcanic SO<sub>2</sub> as a function of days after the eruption. As the cloud  
401 disperses to larger volumes, OH recovers, and the initial e-folding decay time of more than 400

402 days drops over the first month to reach a constant value of  $30.9 \pm 0.5$  days (45 to 59 days after  
403 eruption). In contrast, the volcanic  $\text{SO}_2$  in our SCVOLC simulation, for which OH is prescribed,  
404 decays with a constant e-folding time of  $34.1 \pm 1.4$  days (2 to 21 days after eruption, figures 7  
405 and 9). The specified chemistry simulations show greater variability, particularly as volcanic  
406 minus clean  $\text{SO}_2$  burdens approach zero, because they are not nudged as the SD simulations are,  
407 and therefore have unmatched non-volcanic burdens.

408 Constant e-folding decay times ranging from 23 to 35 days have been derived from observations  
409 for Pinatubo  $\text{SO}_2$  [Bluth *et al.*, 1992; McPeters, 1993; Read *et al.*, 1993; Guo *et al.*, 2004a].  
410 These constant e-folding times do not distinguish the rapid initial removal of  $\text{SO}_2$  on sedimenting  
411 ash and ice in the initial days after the eruption from the variable chemical oxidation rate due to  
412 OH depletion. Our simulations with interactive OH chemistry find a similar constant terminal e-  
413 folding time, but we also find the "average initial e-folding time" for oxidation to be 47 days,  
414 calculated as the time for the initial 10 Tg of Pinatubo  $\text{SO}_2$  to be reduced by 1/e to 3.7 Tg. Pinto  
415 *et al.* [1989] examined the effects of large stratospheric  $\text{SO}_2$  injections on the e-folding time for  
416 loss of  $\text{SO}_2$  by OH. Using a one-dimensional model that accounted for the horizontal dispersion  
417 and expansion of volcanic  $\text{SO}_2$  clouds, they calculated that an injection of 10 Tg of  $\text{SO}_2$  should  
418 increase the e-folding time from 1.3 to 1.8 months, which is consistent with our calculations.  
419 Bekki [1995] found the reduction in OH oxidation to be significant for a much larger 200-Tg  
420 injection, but concluded that the effects of Pinatubo's ~20-Tg  $\text{SO}_2$  injection on OH "would have  
421 been too modest to have had a noticeable effect on the global  $\text{SO}_2$  removal rate." That  
422 assessment, however, relied on a coarse zonally-averaged two-dimensional model with very  
423 large grid cells (~10° latitude x 360° longitude), which could not account for the local OH  
424 depletion within the  $\text{SO}_2$  cloud in three spatial dimensions.

425 Several studies examined the effects of dense SO<sub>2</sub> clouds and volcanic aerosols on OH levels due  
426 to absorption and scattering of sunlight, which affects photolysis rates [*Pinto et al.*, 1989; *Bekki*,  
427 1995; *Bândă et al.*, 2015; *Pitari et al.*, 2016a]. WACCM does not include such effects on  
428 photolysis rates, which these studies indicate are of lesser significance than reductions in OH due  
429 to sulfur chemistry. *Bândă et al.* [2015] found no significant effect of SO<sub>2</sub> absorption on the e-  
430 folding time of SO<sub>2</sub> from the 1991 Pinatubo eruption. Impacts of stratospheric sulfate  
431 geoengineering on photolysis rates and the oxidation capacity of the troposphere might be more  
432 significant [*Pitari et al.*, 2014; *Visioni et al.*, 2017]. Our studies with WACCM focus on middle  
433 atmosphere chemistry, which would be less affected by such effects than the troposphere.

434 Our results show that interactive OH chemistry is essential to accurately calculating oxidation  
435 and dispersal following the input of 10 Tg or more of SO<sub>2</sub> into the stratosphere. Studies of  
436 interactive stratospheric aerosols in earth system models that use invariant prescribed OH values  
437 calculated constant e-folding times for Pinatubo SO<sub>2</sub> of 29-33 days [*Niemeier et al.*, 2009; *Aquila*  
438 *et al.*, 2012], leading to faster initial oxidation. *Bekki and Pyle* [1994] used a two-dimension  
439 model that neglected feedbacks between SO<sub>2</sub> photochemistry and other chemical species, and  
440 calculated a longer e-folding time of 40 days, which they account for by stating: "Since SO<sub>2</sub> is  
441 only significantly removed by OH, this small difference is probably because the modeled OH  
442 levels are low compared to reality in the region of the volcanic cloud." *Sekiya et al.* [2016]  
443 calculated an e-folding time of 38-40 days, using a general circulation model with interactive  
444 OH.

445 The University of L'Aquila Composition-Climate Coupled Model (ULAQ-CCM), which also  
446 includes interactive OH, found a 31-day e-folding time for a 20-Tg SO<sub>2</sub> Pinatubo eruption based  
447 on exponential decay between days 45 and 165 after the eruption, but did not report on variations

448 within the first month [Pitari *et al.*, 2016b]. They also found a much shorter 18-day e-folding  
449 time for a much smaller 1.2-Tg SO<sub>2</sub> eruption, which they account to "more abundant OH due to  
450 an inefficient sink by sulfur dioxide" in the smaller volcanic cloud. The 19-day and 43-day e-  
451 folding times that they find for eruptions of 7 and 12 Tg SO<sub>2</sub>, respectively, suggest factors other  
452 than the mass of SO<sub>2</sub> erupted also affected their calculations.

453 Inclusion of interactive OH chemistry in WACCM is key to understanding variable oxidation  
454 and its importance for the subsequent size and latitudinal distribution of stratospheric aerosol.  
455 We found significantly greater self-lofting of volcanic aerosol in our SCVOLC simulation than  
456 in the FRVOLC ensemble for the two largest eruptions simulated (El Chichón 1982 and Pinatubo  
457 1991). This is due to radiative interaction with the dense aerosol clouds that result from rapid  
458 oxidation of the volcanic SO<sub>2</sub> before it disperses (not shown).

459 *Mills et al.* [2016] presented validations of volcanic aerosol properties in WACCM with CAM5  
460 physics, using half the horizontal resolution used in this study. Here we present similar  
461 validations before examining radiative impacts. In Figure 10, we compare stratospheric aerosol  
462 optical depth (SAOD) at 550 nm for 1990-1998 measured by lidars at 3 locations (black circles)  
463 to 5-day average values calculated at the same locations in our SDVOLC (red dots), FRVOLC  
464 (blue dots, ensemble average), and SCVOLC (orange dots) simulations. The reduced SAOD in  
465 the FRVOLC ensemble compared to SDVOLC reflects a lower stratospheric aerosol burden, and  
466 a shorter aerosol lifetime. This relates to faster circulation and higher (~0.5 km) tropical  
467 tropopause altitudes in FR-WACCM. It also relates to the phase of the QBO at the time of the  
468 eruption, which agrees with observations in SDVOLC (easterlies above 26 km overlying  
469 westerlies below), but which is variable in the FRVOLC ensemble. The easterly shear in  
470 SDVOLC is associated with lofting of the Pinatubo aerosols in the tropics, while the westerly

471 shear is associated with descent and transport to higher latitudes [*Trepte and Hitchman, 1992*],  
472 which we see in our simulations. This difference illustrates a mode of variability affecting  
473 volcanic aerosol evolution in the atmosphere compared to models with interactive QBOs.  
474 Comparison of SAOD in SCVOLC to FRVOLC shows the impact of the enhanced self-lofting of  
475 Pinatubo aerosols when OH depletion is not accounted for.

476 Figure 10a compares to a newly available lidar record from Tomsk, Siberia [*Zuev et al., 2016*].  
477 While the tropopause altitude is generally 11-13 km, the lidar backscatter was integrated between  
478 15-30 km. We have converted the integrated backscatter to aerosol optical depth (AOD) using a  
479 lidar ratio (integrated extinction/backscatter) of 50, which has been found to be appropriate for  
480 the stratosphere within 20% [*Jäger and Deshler, 2002; 2003; Ridley et al., 2014*]. Our calculated  
481 AOD is integrated above the tropopause, yielding slightly higher values in SDVOLC over the  
482 Pinatubo period than the integrated backscatter, which excludes the lowermost 2-4 km. As in  
483 *Mills et al. [2016]*, our SDVOLC calculations show excellent agreement over the Pinatubo  
484 period with the lidars at Geestacht, Germany [*Ansmann et al., 1997*] (Figure 10b), and Mauna  
485 Loa, Hawaii [*Hofmann et al., 2009; Ridley et al., 2014*] (Figure 10c), both of which are  
486 integrated above the tropopause.

### 487 3.4 Volcanic aerosol radiative forcing

488 In Figure 11, we show global mean all-sky net radiative fluxes at the top of the model, compared  
489 to the observed global mean time series from the merged ERBS data [*Allan et al., 2014; Liu et*  
490 *al., 2015*]. Monthly mean net fluxes are shown for January 1991 to December 1995, normalized  
491 and de-seasonalized by subtracting the corresponding flux for each month from the volcanically  
492 quiescent year 1999. Figure 11a shows the de-seasonalized anomaly in the absorbed solar

493 radiation (ASR, positive for downward fluxes), measured as incident minus reflected shortwave  
494 radiation. Following the Pinatubo eruption, observations show a dramatic reduction in ASR, due  
495 to increased scattering of sunlight to space from volcanic aerosols, not fully recovering until  
496 mid-1994. Our SDVOLC simulation calculates a remarkably similar reduction and recovery in  
497 ASR. The FRVOLC ensemble average shows a similar reduction in ASR similar, and the  
498 ensemble range shows the role of other unconstrained climate variables, including ocean states.  
499 The SDVC simulation reveals the effects of constrained sea surface and atmospheric  
500 temperatures on ASR variability without volcanic forcing.

501 Figure 11b shows the de-seasonalized anomaly in the outgoing longwave radiation (OLR,  
502 positive for upward fluxes). Pinatubo aerosols reduced OLR by both direct absorption of  
503 longwave radiation, and by reducing temperatures in the troposphere and at the Earth's surface.  
504 The SDVC simulation includes the latter effect, as it is nudged and driven by observed  
505 tropospheric and sea surface temperatures, which include this cooling. This cooling reduces OLR  
506 by up to  $1.5\text{-}2.0\text{ W m}^{-2}$  by August 1992. Inclusion of volcanic aerosols in our SDVOLC  
507 simulation, however, is necessary to match the observed reduction of  $2.5\text{-}3.0\text{ W m}^{-2}$ . The  
508 differences in the OLR between the SDVC and SDVOLC simulations are due to aerosol  
509 longwave absorption, secondary effects of aerosols on clouds, and cooling of land surface  
510 temperatures.

511 Figure 11c shows net radiative flux (ASR-OLR, positive for downward fluxes), a measure of the  
512 radiative energy imbalance forcing the Earth's climate. The SDVC case shows natural  
513 variability, with a slight upward trend due to increases in greenhouse gases. The SDVOLC  
514 shows a drop in the net flux following the Pinatubo eruption which generally matches well the  
515 observations. The FRVOLC ensemble shows a similar average reduction, and significant

516 variability. In general, the radiative response to Pinatubo in WACCM is a significant  
517 improvement over previous models using prescribed volcanic forcing [*Driscoll et al.*, 2012;  
518 *Neely et al.*, 2016]. This is due to both the limitations of prescribed stratospheric aerosol  
519 climatologies derived from satellite observations [*Ridley et al.*, 2014; *Mills et al.*, 2016], and to  
520 the neglect of aerosol-cloud interactions.

521 The efficacy of volcanic forcing in climate models is quantified by normalizing changes in all-  
522 sky net radiative fluxes to changes in SAOD. The Fifth Assessment Report of the  
523 Intergovernmental Panel on Climate Change [*Myhre et al.*, 2013] uses the value of  $-25 \text{ W m}^{-2}$  per  
524 unit change in volcanic SAOD, based on fixed sea-surface temperature simulations of the  
525 Pinatubo eruption in GISS Model E with prescribed stratospheric aerosol [*Hansen et al.*, 2005].  
526 ULAQ-CCM, with prognostic volcanic aerosols, calculates volcanic forcing efficiencies for  
527 Pinatubo of  $-15.3 \text{ W m}^{-2} \text{ SAOD}^{-1}$  in all-sky conditions [*Pitari et al.*, 2016b]. We calculate the  
528 efficacy of Pinatubo volcanic forcing in WACCM by linearly regressing the differences between  
529 volcanic and clean simulations in annually averaged top-of-model all-sky net fluxes and global  
530 SAOD for the years 1991-1996. For SDVOLC minus SDVC, we calculate  $-18.3 \pm 1.0 \text{ W m}^{-2}$   
531  $\text{SAOD}^{-1}$ . This indicates a reduced volcanic radiative forcing efficacy in WACCM compared to  
532 *Hansen et al.* [2005], which neglected the interaction of volcanic aerosol with clouds, and a  
533 greater efficacy compared to *Pitari et al.* [2016b]. For SCVOLC minus SCVC, we calculate -  
534  $20.2 \pm 4.6 \text{ W m}^{-2} \text{ SAOD}^{-1}$ , suggesting that interactive chemistry is not a significant factor in  
535 volcanic radiative forcing efficacy in WACCM.



### 536 3.5 Volcanic impacts on stratospheric ozone

537 The Antarctic ozone hole, defined as a region where total column ozone measures less than 220  
538 Dobson Units (DU), has developed each austral spring since it first developed in the early 1980s.  
539 The area of the ozone hole is at its largest in October, when total column ozone over Antarctica  
540 reaches minimum annual values. These minimum values depend on the amount of halogen  
541 loading in the stratosphere, as well as meteorology, with greater ozone loss in colder years. In  
542 addition, ozone loss increases when enhanced sulfate aerosol levels from volcanic eruptions  
543 reach the Antarctic stratosphere, due to the effects of heterogeneous chemistry on halogens  
544 [Portmann *et al.*, 1996; Solomon *et al.*, 2016].

545 Figure 12 shows observations of total column ozone measured from the Solar Backscatter Ultra-  
546 Violet satellite (SBUV), averaged over 63-90°S, from 1980 to 2015. The SBUV record has been  
547 carefully calibrated and compared to observations from ground-based, *in situ*, and other satellite  
548 instruments [McPeters *et al.*, 2013]. The observations show the development of the ozone hole in  
549 the 1980s, with significant interannual variability depending on temperature and volcanic aerosol  
550 loading following the eruptions of El Chichón (1982) and Pinatubo (1991). The FRVOLC  
551 ensemble reproduces the magnitude of the decline in Antarctic ozone in October from 1980 to  
552 the mid-1990s, and indicates significant drops following these two major tropical eruptions.  
553 Ozone loss leveled off after a peak in the late 1990s, and FRVOLC reproduces this general trend  
554 in the observations, although ozone columns are generally biased low throughout this simulation,  
555 consistent with the cold bias in the Antarctic spring stratosphere (see Figure 2).

556 Comparison of the SDVOLC and SDVC simulations, which were both initialized from FRVOLC  
557 on January 1, 1990, shows ~40 DU of ozone loss attributable to the 1991 eruptions of Pinatubo

558 and Cerro Hudson. Both SD simulations match the interannual variability in Antarctic ozone  
559 particularly well in the period of reduced and moderate volcanic aerosols post-2000. Significant  
560 drops in ozone followed the eruptions of Puyehue-Cordón Caulle (2011, 40.6°S, 72.1°W) and  
561 Calbuco (2015, 41.3°S, 72.6°W). While comparison of SDVOLC to SDVC in 2011 and 2015  
562 shows significant effects due to volcanic aerosols, cold temperatures also played a role, as  
563 revealed by significant drops in the SDVC ozone columns in those years. Because ozone heats  
564 the stratosphere, cold stratospheric temperatures are a positive feedback of polar ozone loss, and  
565 thermal and dynamical feedbacks may enhance the loss of polar ozone following volcanic  
566 eruptions [Solomon *et al.*, 2016; Ivy *et al.*, 2017; Pitari *et al.*, 2016a].

## 567 **4 Conclusions**

568 We have described a new version of WACCM with improved horizontal resolution, updated  
569 physics, and an interactive QBO. We have validated the chemistry and climate of WACCM with  
570 detailed comparisons to observations. We have paid particular attention to the evolution and  
571 impacts of volcanic sulfate aerosol, which WACCM derives from emissions of SO<sub>2</sub> gas. The  
572 completeness of the chemistry, dynamics, and aerosol microphysics qualify WACCM for studies  
573 of stratospheric sulfate geoengineering.

574 Our calculations reveal the importance of interactive chemistry to the development of sulfate  
575 aerosol from large inputs of SO<sub>2</sub>. Previous findings of a ~30-day e-folding decay time of SO<sub>2</sub>  
576 from the 1991 eruption of Mt. Pinatubo were based on observations that ignored the rapid initial  
577 losses of SO<sub>2</sub> on ice and ash, and on calculations that either did not include interactive OH  
578 chemistry, or did not discuss the impacts of OH depletion in the first month after the eruption.  
579 We show that the dense SO<sub>2</sub> cloud oxidized much more slowly in the first 2 weeks after the

580 eruption due to the depletion of OH by the SO<sub>2</sub> oxidation itself. We calculate a 47-day average  
581 initial e-folding decay time for the Pinatubo SO<sub>2</sub> that remained aloft after the initial losses on ice  
582 and ash, and show the calculated evolution of Pinatubo SO<sub>2</sub> to be largely consistent with  
583 observations. The evolution of stratospheric AOD following the Pinatubo eruption in WACCM  
584 agrees well with lidar observations from 3 independent locations at mid-latitudes and in the  
585 tropics. The radiative impacts of Pinatubo on ASR and OLR match satellite observations very  
586 well. This is crucial for assessing the impacts of stratospheric SO<sub>2</sub> injections on surface climate  
587 and on stratospheric chemistry and dynamics.

588 We have validated the climate and chemistry in an ensemble of fully coupled WACCM  
589 simulation of the years 1975-2016. The trend in global average surface temperatures over this  
590 period closely matches that derived from observations. Temperatures and winds from the  
591 troposphere through the middle atmosphere agree well with observations. WACCM now  
592 includes an internally generated QBO, which exhibits a period close to that observed. This  
593 feature is important to studies of stratospheric sulfate geoengineering, which has been shown in  
594 other studies to disrupt the QBO. Stratospheric water vapor in WACCM is close to that  
595 observed, as is the seasonal cycle in water vapor mixing ratios entering from the troposphere.  
596 Stratospheric ozone columns in WACCM agree well with global satellite observations, and with  
597 ground-based observations in Antarctica showing the development of the ozone hole over this  
598 period.

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617 simulations, may be obtained from M.J.M. (email: [mmills@ucar.edu](mailto:mmills@ucar.edu)).

618 **Tables**

619 *Table 1: WACCM simulations conducted for this work.*

<b>Simulation name</b>	<b>Dynamics</b>	<b>Chemistry</b>	<b>Ocean/ Sea Ice</b>	<b># of runs</b>	<b>Years per run</b>	<b>Conditions</b>	<b>SO<sub>2</sub> from eruptions</b>
<b>FRPI</b>	Free-running	Interactive	Coupled	1	50	Pre-industrial (1850)	No

<b>FRVOLC</b>	Free-running	Interactive	Coupled	4	42	1975-2016	Yes
<b>SDVOLC</b>	Nudged	Interactive	Data	1	27	1990-2016	Yes
<b>SDVC</b>	Nudged	Interactive	Data	1	27	1990-2016	No
<b>SCVOLC</b>	Free-running	Prescribed	Data	1	21	1979-1999	Yes
<b>SCVC</b>	Free-running	Prescribed	Data	1	10	1990-1999	No

620

## 621 Figure captions

622 Figure 1: Global annual surface temperature anomalies for 1979-2015 from the WACCM  
623 FRVOLC ensemble are compared to HadCRUT4 and GISTEMP reconstructions of observations.  
624 Anomalies are calculated with respect to the 1979-2015 average for each data set. Least squares  
625 linear fit trends are listed and plotted.

626 Figure 2: WACCM temperature differences from ERA-Interim reanalysis for DJF, MAM, JJA,  
627 and SON. Differences are plotted for the 1979-2014 time period from the FRVOLC ensemble  
628 average. The contour interval is 1K.

629 Figure 3: WACCM (top panels) and URAP (bottom panels) zonal mean wind for DJF (left  
630 panels) and for JJA (right panels). WACCM winds are averaged over the years 1980-1999 from  
631 the FRVOLC ensemble average. Contour interval is  $10 \text{ m s}^{-1}$ .

632 Figure 4: Tropical zonal winds ( $2^{\circ}\text{S}$ - $2^{\circ}\text{N}$ ) from 1980 to 2000 for ERA-Interim reanalysis (a) and  
633 WACCM FRVOLC ensemble members (b, c, d, and e). Contours are plotted in intervals of  $5 \text{ m}$   
634  $\text{s}^{-1}$ .

635 Figure 5: (a) Height-time seasonal variations of  $\text{H}_2\text{O}$  mixing ratios (ppmv) averaged between  
636 latitudes  $10^{\circ}\text{N}$ - $10^{\circ}\text{S}$  and years 2004-2014 for (a) the WACCM FRVOLC ensemble average  
637 (color filled white contours) and MLS satellite observations (black contours); (b) percent  
638 difference between the ensemble average and observations,  $100 * (\text{FRVOLC} - \text{MLS}) / \text{MLS}$ ; (c)  
639 2004-2014 mean vertical profiles of  $\text{H}_2\text{O}$  mixing ratios (ppmv) from MLS (black) and FRVOLC  
640 (red).

641 Figure 6: Monthly and zonally averaged stratospheric ozone column (in DU) comparison  
642 between OMI/MLS observations between 2004 and 2010 (black) and the WACCM FRVOLC  
643 ensemble (red) between 2004 and 2010 (for ozone < 150 ppb in the model), for four months.  
644 OMI/MLS error bars show the zonally averaged 2- $\sigma$  six-year root mean square standard error of  
645 the mean at a given grid point, derived from the gridded product [Ziemke *et al.*, 2011]. Model  
646 results are interpolated to the same latitude grid as the observations. Shading indicates the  
647 standard deviation (1- $\sigma$ ) of the interannual variability per latitude interval for the FRVOLC  
648 ensemble.

649 Figure 7: Calculated global volcanic SO<sub>2</sub> burden following the June 15, 1991 eruption of Mt.  
650 Pinatubo is compared to observations. The solid line shows the daily average global burden of  
651 SO<sub>2</sub> calculated in the SDVOLC simulation minus the non-volcanic SO<sub>2</sub> burden calculated in the  
652 SDVC simulation. The dashed red line shows the same for SCVOLC minus SCVC. The SO<sub>2</sub>  
653 burden is less than 10 Tg in the day 0 average because the eruption occurred mid-day. The  
654 additional input of 1.5 Tg SO<sub>2</sub> from the August 12 eruption of Cerro Hudson is noted 60 days  
655 after the Pinatubo eruption. Observations from TOVS (blue circles) and TOMS (red asterisks)  
656 show an initial burden of 13-18 Tg SO<sub>2</sub>, of which 10 Tg remained after loss to sedimenting ice  
657 and ash in the first 7-9 days [Guo *et al.*, 2004a]. Observations from SBUV, aircraft, and MLS are  
658 shown as presented in Read *et al.* [1993]. Because MLS column is integrated above 21 km, the  
659 WACCM column integrated above 50 hPa (dashed black line) is shown for comparison.

660 Figure 8: Maps of calculated daily averaged SO<sub>2</sub> (left column) and OH (right column) volume  
661 mixing ratios (moles/mole air) at 61 hPa on days 3 (top row), 7 (middle row), and 13 (bottom  
662 row) after the June 15, 1991 eruption of Mt. Pinatubo. Calculations are shown from the  
663 SDVOLC simulation.

664 Figure 9: Volcanic SO<sub>2</sub> e-folding time (days) shown as a function of days following the June 15,  
665 1991 eruption of Mt. Pinatubo in the SCVOLC (solid black line) and SCVOLC (dashed red line)  
666 simulations. The e-folding time is derived from the daily change in the global volcanic SO<sub>2</sub>  
667 burden. Volcanic SO<sub>2</sub> is calculated by subtracting the global burdens from volcanically clean  
668 simulations (SDVC and SCVC, respectively).

669 Figure 10: Aerosol optical depth (AOD) measured by lidars at 3 locations (black circles) are  
670 compared to calculated 5-day average AOD above the tropopause in corresponding model  
671 columns from our SDVOLC (red dots), FRVOLC (ensemble average, blue dots), and SCVOLC  
672 (orange dots) simulations. Observations are (a) integrated backscatter from 15-30 km measured  
673 in Tomsk, Siberia [Zuev *et al.*, 2016], converted to AOD using a lidar ratio of 50; (b) AOD above  
674 the tropopause measured in Geestacht, Germany [Ansmann *et al.*, 1997]; and (c) AOD above the  
675 tropopause measured in Mauna Loa, Hawaii [Hofmann *et al.*, 2009; Ridley *et al.*, 2014].

676 Figure 11: Top-of-model all-sky radiative fluxes from our SDVOLC (solid red) and SDVC (solid  
677 blue) simulations are compared to top-of-atmosphere ERBS observations (black) merged with  
678 additional data to provide a global dataset [Allan *et al.*, 2014; Liu *et al.*, 2015]. Monthly mean net  
679 fluxes are shown for January 1991 to December 1995, normalized and de-seasonalized by  
680 subtracting the corresponding flux for each month from 1999, a volcanically quiescent year.  
681 Fluxes from our FRVOLC ensemble average (dashed orange line) and range (yellow shading)  
682 are also shown. (a) Absorbed solar radiation (positive for downward fluxes); (b) outgoing  
683 longwave radiation (positive for upward fluxes); (c) net radiative flux (positive for downward  
684 fluxes).



685 *Figure 12: October monthly average column ozone over the South Polar cap (63-90°S) for years*  
686 *1980-2016 from SBUV satellite observations (black solid line and circles), and in WACCM*  
687 *SDVOLC (red solid line and diamonds), SDVC (blue dashed line and diamonds), and FRVOLC*  
688 *experiments. The orange dashed line shows the ensemble average, and yellow shading shows the*  
689 *ensemble range. Grey dots show monthly averages for individual FRVOLC ensemble members.*

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Figure 1.

Surface Temperature Anomaly (K)

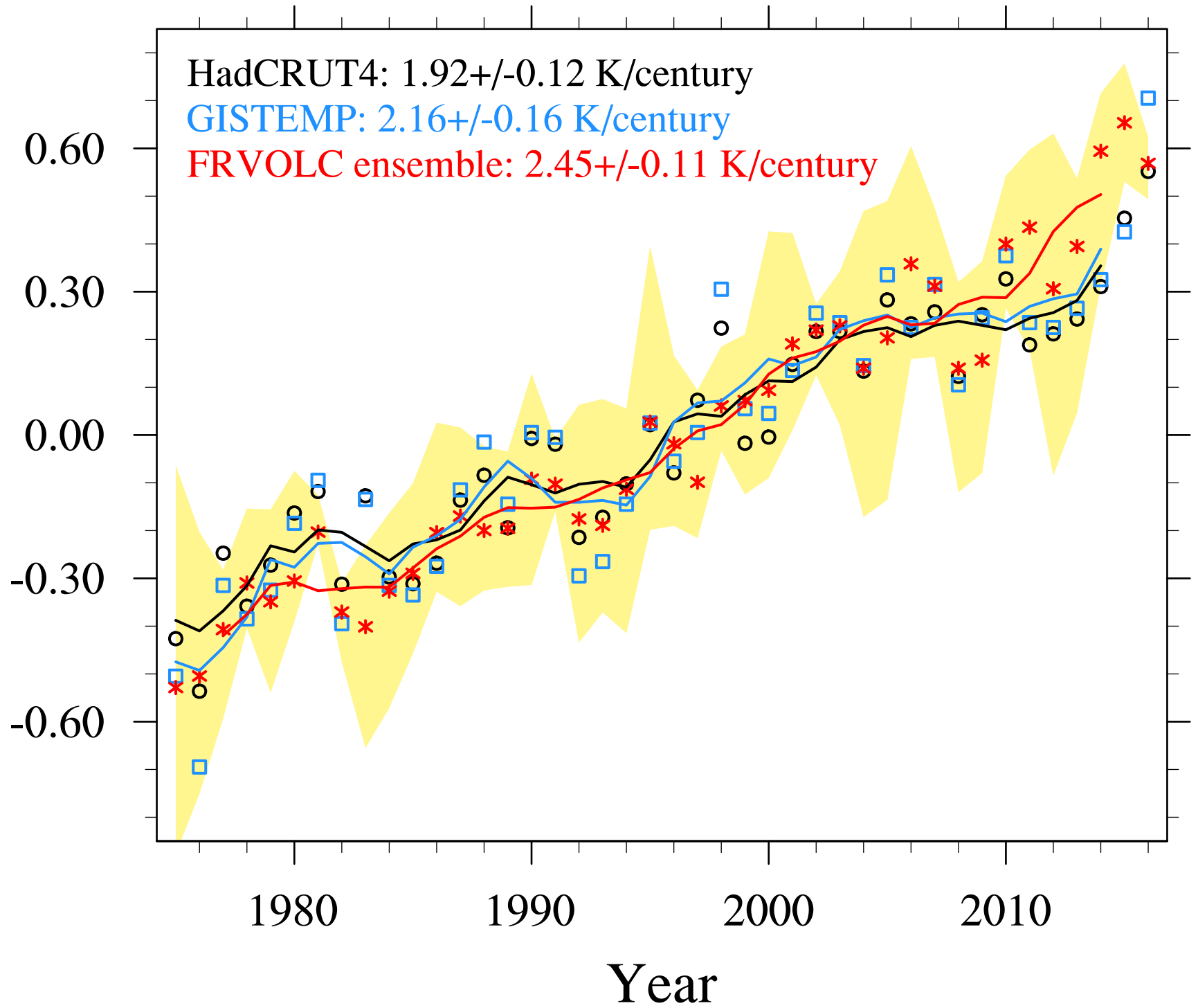


Figure 2.

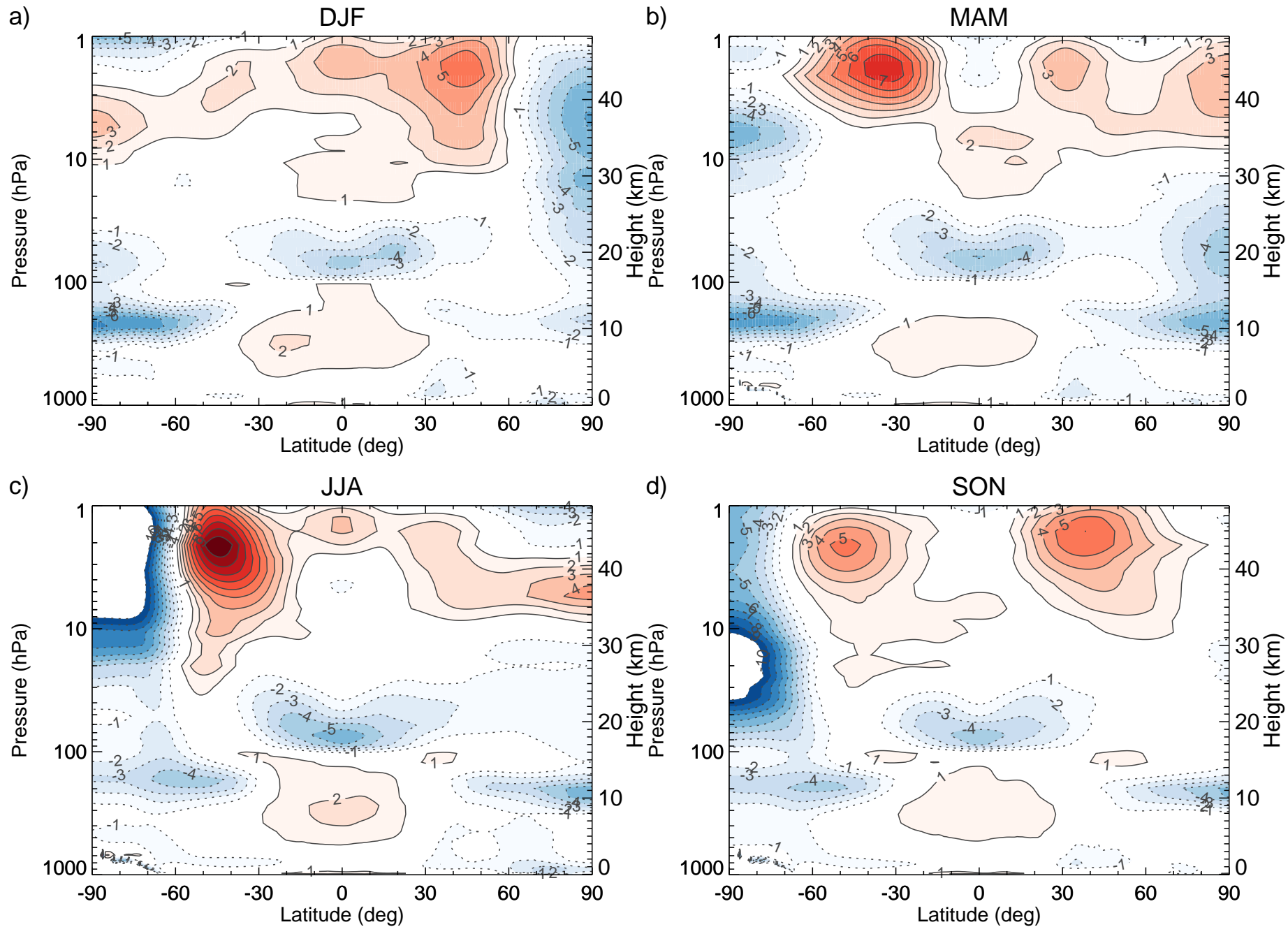


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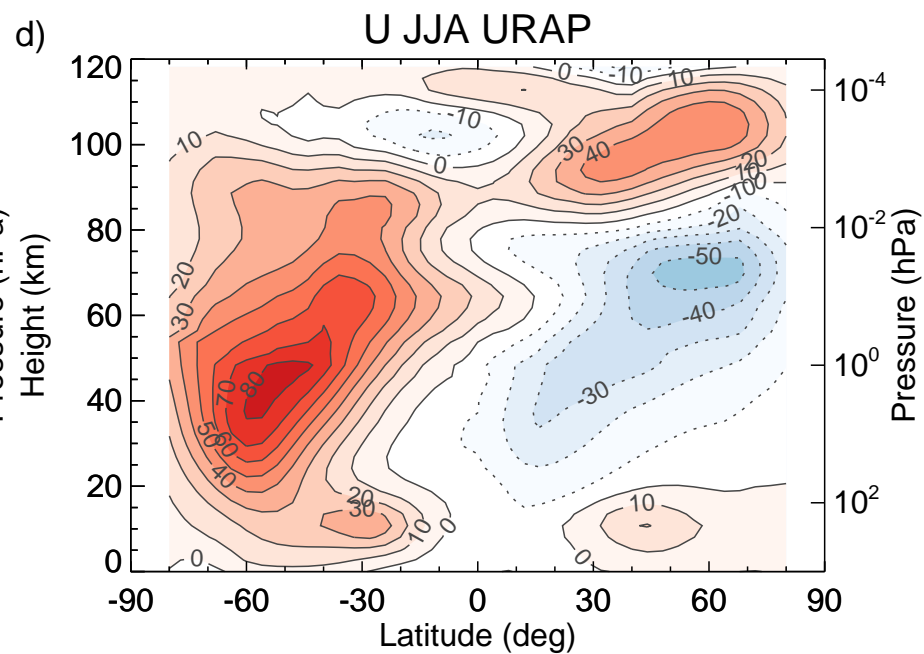
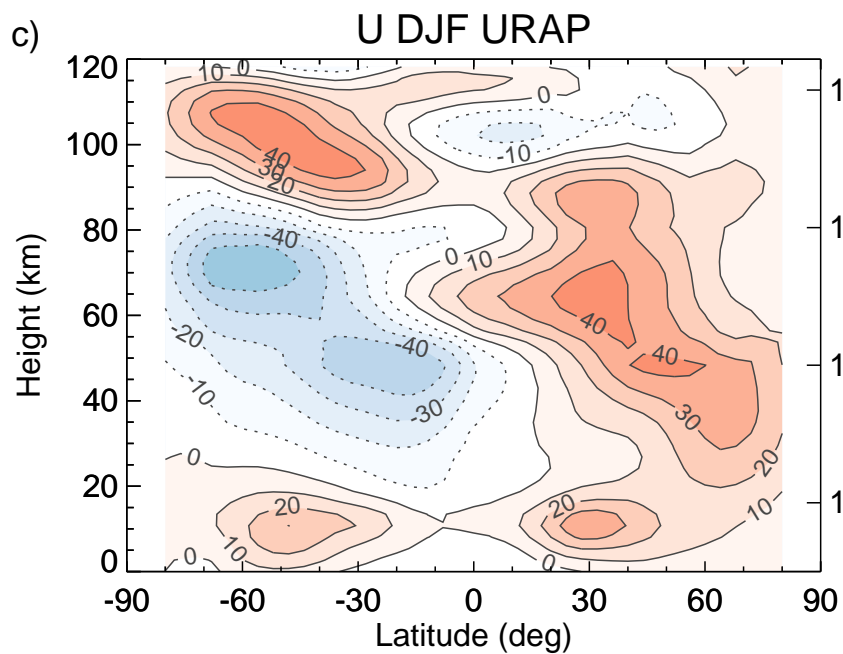
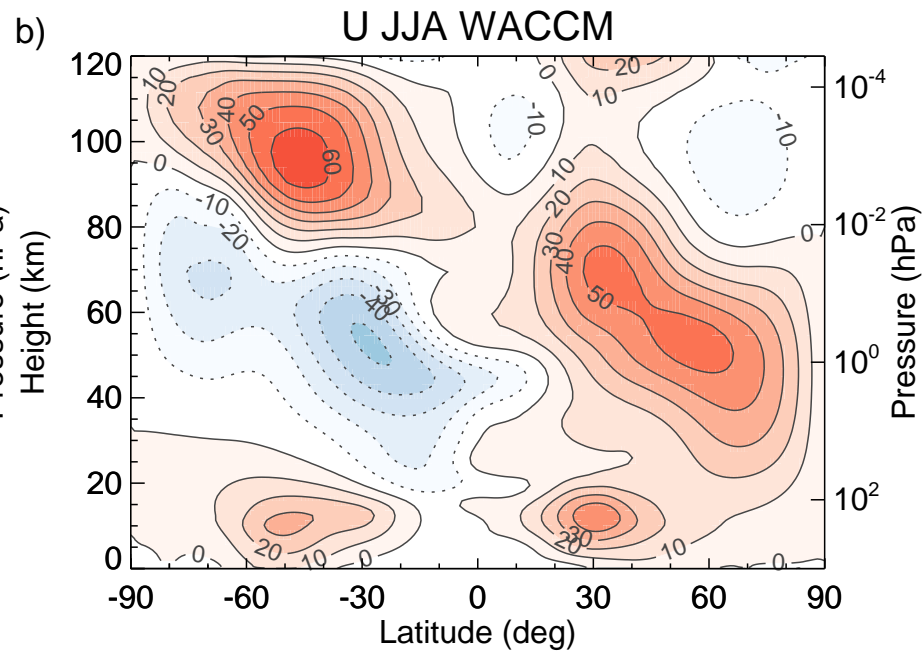
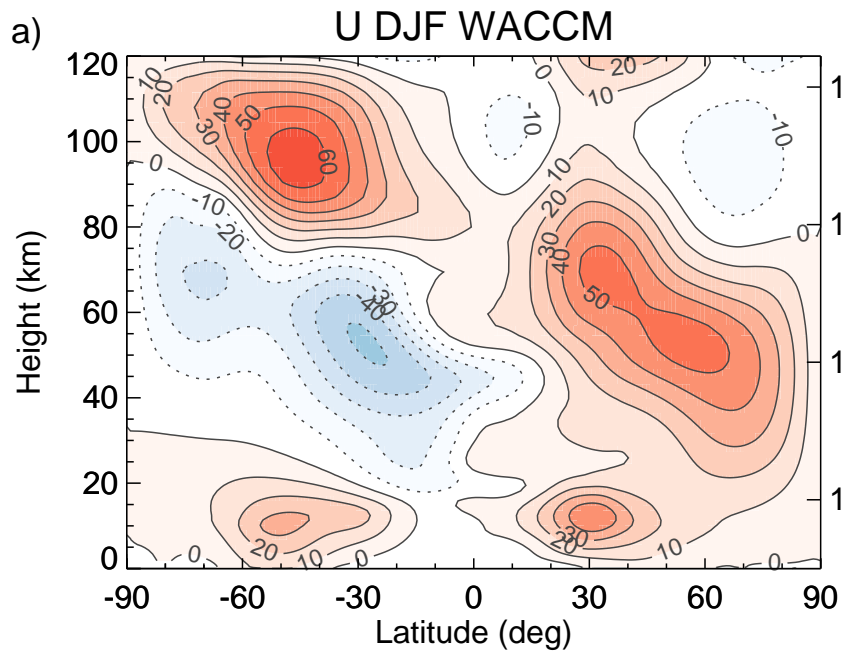
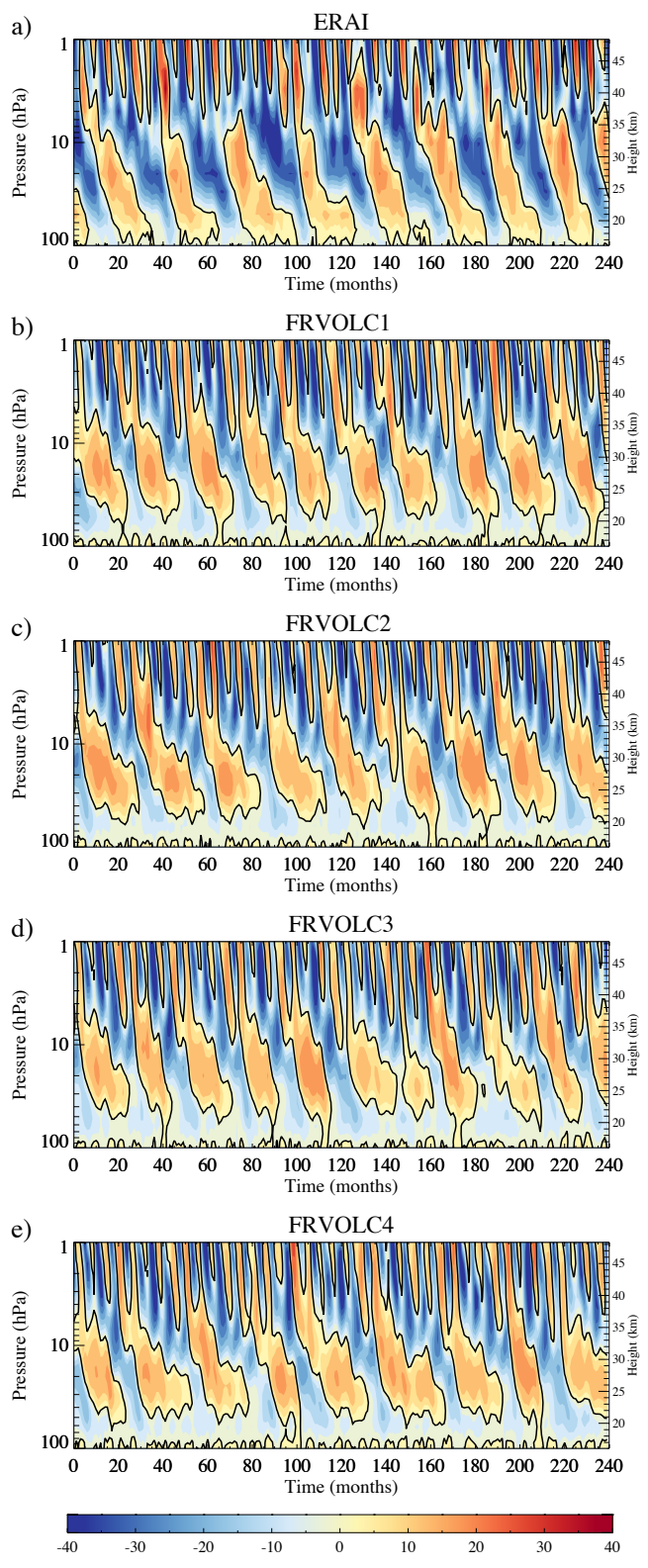


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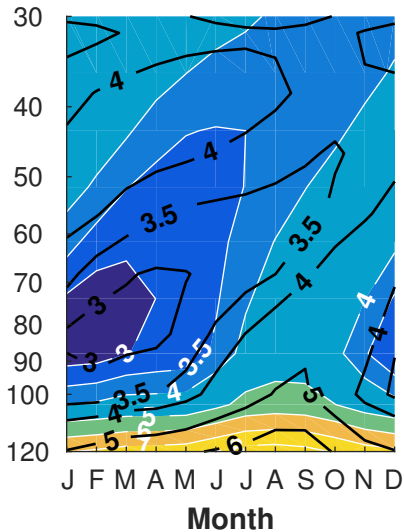




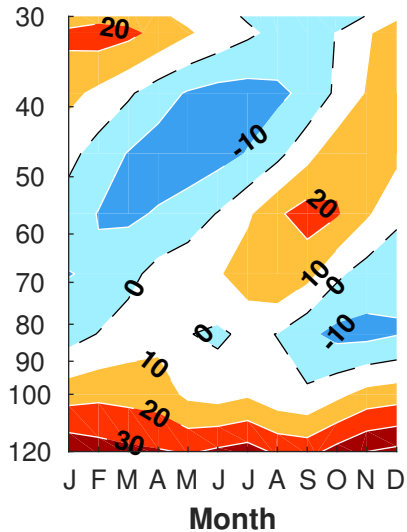
Zonal Wind ( $\text{m s}^{-1}$ )

Figure 5.

**(a) Water Vapor, ppmv**



**(b) Difference, %**



**(c) 2004-2014 Mean**

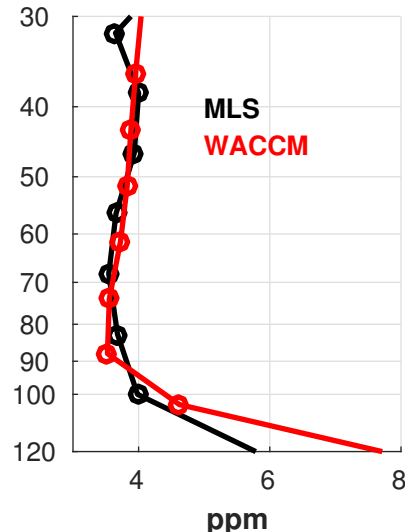
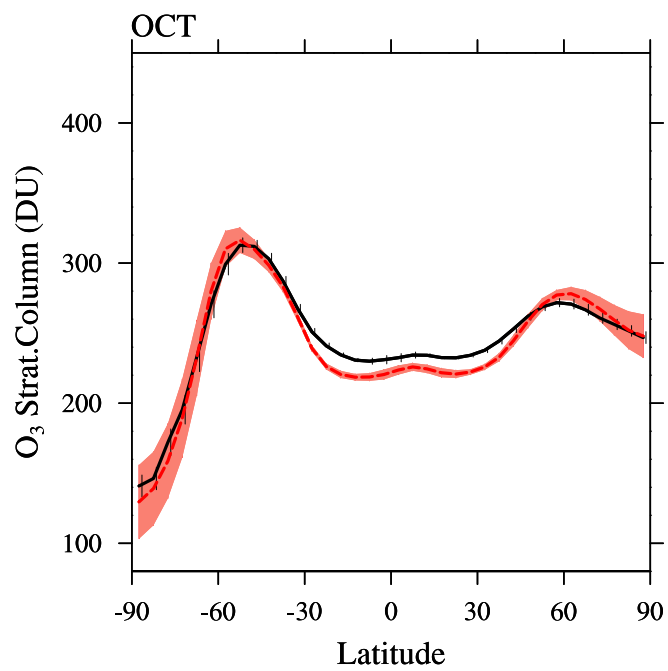
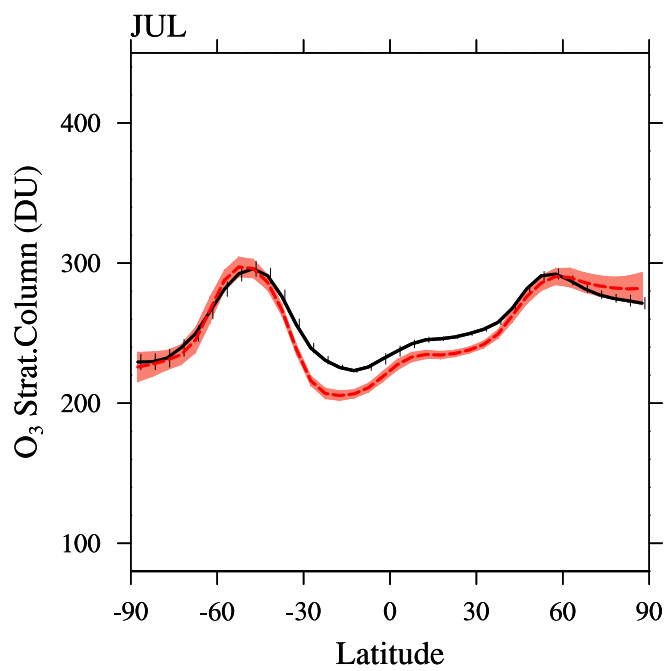
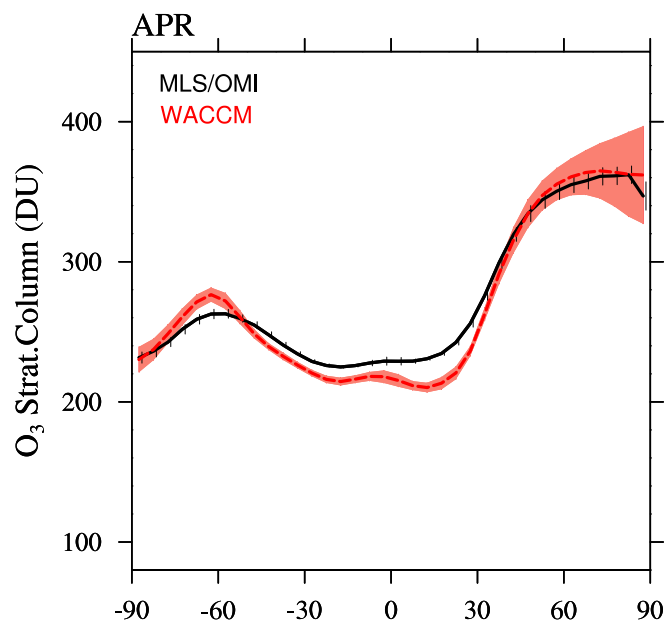
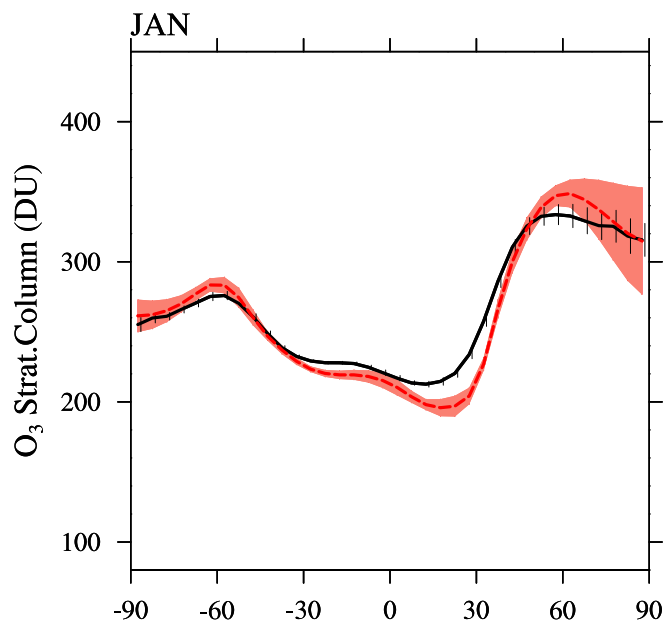


Figure 6.



**Figure 7.**

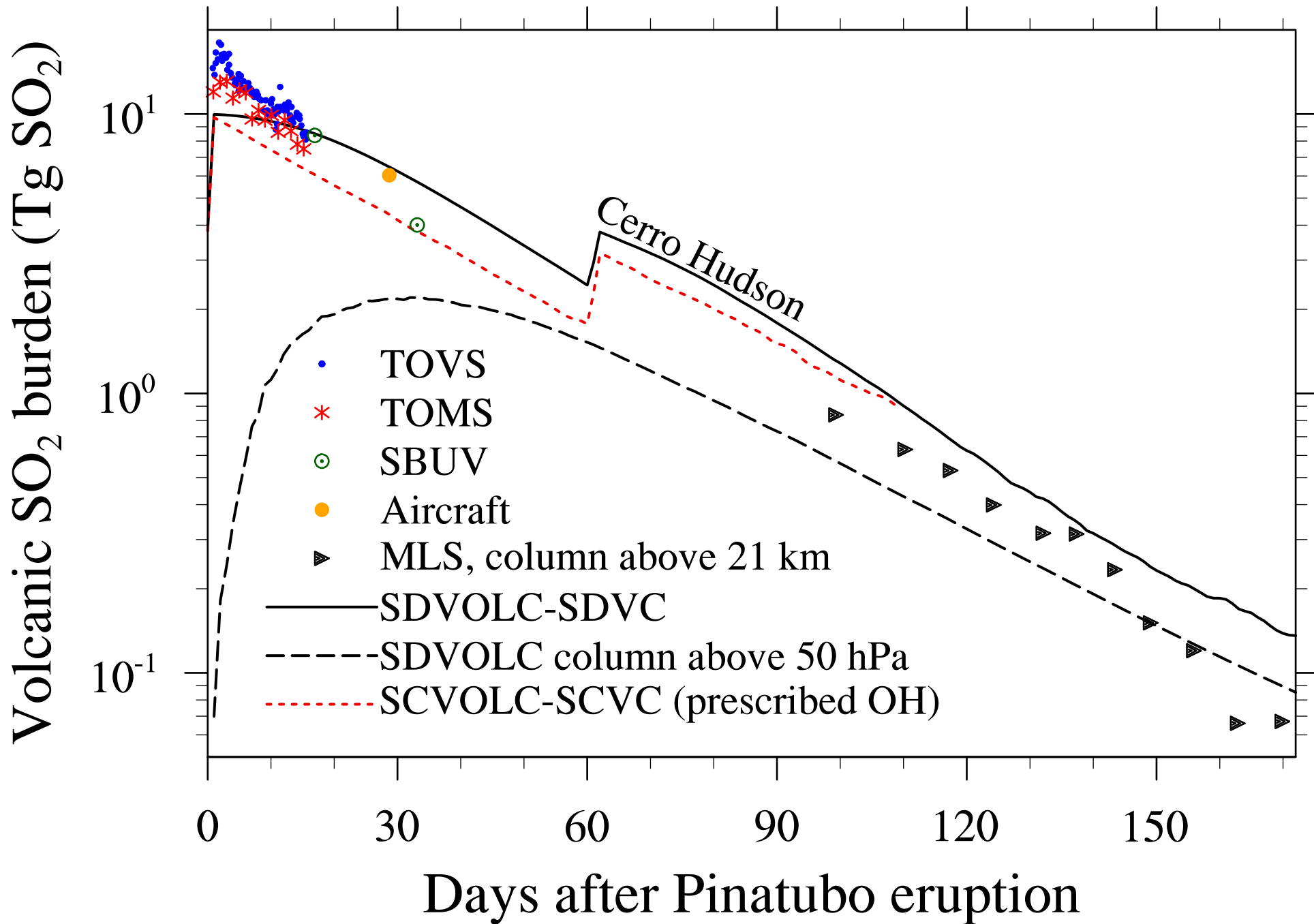


Figure 8.



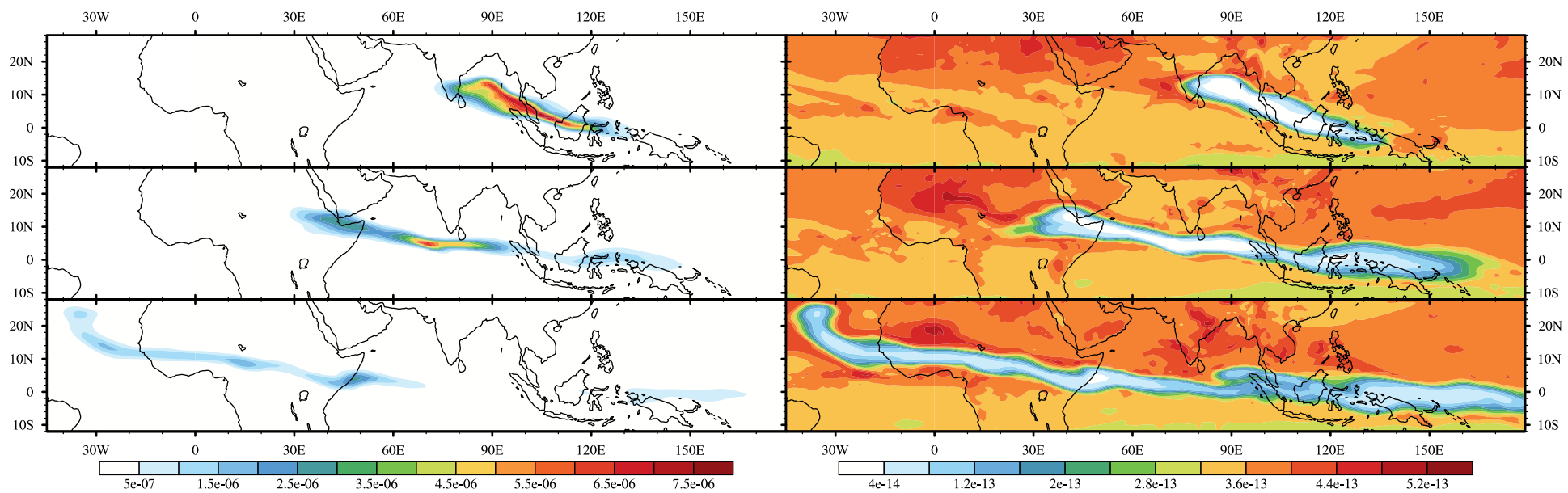


Figure 9.

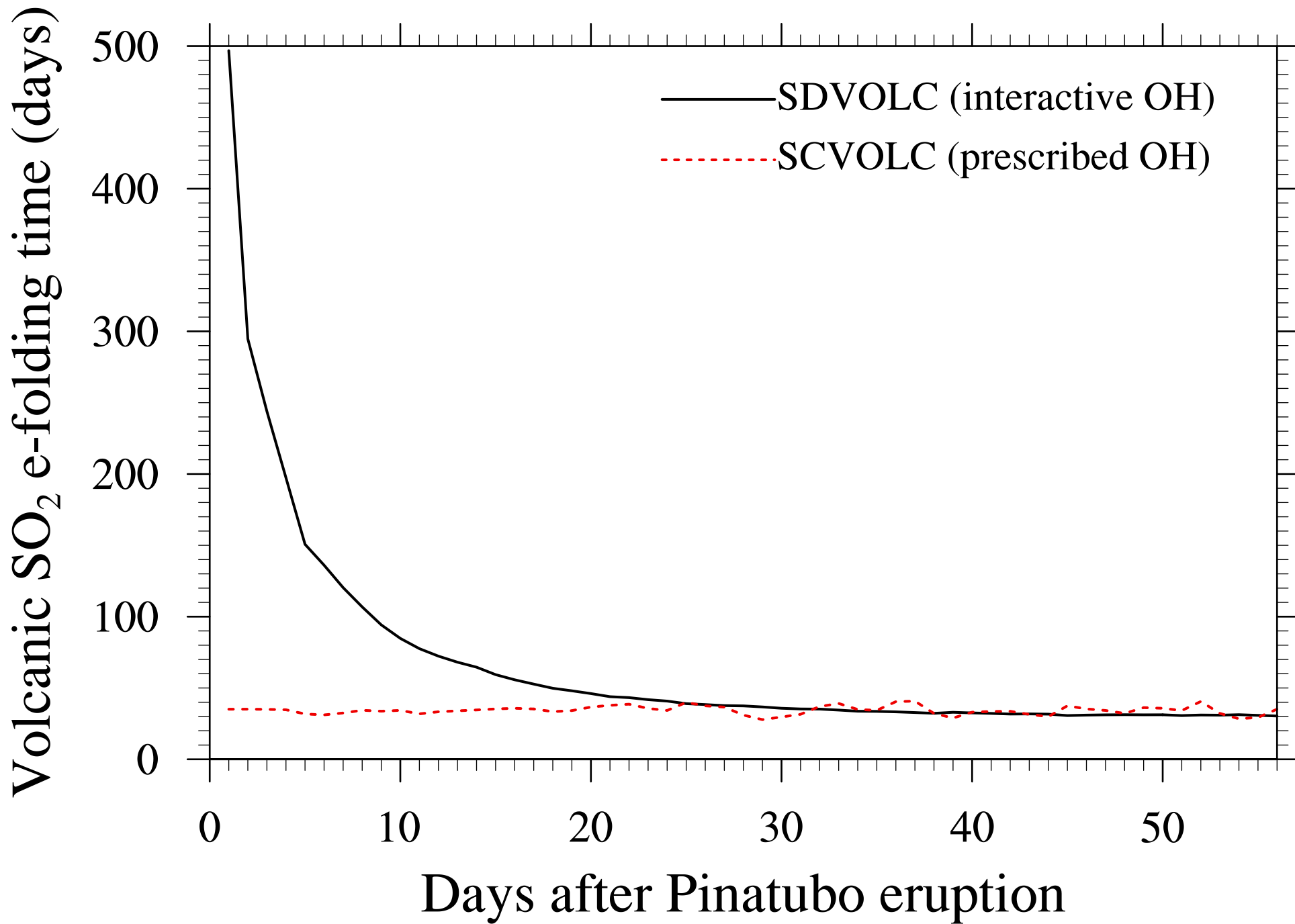


Figure 10.

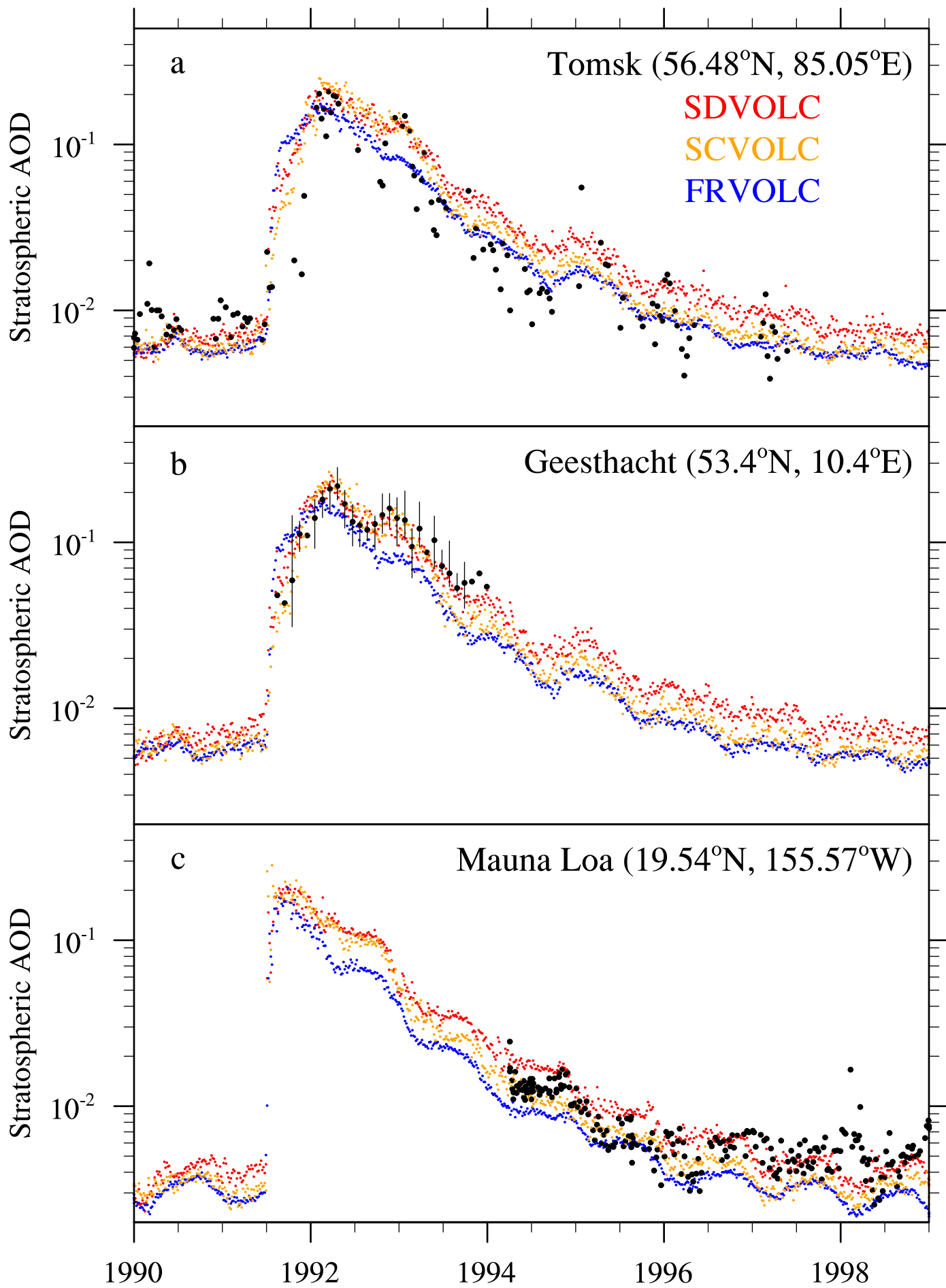


Figure 11.

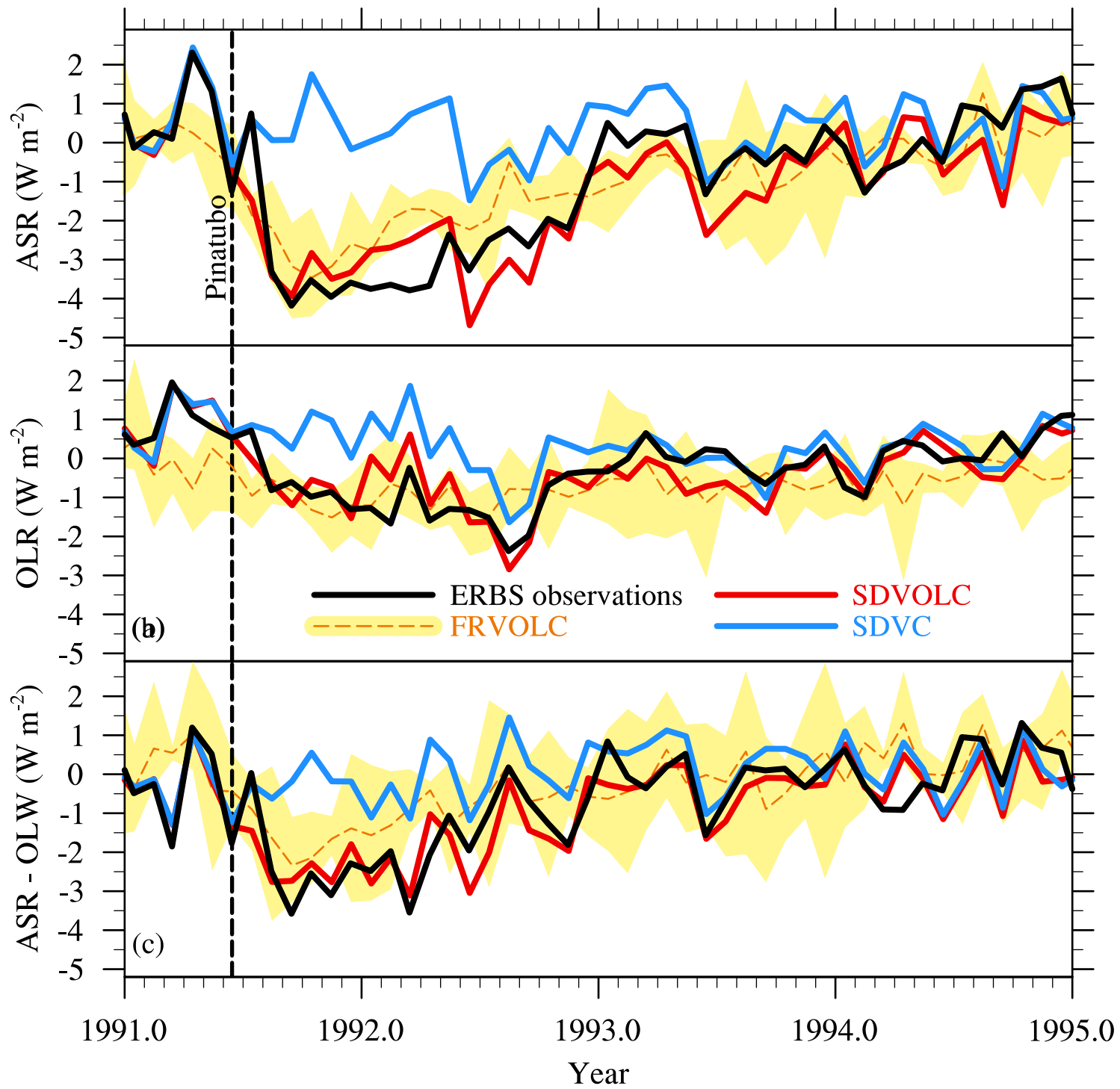


Figure 12.



# October, South Polar Cap Average

