

submitted to *Geophys. J. Int.*

A continuous map of near-surface S-wave attenuation in New Zealand

Chris Van Houtte^{1,2*}, Olga-Joan Ktenidou^{3,4}, Tam Larkin² and Caroline Holden¹

¹ *GNS Science, 1 Fairway Drive, Lower Hutt 5010, New Zealand*

² *Department of Civil Engineering, University of Auckland, 20 Symonds Street, Auckland 1010, New Zealand*

³ *Department of Engineering Science, University of Greenwich, Medway Campus, Central Avenue, Chatham Maritime,*

⁴ *Institut des Sciences de la Terre, ISTerre, Universite de Grenoble 1, CNRS, F-38041 Grenoble, France*

Received 2017 March 1

SUMMARY

Quantifying the near-surface attenuation of seismic waves at a given location can be important for seismic hazard analysis of high-frequency ground motion. This study calculates the site attenuation parameter, κ_0 , at 41 seismograph locations in New Zealand. Combined with results of a previous study, a total of 46 κ_0 values are available across New Zealand. The results compare well with previous t^* studies, revealing high attenuation in the volcanic arc and forearc ranges, and low attenuation in the South Island. However, for site-specific seismic hazard analyses, there is a need to calculate κ_0 at locations away from a seismograph location. For these situations, it is common to infer κ_0 from weak correlations with the shear-wave velocity in the top 30 m, V_{S30} , or to adopt an indicative regional value. This study attempts to improve on this practice. Geostatistical models of the station-specific κ_0 data are developed, and continuous maps are derived using ordinary kriging. The obtained κ_0 maps can provide a median κ_0 and its uncertainty for any location in New Zealand, which may be useful for future site-specific seismic hazard analyses.

Key words: kappa, attenuation, seismic hazard, kriging, quality factor, site effects

2 *C. Van Houtte et al.*

16 **1 INTRODUCTION**

17 The expected level of rock-site ground motion due to a given earthquake rupture is often needed
 18 for use in site-specific hazard studies and to develop building codes. In many cases, an empirical
 19 ground motion model is used to predict the behaviour of an intensity measure from some source,
 20 path, and site response parameters. For the site response, the effect of the soil column has been
 21 extensively studied, however the site response of a rock site can also greatly affect surface ground
 22 motions (Steidl et al. 1996; Boore & Joyner 1997; Ktenidou & Abrahamson 2016). In the case
 23 of soil sites, modern empirical ground motion models such as the Next Generation Attenuation
 24 (NGA) West-1 and West-2 models characterise the site response using the time-averaged shear-
 25 wave velocity in the first 30 m below ground surface (V_{S30}) and the depth to a shear-wave velocity
 26 horizon of 1 km/s ($Z_{1.0}$) as the predictive parameters. For rock-site response, empirical models
 27 typically use a single site factor for all rock sites (e.g., Abrahamson & Silva 1997; McVerry et al.
 28 2006), or more recently, consider the rock and soil site response as a continuous function of V_{S30}
 29 (Abrahamson et al. 2014; Boore et al. 2014; Campbell & Bozorgnia 2014; Chiou & Youngs 2014).

30 However, rock-site response is not currently well-constrained in empirical ground motion mod-
 31 els. Firstly, rock site V_{S30} measurements are often unavailable or unreliable, both for model de-
 32 velopers and for users. Secondly, research by Anderson et al. (1996), Douglas et al. (2009) and
 33 Laurendeau et al. (2013) shows that rock-site response cannot be well-modelled using V_{S30} alone,
 34 as it also depends on the site attenuation parameter, κ_0 , which models the attenuation of approxi-
 35 mately vertically-propagating seismic waves beneath the ground surface. κ_0 is often derived from
 36 the spectral decay parameter, κ (Anderson & Hough 1984), which controls the steepness of the
 37 decay of earthquake Fourier amplitudes in the commonly-used high frequency filter model,

$$A(f) = A_0 \exp(-\pi\kappa f), \quad f > f_e \quad (1)$$

38 where A_0 is a source- and path-dependent Fourier amplitude of acceleration, f is the frequency,
 39 f_e is the frequency above which the decay is approximately linear on a plot of $\log(A)$ against f ,

* Pacific Region Office, GJI

and κ controls the steepness of the spectral decay. In many studies, f_e is chosen at around 10 Hz. General consensus in the seismological community is that the predominant physical mechanism behind κ is attenuation along the wave propagation path, and thus κ is typically modelled as a function of a distance variable, r , and site variable, s (Anderson 1991), i.e.

$$\kappa(r, s) = \kappa_0(s) + \tilde{\kappa}(r) \quad , \quad (2)$$

where $\tilde{\kappa}(r)$ is the distance-dependence of κ and represents attenuation along the ray path before seismic waves enter the near-surface material, and $\kappa_0(s)$ is the site attenuation parameter. Papageorgiou & Aki (1983) initially suggested that κ_0 may be a source effect, rather than site effect, however later Abercrombie (1997) used borehole data to show that near-surface attenuation is likely to be the dominant contributor to $\kappa_0(s)$.

If it is assumed that the spectrum of Fourier amplitudes released by the source, the seismic quality factor Q , and the site amplification are frequency-independent for frequencies greater than f_e , that the observed high-frequency decay is due to the direct wave, and that Q is laterally homogeneous, then equation 2 can be rewritten as,

$$\kappa(r, s) = \int_{path} \frac{dr}{Q(z)V(z)} \quad , \quad (3)$$

(Hough & Anderson 1988), where z is the depth and V is wave velocity. Conditioned on these assumptions, $\kappa(r, s)$, hereafter denoted as κ , can be interpreted in terms of site and path attenuation alone.

The validity of these assumptions is rather contentious. Firstly, the assumption of frequency-independent source spectra above f_e may not be a good representation of reality. An ω -square source model (e.g., Housner 1947; Haskell 1964; Aki 1967; Brune 1970), for which the displacement amplitudes decay according to ω^{-2} for large ω , will result in flat amplitudes of the source spectrum of acceleration at frequencies greater than the source corner frequency. While analyses of the Canterbury earthquakes, for example, have found the ω -square model to be, on average, an excellent representation of sources in that region (Oth & Kaiser 2014), this may not hold for all

4 *C. Van Houtte et al.*

63 events and stations. It has been shown that the falloff rate of displacement spectra can vary strongly
64 across the focal sphere (Madariaga 1976; Boatwright 1980; Kaneko & Shearer 2014, 2015), which
65 means the assumptions of the κ model may not hold for all azimuths and takeoff angles. Despite
66 the uncertain nature of the high frequency shape of source spectra, this study assumes that an
67 ω -square source model applies to all events and stations used to calculate κ values.

68 For the κ model to represent attenuation, it is also necessary to assume that site effects above
69 f_e are negligible (Silva & Darragh 1995; Anderson et al. 1996; Ktenidou et al. 2017). Site effects
70 can be both narrow band, for example from a topographic effect or an impedance contrast in
71 the subsurface geology, or broadband amplification from an impedance gradient. Narrow band
72 amplification can to some extent be identified by spectral ratios, hence this study calculates the
73 horizontal-to-vertical spectral ratio from earthquake recordings (HVSR, Lermo & Chávez-García
74 1993) at each station before computing κ . If high frequency peaks in the HVSR are observed, it was
75 assumed that the model in equations 2 and 3 was not appropriate for the particular site, and the site
76 was subsequently discarded from the analysis. A broadband amplification is much more difficult
77 to detect, because it can typically only be inferred from detailed, site-specific, subsurface data. To
78 apply the κ model in New Zealand, where this information is unavailable, it is therefore necessary
79 to equivocally assume that broadband amplification at New Zealand rock sites is insignificant at
80 high frequencies.

81 Another assumption of the κ model is the frequency-independence of Q above f_e . The belief
82 that Q does depend on frequency has a long history in seismology (Aki 1980; Lay & Wallace
83 1995). It is typically assumed that Q^{-1} (defined as the fractional loss of energy per cycle) can be
84 divided into two parts, an intrinsic component that is related to complex microscopic processes
85 known collectively as ‘internal friction’, and a component arising from scattering of energy due
86 to heterogeneities distributed within a medium (Dainty 1981). It is rare for direct-wave studies to
87 separate the mechanisms, and Q is usually modelled as an effective attenuation operator Q_e^{-1} . It is
88 commonly observed that Q_e depends strongly on frequency, at least between around 1 and 10 Hz.
89 For this reason, Q_e is often modelled using the parametric function

$$Q_e(f) = Q_0 f^\eta, \quad (4)$$

where Q_0 is the value of Q_e at a reference frequency, such as 1 Hz, and η is an exponent controlling the degree of frequency-dependence. Atkinson (2012) and Mereu et al. (2013) compile values of η for numerous attenuation studies in eastern North America, and find that most studies have η values that range from 0.2 to 0.7. This indicates a strong dependence on frequency. However, several studies have also found that the frequency-dependence of Q_e is more complicated than the simple model in equation 4. In particular, coda Q studies using deep borehole data, which are subject to fewer tradeoffs than surface data, have shown that Q_e has a weaker dependence on frequency above 10 Hz (Leary & Abercrombie 1994; Adams & Abercrombie 1998; Abercrombie 1998; Yoshimoto & Okada 2009), although η values are still around 0.3-0.5. While the relationship between coda Q and direct-wave Q is not straight-forward (Yogomida & Benites 1995), it is nevertheless possible that direct-wave Q in the lithosphere has a more mild frequency-dependence above 10 Hz, than for mid-range frequencies. Whether the ‘10 Hz transition’ is related to Q attenuation or geometric attenuation is a matter for debate (Morozov 2008), however regardless of the origin, the consequences of the frequency-dependent attenuation on κ values will be the same. Given the trade-offs between source, path and site effects, testing the sensitivity of κ values (as defined by equation 1) to the value of η is difficult. κ values calculated using equation 1 are sensitive only to spectral shapes and not absolute amplitudes. If a model with $\eta > 0$ is fit to the high-frequency spectrum using equation 1, the fit to the data will not change significantly, but the obtained Q_0 value will differ greatly from an $\eta = 0$ model. However, even if a preferred value of $\eta > 0$ is propagated into the κ_0 calculations, extrapolating the κ_0 results to lower frequencies depends strongly on how one models the 10 Hz transition. Given these complications, this study follows the precedence of Anderson & Hough (1984) and many subsequent researchers, and assumes that $\eta = 0$ for the purposes of deriving κ_0 . The recent results of Perron et al. (2017) suggest the practical implications of this assumption may be minimal. These authors show that calculating κ in high and low frequency bands (on acceleration and displacement spectra respectively) yields different values of $\tilde{\kappa}(r)$ but similar values of hard-rock κ_0 , for two sites in France.

6 *C. Van Houtte et al.*

1
2
3
4
5
6 116 Despite these seemingly limiting assumptions, the κ model has become a popular method
7
8 117 for calculating site attenuation for engineering applications. The κ value itself is not of primary
9
10 118 interest, because it is usually preferred to model path attenuation using Q directly with the seismic
11
12 119 wave velocity V , according to

$$P(f) = A_0 \exp(-\pi f R / QV) , \quad (5)$$

13
14
15
16
17
18 120 where $P(f)$ is the attenuation function, A_0 is a reference amplitude, f is the frequency and R
19
20 121 is the distance along the ray path. However, the site attenuation parameter $\kappa_0(s)$ (from equation
21
22 122 2), subsequently denoted as κ_0 , has become a widely used parameter for modelling near-surface
23
24 123 attenuation at rock sites. κ_0 is used in ground motion simulations using the stochastic method (e.g.,
25
26 124 Boore 2003; Motazedian & Atkinson 2005; Graves & Pitarka 2010), and as a GMPE adjustment
27
28 125 parameter in the host-to-target method, to account for regional differences in rock site attenua-
29
30 126 tion between the host and target regions (Campbell 2003; Cotton et al. 2006; Douglas et al. 2006;
31
32 127 Van Houtte et al. 2011). For simplicity, these applications almost always assume that S-waves are
33
34 128 the dominant wave phase for peak motions, hence κ_0 , as well as Q and V , correspond to atten-
35
36 129 uation, quality factor and velocity for S-waves. Recently, Douglas et al. (2009) and Laurendeau
37
38 130 et al. (2013) corroborate the earlier work of Anderson et al. (1996) and show that the rock site
39
40 131 response is potentially modelled better using κ_0 in conjunction with V_{S30} . Therefore, including κ_0
41
42 132 in empirical ground motion models may improve rock site predictions.

43
44
45 133 Although κ_0 is a widely-used parameter in a multitude of ground-motion prediction applica-
46
47 134 tions, obtaining reliable data for forward prediction is often problematic. For site-specific hazard
48
49 135 studies, κ_0 cannot typically be calculated from earthquake data due to the lack of instrumentation,
50
51 136 and hence is often inferred from weak global $\kappa_0 - V_{S30}$ correlations (Silva & Darragh 1995; Chan-
52
53 137 dler et al. 2006; Van Houtte et al. 2011). While this method may have some benefit for ground
54
55 138 motion prediction compared to excluding κ_0 altogether, the benefits are likely to be minimal. Be-
56
57 139 fore κ_0 can be widely implemented in empirical ground motion prediction for a given region, it is
58
59 140 essential that it is well-defined and understood in the intended area of application. It is therefore
60
141 the objective of this study to provide nationwide κ_0 values for New Zealand, to help facilitate the

uptake of κ_0 as a parameter in empirical ground motion models. κ_0 values are calculated at individual stations across New Zealand, before kriging the data with a spatially-varying mean, to derive a continuous κ_0 map. The map can be used to infer κ_0 values for any rock site in New Zealand, which is likely to be a more reliable method of inference than the use of a $\kappa_0 - V_{S30}$ correlation.

2 DATA

The majority of the data for this study come from GeoNet's New Zealand National Seismic Network (NZNSN, Petersen et al. 2011). At the time of writing, the NZNSN consists of 52 sites with paired strong-motion and broadband instruments, located throughout New Zealand at approximately 100 km spacing. From this network, there are many continuously recording stations on rock sites that have been operational for up to ten years, providing a wealth of data from which κ_0 can be calculated. The dataset for this study is obtained from 33 rock sites in the NZNSN, the locations of which are shown as triangles in Figure 1.

To supplement the network in the North Island, where the station spacing of the NZNSN was deemed insufficient, eight stations from regional networks are also included in this study (shown as squares in Figure 1). The NZNSN sites in this study comprise co-located continuous and strong-motion sensors, however for this study only data from the continuous broadband seismometers are utilised. The regional networks use continuously-recording short-period instruments. All stations in the dataset sample at 100 Hz. The only site information available is NZS1170.5:2004 site classifications (Standards New Zealand 2004), according to which some stations are classified either as hard rock (class A) or soft rock (class B). However, these classifications are inferred, and no direct measurements of site properties exist for any sites in this study. In addition to the 41 stations analysed here, the Christchurch κ_0 results from Van Houtte et al. (2014) are used in the development of the continuous κ_0 map.

In accordance with the minimum magnitude recommendations of Van Houtte et al. (2014), only events with magnitudes greater than 2.5 are considered in this study. The purpose of this limit is to provide sufficient bandwidth above the source corner frequency, f_c , and below the maximum usable frequency, to allow the calculated κ slope to be attributed solely to the site and

1
2
3
4 8 *C. Van Houtte et al.*

5
6 169 regional attenuation without contamination by source effects. Although f_c also depends on the
7
8 170 stress parameter $\Delta\sigma$, which may exhibit systematic regional differences as well as event-to-event
9
10 171 variability, this minimum magnitude requirement is considered sufficient for avoiding strong bi-
11
12 172 ases of f_c on the calculated κ slope. The available data span a magnitude range from $M_w 2.5$ to
13
14 173 $M_w 6.6$. Only crustal events with depths less than 15 km are analysed in this study to ensure that
15
16 174 the wave propagation paths from all events are sampling similar portions of the seismogenic crust.
17
18 175 Additionally, source-to-site distances are limited to less than 150 km, so that path attenuation does
19
20 176 not overshadow the site attenuation effect. With these criteria, a total of 1050 recordings form the
21
22 177 dataset for this study. The magnitudes and epicentral distances for these recordings are shown in
23
24 178 Figure 2.

25 26 27 28 179 **3 METHOD**

29
30 180 The Anderson & Hough (1984) method is adopted to calculate κ_0 (rather than other calculation
31
32 181 methods), where the high frequency slope of the acceleration spectra for recorded earthquake
33
34 182 data is empirically fit with the low-pass filter model in equation 1. The slope usually steepens as
35
36 183 distance from the source increases, and removal of the trend with distance yields a site-specific κ_0 .
37
38 184 This method was denoted by Ktenidou et al. (2014), as $\kappa_{0,AS}$. This choice of method is to retain
39
40 185 consistency between how κ_0 is calculated and how it is intended to be applied, in empirical ground
41
42 186 motion modelling.

43
44
45 187 The details of the data processing are very similar to the method detailed in Van Houtte et al.
46
47 188 (2014), and are only briefly repeated here. For each recording, S-wave windows of five second
48
49 189 duration are used are used to calculate the FAS. The instrument response is removed using the
50
51 190 transfer function in SAC (Goldstein & Snoke 2005). Five second pre-event noise windows are
52
53 191 also calculated, and the recordings are only analysed for frequencies where the signal-to-noise
54
55 192 ratio (SNR) is greater than three. κ values are calculated by fitting the logarithm of the high-
56
57 193 frequency acceleration spectrum, using linear regression between two frequencies, f_e and f_x . f_e
58
59 194 is selected to be well beyond a theoretical source corner frequency (from the Brune (1970) model
60
195 with $\Delta\sigma=2$ MPa), and in general is greater than 10 Hz. Given the 100 Hz sampling of the data,

f_x is selected to always be lower than 40 Hz, although in some cases f_x is limited by the SNR criterion. A minimum value of $f_x - f_e = 10$ Hz is applied to obtain reasonably stable regression.

To identify narrow band site effects that may bias κ measurements, average HVSRs are calculated at each station using multiple earthquakes. If high-frequency amplification was detected at a given station, it was considered to violate the assumptions for calculating κ , and removed from the dataset. An example κ calculation, for a recording with a relatively low κ value, is shown in Figure 3. To obtain a single horizontal κ value per recording, the horizontal component orientation of the recording is rotated at 5° increments through 180° , and κ is calculated as detailed above for each rotation increment. The mean κ value across all orientations is adopted as the κ value for that particular recording. More details on this procedure can be found in Van Houtte et al. (2014).

4 CONSIDERATION OF DISTANCE AND DIRECTION DEPENDENCE

Once κ values are calculated for all events recorded at a station, the site attenuation parameter, κ_0 , is then calculated by extrapolating the trend of κ with epicentral distance to zero distance (i.e. removing the $\tilde{\kappa}(r)$ term from equation 2). The distance-dependence is approximately related to the intrinsic Q for S-waves (Q_S) and V_S structure along the horizontal wave-propagation path from source to site, as shown in equation 3. For a medium with constant Q_S and V_S , κ will increase linearly with distance, and several previous studies adopt this simple form for the distance-dependence (e.g., Hough et al. 1988; Douglas et al. 2010; Van Houtte et al. 2011; Ktenidou et al. 2013). While many alternative forms for $\tilde{\kappa}(r)$ have also been proposed to better represent the geological structure in the region of interest, this study uses a constant Q_S and V_S model,

$$\tilde{\kappa}(r) \approx \frac{R}{Q_S V_S}, \quad (6)$$

where R is the epicentral distance. A more complex form for $\tilde{\kappa}(r)$ was considered, but was limited by depth uncertainties for many recorded events, making it too difficult to reliably trace the Q_S structure being sampled by the ray paths. Therefore, a linear distance-dependence of κ is adopted for simplicity and consistency.

1
2
3
4 10 *C. Van Houtte et al.*

5
6 220 For three stations in the Otago region of New Zealand, Figure 4 shows the locations of events
7
8 221 used to calculate κ , and plots the κ results against epicentral distance. If the κ data at each sta-
9
10 222 tion are fitted using unconstrained linear least-squares regression (labelled as ‘free Q ’ in Figure
11
12 223 4), the Q_S values implied by the $\tilde{\kappa}(r)$ slope range from 900-1600, assuming constant V_S of 3.5
13
14 224 km/s. These Q_S values can be compared to other Q values obtained in an independent study. The
15
16 225 3D South Island attenuation model of Eberhart-Phillips et al. (2008) calculates the frequency-
17
18 226 independent Q structure for P-waves, Q_P , in Otago. In the typical frequency bandwidth for κ cal-
19
20 227 culations, 10-40 Hz, and a similar depth range to the data in this study, less than 10 km, Eberhart-
21
22 228 Phillips et al. (2008) calculate Q_P values ranging from 900-1100. Given that Eberhart-Phillips
23
24 229 et al. (2008) calculate Q_P and this study calculates Q_S , it is not possible to directly compare
25
26 230 the results. Lay & Wallace (1995) and Castro et al. (1999) suggest Q_P/Q_S is around 2 and 1
27
28 231 respectively, although it is not clear how these values should depend on frequency. Yoshimoto
29
30 232 et al. (1993) show that Q_P/Q_S does depend on frequency, and can range between 0.5-2.5 for fre-
31
32 233 quencies greater than 1 Hz. The variability in the ratio means that the absolute Q values are not
33
34 234 comparable, but qualitatively the relative Q_S values from this study compare well with the general
35
36 235 regional trends of the Q_P structure derived by Eberhart-Phillips et al. (2008).

37
38 236 The κ_0 values in Otago inferred from the ‘free Q ’ regression are also indicated in Figure 4.
39
40 237 However, the station-specific κ_0 values may be sensitive to constraints associated with the linear
41
42 238 distance-dependence model (Edwards et al. 2015; Ktenidou et al. 2017). With the assumption that
43
44 239 the ray paths sample similar underlying Q structure for all three Otago stations, their combined
45
46 240 κ data may be regressed together to calculate a better-constrained regional $\tilde{\kappa}(r)$. In this case, the
47
48 241 combined regression yields an ‘average regional Q_S ’ of 1100. If the slope of the linear trend
49
50 242 is constrained at each station to fit this Q model (indicated in Figure 4 as ‘fixed regional Q ’),
51
52 243 the obtained κ_0 values may be different from those of the unconstrained, ‘free Q ’ case. For the
53
54 244 example of the SYZ station, adopting either ‘free Q ’ or ‘fixed regional Q ’ distance-dependence
55
56 245 models causes changes in κ_0 that are nearly a factor of two, but for most stations the difference is
57
58 246 within a factor of two.

59
60 247 For groups of stations where sufficient data have propagation paths sampling similar areas of

the crust, such as the three Otago stations in Figure 4, the overall κ_0 is taken as the average between a ‘free Q ’ and ‘fixed Q ’ distance model. This accounts for some of the epistemic uncertainty in κ_0 due to the adopted distance-dependence model. In addition to the three Otago stations, the following stations were also combined to calculate a regionally-fixed distance dependence:

- (i) DCZ and PYZ, in Fiordland;
- (ii) WEL and BHW, two stations at the southern tip of the North Island; and,
- (iii) WAZ and VRZ, near the west coast of the North Island.

The other regions in New Zealand have insufficient spatial coverage of data and stations to allow the computation of robust regionally-fixed Q models, and for stations in these regions, κ_0 is calculated using only a ‘free Q ’ linear-regression model.

To assist with identifying and interpreting strong Q contrasts that are affecting κ , the results of this study are compared with available crustal Q models for New Zealand (Eberhart-Phillips & Chadwick 2002; Eberhart-Phillips et al. 2005, 2008; Eberhart-Phillips & Bannister 2010; Eberhart-Phillips et al. 2015). Linear regression of the κ data is only applied to events where the wave propagation paths sample relatively constant Q , and events that are likely to be affected by highly heterogeneous Q structures are not considered when calculating κ_0 . For example, Eberhart-Phillips et al. (2008) observe a localised, shallow, low Q region in the central South Island, which they believe is likely to be due to the high seismic activity rate in this part of the brittle crust. Given that this Q anomaly greatly affects the path-dependence of κ , events with ray paths passing through this region are omitted from the κ_0 calculations.

5 CALCULATION OF κ_0 AT STATION LOCATIONS

For 33 stations in the NZNSN, there were sufficient reliable data to furnish κ_0 values. Additionally, κ_0 is calculated at eight stations from regional networks. Combined with five rock-site κ_0 values near the city of Christchurch, calculated in Van Houtte et al. (2014), this gives a total of 46 rock-site κ_0 values throughout the country. Table 1 lists the κ_0 values corresponding to each station, the gradient $d\kappa/dR$ and the number of events used to calculate κ_0 . It also indicates stations for

1
2
3
4 12 *C. Van Houtte et al.*

5
6 274 which κ_0 was calculated as the average of ‘free Q ’ and ‘fixed regional Q ’ regressions. Figures 5
7
8 275 and 6 show the individual station κ data for the South and North Islands respectively, along with
9
10 276 the derived κ_0 values and their 5-95% confidence intervals. Figure 7 shows the residuals between
11
12 277 the New Zealand κ data and the linear distance-dependence model in equation 2. The residuals do
13
14 278 not have a significant trend with distance, which suggests the adopted constant Q assumption is a
15
16 279 reasonable model for this dataset. The histogram and normal Q-Q plots suggest that the residuals
17
18 280 do not deviate significantly from a normal distribution. There is a slight increasing trend between
19
20 281 the κ residuals and magnitude, as shown in Figure 7d. If a linear trend is fit to the residual data, the
21
22 282 p -value of its slope, 0.00575, would be considered statistically significant in traditional statistics,
23
24 283 although perhaps not according to modern thinking (Benjamin et al. 2017; McShane et al. 2017).
25
26 284 While there may be some effect of magnitude on the obtained κ estimates, it is clear that the
27
28 285 effect is very small compared to distance and other factors (the R -squared for a linear model of
29
30 286 $\kappa_{observed} - \kappa_{predicted}$ vs M_L is only 0.007). The trend may also be a result of underlying correlations
31
32 287 in the dataset, for example large earthquakes mostly occurring in regions of high fracture density
33
34 288 and hence low Q .

35
36 289 The calculated κ_0 values are illustrated in Figure 8 over a map of New Zealand’s surface ge-
37
38 290 ology, with the tectonic plate boundary between the Pacific and Australian plates also included
39
40 291 for reference. Dashed lines are the boundaries between regions with distinctly different site atten-
41
42 292 uation characteristics, which are to be discussed in detail. There appears to be a strong regional
43
44 293 dependence of κ_0 , the variation of which is likely due to near-surface geology (i.e. depth less than
45
46 294 5 km) and the tectonic setting. Figure 9 shows the range of κ_0 values observed for each region that
47
48 295 is delineated in Figure 8. These regions are addressed individually hereafter.

51 296 **5.1 Otago**

52
53
54 297 The κ_0 values in the southernmost regions of New Zealand are consistently lower than other re-
55
56 298 gions of the country. The Otago region is geologically characterised by hard rock (schist) sites, and
57
58 299 is identified by Eberhart-Phillips et al. (2008) as a region of high Q_P . κ_0 values from this region
59
60 300 range from 0.006 to 0.012 s, similar to those observed in Eastern North America (Silva & Darragh

1995; Ktenidou & Abrahamson 2016; Ktenidou et al. 2016). This result might be significant for some engineering purposes, as the low κ_0 values are likely to result in stronger high-frequency ground motion in Otago compared to other regions in New Zealand.

5.2 Fiordland

κ_0 values are obtained at two stations in Fiordland, PYZ and DCZ, both of which are hard rock sites. κ_0 for PYZ and DCZ are 0.014 and 0.013 s respectively, which are slightly higher than the values in Otago, but still lower than other regions around New Zealand. The implied Q_S values of the $d\kappa/dR$ slope are also slightly lower than in Otago, a trend that is similarly observed by Eberhart-Phillips et al. (2008) in their 3D South Island Q_P model.

5.3 Canterbury and West Coast

For Canterbury and West Coast regions of the South Island, κ_0 values range from 0.009 to 0.027 s. Three stations in Canterbury, to the east of the Alpine Fault, are located on greywacke rock, and despite station separations of up to 250 km, have similar κ_0 values ranging from 0.016 to 0.022 s. The surface geology is more variable to the west of the Alpine fault, and this is reflected in the larger range of κ values. The $d\kappa/dR$ slopes are reasonably consistent to the east of the Alpine Fault, but there appears to be more variable Q_S structure near the Alpine Fault, and on the West Coast.

5.4 Banks Peninsula

There are several strong motion stations on the Banks Peninsula, which recorded many events in the Canterbury earthquake sequence (Bannister & Gledhill 2012) and have been assigned κ_0 values in Van Houtte et al. (2014). Five of these stations are deemed to be good representations of rock sites, and κ_0 values from these stations are indicated in Figures 8, 9 and Table 1. Observed κ_0 values are larger than other locations in Canterbury, likely due to stations being located on a highly-weathered volcanic outcrop.

5.5 East Coast of the North Island

Stations located on the North Island's East Coast have some of the highest κ_0 values in the country. Sedimentary rock sites located on the coastal ranges and in the forearc basin have κ_0 values ranging from 0.030 to 0.053, similar to observations in Western North America (Anderson & Hough 1984; Silva & Darragh 1995), although the confidence intervals in Figure 6 indicate that these values are poorly constrained in some places. This region was identified in Eberhart-Phillips et al. (2005) as having very low Q_P in the upper crust, due to the high fracture density of material adjacent to the nearby Hikurangi subduction interface. Further west, stations located on the arc ranges have lower κ_0 values than in the forearc, ranging from 0.029 to 0.036 s. A similar increase in shallow Q_P structure in the arc ranges is observed by Eberhart-Phillips et al. (2005).

In general, the $d\kappa/dR$ values are much higher on the North Island's east coast compared to those in the South Island. However, these values are highly variable. This observed variation in $d\kappa/dR$ may be due to strong Q_S variation, or it could be due to uncertainties in the depth of seismicity, as events occurring within the subducting Pacific plate are likely to have much higher κ values compared to events in the overlying Australian plate.

5.6 Taupo Volcanic Zone

The Taupo Volcanic Zone (TVZ), which forms the northern section of the backarc region, has long been recognised as a region of very high attenuation (e.g., Mooney 1970; Cousins et al. 1999; Dowrick & Rhoades 1999; Eberhart-Phillips & McVerry 2003; McVerry et al. 2006). While there are no NZNSN stations located in the TVZ, there are many short-period instruments from volcanic networks. κ_0 values could be estimated at three of these stations. While the results are poorly constrained, particularly at the UTU station, the κ_0 values are all greater than 0.050 s and are the largest in the country. These observations have an important influence on the derivation of the continuous κ_0 map, which will be addressed in the following section.

The TVZ κ_0 results, along with those on the North Island's east coast, compare well with previous Q attenuation studies. The lower North Island Q_P model of Eberhart-Phillips et al. (2005)

351 suggests that either side of the high- Q_P arc ranges, the forearc and volcanic fronts have very low
 352 Q_P values. This pattern has also been observed in subduction regions of Japan (Pei et al. 2009).

353 5.7 Northern Districts

354 Calculating κ_0 values for stations in the northern regions of New Zealand was more difficult than
 355 in the South Island, as these regions typically have low seismicity and hence have fewer data
 356 available. Four stations located on sedimentary rock sites, WAZ, VRZ, HIZ and TLZ, have very
 357 consistent attenuation properties, with κ_0 values ranging from 0.015 to 0.019 s. The four north-
 358 ernmost stations in this study, TGRZ, TOZ, WIAZ and WCZ, have very similar κ_0 values ranging
 359 from 0.024 to 0.030, however for WIAZ and WCZ in particular, these are derived from few events
 360 and thus are not well constrained. The $d\kappa/dR$ slopes are in general lower than the rest of the North
 361 Island, and are similar to South Island values.

362 6 DEVELOPMENT OF A CONTINUOUS κ_0 MAP

363 It is often the case in seismic hazard assessment that ground motion estimates are required at a
 364 site located away from a recording station where κ_0 can be directly computed. To facilitate the
 365 use of the New Zealand κ_0 data in ground motion modelling, and subsequently seismic hazard, a
 366 continuous map is developed. In this section, the discrete station-specific κ_0 values are spatially
 367 smoothed and interpolated, to develop a continuous rock-site κ_0 map for New Zealand. A number
 368 of spatial interpolation schemes were considered as candidates for developing the κ_0 map and ulti-
 369 mately, Gaussian process regression, or ‘kriging’, was preferred due to its flexibility for modelling
 370 the underlying process. Information on kriging is readily available in the literature, and the texts
 371 of Cressie (1993) and Diggle & Ribeiro (2007) are recommended.

372 The problem is initially formulated as follows. Each κ_0 observation is assumed to be a reali-
 373 sation of a random variable, which has a distribution that depends on the value of an underlying
 374 spatially-continuous Gaussian process $S(x)$ at the location x_i , i.e.

$$K_{0,i} = S(x_i) + \epsilon_i \quad : \quad i = 1, \dots, n, \quad (7)$$

1
2
3
4 16 *C. Van Houtte et al.*

5
6 375 where $K_{0,i}$ is a vector of n observed κ_0 values, and ϵ_i are normally distributed errors with mean
7
8 376 of 0 and variance of τ^2 . $S(x)$, known as the ‘signal’, has mean μ , variance σ^2 and a correlation
9
10 377 function $\rho(u)$, where u is the distance between two locations x and x' (Diggle & Ribeiro 2007).
11
12 378 The objective of this study is to provide a model for $S(x)$ based on the New Zealand κ_0 data in
13
14 379 Table 1. Developing this model is a three step process. Firstly, the empirical spatial relationship
15
16 380 of the data is examined, secondly the parameters of a theoretical spatial correlation model are
17
18 381 estimated, and then finally the prediction of $S(x)$.

21 382 **6.1 Empirical spatial behaviour**

22
23
24 383 A useful tool for investigating the behaviour of geospatial data is the semi-variogram, γ , which
25
26 384 describes the variance of the difference between two realisations of the signal $S(x)$ at locations x
27
28 385 and x' ,

$$32 \quad \gamma(x, x') = \frac{1}{2} \text{Var}[S(x) - S(x')] \quad . \quad (8)$$

33
34
35 386 In the case of a stationary Gaussian process, the semi-variogram can also be expressed as

$$36 \quad \gamma(x, x') = \gamma(u) = \tau^2 + \sigma^2[(1 - \rho(u))] \quad . \quad (9)$$

37
38
39 387 A typical function for $\rho(u)$ monotonically decreases as the distance u increases, hence $\gamma(u)$
40
41 388 usually monotonically increases with u . Semi-variograms tend to have the following features.
42
43 389 When $u = 0$, the intercept of the semi-variogram is τ^2 , which is known as the ‘nugget variance’,
44
45 390 or simply ‘nugget’. Also, when $\rho(u) = 0$ (no correlation between observations beyond a given
46
47 391 distance), the semi-variance becomes $\tau^2 + \sigma^2$, called the ‘sill’. The ‘range’ of the semi-variogram
48
49 392 is the distance at which $\gamma(u)$ equals the sill. Figure 10 shows a schematic example of a semi-
50
51 393 variogram with $\rho(u)$ decreasing according to $\exp(-u)$, with a nugget variance of 0.1 and sill of
52
53 394 0.8. Note that, as the correlation function is exponential, it never reaches zero and hence the range
54
55 395 is undefined. It is typical to instead define a ‘practical range’ equal to 0.95 times the sill.

56
57
58
59 396 The solid line in Figure 11 is the empirical semi-variogram for the New Zealand κ_0 data,

397 determined using equation 8, substituting the observations κ_0 and κ'_0 in place of $S(x)$ and $S(x')$.
 398 The data are assumed to be log-normally distributed, as it was found that a logarithmic transform
 399 of the samples brings the data much closer to Gaussian. Equation 7 is then rewritten as

$$\log K_{0,i} = S(x_i) + \epsilon_i \quad : \quad i = 1, \dots, n . \quad (10)$$

400 Semi-variance samples are combined into 50 km bins. There are several points of interest in
 401 this plot. Firstly, there appears to be a small, but non-zero, nugget variance. The nugget variance
 402 has a dual interpretation as representing either the measurement error for each observation, or the
 403 short-distance spatial variation. If the nugget parameter is calculated from these κ_0 semi-variance
 404 data, then its value will be primarily determined by data from the Banks Peninsula, where inter-
 405 station distances are around 5-10 km. The increasing semi-variance plateaus at around 400 km,
 406 before significantly increasing beyond 800 km. The second increase reflects the large differences
 407 between κ_0 values at long inter-station distances, for example between Otago and the TVZ. This
 408 indicates that $S(x)$ is likely to have a mean that depends on location, known as a trend surface,
 409 which is a departure from a stationary Gaussian process.

410 Given that the spatially-varying mean is likely due to complex variation of crustal properties,
 411 this trend surface is difficult to model geostatistically. To minimise the effect of this trend on the
 412 spatial prediction, particularly at short inter-station distances, a simple trend surface is applied to
 413 the spatial κ_0 data. The trend surface consists of a single indicator variable for the Taupo Volcanic
 414 Zone, F_{TVZ} , such that the mean function is

$$\mu(x) = \beta_0 + \beta_1 \cdot F_{TVZ} \quad (11)$$

415 where F_{TVZ} is 1 for sites located within the TVZ, and 0 otherwise. The justification for this
 416 trend is that the TVZ is well-known for having very high near-surface attenuation properties, and
 417 its geographical extent is well-defined in Wilson et al. (1995). The dashed line in Figure 11 shows
 418 the new empirical semi-variogram with this trend surface applied. Not only does this model lower
 419 the semi-variogram at large distances, but the semi-variances are also lower at short distances.

1
2
3 18 *C. Van Houtte et al.*

4
5 420 This is because the trend surface allows stronger correlations between κ_0 from the TVZ sites and
6
7 421 nearby stations outside the volcanic arc.
8
9

10 422 **6.2 Semi-variogram parameter estimation**

11
12
13 423 Before spatial prediction can occur, a model for the empirical data in Figure 11 must be developed.

14
15 424 A first step is to choose a parametric function for $\rho(u)$ that best models the empirical κ_0 semi-
16
17 425 variances. There are many different correlation functions that can be adopted to model different
18
19 426 types of spatial correlation behaviour. In this study, a widely-used family of correlation functions
20
21 427 is adopted, known as the Matérn (1960) family, which have the general form

$$22 \rho(u) = [2^{\theta-1}\Gamma(\theta)]^{-1} \left(\frac{u}{\phi}\right)^{\theta} K_{\theta}\left(\frac{u}{\phi}\right) . \quad (12)$$

23
24
25
26
27
28 428 K_{θ} denotes a modified Bessel function with non-negative order θ , ϕ is a non-negative scale
29
30 429 parameter with units of distance and Γ is the gamma function. As with previous notation, u is
31
32 430 the distance between two locations x and x' . The advantage of this form is that it is very flexible,
33
34 431 because θ dictates the behaviour of the correlation structure where the station separation distance
35
36 432 is small, while ϕ controls the degree of correlation between stations with large separation distance.
37
38 433 Figure 12 gives an example of the effect of different values of θ and ϕ on $\rho(u)$. Note that for $\theta =$
39
40 434 0.5 , $\rho(u)$ reduces to the exponential function $\exp(u/\phi)$.

41
42
43 435 The Matérn parametric form is used to model the spatial correlation of the κ_0 data. Treating
44
45 436 the transformed spatial κ_0 data as Gaussian, the model can be written as

$$46 \log K_0 \sim N(Y\beta, \sigma^2 B(\theta, \phi) + \tau^2 I) , \quad (13)$$

47
48
49 437 where $K_0 = (\kappa_{0,1}, \dots, \kappa_{0,n})$, Y is a matrix of covariates and β the corresponding matrix of
50
51 438 coefficients, B is an n by n correlation matrix, which in this case depends on the Matérn parameters
52
53 439 θ and ϕ , and I is the identity matrix. The corresponding log-likelihood function is

$$\mathcal{L}(\beta, \tau, \sigma, \theta, \phi) = -0.5 \left\{ (n \log(2\pi) + \log |\sigma^2 B(\theta, \phi) + \tau^2 I| + (\log \kappa_0 - Y\beta)^T [\sigma^2 B(\theta, \phi) + \tau^2 I]^{-1} (\log \kappa_0 - Y\beta) \right\}. \quad (14)$$

440 Maximising equation 14 gives the model parameters. However, the Matérn parameters θ and
 441 ϕ tend to be strongly correlated, so a common, alternative method for parameter estimation is to
 442 fix a discrete set of θ , then maximise equation 14 to determine β , τ^2 , σ^2 and ϕ (Diggle & Ribeiro
 443 2007).

444 Using this method, Table 2 shows the maximum likelihood parameters for two candidate mod-
 445 els, and their corresponding log-likelihoods. Both models apply the trend surface from equation
 446 11, but the first model solves for the nugget variance, while the second model fixes the nugget
 447 variance to an indicative value of the station-specific κ_0 uncertainty. Guided by the $\kappa(r)$ regres-
 448 sion results in Figures 5 and 6, an approximate standard deviation value of 50% ($0.18 \log_{10}$ units)
 449 was assumed as the average uncertainty of κ_0 , and hence τ^2 was fixed to equal 0.03. For the
 450 first model, three values of the Matérn order are considered, 0.5, 1.5 and 2.5. These three val-
 451 ues are commonly adopted in geostatistical modelling to represent different smoothness of $S(x)$.
 452 Specifically, they represent processes that are mean-square continuous, once mean-square differ-
 453 entiable and twice mean-square differentiable respectively. For the second model, only a model
 454 with $\theta = 0.5$ is considered, as models with higher values of θ are over-parameterised and give
 455 unstable solutions.

456 The first point from Table 2 is that ϕ usually decreases as θ increases, illustrating the correlation
 457 between the parameters. However, the most important effect of the θ parameter is to increase
 458 the nugget variance τ^2 . The models with different θ values have similar log-likelihoods, which
 459 indicates that τ^2 is not well-constrained by the data. The models with free τ^2 have higher log-
 460 likelihoods than the model with fixed τ^2 .

461 The best-fit theoretical semi-variograms are shown in Figure 13. The best-fit models fit the
 462 empirical semi-variances at small separation distances better than at large distances, which is ex-
 463 pected because equation 14 allows for less precise estimates as distance increases. The inset in
 464 Figure 13a shows the effect of the different θ values at short station-separation distances. The

1
2
3
4 20 *C. Van Houtte et al.*

5
6 465 more complicated form of the $\theta = 1.5$ and $\theta = 2.5$ models is not well justified by the data, hence
7
8 466 the model with $\theta = 0.5$ is preferred in this study.

9
10
11 467 **6.3 Spatial prediction**

12
13
14 468 The final step for developing a continuous site attenuation model is to predict the underlying signal
15
16 469 $S(x)$ from the observed κ_0 realisations and the candidate covariance models. This study uses the
17
18 470 ordinary kriging algorithm, the theory behind which is well described in the statistical literature
19
20 471 and is only briefly repeated here. In short, ordinary kriging, involves minimising the mean square
21
22 472 prediction error

$$23$$

$$24$$

$$25$$

$$26 \quad MSPE[\hat{S}(x)] = \frac{1}{n} \sum_{i=1}^n [(\hat{S}(x) - S(x_i))]^2, \quad (15)$$

$$27$$

28
29 473 where $\hat{S}(x)$ is the prediction of $S(x)$ at an unobserved location. The form of $\hat{S}(x)$ that min-
30
31 474 imises $MSPE[\hat{S}(x)]$ is,

$$32$$

$$33$$

$$34$$

$$35 \quad \hat{S}(x) = \mu + b'V^{-1}(K_0 - \mu), \quad (16)$$

$$36$$

37
38 475 where μ is a vector of mean values of $S(x)$, b is a vector of length n with values corresponding
39
40 476 to the correlation between the unobserved location and the n observed locations, and $V = B +$
41
42 477 $(\tau^2/\sigma^2)I$. In ordinary kriging, μ is assumed to be unknown and is estimated by generalised least
43
44 478 squares, given by

$$45$$

$$46$$

$$47$$

$$48 \quad \hat{\mu} = (\mathbf{1}'V^{-1}\mathbf{1})^{-1}\mathbf{1}'V^{-1}K_0. \quad (17)$$

$$49$$

50
51 The prediction variance is given by

$$52$$

$$53 \quad \text{Var}(S(x)|K_0) = \sigma(1 - b'V^{-1}b), \quad (18)$$

$$54$$

55
56 479 Full derivations of these expressions are readily available in the literature.

57
58 480 As the models contain the F_{TVZ} indicator variable, it is necessary to define the boundaries
59
60 481 of the TVZ in the prediction grid. Wilson et al. (1995) define different boundaries for the TVZ

based on its geological history, and their model for the ‘whole TVZ’ was preferred by Cousins et al. (1999) to represent volcanic path attenuation within empirical ground motion models. Like Cousins et al. (1999), the ‘whole TVZ’ model is selected here as the boundary for the high attenuation volcanic region.

Figure 14 compares the predicted median and standard deviation of κ_0 for models 1 and 2, both with θ equal to 0.5. Due to the larger τ^2 , model 2 predicts a smoother median surface. The effect of increasing τ^2 is essentially to trade in data reproduction for variance, and hence $\hat{S}(x)$ approaches $\bar{\kappa}_0$. The standard deviations in Figure 14 are obtained by taking the square root of the sum of the underlying signal variance and τ^2 . This calculation means that the maps in Figures 14b and 14d provide the prediction standard deviation for an unobserved κ_0 value.

6.4 Consideration of κ_0 ‘measurement uncertainty’

Figure 14 clearly demonstrates that the nugget is an important parameter when considering forward application of the κ_0 map. While the seismograph network in New Zealand is insufficiently dense to derive the nugget parameter from the kriging analysis alone, estimates of κ_0 measurement uncertainty are available from the linear regression of the $\kappa(r)$ data with distance (Figures 5 and 6). This is the purpose of model 2 in Table 2.

The fixed τ^2 value of 0.03 for model 5 is only an indicative value of the κ_0 uncertainty across all stations. A limitation of this assumption is that some of the κ_0 data are better constrained than others, as evidenced by the confidence intervals in Figures 5 and 6. Ideally, a fixed vector of τ^2 with each station’s κ_0 uncertainty could be propagated into the maximum likelihood step. However, this calculation is complicated by the assumption of two different distributions to derive the κ_0 map. The station-specific κ_0 are derived by regression of κ data with distance on a linear scale, while a lognormal distribution was assumed to derive the κ map. The combination of distributions prevents an elegant determination of the spatially-varying nugget, hence the fixed nugget of 0.03 ($\pm 50\%$ of the median value) is adopted for simplicity.

The choice of the preferred model is likely a personal preference, however we are inclined to select model 5. While the AIC of model 2 is lower than model 1, the higher value of the nugget

1
2
3
4 22 *C. Van Houtte et al.*

5
6 509 is more representative of the data and its underlying uncertainty. Model 1 relies on spatially-dense
7
8 510 sampling of κ_0 to constrain the nugget, which are mostly unavailable in New Zealand.

9
10 511 For model 1, the predicted standard deviation on κ_0 is around $0.1 \log_{10}$ units, while for model 2
11
12 512 it is around $0.2 \log_{10}$ units. Note that this 0.2 is only slightly larger than the value that was initially
13
14 513 assumed for the uncertainty in the station-specific κ_0 values, which suggests that the uncertainty in
15
16 514 calculating κ_0 at a given station location is much larger than the uncertainty introduced by kriging.

17
18 515 An additional aspect of uncertainty in the κ_0 map comes from the small dataset. This un-
19
20 516 certainty can be demonstrated using jackknife resampling, also known as leave-one-out cross-
21
22 517 validation. Maps of the jackknife standard deviation for the κ_0 map can be found in Figure S1 of
23
24 518 the supplementary materials, although in general, the jackknife standard deviation is much smaller
25
26 519 than the prediction standard deviation shown in Figure 14.

27 28 29 30 520 **7 SUMMARY AND DISCUSSION**

31
32 521 The overarching objective of this study is to provide the best available information on κ_0 in New
33
34 522 Zealand, to aid efforts to include it in empirical ground-motion prediction, and subsequently in
35
36 523 hazard calculations. Regions within New Zealand are shown to have significantly different atten-
37
38 524 uation properties depending on local geology and tectonic setting, and rock site κ_0 can vary from
39
40 525 0.006 - 0.055 s. The general path attenuation trends observed in this study are consistent with 3D
41
42 526 Q_P attenuation models of New Zealand (Eberhart-Phillips et al. 2015). Many stations have κ_0 val-
43
44 527 ues that are poorly constrained, with 19 values being derived from fewer than 20 recordings each.
45
46 528 While it is acknowledged that these poor constraints limit the level of confidence in the interpre-
47
48 529 tation of the κ_0 map, the proposed model nevertheless represents the best information currently
49
50 530 available on κ_0 in New Zealand. The model is a first attempt at a nationwide mapping of the site
51
52 531 attenuation parameter, and can be updated and improved as more data become available.

53
54 532 Two geostatistical models of the κ_0 data are used to derive continuous maps of κ_0 . The purpose
55
56 533 of presenting two models was to illustrate that very different κ_0 maps can be obtained with differ-
57
58 534 ent kriging assumptions. Our preferred model is model 2. This model accounts for the measure-
59
60 535 ment uncertainty in κ_0 , albeit in a simplified manner. A spatially-varying measurement uncertainty

could not be easily applied, because the selection of the best distribution to represent κ_0 is still an unresolved issue. Hough et al. (1988), Oth et al. (2011), Edwards et al. (2015) and Van Houtte et al. (2014) suggest that κ_0 is normally distributed. However, the assumption of normality generates problems in forward modelling, as it introduces the possibility of unphysical, negative κ_0 values. As such, κ_0 is typically assumed to be either uniformly or lognormally distributed when it is applied in seismic hazard assessments (e.g., EPRI 1993; Silva & Darragh 1995; Atkinson & Boore 2006; Campbell et al. 2014), to prevent consideration of non-physical model parameters. This study supports these assumptions, because the population of κ_0 data in New Zealand are much better represented by a lognormal distribution.

Ktenidou & Abrahamson (2016) suggest that negative κ_0 values may be a result of site amplification at high frequencies, although still find some negative κ_0 values when correcting for generic ‘crustal amplification’ effects. Edwards et al. (2015) suggest that negative κ_0 values are likely to be a result of overly-simplistic assumptions in the derivation of κ_0 , for example the assumption of ω^2 decay of the source displacement spectrum. Should regional earthquakes have displacement spectra that consistently decay with $\omega^{n<2}$, combined with low site attenuation, it is conceivable that a negative κ_0 value can be observed, when deriving κ_0 empirically from the FAS of acceleration (the method used in this study). Such behaviour of source spectra may bias empirical κ_0 observations, like those derived in this study. However, it is our opinion that if negative κ_0 values are being observed in a given region, then the typical assumptions for deriving κ_0 values are being violated, and it is necessary to re-examine the spectral behaviour of regional seismic sources. If near-surface attenuation is to be modelled in seismic hazard analysis, it is essential that the applied values have a physical interpretation. Rather than permitting non-physical attenuation values, and propagating these into seismic hazard calculations, it is preferable to develop a more appropriate physical model for the spectral characteristics.

Given that no negative κ_0 values are observed in this study, it is assumed that the κ_0 data here can be considered positive-only. It is therefore desirable to adopt a distribution for κ_0 that prevents negative values. Unfortunately, this is not a trivial task. A logarithmic transformation of the κ_0 data seems to be an obvious starting point. However, the results in Figure 7 suggest that a normal

1
2
3 24 *C. Van Houtte et al.*

4
5 564 distribution is a good representation of κ data. While beyond the scope of this work, resolving
6
7 565 this issue would be important for future research. When a satisfactory method for modelling the
8
9 566 distribution of κ_0 is available, a more robust κ_0 map and its uncertainty can be derived. This will
10
11 567 make the derived median and standard deviation models more useful in the context of seismic
12
13 568 hazard assessment.

14
15 569 Lastly, this model is intended to represent rock site κ_0 only. Should a κ_0 estimate be required
16
17 570 for a soil site, the continuous κ_0 model in Figure 14 can be used in conjunction with a model for
18
19 571 the dependence of κ_0 with sediment depth. The effect of the sedimentary column on κ_0 has not
20
21 572 yet been widely investigated, at least in terms of observed spectral amplitudes. To the authors'
22
23 573 knowledge, the only available models are those of Campbell (2009) and Ktenidou et al. (2015).
24
25 574 These models are derived using data from Eastern North America and Greece respectively. Given
26
27 575 the very different tectonic environment and sediment types in New Zealand, it is unlikely that these
28
29 576 models are applicable to New Zealand conditions. Therefore, the κ_0 map in this study can only
30
31 577 currently be applied to rock sites, and further research is necessary to determine the effect of the
32
33 578 soil column on the site attenuation in New Zealand.

34 35 36 37 38 579 **ACKNOWLEDGMENTS**

39
40 580 This research was partially funded by a University of Auckland Doctoral Scholarship, and by Insti-
41
42 581 tut des Sciences de la Terre. Peter Stafford, Gail Atkinson, Zach Eilon and an anonymous reviewer
43
44 582 are greatly and genuinely thanked for their encouragement and comments on early versions of this
45
46 583 manuscript, which greatly improved the final model. All data in this study are publicly available
47
48 584 from GeoNet (www.geonet.org.nz), who are funded by the Earthquake Commission (EQC), Land
49
50 585 Information New Zealand (LINZ) and GNS Science.

51 52 53 54 586 **REFERENCES**

55
56
57 587 Abercrombie, R., 1997. Near-surface attenuation and site effects from comparison of surface and deep
58
59 588 borehole recordings, *Bulletin of the Seismological Society of America*, **87**(3), 731–744.

60

- 589 Abercrombie, R., 1998. A summary of attenuation measurements from borehole recordings of earth-
590 quakes: the 10 Hz transition problem, *Pure and Applied Geophysics*, **153**(2-4), 475–487.
- 591 Abrahamson, N. & Silva, W., 1997. Empirical response spectral attenuation relations for shallow crustal
592 earthquakes, *Seismological Research Letters*, **68**(1), 94–127.
- 593 Abrahamson, N., Silva, W., & Kamai, R., 2014. Summary of the ASK14 ground motion relation for active
594 crustal regions, *Earthquake Spectra*, **30**(3), 1025–1055.
- 595 Adams, D. & Abercrombie, R., 1998. Seismic attenuation above 10 Hz in southern California from coda
596 waves recorded in the Cajon Pass borehole, *Journal of Geophysical Research*, **103**(B10), 24257–24270.
- 597 Aki, K., 1967. Scaling law of seismic spectrum, *Journal of Geophysical Research*, **72**(4), 1217–1231.
- 598 Aki, K., 1980. Attenuation of shear-waves in the lithosphere for frequencies from 0.05 to 25 hz, *Physics
599 of the Earth and Planetary Interiors*, **21**(1), 50–60.
- 600 Anderson, J., 1991. A preliminary descriptive model for the distance dependence of the spectral decay
601 parameter in southern California, *Bulletin of the Seismological Society of America*, **81**(6), 2186–2193.
- 602 Anderson, J. & Hough, S., 1984. A model for the shape of the Fourier amplitude spectrum of acceleration
603 at high frequencies, *Bulletin of the Seismological Society of America*, **74**(5), 1969–1993.
- 604 Anderson, J., Lee, Y., Zeng, Y., & Day, S., 1996. Control of strong motion by the upper 30 meters, *Bulletin
605 of the Seismological Society of America*, **86**(6), 1749–1759.
- 606 Atkinson, G., 2012. Evaluation of attenuation models for the northeastern United States / southeastern
607 Canada, *Seismological Research Letters*, **83**(1), 166–178.
- 608 Atkinson, G. & Boore, D., 2006. Earthquake ground-motion prediction equations for eastern North Amer-
609 ica, *Bulletin of the Seismological Society of America*, **96**(6), 2181–2205.
- 610 Bannister, S. & Gledhill, K., 2012. Evolution of the 2010–2012 Canterbury earthquake sequence, *New
611 Zealand Journal of Geology and Geophysics*, **55**(3), 295–304.
- 612 Benjamin, D. J., Berger, J. O., Johannesson, M., Nosek, B. A., Wagenmakers, E.-J., Berk, R., Bollen,
613 K. A., Brems, B., Brown, L., Camerer, C., Cesarini, D., Chambers, C. D., Clyde, M., Cook, T. D.,
614 Boeck, P. D., Dienes, Z., Dreber, A., Easwaran, K., Efferson, C., Fehr, E., Fidler, F., Field, A. P., Forster,
615 M., George, E. I., Gonzalez, R., Goodman, S., Green, E., Green, D. P., Greenwald, A. G., Hadfield, J. D.,
616 Hedges, L. V., Held, L., Ho, T. H., Hoijsink, H., Hruschka, D. J., Imai, K., Imbens, G., Ioannidis, J. P. A.,
617 Jeon, M., Jones, J. H., Kirchler, M., Laibson, D., List, J., Little, R., Lupia, A., Machery, E., Maxwell,
618 S. E., McCarthy, M., Moore, D. A., Morgan, S. L., Munafò, M., Nakagawa, S., Nyhan, B., Parker, T. H.,
619 Pericchi, L., Perugini, M., Rouder, J., Rousseau, J., Savalei, V., Schnbrodt, F. D., Sellke, T., Sinclair, B.,
620 Tingley, D., Zandt, T. V., Vazire, S., Watts, D. J., Winship, C., Wolpert, R. L., Xie, Y., Young, C., Zinman,
621 J., & Johnson, V. E., 2017. Redefine statistical significance, *Nature Human Behaviour*.
- 622 Boatwright, J., 1980. A spectral theory for circular seismic sources; simple estimates of source dimension,
623 dynamic stress drop, and radiated seismic energy, *Bulletin of the Seismological Society of America*, **70**(1),

1
2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40
41
42
43
44
45
46
47
48
49
50
51
52
53
54
55
56
57
58
59
60

624 1–27.

625 Boore, D., 2003. Simulation of ground motion using the stochastic method, *Pure and Applied Geophysics*,
626 **160**(3-4), 635–676.

627 Boore, D. & Joyner, W., 1997. Site amplifications for generic rock sites, *Bulletin of the Seismological*
628 *Society of America*, **87**(2), 327–341.

629 Boore, D., Stewart, J., Seyhan, E., & Atkinson, G., 2014. NGA-West2 equations for predicting PGA,
630 PGV, and 5% damped PSA for shallow crustal earthquakes, *Earthquake Spectra*, **30**(3), 1057–1085.

631 Brune, J., 1970. Tectonic stress and the spectra of seismic shear waves from earthquakes, *Journal of*
632 *Geophysical Research*, **75**(26), 4997–5009.

633 Campbell, K., 2003. Prediction of strong ground motion using the hybrid empirical method and its use
634 in the development of ground-motion (attenuation) relations in eastern North America, *Bulletin of the*
635 *Seismological Society of America*, **93**(3), 1012–1033.

636 Campbell, K., 2009. Estimates of shear-wave Q and κ_0 for unconsolidated and semiconsolidated sediments
637 in eastern North America, *Bulletin of the Seismological Society of America*, **99**(4), 2365–2392.

638 Campbell, K. & Bozorgnia, Y., 2014. NGA-West2 ground motion model for the average horizontal com-
639 ponents of PGA, PGV, and 5% damped linear acceleration response spectra, *Earthquake Spectra*, **30**(3),
640 1087–1115.

641 Campbell, K., Hashash, Y., Kim, B., Kottke, A., Rathje, E., Silva, W., & Stewart, J., 2014. Reference rock
642 site condition for Central and Eastern North America part II-attenuation (kappa) definition.

643 Castro, R., Monachesi, G., Mucciarelli, M., Trojani, L., & Pacor, F., 1999. P- and S-wave attenuation in
644 the region of Marche, Italy, *Tectonophysics*, **302**(1), 123–132.

645 Chandler, A., Lam, N., & Tsang, H., 2006. Near-surface attenuation modelling based on rock shear-wave
646 velocity profile, *Soil Dynamics and Earthquake Engineering*, **26**(11), 1004–1014.

647 Chiou, B. & Youngs, R., 2014. Update of the Chiou and Youngs NGA model for the average horizontal
648 component of peak ground motion and response spectra, *Earthquake Spectra*, **30**(3), 1117–1153.

649 Cotton, F., Scherbaum, F., Bommer, J., & Bungum, H., 2006. Criteria for selecting and adjusting ground-
650 motion models for specific target regions: Application to Central Europe and rock sites, *Journal of Seis-*
651 *mology*, **10**(2), 137–156.

652 Cousins, W. J., Zhao, J., & Perrin, N., 1999. A model for the attenuation of peak ground acceleration in
653 New Zealand earthquakes based on seismograph and accelerograph data, *Bulletin of the New Zealand*
654 *Society for Earthquake Engineering*, **32**(4), 193–220.

655 Cressie, N., 1993. *Statistics for spatial data*, John Wiley & Sons, New York.

656 Dainty, A., 1981. A scattering model to explain seismic Q observations in the lithosphere between 1 and
657 30 Hz, *Geophysical Research Letters*, **8**(11), 1126–1128.

658 Diggle, P. & Ribeiro, P., 2007. *Model-based geostatistics*, Springer Series in Geostatistics, Springer-

Verlag, New York.

Douglas, J., Bungum, H., & Scherbaum, F., 2006. Ground-motion prediction equations for southern Spain and southern Norway obtained using the composite model perspective, *Journal of Earthquake Engineering*, **10**(1), 33–72.

Douglas, J., Gehl, P., Bonilla, L. F., Scotti, O., Régnier, J., Duval, A.-M., & Bertrand, E., 2009. Making the most of available site information for empirical ground-motion prediction, *Bulletin of the Seismological Society of America*, **99**(3), 1502–1520.

Douglas, J., Gehl, P., Bonilla, L. F., & Gélis, C., 2010. A κ model for mainland France, *Pure and Applied Geophysics*, **167**(11), 1303–1315.

Dowrick, D. & Rhoades, D., 1999. Attenuation of Modified Mercalli Intensity in New Zealand earthquakes, *Bulletin of the New Zealand Society for Earthquake Engineering*, **32**(2), 55–89.

Eberhart-Phillips, D. & Bannister, S., 2010. 3-D imaging of Marlborough, New Zealand, subducted plate and strike-slip fault systems, *Geophysical Journal International*, **182**(1), 73–96.

Eberhart-Phillips, D. & Chadwick, M., 2002. Three-dimensional attenuation model of the shallow Hikurangi subduction zone in the Raukumara Peninsula, New Zealand, *Journal of Geophysical Research: Solid Earth*, **107**(B2).

Eberhart-Phillips, D. & McVerry, G., 2003. Estimating slab earthquake response spectra from a 3D Q model, *Bulletin of the Seismological Society of America*, **93**(6), 2649–2663.

Eberhart-Phillips, D., Reyners, M., Chadwick, M., & Chiu, J.-M., 2005. Crustal heterogeneity and subduction processes: 3-D V_p , V_p/V_s and Q in the southern North Island, New Zealand, *Geophysical Journal International*, **162**(1), 270–288.

Eberhart-Phillips, D., Chadwick, M., & Bannister, S., 2008. Three-dimensional attenuation structure of central and southern South Island, New Zealand, from local earthquakes, *Journal of Geophysical Research: Solid Earth*, **113**(B5).

Eberhart-Phillips, D., Reyners, M., & Bannister, S., 2015. A 3D Q_P attenuation model for all of New Zealand, *Seismological Research Letters*, **86**(6), 1655–1663.

Edwards, B., Ktenidou, O.-J., Cotton, F., Abrahamson, N., Van Houtte, C., & Fäh, D., 2015. Epistemic uncertainty and limitations of the κ_0 model for near-surface attenuation at hard rock sites, *Geophysical Journal International*, **202**(3), 1627–1645.

EPRI, 1993. Method and guidelines for estimating earthquake ground motion in eastern North America.

Goldstein, P. & Snoke, A., 2005. SAC availability for the IRIS community, *Incorporated Institutions for Seismology Data Management Center Electronic Newsletter*, **7**.

Graves, R. & Pitarka, A., 2010. Broadband ground-motion simulation using a hybrid approach, *Bulletin of the Seismological Society of America*, **100**(5A), 2095–2123.

Haskell, N., 1964. Total energy and energy spectral density of elastic wave radiation from propagating

- 694 faults, *Bulletin of the Seismological Society of America*, **54**(6A), 1811–1841.
- 695 Hough, S. & Anderson, J., 1988. High-frequency spectra observed at Anza, California: implications for Q
696 structure, *Bulletin of the Seismological Society of America*, **78**(2), 692–707.
- 697 Hough, S., Anderson, J., Brune, J., Vernon, F., Berger, J., Fletcher, J., Haar, L., Hanks, L., & Baker, L.,
698 1988. Attenuation near Anza, California, *Bulletin of the Seismological Society of America*, **78**(2), 672–
699 691.
- 700 Housner, G., 1947. Characteristics of strong-motion earthquakes, *Bulletin of the Seismological Society of*
701 *America*, **37**(1), 19–31.
- 702 Kaneko, Y. & Shearer, P., 2014. Seismic source spectra and estimated stress drop derived from cohesive-
703 zone models of circular subshear rupture, *Geophysical Journal International*, **197**(2), 1002–1015.
- 704 Kaneko, Y. & Shearer, P., 2015. Variability of seismic source spectra, estimated stress drop, and radi-
705 ated energy, derived from cohesive-zone models of symmetrical and asymmetrical circular and elliptical
706 ruptures, *Journal of Geophysical Research: Solid Earth*, **120**(2), 1053–1079.
- 707 Ktenidou, O.-J. & Abrahamson, N., 2016. Empirical estimation of high-frequency ground motion on hard
708 rock, *Seismological Research Letters*, **87**(6), 1465–1478.
- 709 Ktenidou, O.-J., Gélis, C., & Bonilla, L.-F., 2013. A study on the variability of kappa (κ) in a borehole:
710 Implications of the computation process, *Bulletin of the Seismological Society of America*, **103**(2A),
711 1048–1068.
- 712 Ktenidou, O.-J., Cotton, F., Abrahamson, N., & Anderson, J., 2014. Taxonomy of κ : A review of defini-
713 tions and estimation approaches targeted to applications, *Seismological Research Letters*, **85**(1), 135–146.
- 714 Ktenidou, O.-J., Abrahamson, N., Drouet, S., & Cotton, F., 2015. Understanding the physics of kappa (κ):
715 insights from a downhole array, *Geophysical Journal International*, **203**(1), 678–691.
- 716 Ktenidou, O.-J., Abrahamson, N., Darragh, R., & Silva, W., 2016. A methodology for the estimation
717 of kappa (κ) for large datasets: example application to rock sites in the NGA-East database, *Pacific*
718 *Earthquake Engineering Research Center Report 2016/01, Berkeley CA, United States*.
- 719 Ktenidou, O.-J., Silva, W., Darragh, R., Abrahamson, N., & Kishida, T., 2017. Squeezing kappa (κ) out of
720 the Transportable Array: a strategy for using bandlimited data in regions of sparse seismicity, *Bulletin of*
721 *the Seismological Society of America*, **107**(1), 256–275.
- 722 Laurendeau, A., Cotton, F., Ktenidou, O.-J., Bonilla, L.-F., & Hollender, F., 2013. Rock and stiff soil
723 site amplification: dependency on v_{S30} and kappa (κ_0), *Bulletin of the Seismological Society of America*,
724 **103**(6).
- 725 Lay, T. & Wallace, T., 1995. *Modern global seismology*, Academic Press.
- 726 Leary, P. & Abercrombie, R., 1994. Frequency-dependent crustal scattering and absorption at 5-160 Hz
727 from coda decay observed at 2.5 km depth, *Geophysical Research Letters*, **21**(11), 971–974.
- 728 Lermo, J. & Chávez-García, F., 1993. Site effect evaluation using spectral ratios with only one station,

- 729 *Bulletin of the Seismological Society of America*, **83**(5), 1574–1594.
- 730 Madariaga, R., 1976. Dynamics of an expanding circular fault, *Bulletin of the Seismological Society of*
731 *America*, **66**(3), 639–666.
- 732 Matérn, B., 1960. Spatial variation, *Technical Report, Statens Skogsforsningsinstitut, Stockholm*.
- 733 McShane, B., Gal, D., Gelman, A., Robert, C., & Tackett, J., 2017. Abandon statistical significance,
734 *arXiv:1709.07588*.
- 735 McVerry, G., Zhao, J., Abrahamson, N., & Somerville, P., 2006. New Zealand acceleration response
736 spectrum attenuation relations for crustal and subduction zone earthquakes, *Bulletin of the New Zealand*
737 *Society for Earthquake Engineering*, **39**(1), 1–58.
- 738 Mereu, R., Dineva, S., & Atkinson, G., 2013. The application of velocity spectral stacking to extract infor-
739 mation on source and path effects for small-to-moderate earthquakes in Southern Ontario with evidence
740 for constant-width faulting, *Seismological Research Letters*, **84**(5), 899–916.
- 741 Mooney, H., 1970. Upper mantle inhomogeneity beneath New Zealand: seismic evidence, *Journal of*
742 *Geophysical Research*, **75**(2), 285–309.
- 743 Morozov, I., 2008. Geometrical attenuation, frequency dependence of Q, and the absorption band problem,
744 *Geophysical Journal International*, **175**(1), 239–252.
- 745 Motazedian, D. & Atkinson, G., 2005. Stochastic finite-fault modeling based on a dynamic corner fre-
746 quency, *Bulletin of the Seismological Society of America*, **95**(3), 995–1010.
- 747 Oth, A. & Kaiser, A. E., 2014. Stress release and source scaling of the 2010–2011 Canterbury, New
748 Zealand earthquake sequence from spectral inversion of ground motion data, *Pure and Applied Geo-*
749 *physics*, **171**(10), 2767–2782.
- 750 Oth, A., Bindi, D., Parolai, S., & Di Giacomo, D., 2011. Spectral analysis of K-NET and KiK-net data in
751 Japan, part II: On attenuation characteristics, source spectra, and site response of borehole and surface
752 stations, *Bulletin of the Seismological Society of America*, **101**(2), 667–687.
- 753 Papageorgiou, A. & Aki, K., 1983. A specific barrier model for the quantitative description of inhom-
754 geneous faulting and the prediction of strong ground motion. I. Description of the model, *Bulletin of the*
755 *Seismological Society of America*, **73**(3), 693–722.
- 756 Pei, S., Cui, Z., Sun, Y., Toksöz, M. N., Rowe, C. A., Gao, X., Zhao, J., Liu, H., He, J., & Morgan,
757 F. D., 2009. Structure of the upper crust in Japan from S-wave attenuation tomography, *Bulletin of the*
758 *Seismological Society of America*, **99**(1), 428–434.
- 759 Perron, V., Hollender, F., Bard, P.-Y., Gélis, C., Guyonnet-Benaize, C., Hernandez, B., & Ktenidou, O.-J.,
760 2017. Robustness of kappa (κ) measurement in low-to-moderate seismicity areas: insight from a site-
761 specific study in Provence, France, *Bulletin of the Seismological Society of America*, **107**(5), 2272–2292.
- 762 Petersen, T., Gledhill, K., Chadwick, M., Gale, N., & Ristau, J., 2011. The New Zealand national seismo-
763 graph network, *Seismological Research Letters*, **82**(1), 9–20.

30 *C. Van Houtte et al.*

- 764 Silva, W. & Darragh, R., 1995. Engineering characterization of earthquake strong ground motion recorded
 765 at rock sites, *Electric Power Research Institute, Report TR-102261, Palo Alto, California.*
- 766 Standards New Zealand, 2004. Structural Design Actions Part 5: Earthquake Actions, *Department of*
 767 *Building and Housing, Wellington, New Zealand.*
- 768 Steidl, J., Tumarkin, A., & Archuleta, R., 1996. What is a reference site?, *Bulletin of the Seismological*
 769 *Society of America*, **86**(6), 1733–1748.
- 770 Van Houtte, C., Drouet, S., & Cotton, F., 2011. Analysis of the origins of κ (kappa) to compute hard
 771 rock to rock adjustment factors for GMPEs, *Bulletin of the Seismological Society of America*, **101**(6),
 772 2926–2941.
- 773 Van Houtte, C., Ktenidou, O., Larkin, T., & Holden, C., 2014. Hard-site κ_0 (kappa) calculations for
 774 Christchurch, and comparison with local ground-motion prediction models, *Bulletin of the Seismological*
 775 *Society of America*, **104**(4), 1899–1913.
- 776 Wilson, C., Houghton, B., McWilliams, M., Lanphere, M., Weaver, S., & Briggs, R., 1995. Volcanic
 777 and structural evolution of Taupo Volcanic Zone, New Zealand: a review, *Journal of Volcanology and*
 778 *Geothermal Research*, **68**(1), 1–28.
- 779 Yogomida, K. & Benites, R., 1995. Relation between direct wave Q and coda Q: a numerical approach,
 780 *Geophysical Journal International*, **123**(2), 471–483.
- 781 Yoshimoto, K. & Okada, M., 2009. Frequency-dependent attenuation of S-waves in the Kanto region,
 782 *Japan, Earth, Planets and Space*, **61**(9), 1067–1075.
- 783 Yoshimoto, K., Sato, H., & Ohtake, M., 1993. Frequency-dependent attenuation of P and S waves in the
 784 Kanto area, Japan, based on the coda-normalization method, *Geophysical Journal International*, **114**(1),
 785 165–174.

Table 1: Summary of κ data from New Zealand. Refer to Figure 1 for station locations.

Station	κ_0 (s)	Slope $d\kappa/dR$	Number of recordings
<i>Fiordland</i>			
DCZ	0.013	3.62×10^{-4}	41
		$3.34 \times 10^{-4} *$	
PYZ	0.014	3.11×10^{-4}	22
		$3.34 \times 10^{-4} *$	
<i>Otago</i>			
WHZ	0.009	1.79×10^{-4}	28
SYZ	0.012	1.72×10^{-4}	18
		$2.51 \times 10^{-4} *$	
TUZ	0.006	2.54×10^{-4}	50
		$2.51 \times 10^{-4} *$	
OPZ	0.010	3.15×10^{-4}	24
		$2.51 \times 10^{-4} *$	
<i>Canterbury</i>			
EAZ	0.016	3.33×10^{-4}	50

Table 1: continued.

Station	κ_0 (s)	Slope $d\kappa/dR$	Number of recordings
JCZ	0.027	2.94×10^{-4}	45
LBZ	0.022	2.14×10^{-5}	37
RPZ	0.021	2.93×10^{-4}	26
LTZ	0.016	3.85×10^{-4}	22
KHZ	0.031	5.16×10^{-4}	44
<i>Banks Peninsula[†]</i>			
CRLZ	0.032	-	-
MQZ	0.030	-	-
AKSS	0.033	-	-
D14C	0.025	-	-
MTPS	0.039	-	-
<i>West Coast</i>			
FOZ	0.019	1.63×10^{-4}	30
WVZ	0.009	4.08×10^{-4}	19
INZ	0.022	3.91×10^{-4}	16
DSZ	0.023	1.91×10^{-4}	28
THZ	0.012	4.18×10^{-4}	50
NNZ	0.021	5.78×10^{-4}	21
<i>Forearc / East Coast Ranges</i>			
PAWZ	0.052	5.93×10^{-4}	11
BFZ	0.050	5.27×10^{-4}	11
PXZ	0.044	3.87×10^{-4}	9
KNZ	0.042	9.67×10^{-4}	9
CKHZ	0.053	7.35×10^{-4}	15
PUZ	0.048	1.61×10^{-3}	39
MWZ	0.050	2.82×10^{-4}	18
MXZ	0.030	1.49×10^{-3}	15
<i>Arc ranges</i>			
WEL	0.032	4.10×10^{-4} $3.73 \times 10^{-4} *$	48
BHW	0.034	4.77×10^{-4} $3.73 \times 10^{-4} *$	13
BKZ	0.029	3.38×10^{-4}	29
URZ	0.030	3.64×10^{-4}	25
<i>Taupo Volcanic Zone</i>			
WHTZ	0.055	4.34×10^{-4}	12
UTU	0.053	4.46×10^{-4}	10
MARZ	0.052	6.59×10^{-4}	17
<i>Northern Districts</i>			
WAZ	0.019	1.88×10^{-4} $2.10 \times 10^{-4} *$	76
VRZ	0.017	3.13×10^{-4} $2.10 \times 10^{-4} *$	14
HIZ	0.015	2.38×10^{-4}	19
TLZ	0.016	3.24×10^{-4}	14

32 *C. Van Houtte et al.*

Table 2. Results for the two semi-variogram models described in the text. *LL* represents the model log-likelihood and *AIC* represents the Akaike Information Criterion.

<i>Model 1 - trend surface for TVZ and free τ^2</i>							
Order	β_0	β_1	σ^2	τ^2	ϕ	<i>LL</i>	<i>AIC</i>
$\theta = 0.5$	-1.630	0.355	0.055	0.003	274.2	21.81	-33.62
$\theta = 1.5$	-1.624	0.350	0.045	0.008	105.5	21.80	-33.60
$\theta = 2.5$	-1.622	0.344	0.044	0.009	76.5	21.82	-33.64
<i>Model 2 - trend surface for TVZ and fixed $\tau^2 = 0.03$</i>							
Order	β_0	β_1	σ^2		ϕ	<i>LL</i>	<i>AIC</i>
$\theta = 0.5$	-1.654	0.294	0.045		646.0	14.31	-20.62

Table 1: continued.

Station	κ_0 (s)	Slope $d\kappa/dR$	Number of recordings
TGRZ	0.024	2.80×10^{-4}	29
TOZ	0.027	1.63×10^{-4}	26
WIAZ	0.030	2.20×10^{-4}	13
WCZ	0.029	5.39×10^{-4}	7

* Regionally fixed path dependence.

† Values calculated in Van Houtte et al. (2014).

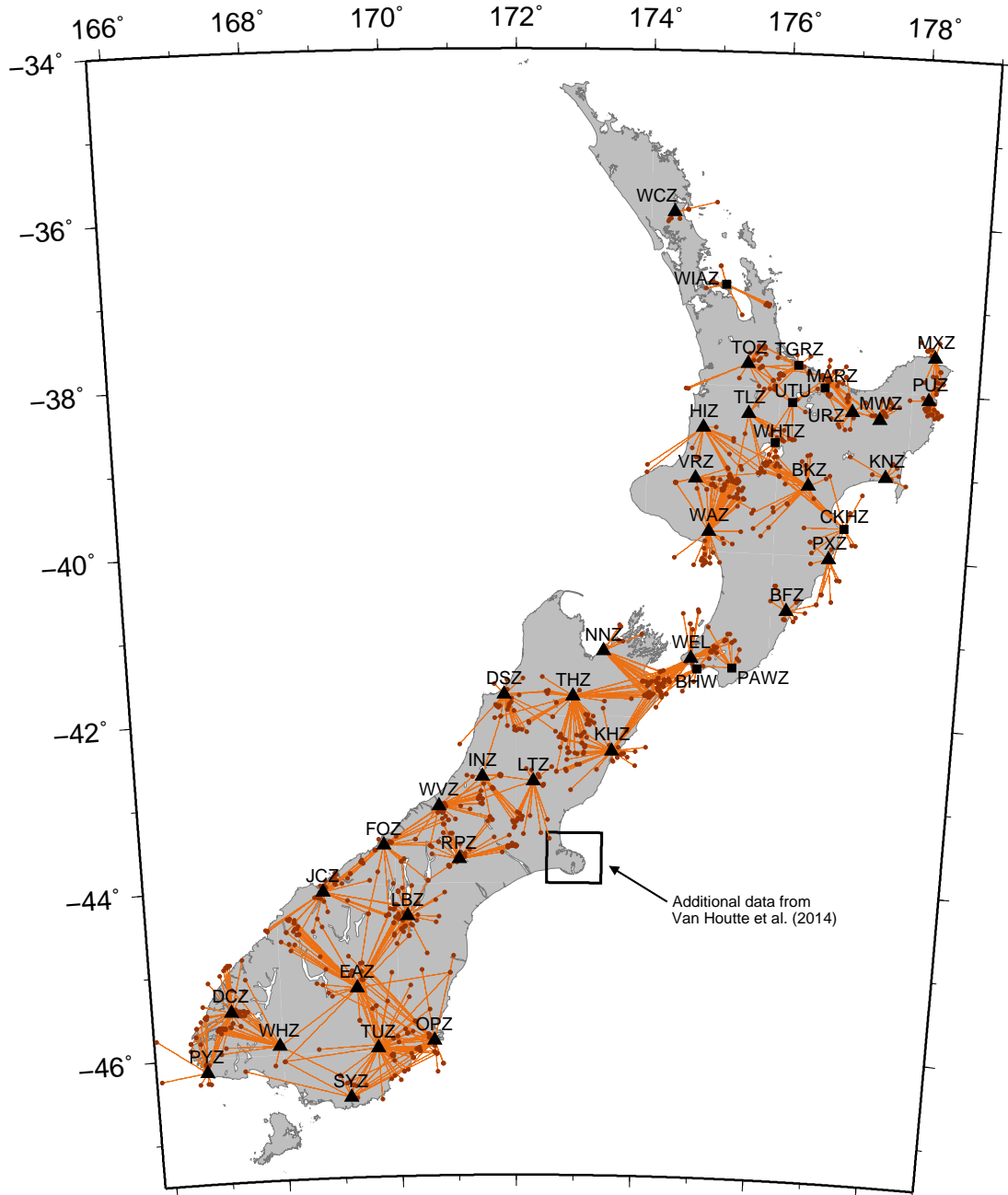


Figure 1. Map of continuous recording stations utilised from the New Zealand National Seismic Network (triangles), and supplementary stations from volcanic or regional networks (squares) that were also used in this study. Circles and lines indicate events and event-station paths respectively.

34 *C. Van Houtte et al.*

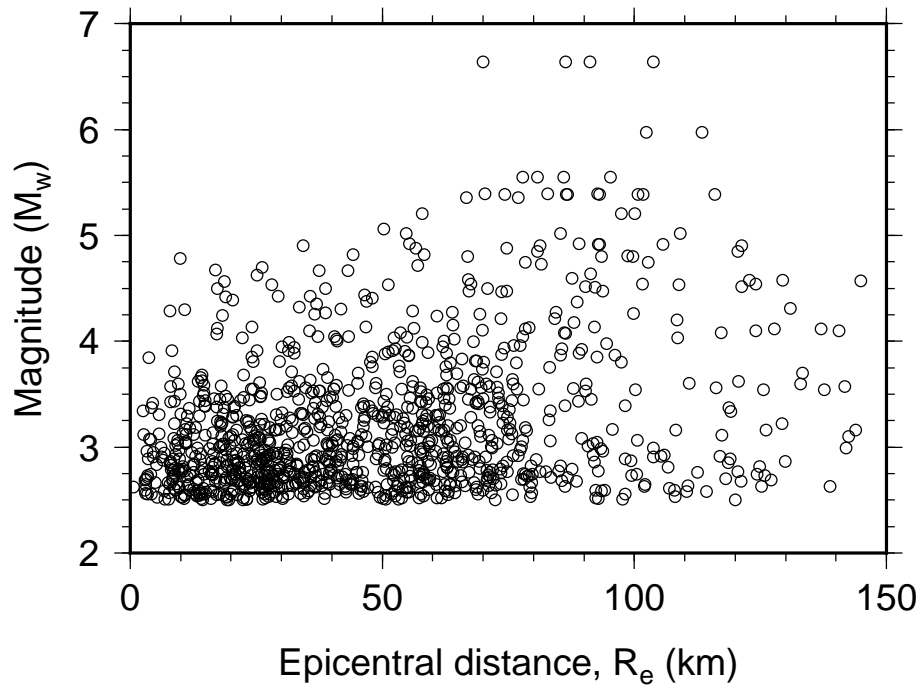


Figure 2. Magnitude and distance distribution of earthquakes that form the dataset for this study.

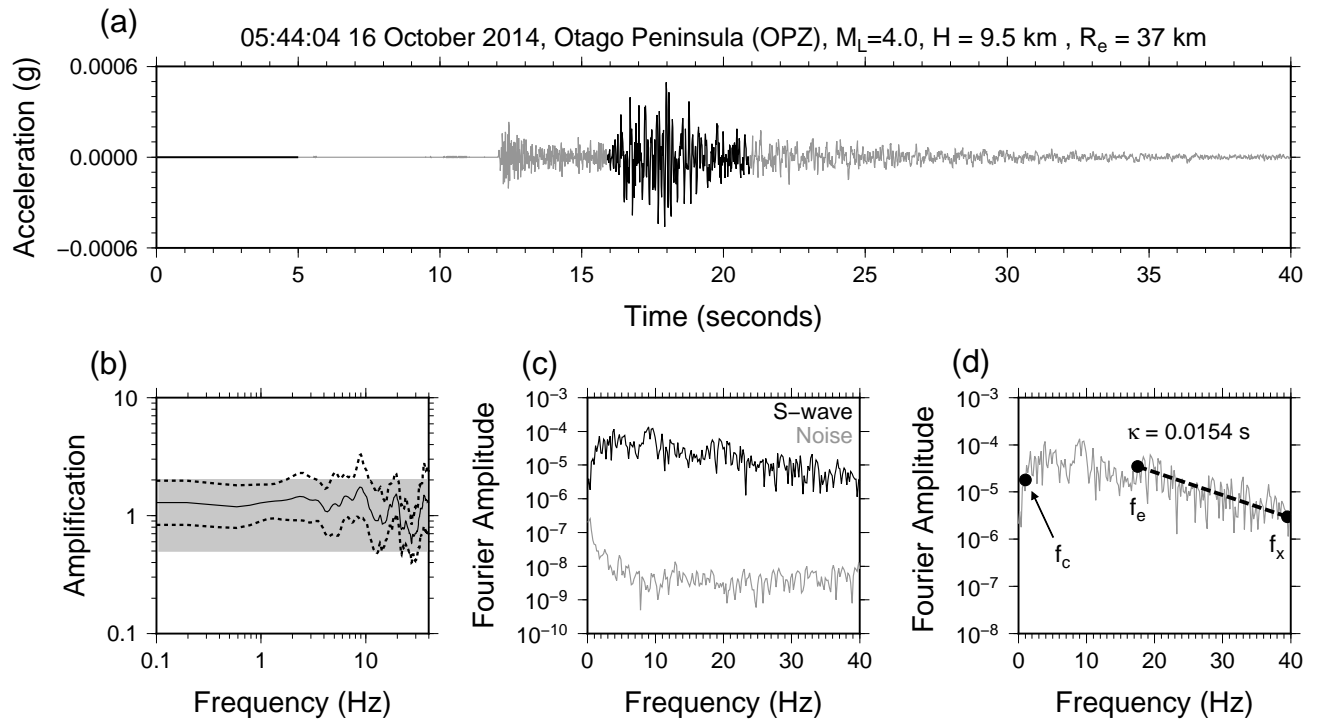


Figure 3. Example calculation of a low κ value. (a) S-wave and noise windows from an earthquake recorded by the OPZ station, (b) the HVSR for OPZ averaged over 23 events, where dashed lines represent ± 1 standard deviation from the median (solid line) and shaded area indicates the region where the amplification can be considered insignificant, (c) the FAS for the noise and S-wave windows in (a), and (d) f_c , f_e , and f_x picks for calculating κ .

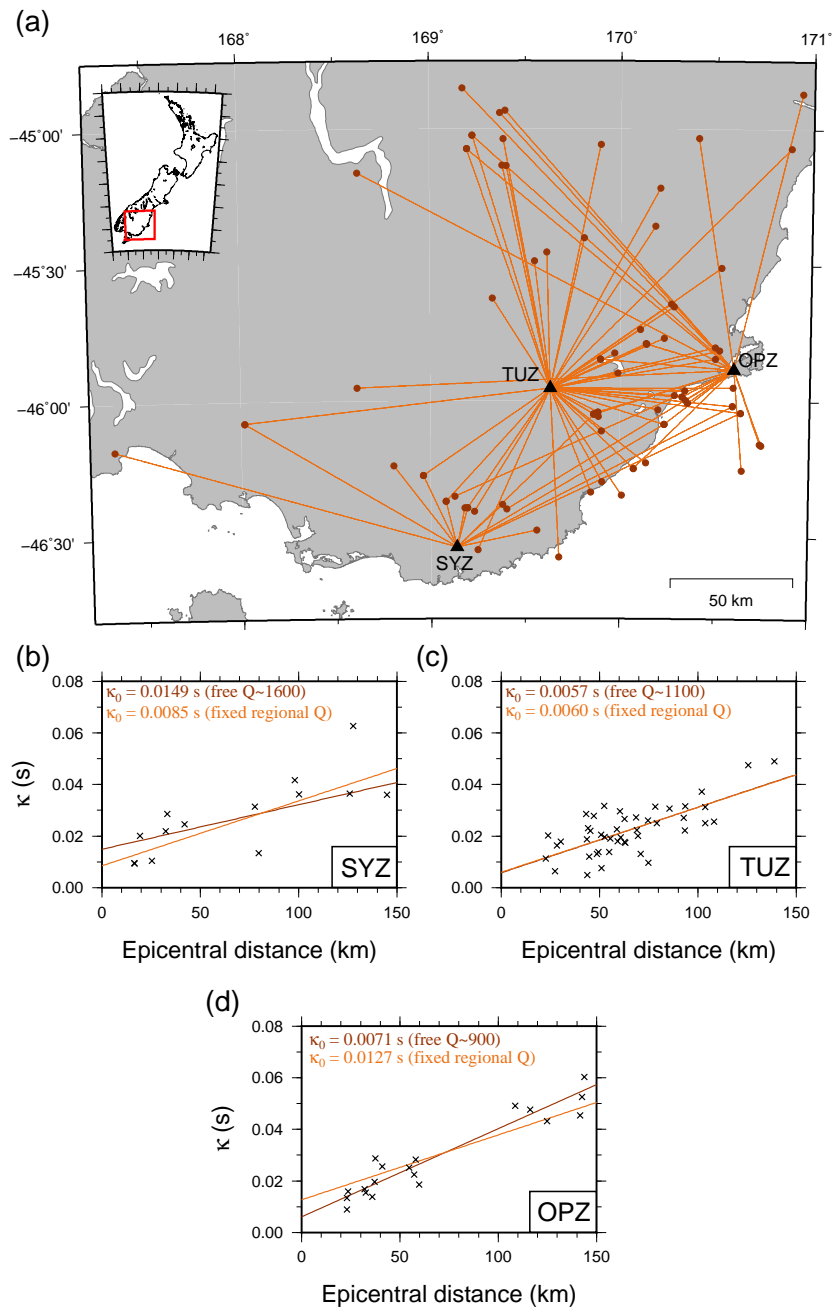
36 *C. Van Houtte et al.*

Figure 4. (a) Locations of three stations in Otago, with events and event-station paths indicated by circles and lines respectively, and (b)-(d) their κ data. Dark lines in (b)-(d) indicate a ‘free Q ’ distance-dependence model derived separately for each individual station, while lighter lines indicate a regionally-fixed $Q=1100$ model, common for all stations. κ_0 calculated based on free and fixed Q assumptions are indicated on the plots.

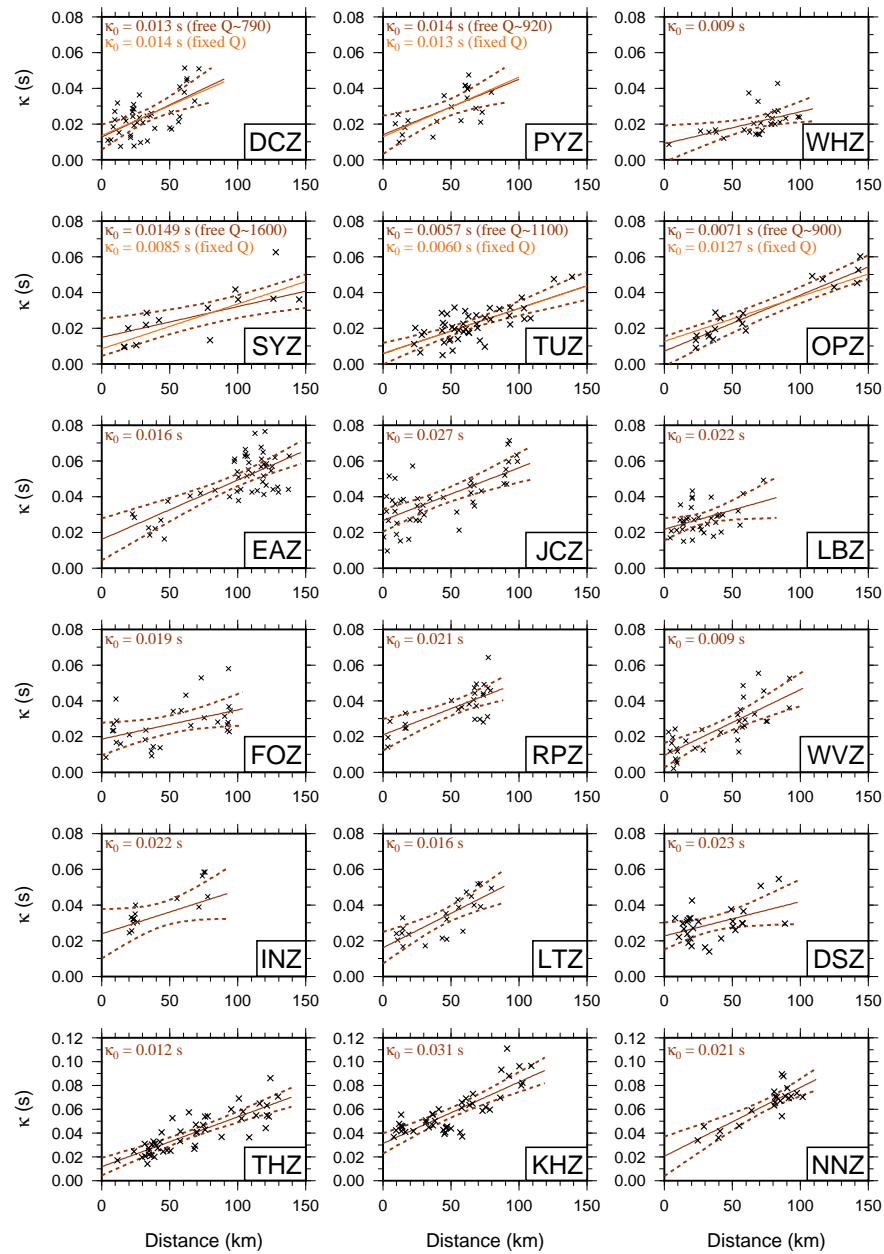


Figure 5. κ data for South Island stations in this study, ordered roughly south-to-north, and by region (refer to Figure 1). Dark lines indicate an unconstrained regression per station (free Q), while lighter lines indicate fixed regional Q slopes calculated using κ data from multiple neighbouring stations. Calculated κ_0 values are also indicated in the plots.

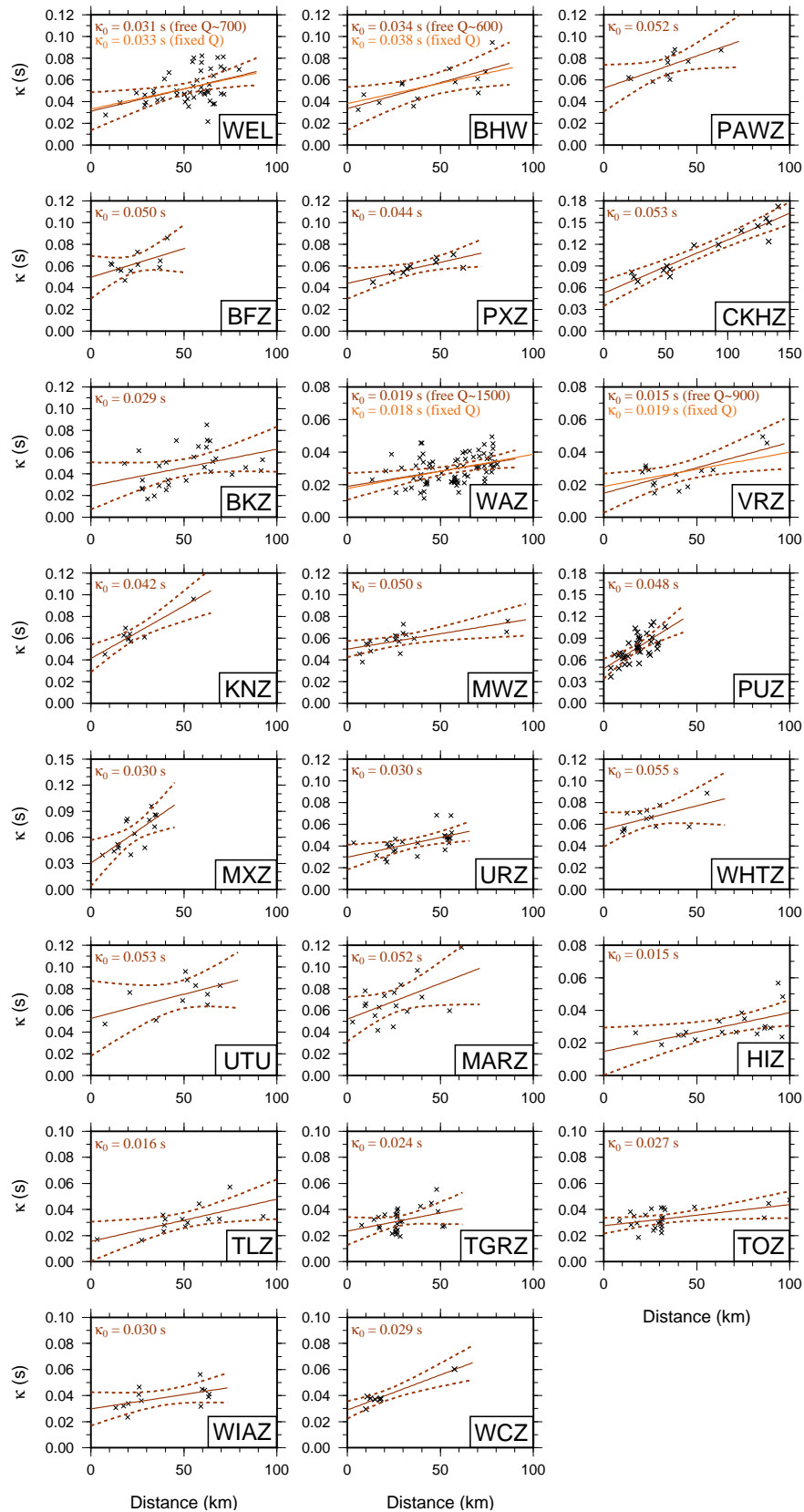


Figure 6. κ data for North Island stations in this study. Refer to Figure 1 for station locations. As with Figure Dark lines indicate an unconstrained regression (free Q), while light lines indicate fixed regional Q slopes calculated using κ data from adjacent stations. Dashed lines represent 5-95% confidence intervals from the unconstrained regression. Calculated κ_0 values are indicated in the plots.

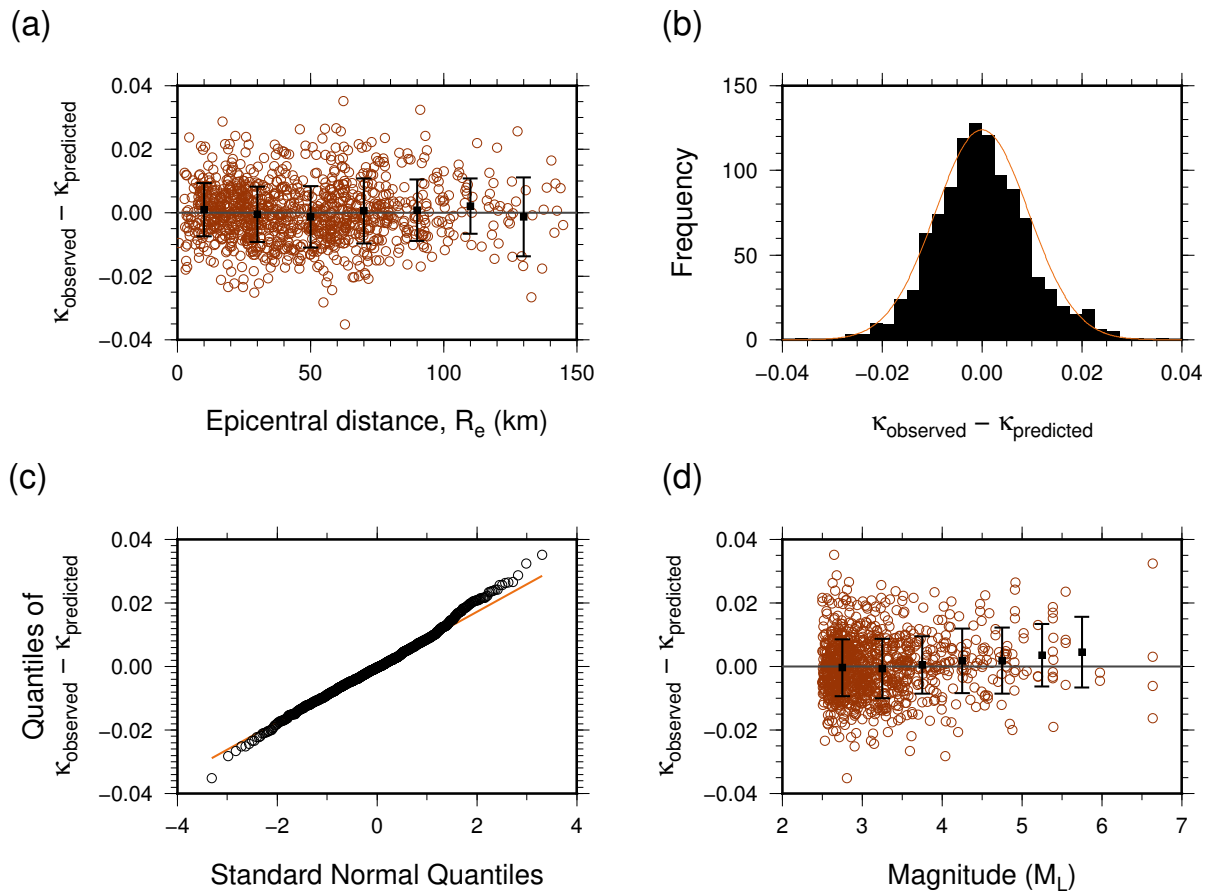


Figure 7. (a) The residuals between the κ data and the model in equation 2, including mean and standard deviation of 20 km distance bins, indicated as squares. (b) A histogram of the residuals, with the solid line indicating the corresponding Gaussian function. (c) A comparison of the residuals with theoretical normal quantiles. (d) Residuals plotted against magnitude.

40 *C. Van Houtte et al.*

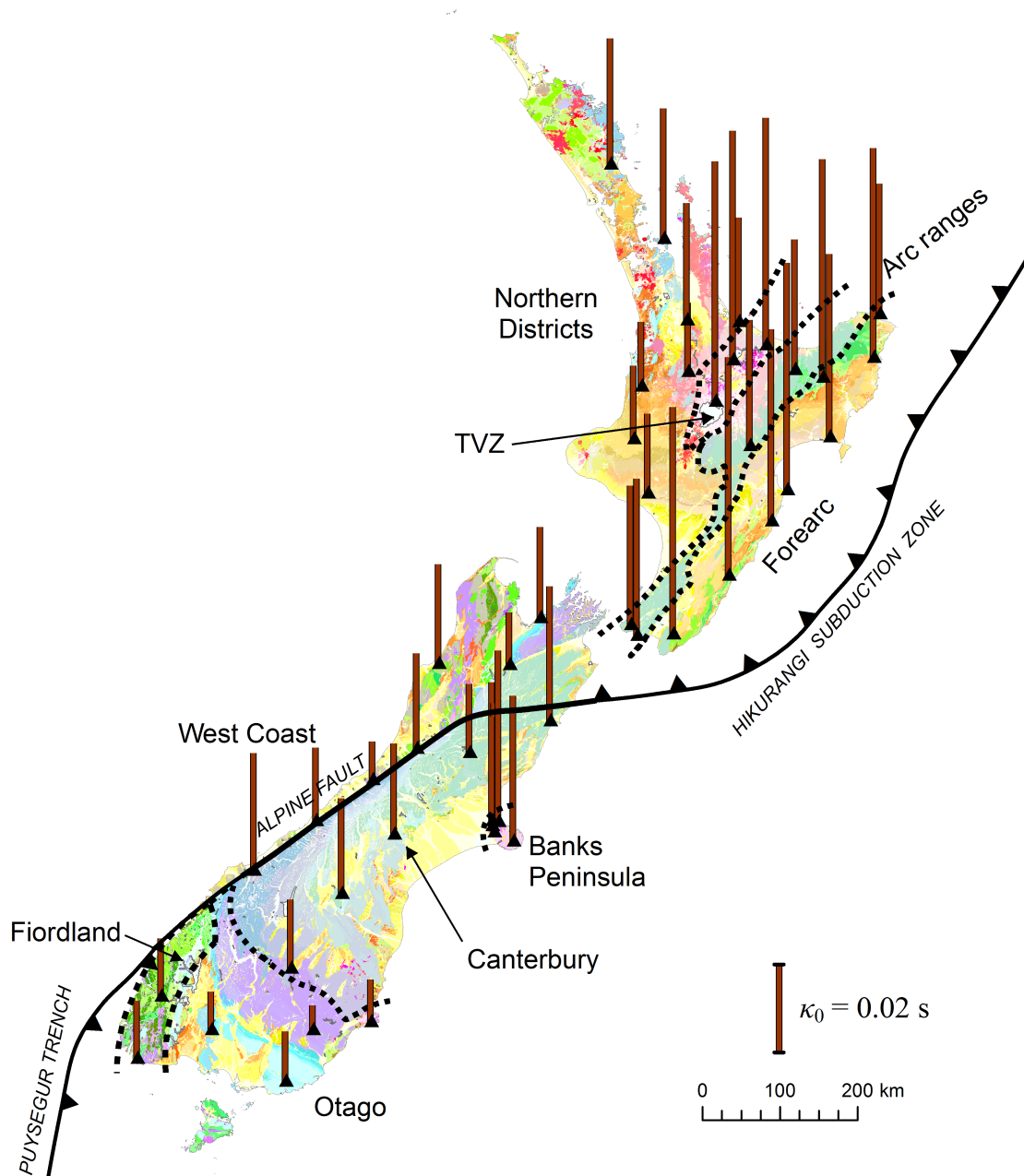


Figure 8. Calculated κ_0 values for the stations analysed in this study, superimposed on a geological map. The length of the bars is proportional to the κ_0 values, with a scale included for reference. Dashed lines indicate boundaries of regions with similar κ_0 values (discussed in detail within the text), while the solid black line is the boundary between the Australian and Pacific tectonic plates.

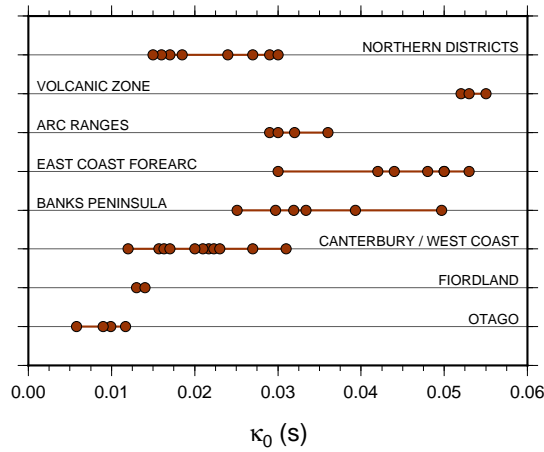


Figure 9. Regional differences in κ_0 . Each point corresponds to an individual station κ_0 value, grouped into the regions indicated by Figure 8.

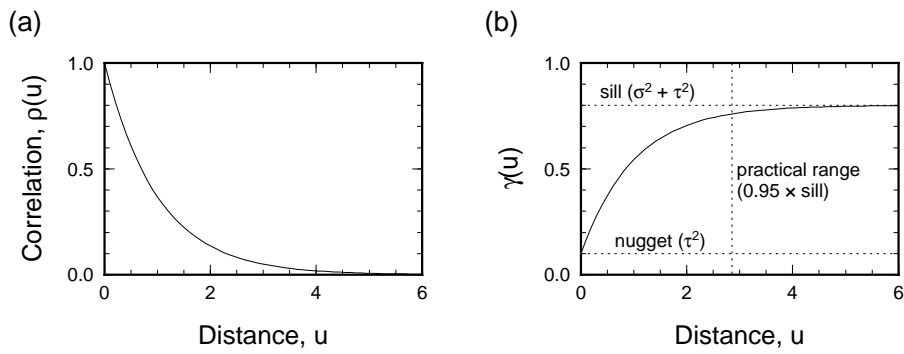


Figure 10. An example spatial correlation function (a), and its subsequent semi-variogram (b). Explanation of functions and terms is included within the text.

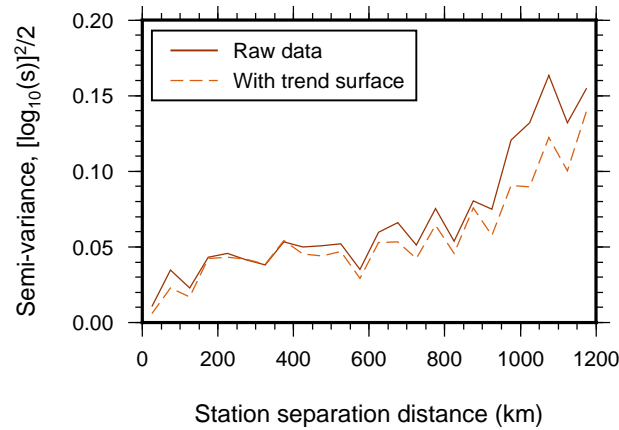


Figure 11. Empirical semi-variograms of the logarithm of the New Zealand κ_0 data, in 50 km distance bins. Solid line represents the semi-variances calculated based on inter-station distance alone, while dashed line is calculated using a spatially-varying mean, using an indicator variable for stations in the Taupo Volcanic Zone.

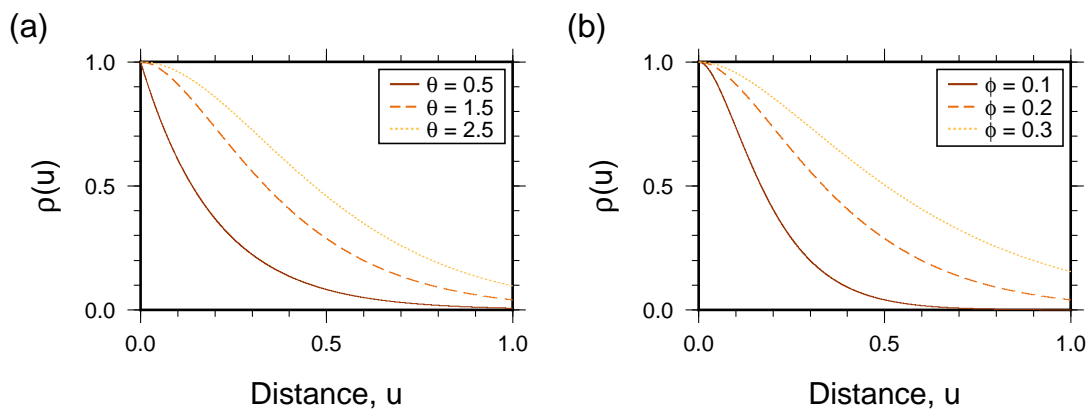


Figure 12. (a) The effect of θ on the Matérn correlation function, with a fixed value of $\phi = 0.2$. (b) The effect of ϕ , for a θ value of 1.5.

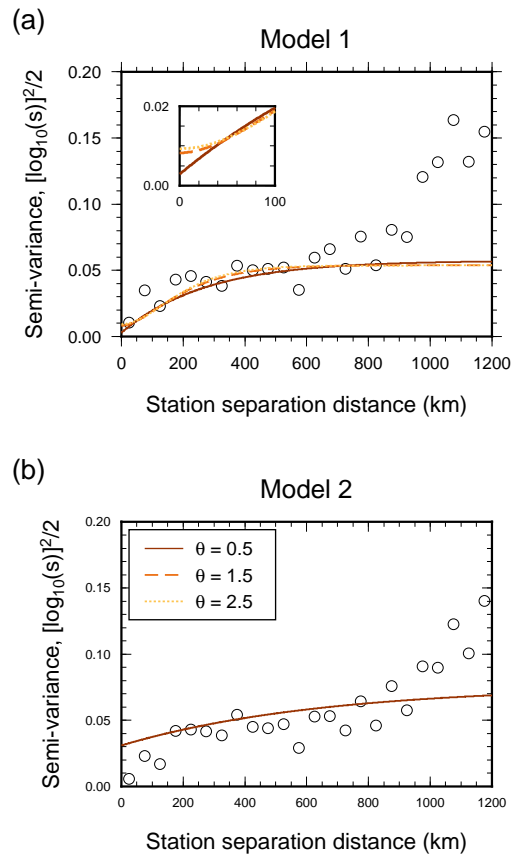


Figure 13. The binned empirical variogram for the logarithm of the New Zealand κ_0 data, corresponding to (a) ‘Model 1’ and (b) ‘Model 2’. Solid, dashed and dotted lines correspond to models with Matérn orders of 0.5, 1.5 and 2.5 respectively, and other parameters determined by maximum likelihood estimation. Circles in (a) and (b) correspond to the dashed lines in Figure 11. The inset in (a) is a more detailed view of the short distance behaviour of Model 1 for different values of the Matérn order.

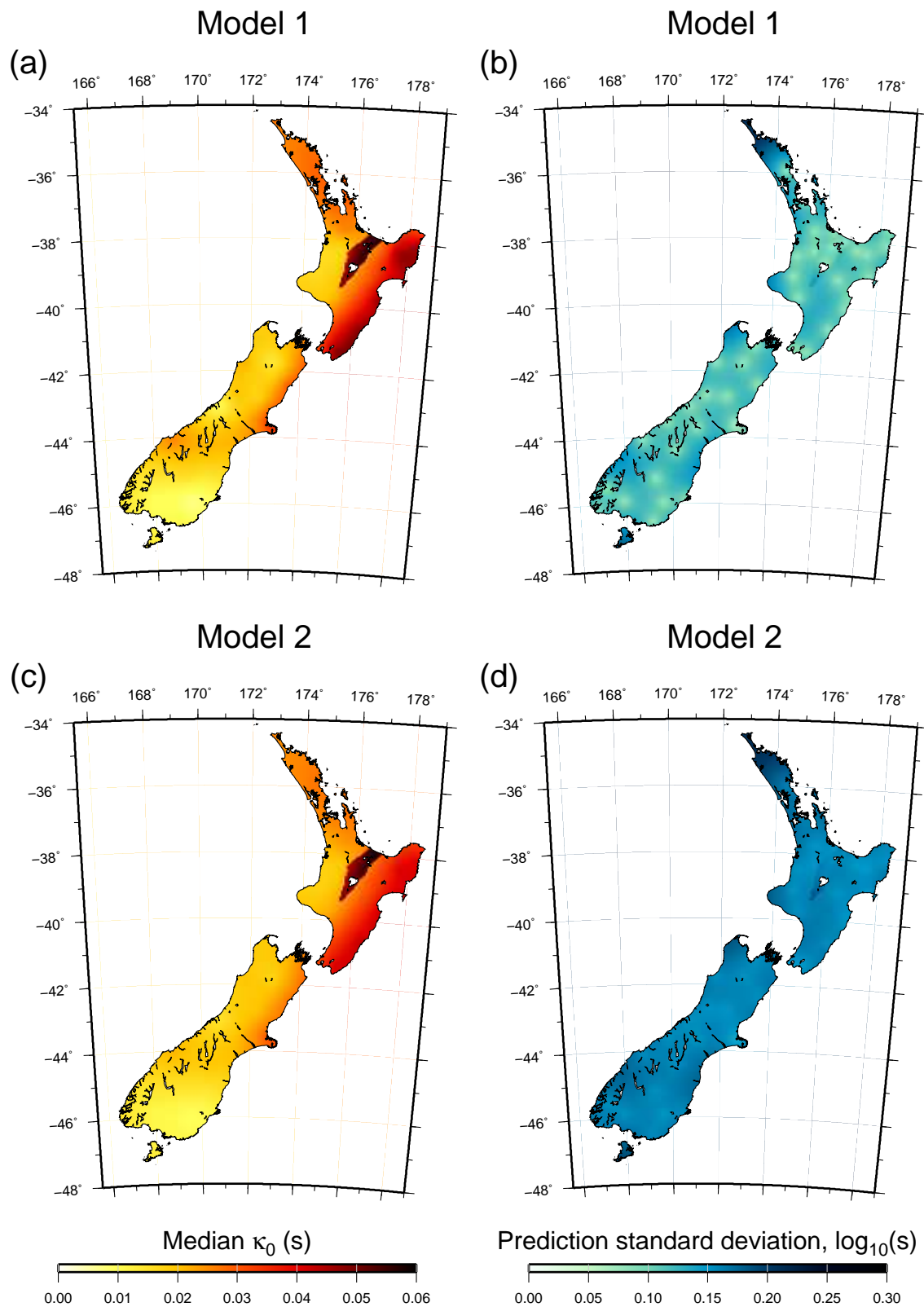
44 *C. Van Houtte et al.*

Figure 14. (a) Median κ_0 map for model 1 with $\theta = 0.5$ and (b) its standard deviation in $\log_{10}(s)$. (c) The median model for model 2 with $\theta = 0.5$ and (d) its standard deviation.

submitted to *Geophys. J. Int.*

Supporting Information for "A continuous near-surface S-wave attenuation map of New Zealand"

Chris Van Houtte^{1,2*}, Olga-Joan Ktenidou^{3,4}, Tam Larkin² and Caroline Holden¹

¹ *GNS Science, 1 Fairway Drive, Lower Hutt 5010, New Zealand*

² *Department of Civil Engineering, University of Auckland, 20 Symonds Street, Auckland 1010, New Zealand*

³ *Department of Engineering Science, University of Greenwich, Medway Campus, Central Avenue, Chatham Maritime,*

⁴ *Institut des Sciences de la Terre, ISTerre, Universite de Grenoble 1, CNRS, F-38041 Grenoble, France*

31 July 2017

Contents of this file

(i) Figure S1.

Details of supplementary material

The supplementary material provides the results of leave-one-out cross-validation. The results are presented using the jackknife standard deviation, for both model 1 and model 2, in Figure S1. The jackknife standard deviation is smaller for model 2 than model 1 and is more spatially homogeneous. For both models the jackknife variance is much smaller than the κ_0 prediction variance.

2 *C. Van Houtte et al.*

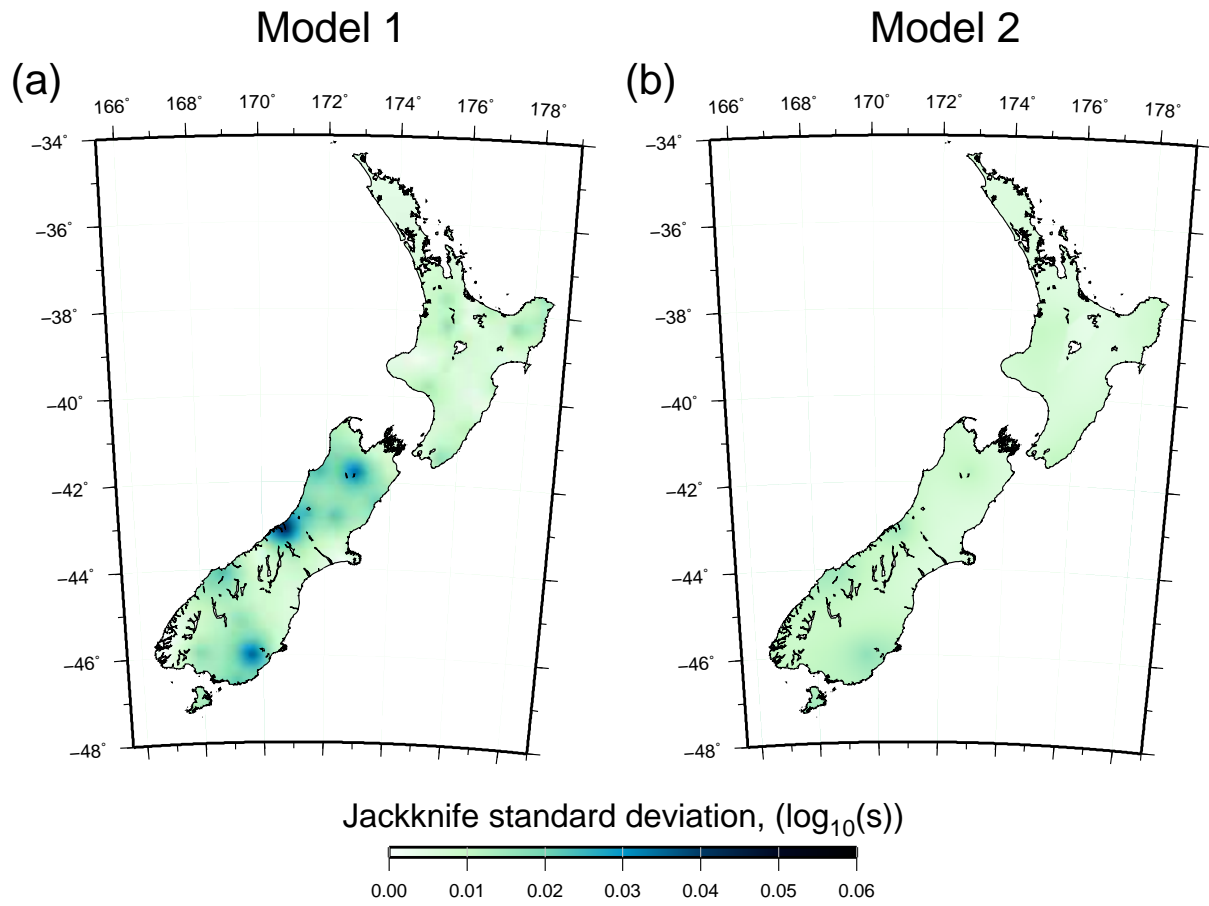


Figure S1. Jackknife standard deviation for Model 1 (a) and Model 2 (b).