

Vertical density gradient in the eastern North Atlantic during the last 30,000 years.

Rogerson, M¹, Bigg, G.R.², Rohling, E.J.³ and Ramirez, J¹.

¹ Geography Department, University of Hull, Cottingham Road, Hull, HU6 7RX, UK, m.rogerson@hull.ac.uk

² Department of Geography, University of Sheffield, Winter Street, S10 2TN, grant.bigg@sheffield.ac.uk

³ School of Ocean and Earth Science, University of Southampton, National Oceanography Centre Southampton, SO14 3ZH, e.rohling@noc.soton.ac.uk

Abstract.

Past changes in the density and momentum structure of oceanic circulation are an important aspect of changes in the Atlantic Meridional Overturning Circulation and consequently climate. However, very little is known about past changes in the vertical density structure of the ocean, even very extensively studied systems such as the North Atlantic. Here we exploit the physical controls on the settling depth of the dense Mediterranean water plume derived from the Strait of Gibraltar to obtain the first robust, observations-based, probabilistic reconstruction of the vertical density gradient in the eastern North Atlantic during the last 30,000 years. We find that this gradient was weakened by more than 50%, relative to the present, during the last Glacial Maximum, and that changes in general are associated with reductions in AMOC intensity. However, we find only a small change during Heinrich Event 1 relative to the Last Glacial Maximum, despite strong evidence that overturning was substantially altered. This implies that millennial-scale changes may not be reflected in vertical density structure of the ocean, which may be limited to responses on an ocean-overturning timescale or longer. Regardless, our novel reconstruction of Atlantic density structure can be used as the basis for a dynamical measure for validation of model-based AMOC reconstructions. In addition, our general approach is transferrable to other marginal sea outflow plumes, to provide estimates of oceanic vertical density gradients in other locations.

38 Atlantic Ocean: Meridional Overturning: Overflow Physics: Mediterranean Outflow:
39 Iberian Margin: Paleoceanography

40 **Introduction**

41

42 Progress in the reconstruction of past Atlantic Meridional Overturning Circulation
43 (AMOC) changes (McManus et al. 2004; Robinson et al. 2005; Hall et al. 2006;
44 Stanford et al. 2006; Lynch-Stieglitz et al. 2007) has revealed that AMOC reductions
45 coincided with colder episodes within the Last Glacial, especially Heinrich Events
46 (McManus et al. 2004; Robinson et al. 2005). Also, a prominent chemocline has been
47 identified at around 2000m depth in the North Atlantic during the Last Glacial
48 Maximum (LGM) and Heinrich Event 1 (H1) (Robinson et al. 2005; Grootes et al.
49 2008), which suggests an altered deep-water circulation state. However, so far hardly
50 anything is known about the past subsurface density structure in the North Atlantic.
51 As this structure is fundamental for understanding deep-water circulation (Lynch-
52 Stieglitz et al. 2007), it is critically important that new means are established for
53 assessing changes in oceanic vertical density structures. We present new insight into
54 this structure in the eastern North Atlantic from a novel approach that centres on the
55 physical controls on depth-changes of the Mediterranean Outflow plume through
56 time.

57

58 The settling depth of the plume of dense water that results from subsurface outflow
59 from the Mediterranean through the Strait of Gibraltar (Atlantic Mediterranean Water;
60 AMW) is controlled by: (1) the density anomaly of pure Mediterranean Outflow
61 Water (MOW) at its exit from the Strait of Gibraltar (SoG); (2) the degree of mixing
62 this water experiences as it descends down the slope of the Gulf of Cadiz; and (3) the
63 density structure of the Atlantic Ocean to the west of the Gulf of Cadiz (O'Neill-
64 Baringer and Price 1997). The density anomaly of pure MOW (relative to ambient
65 North Atlantic waters) has varied through time, mainly due to hydraulic controls
66 imposed by changes in eustatic sea level on the geometry of the SoG (Rogerson et al.
67 2005), with second-order influences from regional climate changes (Voelker et al.
68 2006; Rogerson et al. 2010).

69

70 Globally cold periods such as the LGM and H1 coincide with generally low sea levels
71 (130-60 m below present) (Siddall et al. 2003), high buoyancy loss from the
72 northwest Mediterranean (Hayes et al. 2005; Kuhlemann et al. 2008), and
73 consequently very dense MOW (Rogerson et al. 2005). During those times, the AMW
74 plume settled at a much greater depth than today, as evidenced by a withdrawal of
75 flow from the upper slope of the Gulf of Cadiz and enhanced flow on the lower slope
76 (Rogerson et al. 2006; Voelker et al. 2006). A prominent sedimentary unconformity,
77 significant changes in benthic foraminiferal isotopes, and distinct palaeoecological
78 changes (Schönfeld and Zahn 2000) on the Portuguese margin between 1700 and
79 2000 m depth, indicate that AMW influences reached to at least those depths during
80 the LGM. Stable isotope data from the Portuguese margin suggest that AMW
81 influences even reached down to ~2600m during Heinrich Stadials (Skinner and
82 Elderfield 2007). In contrast, the core of the AMW plume today resides at ~800m,
83 with the very deepest AMW influences around 1700m (O'Neill-Baringer and Price
84 1999).

85

86 Overall, to satisfy the whole range of evidence from sedimentary records, it appears
87 that the core of the AMW resided at least 900m deeper during glacial times than today
88 (Rogerson et al. 2005) (Fig. 1a). Although the intuitive expectation would be that
89 denser MOW produces a denser AMW plume that settles at greater depths, this
90 expectation is incorrect. Instead, it has been well established that denser MOW leads
91 to a higher plume velocity, which in turn drives enhanced ambient water entrainment
92 during settling. This is a negative feedback process that is ubiquitous in overflow
93 plumes (Price and O'Neill-Baringer 1994). Thus, an enhanced glacial density contrast
94 at the SoG would result in either a negligible change in the AMW plume settling
95 depth, or even a reduction (Price and O'Neill-Baringer 1994). Here, we quantitatively
96 evaluate the controls on past AMW settling depth to explain the apparent
97 contradiction between theoretical and observational constraints on the system.

98 **Methods**

99

100 We quantify the Mediterranean Outflow and entrainment system concept using a
101 widely accepted theory of marginal sea overflow mixing (Price and O'Neill-Baringer
102 1994), coupled with a somewhat modified version of a model for the Gibraltar

103 Exchange (Bryden and Kinder 1991). We will refer to these two models as "PO94"
 104 and "BK91" respectively. Because glacial conditions cannot be specified without
 105 considerable uncertainty, we instead present our analysis of controls on the AMW
 106 settling depth in a Monte Carlo-style approach across a broad parameter space of
 107 possible conditions. The Mediterranean excess of evaporation over total freshwater
 108 input (X_{Med}), temperature loss during conversion of Atlantic Inflow to Mediterranean
 109 Intermediate and Deep Waters and sea level are all allowed to vary independently,
 110 simultaneously and randomly, with the only constraint being that deep water in the
 111 Western Mediterranean must be of higher density (i.e. lower temperature) than
 112 intermediate water. Each experiment comprises 5,000 iterations of the model.

113

114 The PO94 model assumes that entrainment of ambient water occurs as a single event
 115 around 100km downstream of the Camarinal sill, mixing a single type of
 116 Mediterranean water and a single type of ambient water to produce a single type of
 117 product water (Price and O'Neill-Baringer 1994). It does not account for differential
 118 mixing due to Ekman veering (O'Neill-Baringer and Price 1997). Consequently, this
 119 model provides only a single product water density, which in essence identifies the
 120 mean isopycnal on which the final AMW will settle. The critical parameter for PO94
 121 is the "proportional mixing coefficient"

122

$$123 \quad \Phi = 1 - (B_{geo}^{1/3} / U_{geo})$$

124

125 where B_{geo} is the geostrophic buoyancy flux of the AMW plume and U_{geo} is the
 126 geostrophic velocity. These parameters are given by

127

$$128 \quad U_{geo} = (g' \pi / f)$$

$$129 \quad B_{geo} = (H_{src} U_{src} g') / (1 + 2K_{geo} x / W_{src}),$$

130

131 where π is the gradient of the continental slope, f is the Coriolis parameter (0.000084),
 132 x is the distance downstream from Gibraltar (100km), and g' is the reduced gravity of
 133 pure MO water ($g' = (g \rho_{MO} - \rho_{ATL}) / \rho_{MO}$) where ρ_{MO} and ρ_{ATL} are the densities of
 134 Mediterranean Outflow and Atlantic water). H_{src} , W_{src} , U_{src} are the height, width and
 135 velocity of the flow at the Camarinal Sill, g is the acceleration due to gravity, and K_{geo}

136 is the geostrophic Ekman number that is generally ≤ 0.2 (PO94). Today, $\Phi = 0.58$,
 137 meaning that AMW comprises 42% Mediterranean Water (PO94), which is of similar
 138 order as estimates from direct measurement ($\sim 33\%$) (O'Neill-Baringer and Price
 139 1999). The difference between the observed and modelled values for entrainment
 140 results in the AMW plume in the model having density enhanced by $\sim 1.27 \times 10^{-4} \text{ kg}$
 141 m^{-3} , which is $\sim 7.5\%$ of the initial density anomaly of pure MOW and therefore
 142 negligible. The necessary input parameters for assessment of past changes in Φ
 143 therefore relate to the geometry of the flow in the SoG (H_{src} , U_{src} and W_{src}) and the
 144 reduced density of pure Mediterranean Outflow water (g').

145

146 The required geometric parameters represent conditions at the shallowest point in the
 147 Strait of Gibraltar, the Camarinal sill, which is the location of hydraulic control on the
 148 outflow (Armi and Farmer 1988). W_{src} is altered by changes in water depth above the
 149 sill and thus is a function of global sea level. W_{src} is estimated on the basis of sea-level
 150 influences on the triangular cross section of the sill section with a width of 20km and
 151 a depth of 284m (Bryden and Kinder 1991). U_{src} depends on the sea floor gradient in
 152 the direction of flow over the shallowest part of the sill, the g' of MO water, and (on
 153 short time scales) on a range of tidal and subinertial forces. Given that the average
 154 variability in U_{src} on timescales above decadal is not sensitive to the latter's short-
 155 term influences (Gomis et al. 2006), changes in U_{src} may be viewed as forced only by
 156 changes in g' . This perspective is further enhanced by the observation that flow over
 157 the westernmost sill within the Gibraltar Strait system (the Spartel Sill) is almost
 158 constant throughout the tidal cycle (Bryden et al. 1988; Garcia-Lafuente et al. 2009).
 159 For any g' , H_{src} can therefore be estimated from the relationship of flux (Q_{MO}) to
 160 velocity and cross sectional area. Consequently, the only boundary conditions we
 161 need to supply the PO94 model with are global sea level change, flux (Q_{MO}) and
 162 reduced density of pure Mediterranean water (g'). When these statements are
 163 combined, settling depth of the AMW plume is a simple estimate from:-

164

$$165 \quad D_{\text{settling}} = (D_s - H') + \frac{(\Delta\rho_{\text{MO}} * \Phi)}{(\partial\rho/\partial z)}$$

166

167

168 Where $D_{setling}$ is the PO94 estimate of mean settling depth for the AMW plume in the
 169 Atlantic. This rearranges to

170

171

$$172 \quad \partial\rho / \partial z = \frac{(\Delta\rho_{MO} * \Phi)}{D_{setling} - (D_s - H')}$$

174

175 providing a simple representation of eastern Atlantic vertical density gradient. Density
 176 is a property that varies quite smoothly in the ocean interior, especially westward of
 177 Gibraltar where the oceanography is essentially zonal due to the structure imposed by
 178 the Azores Front (Gould 1985; New et al. 2001; Alves et al. 2002). Consequently, we
 179 anticipate that the Gulf of Cadiz vertical density gradient will be representative at
 180 least of a region extending west to the Mid Atlantic ridge and north to the Bay of
 181 Biscay.

182

183 Following previous studies (Rogerson et al. 2005; Rogerson et al. 2006), we estimate
 184 the flux of outflowing Mediterranean Water at the Camarinal Sill (Q_{MO}), the vertical
 185 density difference at the sill ($\Delta\rho_{MW}$) and the initial reduced density of water mixing in
 186 the Gulf of Cadiz (g') using the model of Gibraltar exchange of Bryden and Kinder
 187 (1991; BK91). It expresses hydraulic control on flow through the sill and narrows
 188 sections of Gibraltar with mass and salt conservation (Bryden and Kinder 1991). We
 189 modify this model here in one important aspect, in that we include sensitivity of g' to
 190 changes in regional sea surface temperature gradients (BK91 model considered only
 191 salinity effects), as these are known to be variable on glacial-interglacial timescales
 192 (Kuhlemann et al. 2008). The BK91 requires an iterative solution for the relationship
 193 between ΔS_{gib} and Q_{total} , and we approach this by exploiting a convergent solution
 194 within the paired equations following the simplification provided by Mikolajewicz
 195 (Mikolajewicz 2010).

196

$$197 \quad \Delta S_{gib} = S_{atl} (X_{med} / (0.5Q_{total} - X_{med}))$$

198

$$199 \quad Q_{total} = ((C ((W_{src} D_s) / 2)) * \sqrt{((\beta \Delta S_{gib} + \alpha \Delta T_{gib}) g D_s / \rho_{MO})})$$

200

201 where C is a geometric coefficient representing the shape of the Strait of Gibraltar, β
 202 and α are the salinity contraction and thermal expansion coefficients respectively and
 203 ΔT_{gib} is the temperature difference between inflowing and outflowing water. Salinity
 204 of the inflowing Atlantic water is calculated directly from the proportional loss of
 205 global ocean volume due to eustatic sea level change ($S_{\text{atl}} = S_{\text{atl-pres}} H / [H - h']$). The
 206 temperature of the deep and intermediate water layers of Mediterranean water that
 207 pass over the sill (Mediterranean Dense Water {MDW} and Mediterranean
 208 Intermediate Water {MIW} respectively), respectively (Millot 2009)) are provided by
 209 arbitrary offsets (between 0 and 6°C) from winter sea surface temperature in the Gulf
 210 of Cadiz. These offsets are randomly generated via a Monte Carlo approach and this
 211 variability is propagated through the rest of the model (see Table 1). Potential density
 212 of the Atlantic, MIW, and MDW watermasses are calculated from the model output
 213 salinity and temperature data and the Levitus and Isayev polynomial approximation of
 214 the equation of state for seawater (Levitus and Isayev 1992).

215

216 To simulate the impact of entraining sediment into the Mediterranean Outflow plume,
 217 which may be relevant to past changes in plume density, we exploit the relationship
 218 between bottom velocity and sediment entrainment, using a simple parameterisation
 219 (Karim and Kennedy 1990) known to be relatively insensitive to errors in estimation
 220 of velocity and grainsize (Pinto et al. 2006). The equation for q_s , the flux of sediment
 221 entrained, is

222

$$223 \quad q_s = 10^{-2.821+33.69 \log(x_1)+0.840 \log(x_2)} \sqrt{((s_d - 1) g d_{50}^3)}$$

224

225 where $x_1 = U_{\text{geo}} / \sqrt{((s_d - 1) g d_{50})}$ and $x_2 = (U^* - U^*_{\text{crit}}) / \sqrt{((s_d - 1) g d_{50})}$. Here, s_d is
 226 the density of the sediment (2.65 kg m⁻³), U^* is the bottom shear velocity ($U^* =$
 227 $\sqrt{(0.8 U_{\text{geo}}^2)}$) and U^*_{crit} is the critical shear velocity for $d_{50} = 2.4$, which is taken from
 228 the Shields diagram (2.4) (Julien 1998). d_{50} is the median grainsize of the sediment
 229 available on the slope. We use 2.4µm for this parameter, which is taken from core top
 230 data (Rogerson et al. 2011) as there is insufficient data available for the sediment
 231 grainsize distribution on the slope during the LGM. As the MOW plume would have
 232 been smaller in the past (Rogerson et al. 2005), it is likely that, on the scale of the
 233 whole slope, the sediment available for erosion was slightly finer. As this would make

234 the sediment more cohesive, raising U^*_{crit} , our assessment of the impact of sediment
235 entrainment should be seen as representing the maximum likely value.

236 **Results and Discussion**

237 The results of our analysis (Fig. 1) reveal a significant positive relationship between
238 AMW settling depth and density increase (buoyancy decrease) in the Mediterranean
239 due to evaporation (Fig. 1b; $r^2 = 0.85$) but only a weak positive relationship with
240 cooling (Fig. 1c, d) or sea level change (Fig. 1e). However, even within the generous
241 bounds of parameter space investigated, changes in these parameters cannot make
242 AMW settle more than ~200m below its present depth, despite the potential for up to
243 130m of displacement directly from sea level change alone.

244

245 Sediment entrainment may provide a mechanism for secondary enhancement of
246 AMW density and consequently might promote greater settling depth, given that the
247 flow's estimated shear velocities (Fig. 1f) mostly exceed the critical value (2.4 m s^{-1})
248 needed to allow strong sediment entrainment of fine silt grade sediment (McCave and
249 Hall 2006). Indeed, we found that for shear velocities higher than $\sim 6.8 \text{ m s}^{-1}$ sediment
250 entrainment resulted in concentrations $>5\%$ by volume (grey area, Fig. 1f). This level
251 of suspension is generally considered as being "hyperconcentration" and such flows
252 are known to exhibit reduced entrainment of sediment, possibly resulting in net
253 deposition, in addition to having different mixing and flow hydraulics to "normal"
254 flows of suspended sediment (Julien 1998). The physics represented in our model will
255 thus overestimate continued entrainment beyond the 5% level, and as it is unlikely
256 that such a flow would be capable of forming a geostrophic current (instead being
257 more likely to behave like a conventional sedimentary turbidity current) we therefore
258 enforce a maximum entrainment level of 5% for the fastest flows. Unless some form
259 of hyperconcentrated flow is permitted, we again find that the large glacial settling
260 depth increase (Fig. 1a) cannot be explained via sediment entrainment (Fig. 2a-c). In
261 our simulations, the maximum impacts of sea-level, temperature and salinity changes
262 within the Mediterranean, and suspended sediment influences, are an increase in
263 settling depth of the AMW core of $\sim 150\text{m}$, which is insufficient to take even the
264 lowermost AMW down to the depths observed for the LGM. A major additional
265 control on settling depth is clearly dominating the system.

266

267 This leads us to the only remaining potential mechanism for explaining the deep
268 glacial settling of the AMW core at depths of ~1700m or more, namely reduction of
269 the vertical density gradient in the eastern North Atlantic. We investigate the
270 influence of this parameter by varying it in our model between 0.1 and 2 times its
271 modern value of $1.1 \times 10^{-4} \text{ kg m}^{-4}$ (Fig 2 d, e, f). This reveals that the observations of
272 AMW influences at 2000m (with a core depth ~1700m) during the LGM (Schönfeld
273 and Zahn 2000), and potentially even deeper during Heinrich Stadials (Skinner and
274 Elderfield 2007), require a glacial vertical density gradient in the eastern North
275 Atlantic that was reduced to less than half its present-day value (Fig. 2f).

276

277 To gauge whether such a change would be physically plausible, we have extracted
278 glacial-interglacial changes in vertical density gradients in the region between 27-
279 37°N and 10-20°W from existing simulations of the ocean part of the ‘Frugal’ climate
280 model (Levine and Bigg 2008), one of the few climate models to actually represent
281 the SoG as a proper strait. Area-mean σ_θ profiles across all model levels between 0
282 and 3000m depth systematically show somewhat lower absolute values than observed
283 from hydrographic data, but the vertical structure agrees well in a relative sense (Fig.
284 3). The vertical density gradients for the LGM, and for a Heinrich Event (where 0.4Sv
285 of freshwater equivalent in the form of icebergs was released from Hudson Strait;
286 (Levine and Bigg 2008), are less than half the present-day gradient; over the 300-
287 1000m depth range, where most of the conversion of MOW to AMW occurs (Price et
288 al. 1993), the vertical gradient is 4.2×10^{-4} and $4.1 \times 10^{-4} \text{ kg m}^{-4}$ in the LGM and
289 Heinrich simulations respectively, relative to $11 \times 10^{-4} \text{ kg m}^{-4}$ in the present-day
290 simulation (Levine and Bigg 2008). These results demonstrate that our inference of a
291 roughly 50% reduced vertical density gradient in the eastern North Atlantic is
292 physically plausible. Moreover, this analysis confirms that the mechanism causing
293 vertical density gradient reduction is linked to weakened AMOC transport. Given that
294 AMOC and Atlantic vertical density gradients are both reflections of the buoyancy
295 budget of the Atlantic, this relationship is not in itself surprising. High AMOC
296 transport must physically coincide with strong buoyancy loss in the Nordic Seas
297 region, and thus with very dense Atlantic interior waters and consequent very strong
298 vertical density gradients. However, our new approach allows this relationship to be

299 investigated directly from observations, providing powerful dynamical constraints to
300 modelling studies of AMOC change.

301

302 The relative insensitivity of the AMW settling depth to parameters other than the
303 Atlantic vertical density gradient allows us to assess the minimum change in density
304 gradient that is compatible with the observed glacial-interglacial AMW settling-depth
305 changes. Figure 4 shows a probabilistic assessment of the density gradient values
306 necessary to achieve AMW flow at depths reported at selected periods over the last
307 30ka (Table 2). This reveals that the most likely density gradient during the LGM was
308 $\sim 3.1 \times 10^{-4} \text{ kg m}^{-4}$ with a $\pm 2\sigma$ range between 2.08×10^{-4} and $4.27 \times 10^{-4} \text{ kg m}^{-4}$. The
309 relative change from today equates to a 52-77% reduction, which encompasses the
310 fractional reduction from the GCM simulations. Also shown in Figure 4 is the same
311 assessment where sediment entrainment is considered (as above, hyperconcentrated
312 flows are excluded) which indicates only minor modification of the Atlantic density
313 gradients required, although the upper limit is higher ($5.47 \times 10^{-4} \text{ kg m}^{-4}$) implying a
314 somewhat smaller reduction relative to the present than in the case without sediment
315 entrainment. Also shown are scenarios for Heinrich Event 1, the Younger Dryas,
316 Bölling-Alleröd and Holocene (see parameterisation in Table 2). There is a clear
317 relationship with northern hemisphere climate, with colder periods exhibiting lower
318 vertical density gradients. However, the Heinrich Event 1 scenario differs only
319 marginally from the LGM scenarios, which stands in contrast to the considerable
320 inferred change in AMOC (McManus et al. 2004; Robinson et al. 2005; Hall et al.
321 2006; Stanford et al. 2006; Lynch-Stieglitz et al. 2007). This implies that, unlike
322 AMOC, the density structure cannot respond to millennial-scale forcing and that its
323 response is limited to ocean-overturning timescales. This, in turn, implies that AMOC
324 is not strictly tied to the Atlantic stratification on short ($<1 \text{ ka}$) timescales, despite the
325 common forcing outlined above.

326

327 Nevertheless, our reconstructions largely compare well with concepts of past AMOC
328 change derived entirely from independent sources of palaeoceanographic data and
329 general circulation models. Moreover, they provide a pivotal deep-water validation to
330 concepts of the dynamically determined structure in the glacial eastern North Atlantic
331 (Levine and Bigg 2008), and so of the glacial AMOC. Previously, model validation

332 largely relied on the distribution of water-mass properties such as surface or benthic
333 temperatures. Our reconstruction of the vertical density gradient – based on robust
334 measurements and quantified using physical relationships – for the first time provides
335 a critical reconstruction of dynamical structure within the ocean interior, for testing
336 models of large-scale ocean circulation.

337 **Conclusions**

- 338 1) The settling depth of the Mediterranean Outflow plume is largely insensitive to
339 changes in watermass properties in the Mediterranean Sea, even on glacial-
340 interglacial timescales. Increased settling depth also cannot be related to sediment
341 entrainment effects, unless the flow becomes super-concentrated, which would
342 disagree with the strong geostrophic nature of the flow.
- 343 2) AMW settling depth therefore depends predominantly on the vertical density
344 gradient in the eastern North Atlantic.
- 345 3) The eastern North Atlantic vertical density gradient is found to be reduced by more
346 than 50% during the Last Glacial Maximum, compared to the present, which agrees
347 well with previous reconstructions of AMOC intensity changes.
- 348 4) Little difference is found between the LGM and Heinrich Event 1. This implies
349 there is a limitation to the speed of response of this parameter, which does not seem to
350 alter on timescales lower than millennial-scale ocean overturning rate.
- 351 5) We have elaborated our plume-control approach in a case specific to the
352 Mediterranean outflow, but it is in a general sense transferrable to other marginal
353 seas. Hence, the general approach offers exciting opportunities for estimating oceanic
354 vertical density gradients in many other sites where strait exchange can be modelled
355 and outflow plume height changes through time can be reconstructed from shallow
356 seismics and borehole studies.

357

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479 **Table 1. Monte Carlo variables used in modified PO94 model for Figures 1 and**
 480 **2.**

Parameter	Modern value	Monte Carlo parameterisation
Sea Level	0m	Monotonic random value between 0 and -130.
X_{Med}	0.05	Monotonic random value between 0.025 and 0.1
Winter Sea Surface Temperature (wSST) in Atlantic inflow.	16°C	Monotonic random value between 7 and 20°C.
Temperature difference between MIW and wSST.	~4°C	Monotonic random value between 0 and 6°C.
Temperature difference between MDW and wSST.	~4.5°C	Monotonic random value between 0 and 6°C (must exceed difference for MIW so that MDW is the coldest Mediterranean watermass)
Proportional admixture of MDW in MO	~0.3	Monotonic random value between 0 and 1.
Density gradient in mid-latitude eastern North Atlantic	0.0009 kg m ⁻⁴	Monotonic random value between 0.00009 and 0.0018.

481

Table 2. Boundary conditions for scenario simulations (Figure 4).

Parameter	Upper limit	Lower limit
Holocene		
AMW settling depth (m)	500	1100
Sea Level (m)	0	-10
Surface Temperature (°C)	18	16
Western Med. Cooling (°C)	3	0
Eastern Med. Cooling (°C)	3	0
X_{Med} (Sv)	0.05	0.04
Bølling-Allerød		
AMW settling depth (m) (Schönfeld and Zahn 2000)	800	1000
Sea Level (m)	-70	-90
Surface Temperature (°C)	16	12
Western Med. Cooling (°C)	3	0
Eastern Med. Cooling (°C)	3	0
X_{Med} (Sv)	0.05	0.04
Younger Dryas		
AMW settling depth (m) (Schönfeld and Zahn 2000)	1300	1500
Sea Level (m)	-50	-70
Surface Temperature (°C)	14	10
Western Med. Cooling (°C)	6	0
Eastern Med. Cooling (°C)	6	0
X_{Med} (Sv)	0.05	0.04
LGM		
AMW settling depth (m) (Schönfeld and Zahn 2000)	1400	2000
Sea Level (m)	-100	-130
Surface Temperature (°C)	14	10
Western Med. Cooling (°C)	6	0
Eastern Med. Cooling (°C)	6	0
X_{Med} (Sv)	0.05	0.04
Heinrich Event 1		
AMW settling depth(m) (Skinner and Elderfield 2007)	2000	2600
Sea Level (m)	-100	-130
Surface Temperature (°C)	9	6
Western Med. Cooling (°C)	6	0

Eastern Med. Cooling (°C)	6	0
\bar{X}_{Med} (Sv)	0.1	0.025

483

484

485

486 **Table 3. List of parameters used in this study.**

487

Symbol	Parameter	Value (if constant)	Units
Φ	Mixing coefficient		
B_{geo}	Geostrophic Buoyancy Flux		$m^3 s^{-3}$
U_{geo}	Geostrophic Velocity		$m s^{-1}$
g'	Reduced gravity		
π	Bottom gradient		$^{\circ}$
f	Coriolis parameter	0.000084	
H_{src}	Height of MOW plume at source		m
K_{geo}	Geostrophic Ekman number		
x	Distance from source of entrainment "event"	100,000	m
W_{src}	Width of MOW plume at source		m
g	Acceleration due to gravity	9.81	$m s^{-2}$
ρ_{MO}	Density of Mediterranean Water		$kg m^{-3}$
ρ_{ATL}	Density of inflowing Atlantic Water		$kg m^{-3}$
$D_{setling}$	Mean settling depth of AMW		m
$\Delta\rho_{MO}$	Density difference of Mediterranean and Atlantic water		$kg m^{-3}$
D_s	Depth of water at the Camarinal Sill.		m
H'	Global Sea Level change		m
$\partial\rho/\partial z$	Atlantic vertical density gradient		$kg m^{-4}$
Q_{MO}	Flux of MOW		Sv
ΔS_{gib}	Salinity difference between Atlantic and Mediterranean Water		PSU
S_{atl}	Salinity of inflowing Atlantic water		PSU
X_{med}	Mediterranean net freshwater export flux		Sv
C	Geometric coefficient for Strait of Gibraltar	0.283	
Q_{total}	Total, two-layer export at Gibraltar		Sv
β	Coefficient of saline contraction	0.00077	$kg m^{-3} PSU^{-1}$
α	Coefficient of thermal expansion	0.0002	$kg m^{-3} ^{\circ}C^{-1}$
ΔT_{gib}	Temperature difference between Atlantic and Mediterranean Water		$^{\circ}C$
q_s	Sediment flux		$kg s^{-1}$
$x1$	First entrainment coefficient		
$x2$	Second entrainment coefficient		
s_d	Sediment density	2.65	$g cm^{-3}$
d_{50}	median grainsize of sediment	2.4	μm
U^*	Shear velocity		$m s^{-1}$
U^*_{crit}	Critical Shear Velocity	2.4	$m s^{-1}$

488

489

490

491 **Figure Captions**

492

493 Figure 1: a; Vertical density structure of the North Atlantic in the region between 12
494 and 8°W and 33 and 38°N. Observations of past AMW activity (Schönfeld and Zahn
495 2000; Skinner and Elderfield 2007) are also shown. b-f show the estimated impact on
496 settling depth of the AMW plume from the Monte Carlo-like model derived from: b,
497 net freshwater flux from the Mediterranean (X_{Med}); c, cooling effects in the Gulf of
498 Lions; d, cooling effects in the eastern Mediterranean; e, sea level change. 1f, shows
499 the range of shear velocities (U^*) for the range of settling depths produced by the
500 model. Shear velocities exceeding 2.4 m s^{-1} (shown by the vertical line) are capable of
501 entraining sediment and those exceeding 6.8 m s^{-1} imply hyperconcentration of
502 sediment. For these simulations, sediment concentration is capped at 5% by volume.
503

504 Figure 2 a-c; Modelled increase in AMW settling depth due to sediment entrainment
505 (see Fig. 1f): a, relation to sea level change; b, relation to bottom water temperatures;
506 c, relation to reduced density of the AMW plume (g'). Fig. 2d-f, output of the Monte
507 Carlo-like model when the Atlantic vertical density is allowed to vary randomly
508 between 0.5 and 2 times the modern value: d, control of settling depth by net
509 freshwater flux from the Mediterranean (X_{Med}); e, control from sea level change; f,
510 control from Atlantic vertical density gradient ($\partial\rho / \partial z$).

511

512 Figure 3; Hydrographic and GCM (Levine and Bigg 2008) output showing vertical
513 density gradient in the eastern North Atlantic in a 10° box with its northeast corner at
514 Cape St Vincent. Grey points are hydrographic data (MEDATLAS 2002). Red circles
515 are GCM output for present day, blue circles are GCM output for LGM and black
516 stars are GCM output during a Hudson Strait Heinrich iceberg flux experiment. GCM
517 data have been adjusted so that density at 300m is compatible with hydrographic data.
518

519 Figure 4; a: Probabilistic assessment of Atlantic vertical density gradient during the
520 late Holocene (unfilled, broken lines), Bölling-Allerod (grey fill), Younger Dryas (red
521 fill), LGM (black fill), LGM with sediment entrainment impact (unfilled, green line)
522 and Heinrich Event 1 (blue fill). Boundary conditions for simulations presented in
523 Supporting Table 2. Fig 4b: Black data is simulation hydrographic data for LGM

524 scenario, grey data show simulations incorporating the impact of sediment
525 entrainment.
526

Figure 1

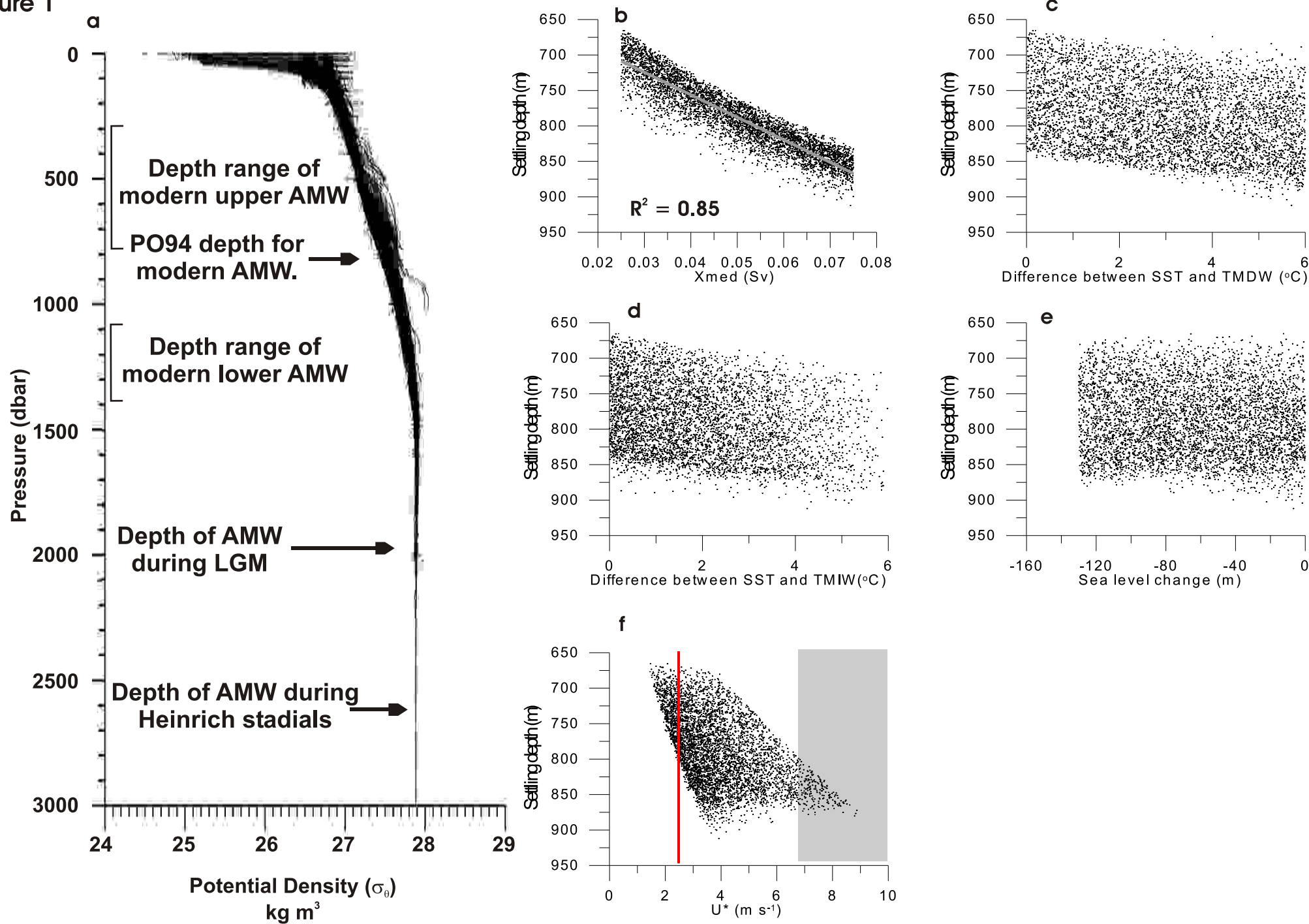


Figure 2

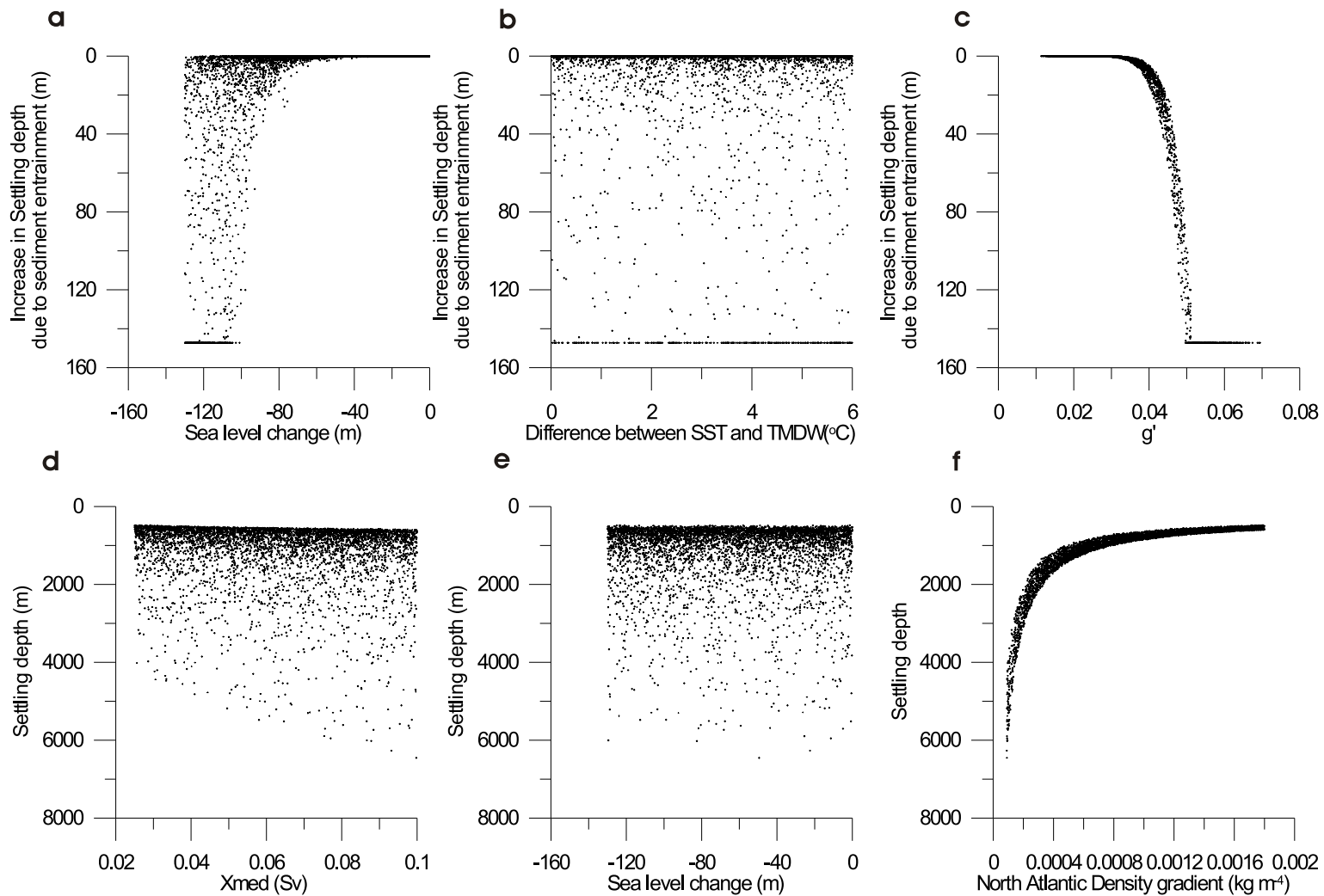


Figure 3

