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# Absolute Motion and Gravitational Effects

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The new information-theoretic Process Physics provides an explanation of space as a quantum foam system in which gravity is an inhomogeneous flow of the quantum foam into matter. Here an analysis of data from seven experiments demonstrates that absolute motion relative to space has been observed by Michelson and Morley (1887), Miller (1925/26), Illingworth (1927), Joos (1930), Jaseja et al (1963), Torr and Kolen (1981), and by DeWitte (1991). The Dayton Miller data also reveals the in-flow of space into the sun. The data reveals a new form of gravitational waves, predicted by the new theory of gravity in the accompanying paper ‘Gravity as Quantum Foam In-Flow’.

*Keywords:* Process Physics, quantum foam, spatial in-flow, absolute motion, gravitational waves, galactic in-flow.

# Contents

<b>1</b>	<b>Introduction</b>	<b>54</b>
<b>2</b>	<b>Detection of Absolute Motion</b>	<b>55</b>
2.1	Space and Absolute Motion . . . . .	55
2.2	Theory of the Michelson Interferometer . . . . .	56
2.3	The Michelson-Morley Experiment: 1887 . . . . .	63
2.4	The Miller Interferometer Experiment: 1925-1926 . . . . .	70
2.5	Gravitational In-flow from the Miller Data . . . . .	75
2.6	The Illingworth Experiment: 1927 . . . . .	81
2.7	The Joos Experiment: 1930 . . . . .	82
2.8	The New Bedford Experiment: 1963 . . . . .	83
2.9	The DeWitte Experiment: 1991 . . . . .	88
2.10	The Torr-Kolen Experiment: 1981 . . . . .	95
2.11	Galactic In-flow and the CMB Frame . . . . .	97
2.12	In-Flow Turbulence and Gravitational Waves . . . . .	99
2.13	Vacuum Michelson Interferometers . . . . .	103
2.14	Solid-State Michelson Interferometers . . . . .	104
<b>3</b>	<b>Conclusions</b>	<b>106</b>
<b>4</b>	<b>Acknowledgements</b>	<b>107</b>
<b>5</b>	<b>References</b>	<b>107</b>

# 1 Introduction

The new information-theoretic Process Physics [1, 2, 3, 4, 5, 6, 7, 8, 9, 10] provides a new explanation of space as a quantum foam system in which gravity is an inhomogeneous flow of the quantum foam into matter. Here an analysis of data from seven experiments demonstrates that absolute motion relative to space has been observed by Michelson and Morley (1887), Miller (1925/6), Illingworth (1927), Joos (1930), Jaseja *et al* (1963), Torr and Kolen (1981), and by DeWitte (1991), contrary to common belief within physics that absolute motion has never been observed. The Dayton Miller data also reveals the in-flow of space into the sun which manifests also as the gravitational ‘attraction’ of the earth by the sun. The data also reveals the in-flow into the Milky Way and local galactic cluster. The experimental data suggests that the in-flow manifests turbulence, which amounts to the observation of a gravitational wave phenomena, predicted by the new theory of gravity in [2].

As explained in [2] absolute motion is consistent with special relativistic effects, which are caused by actual dynamical effects of absolute motion through the quantum foam. The Lorentzian interpretation of special relativistic effects is seen to be essentially correct.

The detection of absolute motion and the related gravitational in-flow implies that space has structure, and this structure is the first experimental evidence of quantum gravity, though the data does yet reveal the scale of that structure.

## 2 Detection of Absolute Motion

### 2.1 Space and Absolute Motion

Absolute motion is motion relative to space itself. It turns out that Michelson and Morley in their historic experiment of 1887 did detect absolute motion, but rejected their own findings because using their method of analysis of the observed fringe shifts the determined speed of some 8 km/s was less than the 30 km/s orbital speed of the earth. The data was clearly indicating that the theory for the operation of the Michelson interferometer was not adequate. Rather than reaching this conclusion Michelson and Morley came to the incorrect conclusion that their results amounted to the failure to detect absolute motion. This had an enormous impact on the development of physics, for as is well known Einstein adopted the absence of absolute motion effects as one of his fundamental assumptions. By the time Miller had finally figured out how to work around the lack of a viable theory for the operation of the Michelson interferometer, and properly analyse data from his own Michelson interferometer, absolute motion had become a forbidden concept within physics, as it still is at present. The experimental observations by Miller and others of absolute motion has continued to be scorned and rejected by the physics community. Fortunately as well as revealing absolute motion the experimental data also reveals evidence in support of a new theory of gravity.

## 2.2 Theory of the Michelson Interferometer

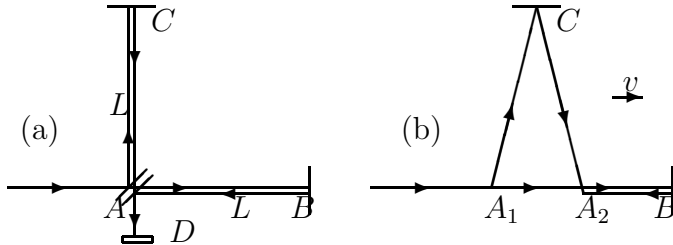


Figure 1: Schematic diagrams of the Michelson Interferometer, with beamsplitter/mirror at  $A$  and mirrors at  $B$  and  $C$ , on equal length arms when parallel, from  $A$ .  $D$  is a quantum detector (not drawn in (b)) that causes localisation of the photon state by a collapse process. In (a) the interferometer is at rest in space. In (b) the interferometer is moving with speed  $v$  relative to space in the direction indicated. Interference fringes are observed at the quantum detector  $D$ . If the interferometer is rotated in the plane through  $90^\circ$ , the roles of arms  $AC$  and  $AB$  are interchanged, and during the rotation shifts of the fringes are seen in the case of absolute motion, but only if the apparatus operates in a gas. By counting fringe changes the speed  $v$  may be determined.

We now show for the first time in over 100 years how three key effects together permit the Michelson interferometer [11] to reveal the phenomenon of absolute motion when operating in the presence of a gas, with the third effect only discovered in 2002 [7]. The main outcome is the derivation of the origin of the Miller  $k^2$  factor in the expression for the time difference for

light travelling via the orthogonal arms,

$$\Delta t = k^2 \frac{L|\mathbf{v}_P|^2}{c^3} \cos(2(\theta - \psi)). \quad (1)$$

Here  $\mathbf{v}_P$  is the projection of the absolute velocity  $\mathbf{v}$  of the interferometer through the quantum-foam onto the plane of the interferometer, where the projected velocity vector  $\mathbf{v}_P$  has azimuth angle  $\psi$  relative to the local meridian, and  $\theta$  is the angle of one arm from that meridian. The  $k^2$  factor is  $k^2 = n(n^2 - 1)$  where  $n$  is the refractive index of the gas through which the light passes,  $L$  is the length of each arm and  $c$  is the speed of light relative to the quantum foam. This expression follows from three key effects: (i) the difference in geometrical length of the two paths when the interferometer is in absolute motion, as first realised by Michelson, (ii) the Fitzgerald-Lorentz contraction of the arms along the direction of motion, and (iii) that these two effects precisely cancel in vacuum, but leave a residual effect if operated in a gas, because the speed of light through the gas is reduced to  $V = c/n$ , ignoring here for simplicity any Fresnel-drag effects. This is one of the aspects of the quantum foam physics that distinguishes it from the Einstein formalism. The time difference  $\Delta t$  is revealed by the fringe shifts on rotating the interferometer. In Newtonian physics, that is with no Fitzgerald-Lorentz contraction,  $k^2 = n^3$ , while in Einsteinian physics  $k = 0$  reflecting the fundamental assumption that absolute motion is not measurable and indeed has no meaning. So the experimentally determined value of  $k$  is a key test of fundamental physics. For air  $n = 1.00029$ , and so for process physics  $k = 0.0241$  and  $k^2 = 0.00058$ , which is close to the Einsteinian value of  $k = 0$ , particularly in comparison to the Newtonian

value of  $k = 1.0$ . This small but non-zero  $k$  value explains why the Michelson interferometer experiments gave such small fringe shifts. Fortunately it is possible to check the  $n$  dependence of  $k$  as two experiments [14, 15] were done in helium gas, and this has an  $n^2 - 1$  value significantly different from that of air.

In deriving (2) in the new physics it is essential to note that space is a quantum-foam system which exhibits various subtle features<sup>1</sup>. In particular it exhibits real dynamical effects on clocks and rods. In this physics the speed of light is only  $c$  relative to the quantum-foam, but to observers moving with respect to this quantum-foam the speed appears to be still  $c$ , but only because their clocks and rods are affected by the quantum-foam. As shown in [2] such observers will find that records of observations of distant events will be described by the Einstein spacetime formalism, but only if they restrict measurements to those achieved by using clocks, rods and light pulses, that is using the Einstein measurement protocol. However if they use an absolute motion detector then such observers can correct for these effects.

It is simplest in the new physics to work in the quantum-foam frame of reference. If there is a gas present at rest in this frame, such as air, then the speed of light in this frame is  $V = c/n$ . If the interferometer and gas are moving with respect to the quantum foam, as in the case of an interferometer attached to the earth, then the speed of light relative to the quantum-foam is still  $V = c/n$  up to corrections due to drag effects. Hence this

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<sup>1</sup>In [9] it is shown that the quantum-foam in-flow theory of gravity [2] explains the borehole  $g$  anomaly and the rotation velocity curves of spiral galaxies, and analysis of the data reveals that the in-flow theory involves the fine structure constant  $\alpha = e^2/\hbar c$ .



new physics requires a different method of analysis from that of the Einstein physics. With these cautions we now describe the operation of a Michelson interferometer in this new physics, and show that it makes predictions different to that of the Einstein physics. Of course experimental evidence is the final arbiter in this conflict of theories.

As shown in Fig.2 the beamsplitter/mirror when at  $A$  sends a photon  $\psi(t)$  into a superposition  $\psi(t) = \psi_1(t) + \psi_2(t)$ , with each component travelling in different arms of the interferometer, until they are recombined in the quantum detector which results in a localisation process, and one spot in the detector is produced. Repeating with many photons reveals that the interference between  $\psi_1$  and  $\psi_2$  at the detector results in fringes. These fringes actually only appear if the mirrors are not quite orthogonal, otherwise the screen has a uniform intensity and this intensity changes as the interferometer is rotated, as shown in the analysis by Hicks [16]. To simplify the analysis here assume that the two arms are constructed to have the same lengths  $L$  when they are physically parallel to each other and perpendicular to  $v$ , so that the distance  $BB'$  is  $L \sin(\theta)$ . The Fitzgerald-Lorentz effect in the new physics is that the distance  $SB'$  is  $\gamma^{-1}L \cos(\theta)$  where  $\gamma = 1/\sqrt{1 - v^2/c^2}$ . The various other distances are  $AB = Vt_{AB}$ ,  $BC = Vt_{BC}$ ,  $AS = vt_{AB}$  and  $SC = vt_{BC}$ , where  $t_{AB}$  and  $t_{BC}$  are the travel times. Applying the Pythagoras theorem to triangle

$ABB'$  we obtain

$$t_{AB} = \frac{2v\gamma^{-1}L \cos(\theta)}{2(V^2 - v^2)} + \frac{\sqrt{4v^2\gamma^{-2}L^2 \cos^2(\theta) + 4L^2(1 - \frac{v^2}{c^2} \cos^2(\theta))(V^2 - v^2)}}{2(V^2 - v^2)}. \quad (2)$$

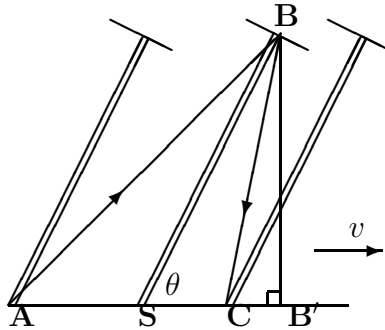


Figure 2: One arm of a Michelson Interferometer travelling at angle  $\theta$  and velocity  $\mathbf{v}$ , and shown at three successive times: (i) when photon leaves beamsplitter at  $A$ , (ii) when photon is reflected at mirror  $B$ , and (iii) when photon returns to beamsplitter at  $C$ . The line  $BB'$  defines right angle triangles  $ABB'$  and  $SBB'$ . The second arm is not shown but has angle  $\theta + 90^\circ$  to  $\mathbf{v}$ . Here  $\mathbf{v}$  is in the plane of the interferometer for simplicity, and the azimuth angle  $\psi = 0$ .

The expression for  $t_{BC}$  is the same except for a change of sign

of the  $2v\gamma^{-1}L \cos(\theta)$  term, then

$$\begin{aligned} t_{ABC} &= t_{AB} + t_{BC} \\ &= \frac{\sqrt{4v^2\gamma^{-2}L^2 \cos^2(\theta) + 4L^2(1 - \frac{v^2}{c^2} \cos^2(\theta))(V^2 - v^2)}}{(V^2 - v^2)}. \end{aligned} \quad (3)$$

The corresponding travel time  $t'_{ABC}$  for the orthogonal arm is obtained from (3) by the substitution  $\cos(\theta) \rightarrow \cos(\theta + 90^\circ) = -\sin(\theta)$ . The difference in travel times between the two arms is then  $\Delta t = t_{ABC} - t'_{ABC}$ . Now trivially  $\Delta t = 0$  if  $v = 0$ , but also  $\Delta t = 0$  when  $v \neq 0$  but only if  $V = c$ . This then would result in a null result on rotating the apparatus. Hence the null result of Michelson interferometer experiments in the new physics is only for the special case of photons travelling in vacuum for which  $V = c$ . However if the interferometer is immersed in a gas then  $V < c$  and a non-null effect is expected on rotating the apparatus, since now  $\Delta t \neq 0$ . It is essential then in analysing data to correct for this refractive index effect. The above  $\Delta t$  is the change in travel time when one arm is moved through angle  $\theta$ . The interferometer operates by comparing the change in the difference of the travel times between the arms. Then for  $V = c/n$  we find for  $v \ll V$  that

$$\Delta t = Ln(n^2 - 1) \frac{v^2}{c^3} \cos(2\theta) + O(v^4), \quad (4)$$

that is  $k^2 = n(n^2 - 1)$ , which gives  $k = 0$  for vacuum experiments ( $n = 1$ ). So the Miller phenomenological parameter  $k$  is seen to accommodate both the Fitzgerald-Lorentz contraction effect

and the dielectric effect, at least for gases. This is very fortunate since being a multiplicative parameter a re-scaling of old analyses is all that is required.  $\Delta t$  is non-zero when  $n \neq 1$  because the refractive index effect results in incomplete cancellation of the geometrical effect and the Fitzgerald-Lorentz contraction effect. Of course it was this cancellation effect that Fitzgerald and Lorentz actually used to arrive at the length contraction hypothesis, but they failed to take the next step and note that the cancellation would be incomplete in the air operated Michelson-Morley experiment. In a bizarre development modern Michelson interferometer experiments, which use resonant cavities rather than interference effects, but for which the analysis here is easily adapted, and with the same consequences, are operated in vacuum mode. That denies these experiments the opportunity to see absolute motion effects. Nevertheless the experimentalists continue to misinterpret their null results as evidence against absolute motion. Of course these experiments are therefore restricted to merely checking the Fitzgerald-Lorentz contraction effect, and this is itself of some interest.

All data from gas-mode interferometer experiments, except for that of Miller, has been incorrectly analysed using only the first effect as in Michelson's initial theoretical treatment, and so the consequences of the other two effects have been absent. Repeating the above analysis without these two effects we arrive at the Newtonian-physics time difference which, for  $v \ll V$ , is

$$\Delta t = Ln^3 \frac{v^2}{c^3} \cos(2\theta) + O(v^4), \quad (5)$$

that is  $k^2 = n^3$ . The value of  $\Delta t$ , which is typically of order  $10^{-17}s$  in gas-mode interferometers corresponding to a fractional

fringe shift, is deduced from analysing the fringe shifts, and then the speed  $v_M$  has been extracted using (5), instead of the correct form (4) or more generally (2). However it is very easy to correct for this oversight. From (4) and (5) we obtain for the corrected absolute (projected) speed  $v_P$  through space, and for  $n \approx 1^+$ ,

$$v_P = \frac{v_M}{\sqrt{n^2 - 1}}. \quad (6)$$

For air the correction factor in (6) is significant, and even more so for helium.

## 2.3 The Michelson-Morley Experiment: 1887

Michelson and Morley reported that their interferometer experiment in 1887 gave a ‘null-result’ which since then, with rare exceptions, has been claimed to support the Einstein assumption that absolute motion has no meaning. However to the contrary the Michelson-Morley published data [12] shows non-null effects, but much smaller than they expected. They made observations of thirty-six  $360^\circ$  turns using an  $L = 11$  meter length interferometer operating in air in Cleveland (Latitude  $41^\circ 30' N$ ) with six turns near 12:00 hrs (7:00 hrs ST) on each day of July 8, 9 and 11, 1887 and similarly near 18:00 hrs (13:00 hrs ST) on July 8, 9 and 12, 1887. Each turn took approximately 6 minutes as the interferometer slowly rotated floating on a tank of mercury. They published and analysed the average of each of the 6 data sets. The fringe shifts were extremely small but within their observational capabilities.

Table 2 shows examples of the averaged fringe shift micrometer readings every  $22.5^\circ$  of rotation of the Michelson-Morley

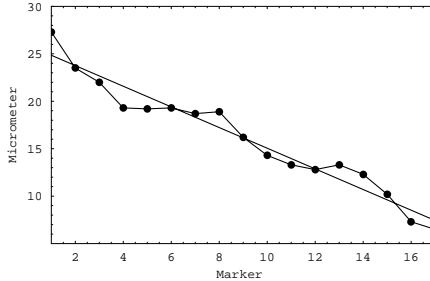


Figure 3: Plot of micrometer readings for July 11 12:00 hr (7:00 ST) showing the absolute motion induced fringe shifts superimposed on the uniform temperature induced fringe drift.

interferometer [12] for July 11 12:00 hr local time and also for July 9 18:00 hr local time. The orientation of the stone slab base is indicated by the marks 16, 1, 2, ... North is mark 16. The dominant effect was a uniform fringe drift caused by temporal temperature effects on the length of the arms, and imposed upon that are the fringe shifts corresponding to the effects of absolute motion, as shown in Fig.3.

local time	16 8	1 9	2 10	3 11	4 12	5 13	6 14	7 15	16
12:00hr July 11	27.3 16.2	23.5 14.3	22.0 13.3	19.3 12.8	19.2 13.3	19.3 12.3	18.7 10.2	18.9 7.3	6.5
18:00hr July 9	26.0 33.0	26.0 35.8	28.2 36.5	29.2 37.3	31.5 38.8	32.0 41.0	31.3 42.7	31.7 43.7	44.0

Table 2. Examples of Michelson-Morley fringe-shift micrometer readings. The readings for July 11 12:00 hr are plotted in Fig.3.

This temperature effect can be removed by subtracting from the data in each case a best fit to the data of  $a + bk$ ,  $\{k = 0, 1, 2, \dots, 8\}$  for the first  $0^\circ$  to  $180^\circ$  part of each rotation data set. Then multiplying by 0.02 for the micrometer thread calibration gives the fringe-shift data points in Fig.5. This factor of 0.02 converts the micrometer readings to fringe shifts expressed as fractions of a wavelength. Similarly a linear fit has been made to the data from the  $180^\circ$  to  $360^\circ$  part of each rotation data set. Separating the full  $360^\circ$  rotation into two  $180^\circ$  parts reduces the effect of the temperature drift not being perfectly linear in time.

In the new physics there are four main velocities that contribute to the total velocity:

$$\mathbf{v} = \mathbf{v}_{cosmic} + \mathbf{v}_{tangent} - \mathbf{v}_{in} - \mathbf{v}_E. \quad (7)$$

Here  $\mathbf{v}_{cosmic}$  is the velocity of the solar system relative to some cosmologically defined galactic quantum-foam system (discussed later) while the other three are local effects: (i)  $\mathbf{v}_{tangent}$  is the tangential orbital velocity of the earth about the sun, (ii)  $\mathbf{v}_{in}$  is a quantum-gravity radial in-flow of the quantum foam past the earth towards the sun, and (iii) the corresponding quantum-foam in-flow into the earth is  $\mathbf{v}_E$  and makes no contribution to a horizontally operated interferometer, assuming the velocity superposition approximation, and also that the turbulence associated with that flow is not significant, discussed in [2]. The minus signs in (7) arise because, for example, the in-flow towards the sun requires the earth to have an outward directed velocity against that in-flow in order to maintain a fixed distance from the sun, as shown in Fig.4. For circular orbits and using in-flow form of Newtonian gravity the speeds  $v_{tangent}$  and  $v_{in}$  are given

by

$$v_{\text{tangent}} = \sqrt{\frac{GM}{R}}, \quad (8)$$

$$v_{\text{in}} = \sqrt{\frac{2GM}{R}}, \quad (9)$$

while the net speed  $v_R$  of the earth from the vector sum  $\mathbf{v}_R = \mathbf{v}_{\text{tangent}} - \mathbf{v}_{\text{in}}$  is

$$v_R = \sqrt{\frac{3GM}{R}}, \quad (10)$$

where  $M$  is the mass of the sun,  $R$  is the distance of the earth from the sun, and  $G$  is Newton's gravitational constant.  $G$  is essentially a measure of the rate at which matter effectively 'dissipates' the quantum-foam. The gravitational acceleration arises from inhomogeneities in the flow. These expressions give  $v_{\text{tangent}} = 30\text{km/s}$ ,  $v_{\text{in}} = 42.4\text{km/s}$  and  $v_R = 52\text{km/s}$ .

Fig.5 shows all the data for the 1887 Michelson-Morley experiment for the fringe shifts after removal of the temperature drift effect for each averaged 180 degree rotation. The dotted curves come from the best fit of  $\frac{0.4}{30^2} k_{\text{air}}^2 v_P^2 \cos(2(\theta - \psi))$  to the data. The coefficient  $0.4/30^2$  arises as the apparatus would give a 0.4 fringe shift, as a fraction of a wavelength, with  $k = 1$  if  $v_P = 30 \text{ km/s}$  [12]. Shown in each figure is the resulting value of  $v_P$ . In some cases the data does not have the expected  $\cos(2(\theta - \psi))$  form, and so the corresponding values for  $v_P$  are not meaningful. The remaining fits give  $v_P = 331 \pm 30 \text{ km/s}$  for



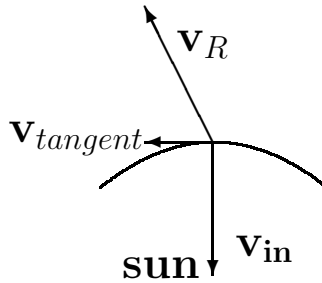


Figure 4: Orbit of earth about the sun defining the plane of the ecliptic with tangential orbital velocity  $\mathbf{v}_{tangent}$  and quantum-foam in-flow velocity  $\mathbf{v}_{in}$ . Then  $\mathbf{v}_R = \mathbf{v}_{tangent} - \mathbf{v}_{in}$  is the velocity of the earth relative to the quantum foam, after subtracting  $\mathbf{v}_{cosmic}$ .

the 7:00 hr (ST) data, and  $v_P = 328 \pm 50$  km/s for the 13:00 hr (ST) data. For comparison the full curves show the predicted form for the Michelson-Morley data, computed for the latitude of Cleveland, using the Miller direction (see later) for  $\mathbf{v}_{cosmic}$  of Right Ascension and Declination ( $\alpha = 4^{hr}54'$ ,  $\delta = -70^{\circ}30'$ ) and incorporating the tangential and in-flow velocity effects for July. The magnitude of the theoretical curves are in general in good agreement with the magnitudes of the experimental data, excluding those cases where the data does not have the sinusoidal form. However there are significant fluctuations in the azimuth angle. These fluctuations are also present in the Miller data, and together suggest that this is a real physical phenomenon, and not solely due to difficulties with the operation of the interferometer.

The Michelson-Morley interferometer data clearly shows the

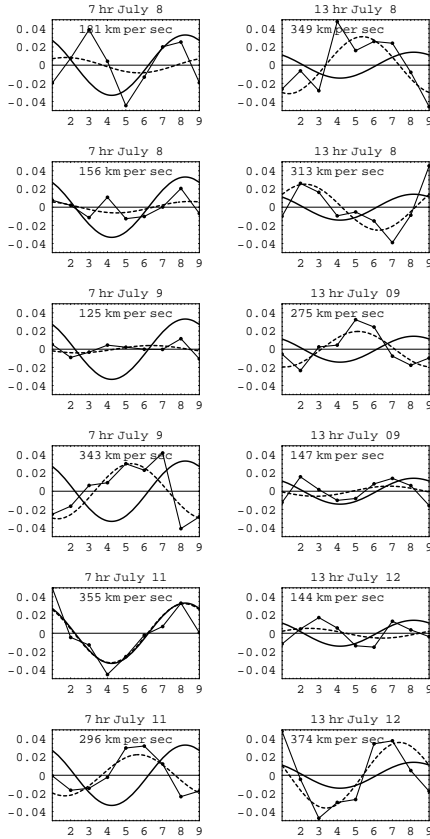


Figure 5: Shows all the Michelson-Morley 1887 data after removal of the temperature induced fringe drifts. The data for each  $360^{\circ}$  full turn (the average of 6 individual turns) is divided into the 1st and 2nd  $180^{\circ}$  parts and plotted one above the other. The dotted curve shows a best fit to the data, while the full curves show the expected forms using the Miller direction for  $\mathbf{v}_{cosmic}$ .

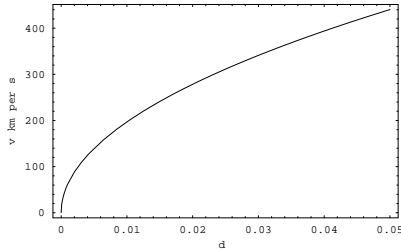


Figure 6: Speed calibration for Michelson-Morley experiment. This shows the value of  $v_P$  in km/s for values of the fringe shifts,  $d$ , expressed as a fraction of one wavelength of the light used, as shown in Fig.5

characteristic sinusoidal form with period  $180^\circ$  together with a large speed. Ignoring the effect of the refractive index, namely using the Newtonian value of  $k = 1$ , gives speeds reduced by the factor  $k_{air}$ , namely  $k_{air}v_P = 0.0241 \times 330\text{km/s} = 7.9 \text{ km/s}$ . Michelson and Morley reported speeds in the range  $5\text{km/s} - 7.5\text{km/s}$ . These slightly smaller speeds arise because they averaged all the 7:00 hr (ST) data, and separately all the 13:00 hr (ST) data, whereas here some of the lower quality data has not been used. Michelson was led to the false conclusion that because this speed of some  $8 \text{ km/s}$  was considerably less than the orbital speed of  $30 \text{ km/s}$  the interferometer must have failed to have detected absolute motion, and that the data was merely caused by experimental imperfections. This was the flawed analysis that led to the incorrect conclusion by Michelson and Morley that the experiment had failed to detect absolute motion. The consequences for physics were extremely damaging, and are only now being rectified after some 115 years.

## 2.4 The Miller Interferometer Experiment: 1925-1926

Dayton Miller developed and operated a Michelson interferometer for over twenty years, see Fig.7, with the main sequence of observations being on Mt.Wilson in the years 1925-1926, with the results reported in 1933 by Miller [13]. Accounts of the Miller experiments are available in Swenson [22]. Miller developed his huge interferometer over the years, from 1902 to 1906 in collaboration with Morley, and later at Mt.Wilson where the most extensive interferometer observations were carried out. Miller was meticulous in perfecting the operation of the interferometer and performed many control experiments. The biggest problem to be controlled was the effect of temperature changes on the lengths of the arms. It was essential that the temperature effects were kept as small as possible, but so long as each turn was performed sufficiently quickly, any temperature effect could be assumed to have been linear with respect to the angle of rotation. Then a uniform background fringe drift could be removed, as in the Michelson-Morley data analysis (see Fig.3).

In all some 200,000 readings were taken during some 12,000 turns of the interferometer<sup>2</sup>. Analysis of the data requires the extraction of the speed  $v_M$  and the azimuth angle  $\psi$  by effec-

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<sup>2</sup>In a remarkable development in 2002 as a result of a visit by James DeMeo to Case Western Reserve University the original Miller data was located, some 61 years after Miller's death in 1941. Until then it was thought that the data had been destroyed. Analysis of that data by the author of this article has confirmed the accuracy of Miller's analysis. Using more thorough computer based techniques the data is now being re-analysed.



Figure 7: Miller's interferometer with an effective arm length of  $L = 32\text{m}$  achieved by multiple reflections, as shown in Fig.8. Used by Miller on Mt.Wilson to perform the 1925-1926 observations of absolute motion. The steel arms weighed 1200 kilograms and floated in a tank of 275 kilograms of Mercury. Analysis of the extensive fringe-shift data from this momentous experiment now reveal a spatial flow past the earth into the sun, together with gravitational waves associated with that flow. From Case Western Reserve University Archives.

tively fitting the observed time differences, obtained from the observed fringe shifts, using (2), but with  $k = 1$ . Miller was of course unaware of the full theory of the interferometer and so he assumed the Newtonian theory, which neglected both the Fitzgerald-Lorentz contraction and air effects.

Miller performed this analysis of his data by hand, and the results for April, August and September 1925 and February 1926 are shown in Fig.9. The speeds shown are the Michelson speeds  $v_M$ , and these are easily corrected for the two neglected

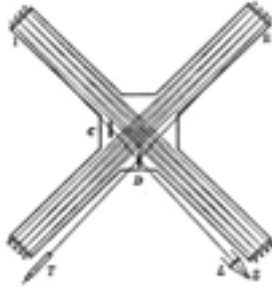


Figure 8: The effective arm length of  $L = 32\text{m}$  was achieved by multiple reflections.

effects by dividing these  $v_M$  by  $k_{air} = \sqrt{(n^2 - 1)} = 0.0241$ , as in (6). Then for example a speed of  $v_M = 10\text{km/s}$  gives  $v_P = v_M/k_{air} = 415\text{km/s}$ . However this correction procedure was not available to Miller. He understood that the theory of the Michelson interferometer was not complete, and so he introduced the phenomenological parameter  $k$  in (2). We shall denote his values by  $\bar{k}$ . Miller noted, in fact, that  $\bar{k}^2 \ll 1$ , as we would now expect. Miller then proceeded on the assumption that  $\mathbf{v}$  should have only two components: (i) a cosmic velocity of the solar system through space, and (ii) the orbital velocity of the earth about the sun. Over a year this vector sum would result in a changing  $\mathbf{v}$ , as was in fact observed, see Fig.9. Further, since the orbital speed was known, Miller was able to extract from the data the magnitude and direction of  $\mathbf{v}$  as the orbital speed offered an absolute scale. For example the dip in the  $v_M$  plots for sidereal times  $\tau \approx 16^{\text{hr}}$  is a clear indication of the direction of  $\mathbf{v}$ , as the dip arises at those sidereal times when the projection  $v_P$  of  $\mathbf{v}$  onto the plane of the interferometer is at

a minimum. During a 24hr period the value of  $v_P$  varies due to the earth's rotation. As well the  $v_M$  plots vary throughout the year because the vectorial sum of the earth's orbital velocity  $\mathbf{v}_{tangent}$  and the cosmic velocity  $\mathbf{v}_{cosmic}$  changes. There are two effects here as the direction of  $\mathbf{v}_{tangent}$  is determined by both the yearly progression of the earth in its orbit about the sun, and also because the plane of the ecliptic is inclined at  $23.5^\circ$  to the celestial plane. Figs.11 and 13 show the expected theoretical variation of both  $v_P$  and the azimuth  $\psi$  during one sidereal day in the months of April, August, September and February. These plots show the clear signature of absolute motion effects as seen in the actual interferometer data of Fig.9.

Note that the above corrected Miller projected absolute speed of approximately  $v_P = 415\text{km/s}$  is completely consistent with the corrected projected absolute speed of some  $330\text{km/s}$  from the Michelson-Morley experiment, though neither Michelson nor Miller were able to apply this correction. The difference in magnitude is completely explained by Cleveland having a higher latitude than Mt. Wilson, and also by the only two sidereal times of the Michelson-Morley observations. So from his 1925-1926 observations Miller had completely confirmed the true validity of the Michelson-Morley observations and was able to conclude, contrary to their published conclusions, that the 1887 experiment had in fact detected absolute motion. But it was too late. By then the physicists had incorrectly come to believe that absolute motion was inconsistent with various 'relativistic effects' that had by then been observed. This was because the Einstein formalism had been 'derived' from the assumption that absolute motion was without meaning and so unobservable in principle.

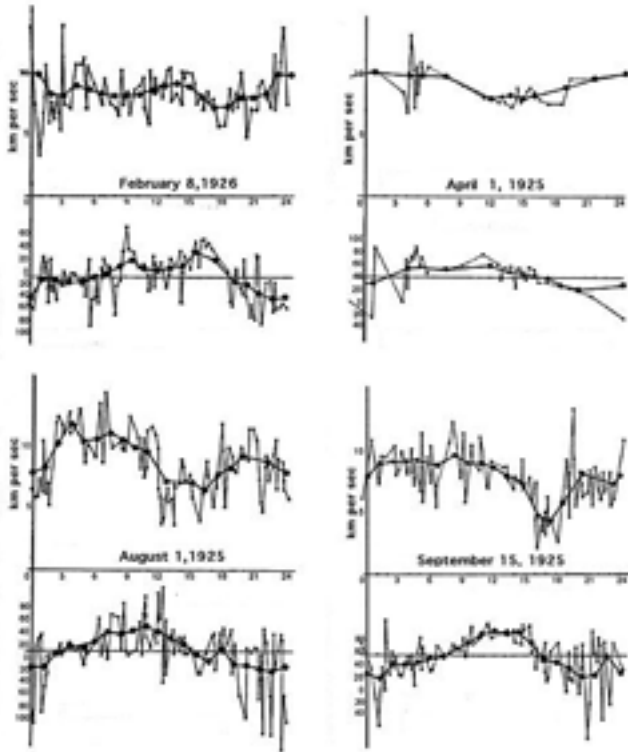


Figure 9: Miller's results from the 1925-1926 observations of absolute motion showing the projected 'Michelson' speed  $v_M$  in km/s and azimuth angle  $\psi$  in degrees plotted against sidereal time in hours. The smoother line is a running time average computed by Miller. The fluctuations in both  $v_M$  and  $\psi$  appear to be a combination of apparatus effects and genuine physical phenomena caused by turbulence in the gravitational in-flow of space towards the sun. Each data point arises from analysis of the average of twenty full rotations of the interferometer.



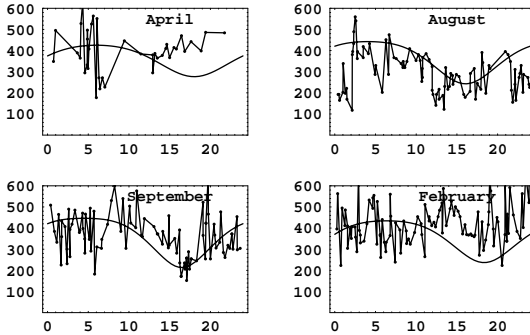


Figure 10: Miller interferometer projected speeds  $v_P$  in km/s showing both data and best fit of theory giving  $v_{cosmic} = 433$  km/s in the direction ( $\alpha = 5.2^{hr}$ ,  $\delta = -67^0$ ), and using  $n = 1.000226$  appropriate for the altitude of Mt. Wilson

## 2.5 Gravitational In-flow from the Miller Data

As already noted Miller was led to the conclusion that for reasons unknown the existing theory of the Michelson interferometer did not reveal true values of  $v_P$ , and for this reason he introduced the parameter  $k$ , with  $\bar{k}$  indicating his numerical values. Miller had reasoned that he could determine both  $\mathbf{v}_{cosmic}$  and  $\bar{k}$  by observing the interferometer determined  $v_P$  and  $\psi$  over a year because the known orbital velocity of the earth about the sun would modulate both of these observables, and by a scaling argument he could determine the absolute velocity of the solar system. In this manner he finally determined that  $|\mathbf{v}_{cosmic}| = 208$  km/s in the direction ( $\alpha = 4^{hr}54^m$ ,  $\delta = -70^033'$ ). However now that the theory of the Michelson interferometer has been revealed

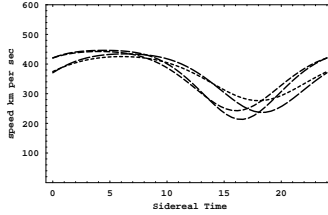


Figure 11: Expected theoretical variation of the projected velocity  $v_P$  during one sidereal day in the months of April, August, September and February, labelled by increasing dash length for cosmic speed of 433km/s in the direction ( $\alpha = 5.2^{hr}$ ,  $\delta = -67^0$ ). This shows the signature of the earth's orbital rotation.

an anomaly becomes apparent. Table 3 shows  $v = v_M/k_{air}$  for each of the four epochs, giving speeds consistent with the revised Michelson-Morley data. However Table 3 also shows that  $\bar{k}$  and the speeds  $\bar{v} = v_M/\bar{k}$  determined by the scaling argument are considerably different. Here the  $v_M$  values arise after taking account of the projection effect. That  $\bar{k}$  is considerably larger than the value of  $k_{air}$  indicates that another velocity component has been overlooked. Miller of course only knew of the tangential orbital speed of the earth, whereas the new physics predicts that as-well there is a quantum-gravity radial in-flow  $\mathbf{v}_{in}$  of the quantum foam. We can re-analyse Miller's data to extract a first approximation to the speed of this in-flow component. Clearly it is  $v_R = \sqrt{v_{in}^2 + v_{tangent}^2}$  that sets the scale and not  $v_{tangent}$ , and because  $\bar{k} = v_M/v_{tangent}$  and  $k_{air} = v_M/v_R$  are the scaling

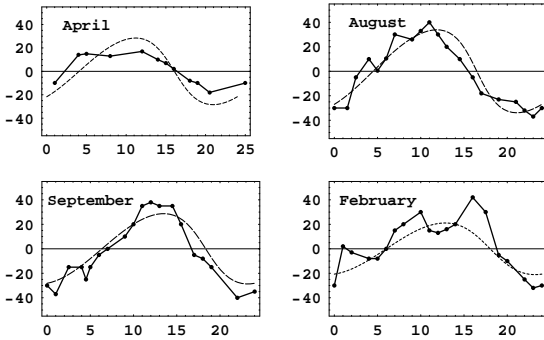


Figure 12: Miller azimuths  $\psi$ , measured from south, showing both data and best fit of theory giving  $v_{cosmic} = 433$  km/s in the direction ( $\alpha = 5.2^{hr}, \delta = -67^{\circ}$ ), and using  $n = 1.000226$  appropriate for the altitude of Mt. Wilson.

relations, then

$$\begin{aligned}
 v_{in} &= v_{tangent} \sqrt{\frac{v_R^2}{v_{tangent}^2} - 1}, \\
 &= v_{tangent} \sqrt{\frac{\bar{k}^{-2}}{k_{air}^2} - 1}.
 \end{aligned}
 \tag{11}$$

Using the  $\bar{k}$  values in Table 3 and the value<sup>3</sup> of  $k_{air}$  we obtain the  $v_{in}$  speeds shown in Table 3, which give an average speed of

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<sup>3</sup>We have not modified this value to take account of the altitude effect or temperatures atop Mt. Wilson. This weather information was not recorded by Miller. The temperature and pressure effect is that  $n = 1.0 + 0.00029 \frac{P}{P_0} \frac{T_0}{T}$ , where  $T$  is the temperature in  $^{\circ}\text{K}$  and  $P$  is the pressure in atmospheres.  $T_0 = 273\text{K}$  and  $P_0 = 1\text{atm}$ .

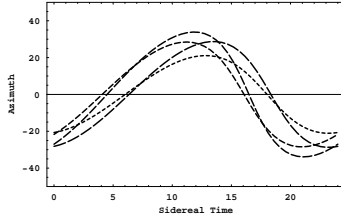


Figure 13: Expected theoretical variation of the azimuths  $\psi$ , measured from south, during one sidereal day in the months of April, August, September and February, labelled by increasing dash length, for a cosmic speed of 433km/s in the direction ( $\alpha = 5.2^{hr}$ ,  $\delta = -67^0$ ). This shows the signature of the earth's orbital rotation.

54 km/s, compared to the 'Newtonian' in-flow speed of 42 km/s. Note that the in-flow interpretation of the anomaly predicts that  $\bar{k} = (v_R/v_{tangent}) k_{air} = \sqrt{3} k_{air} = 0.042$ . Of course this simple re-scaling of the Miller results is not completely valid because (i) the direction of  $\mathbf{v}_R$  is of course different to that of  $\mathbf{v}_{tangent}$ , and also not necessarily orthogonal to  $\mathbf{v}_{tangent}$  because of turbulence, and (ii) also because of turbulence we would expect some contribution from the in-flow effect of the earth itself, namely that it is not always perpendicular to the earth's surface, and so would give a contribution to a horizontally operated interferometer.

An analysis that properly searches for the in-flow velocity effect clearly requires a complete re-analysis of the Miller data, and this is now possible and underway at Flinders University as the original data sheets have been found. It should be noted that the direction approximately diametrically opposite ( $\alpha = 4^{hr} 54^m$ ,  $\delta = -70^0 33'$ ), namely ( $\alpha = 17^{hr}$ ,  $\delta = +68'$ ) was at one

Epoch	$v_M$	$\bar{k}$	$v = v_M/k_{air}$	$\bar{v} = v_M/\bar{k}$	$v = \sqrt{3\bar{v}}$	$v_{in}$
February	9.3	0.048	385.9	193.8	335.7	51.7
April	10.1	0.051	419.1	198.0	342.9	56.0
August	11.2	0.053	464.7	211.3	366.0	58.8
September	9.6	0.046	398.3	208.7	361.5	48.8

Table 3. The  $\bar{k}$  anomaly,  $\bar{k} \gg k_{air} = 0.0241$ , as the gravitational in-flow effect. Here  $v_M$  and  $\bar{k}$  come from fitting the interferometer data, while  $v$  and  $\bar{v}$  are computed speeds using the indicated scaling. The average of the in-flow speeds is  $v_{in} = 54 \pm 5$  km/s, compared to the ‘Newtonian’ in-flow speed of 42 km/s. From column 4 we obtain the average  $v = 417 \pm 40$  km/s. All speeds in table in km/s.

stage considered by Miller as being possible. This is because the Michelson interferometer, being a 2nd-order device, has a directional ambiguity which can only be resolved by using the seasonal motion of the earth. However as Miller did not include the in-flow velocity effect in his analysis it is possible that a re-analysis might give this northerly direction as the direction of absolute motion of the solar system.

Hence not only did Miller observe absolute motion, as he claimed, but the quality and quantity of his data has also enabled the confirmation of the existence of the gravitational in-flow effect [2]. This is a manifestation of a new theory of gravity and one which relates to quantum gravitational effects via the unification of matter and space. As well the persistent evidence that this in-flow is turbulent indicates that this theory of gravity involves self-interaction of space itself.

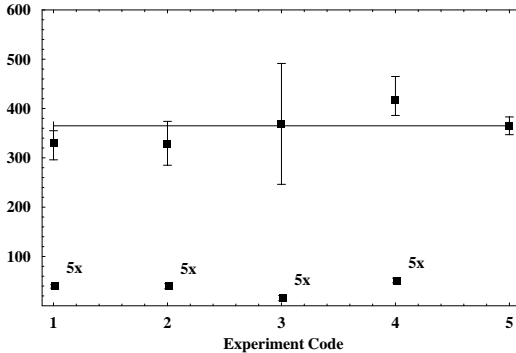


Figure 14: Speeds  $v$  in km/s determined from various Michelson interferometer experiments (1)-(4) and CMB (5): (1) Michelson-Morley (noon observations) and (2) (18<sup>h</sup> observations) see Sect.2.3, (3) Illingworth [14], (4) Miller, Mt. Wilson [13], and finally in (5) the speed from observations of the CMB spectrum dipole term [18]. The results (1)-(3) are not corrected for the  $\pm 30$  km/s of the orbital motion of the earth about the sun or for the gravitational in-flow speed, though these corrections were made for (4) with the speeds from Table 3. The horizontal line at  $v = 369$  km/s is to aid comparisons with the CMB frame speed data. The Miller direction is different to the CMB direction. Due to the angle between the velocity vector and the plane of interferometer the results (1)-(3) are less than or equal to the true speed, while the result for (4) is the true speed as this projection effect was included in the analysis. These results demonstrate the remarkable consistency between the three interferometer experiments. The Miller speed agrees with the speed from the DeWitte non-interferometer experiment, in Sect.2.9. The lower data, magnified by a factor of 5, are the original speeds  $v_M$  determined from fringe shifts using (1) with  $k = 1$ . This figure updates the corresponding figure in Ref.[7].

## 2.6 The Illingworth Experiment: 1927

In 1927 Illingworth [14] performed a Michelson interferometer experiment in which the light beams passed through the gas helium,

*...as it has such a low index of refraction that variations due to temperature changes are reduced to a negligible quantity.*

For helium at STP  $n = 1.000036$  and so  $k_{He}^2 = 0.00007$ , which results in an enormous reduction in sensitivity of the interferometer. Nevertheless this experiment gives an excellent opportunity to check the  $n$  dependence in (6). Illingworth, not surprisingly, reported no “ether drift to an accuracy of about one kilometer per second”. Múnera [17] re-analysed the Illingworth data to obtain a speed  $v_M = 3.13 \pm 1.04\text{km/s}$ . The correction factor in (6),  $1/\sqrt{n_{He}^2 - 1} = 118$ , is large for helium and gives  $v = 368 \pm 123\text{km/s}$ . As shown in Fig.14 the Illingworth observations now agree with those of Michelson-Morley and Miller, though they would certainly be inconsistent without the  $n$ -dependent correction, as shown in the lower data points (shown at  $5\times$  scale).

So the use by Illingworth of helium gas, and also by Joos, has turned out have offered a fortuitous opportunity to confirm the validity of the refractive index effect, though because of the insensitivity of this experiment the resulting error range is significantly larger than those of the other interferometer observations. So finally it is seen that the Illingworth experiment detected absolute motion with a speed consistent with all other observations.

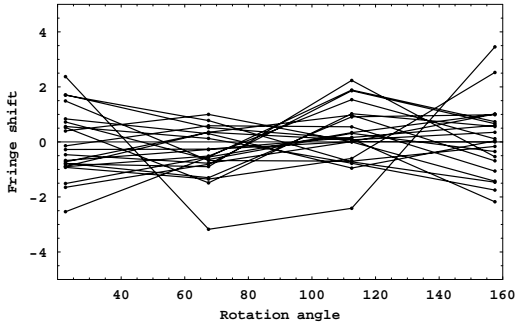


Figure 15: The Joos fringes shifts in  $\lambda/1000$  recorded on May 30, 1930 from a Michelson interferometer using helium. Only one of the rotations produced a clean signal of the form expected.

## 2.7 The Joos Experiment: 1930

Joos set out to construct and operate a large vacuum Michelson interferometer at the Zeiss Works in Jena, Germany 1930 [15]. This interferometer had an effective arm length of 21m achieved using multiple reflections in each arm. The vacuum sealing was ineffective and the penetration of air into the vacuum vessel caused problematic vibrations. Subsequently Joos used helium, assuming apparently that helium could be considered as a substitute for a true vacuum<sup>4</sup>. The use of helium is not mentioned in the Joos paper [15], but is mentioned by Swenson [22]. Joos recorded the fringe shifts photographically, and subsequently analysed the images using a photometer. The data for 22 rotations throughout the day of May 30, 1930 are shown in

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<sup>4</sup>Thanks to Dr Lance McCarthy for pointing out the use of helium in this experiment and in extracting the data from the Joos paper.



Fig.15, and are reproduced from Fig.11 of [15]. From that data Joos concluded, using an analysis that did not take account of the special relativistic length contraction effect, that the fringe shifts corresponded to a speed of only 1.5 km/s. However as previously noted such an analysis is completely flawed. As well the data in Fig.15 shows that for all but one of the rotations the fringe shifts were poorly recorded. Only in the one rotation, at 11 23<sup>58</sup>, does the data actually look like the form expected. This is probably not accidental as the maximum fringe shift was expected at that time, based on the Miller direction of absolute motion, and the sensitivity of the device was  $\pm 1$  thousandth of a fringe shift. In Fig.16 that one rotation data are compared with the form expected for Jena on May 30 using the Miller speed and direction together with the new refractive index effect, and using the refractive index of helium. The agreement is quite remarkable. So again, contrary the Joos paper and to subsequent commentators, Joos did in fact detect a very large velocity of absolute motion.

## 2.8 The New Bedford Experiment: 1963

In 1964 from an absolute motion detector experiment at New Bedford, latitude 42<sup>0</sup>N, Jaseja *et al* [19] reported yet another 'null result'. In this experiment two He-Ne masers were mounted with axes perpendicular on a rotating table, see Fig.17. Rotation of the table through 90<sup>0</sup> produced repeatable variations in the frequency difference of about 275kHz, an effect attributed to magnetorestriction in the Invar spacers due to the earth's magnetic field. Observations over some six consecutive hours on Jan-

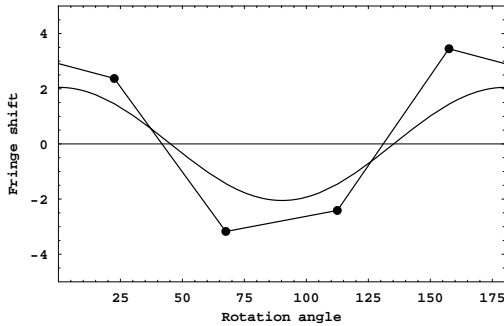


Figure 16: Comparison of the Joos data for the one good rotation at 11 23<sup>58</sup> with the theoretical prediction using the speed and direction from the Miller experiment, together with the length contraction and refractive index effects. The device sensitivity was  $\pm 1$ .

uary 20, 1963 from 6:00 am to 12:00 noon local time did produce a ‘dip’ in the frequency difference of some 3kHz superimposed on the 275kHz effect, as shown in Fig.18 in which the local times have been converted to sidereal times. The most noticeable feature is that the dip occurs at approximately 17 – 18:00<sup>hr</sup> sidereal time (or 9 – 10:00 hrs local time), which agrees with the direction of absolute motion observed by Miller and also by DeWitte (see Sect.2.9). It was most fortunate that this particular time period was chosen as at other times the effect is much smaller, as shown for example for the February data in Fig.9 which shows the minimum at 18:00<sup>hr</sup> sidereal time. The local times were chosen by Jaseja *et al* such that if the only motion was due to the earth’s orbital speed the maximum frequency difference, on rotation, should have occurred at 12:00hr local time, and the

minimum frequency difference at 6:00 hr local time, whereas in fact the minimum frequency difference occurred at 9:00 hr local time.

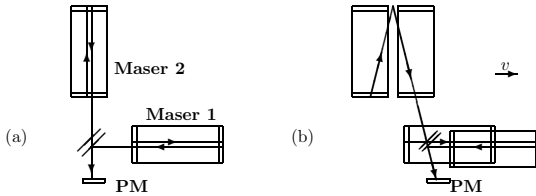


Figure 17: Schematic diagram for recording the variations in beat frequency between two optical masers: (a) when at absolute rest, (b) when in absolute motion at velocity  $\mathbf{v}$ . PM is the photomultiplier detector. The apparatus was rotated back and forth through  $90^\circ$ .

As for the Michelson-Morley experiment the analysis of the New Bedford experiment was also bungled. Again this apparatus can only detect the effects of absolute motion if the cancellation between the geometrical effects and Fitzgerald-Lorentz length contraction effects is incomplete. This occurs only when the radiation travels in a gas, here the He-Ne gas present in the maser.

This double maser apparatus is essentially equivalent to a Michelson interferometer. Then the resonant frequency  $\nu$  of each maser is proportional to the reciprocal of the out-and-back travel time. For maser 1

$$\nu_1 = m \frac{V^2 - v^2}{2LV \sqrt{1 - \frac{v^2}{c^2}}}, \quad (12)$$

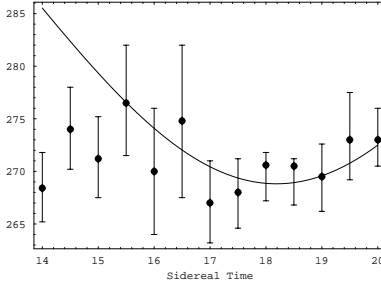


Figure 18: Frequency difference in kHz between the two masers in the 1963 New Bedford experiment after a  $90^0$  rotation. The 275kHz difference is a systematic repeatable apparatus effect, whereas the superimposed ‘dip’ at 17 – 18:00<sup>hr</sup> sidereal time of approximately 3kHz is a real time dependent frequency difference. The full curve shows the theoretical prediction for the time of the ‘dip’ for this experiment using the Miller direction for  $\hat{\mathbf{v}}$  ( $\alpha = 5.2^{hr}$ ,  $\delta = -67^0$ ) with  $|\mathbf{v}| = 433\text{km/s}$  and including the earth’s orbital velocity and sun gravitational in-flow velocity effects for January 20, 1963. The absolute scale of this theoretical prediction was not possible to compute as the refractive index of the He-Ne gas mixture was unknown.

for which a Fitzgerald-Lorentz contraction occurs, while for maser  
2

$$\nu_2 = m \frac{\sqrt{V^2 - v^2}}{2L}. \quad (13)$$

Here  $m$  refers to the mode number of the masers. When the apparatus is rotated the net observed frequency difference is  $\delta\nu = 2(\nu_2 - \nu_1)$ , where the factor of ‘2’ arises as the roles of the two masers are reversed after a  $90^0$  rotation. Putting  $V = c/n$  we find for  $v \ll V$  and with  $\nu_0$  the at-rest resonant frequency,

that

$$\delta\nu = (n^2 - 1)\nu_0 \frac{v^2}{c^2} + O\left(\frac{v^4}{c^4}\right). \quad (14)$$

If we use the Newtonian physics analysis, as in Jaseja *et al* [19], which neglects both the Fitzgerald-Lorentz contraction and the refractive index effect, then we obtain  $\delta\nu = \nu_0 v^2/c^2$ , that is without the  $n^2 - 1$  term, just as for the Newtonian analysis of the Michelson interferometer itself. Of course the very small magnitude of the absolute motion effect, which was approximately 1/1000 that expected assuming only an orbital speed of  $v = 30$  km/s in the Newtonian analysis, occurs simply because the refractive index of the He-Ne gas is very close to one<sup>5</sup>. Nevertheless given that it is small the sidereal time of the obvious 'dip' coincides almost exactly with that of the other observations of absolute motion.

The New Bedford experiment was yet another missed opportunity to have revealed the existence of absolute motion. Again the spurious argument was that because the Newtonian physics analysis gave the wrong prediction then Einstein relativity must be correct. But the analysis simply failed to take account of the Fitzgerald-Lorentz contraction, which had been known since the end of the 19<sup>th</sup> century, and the refractive index effect which had an even longer history. As well the authors failed to convert their local times to sidereal times and compare the time for the 'dip' with Miller's time<sup>6</sup>.

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<sup>5</sup>It is possible to compare the refractive index of the He-Ne gas mixture in the maser with the value extractable from this data:  $n^2 = 1 + 30^2/(1000 \times 400^2)$ , or  $n = 1.0000028$ .

<sup>6</sup>There is no reference to Miller's 1933 paper in Ref.[19].

## 2.9 The DeWitte Experiment: 1991

The Michelson-Morley, Illingworth, Miller, Joos and New Bedford experiments all used Michelson interferometers or its equivalent in gas mode, and all revealed absolute motion. The Michelson interferometer is a 2nd-order device meaning that the time difference between the 'arms' is proportional to  $(v/c)^2$ . There is also a factor of  $n^2 - 1$  and for gases like air and particularly helium or helium-neon mixes this results in very small time differences and so these experiments were always very difficult. Of course without the gas the Michelson interferometer is incapable of detecting absolute motion<sup>7</sup>, and so there are fundamental limitations to the use of this interferometer in the study of absolute motion and related effects.

In a remarkable development in 1991 a research project within Belgacom, the Belgium telecommunications company, stumbled across yet another detection of absolute motion, and one which turned out to be 1st-order in  $v/c$ . The study was undertaken by Roland DeWitte [21]. This organisation had two sets of atomic clocks in two buildings in Brussels separated by 1.5 km and the research project was an investigation of the task of synchronising these two clusters of atomic clocks. To that end 5MHz radiofrequency signals were sent in both directions through two buried coaxial cables linking the two clusters. The atomic clocks were caesium beam atomic clocks, and there were three in each cluster. In that way the stability of the clocks could be established and monitored. One cluster was in a building on Rue du

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<sup>7</sup>So why not use a transparent solid in place of the gas? See Sect.2.14 for the discussion.

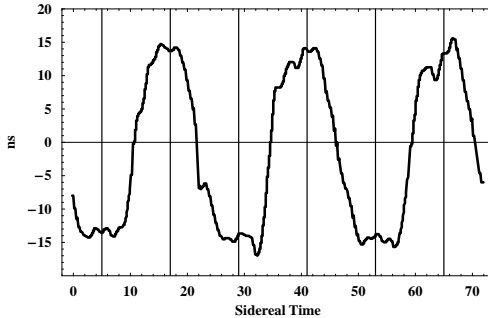


Figure 19: Variations in twice the one-way travel time, in ns, for an RF signal to travel 1.5 km through a coaxial cable between Rue du Marais and Rue de la Paille, Brussels. An offset has been used such that the average is zero. The definition of the sign convention for  $\Delta t$  used by DeWitte is unclear. The cable has a North-South orientation, and the data is  $\pm$  difference of the travel times for NS and SN propagation. The sidereal time for maximum effect of  $\sim 17$ hr (or  $\sim 5$ hr) (indicated by vertical lines) agrees with the direction found by Miller and also by Jaseja *et al*, but because of the ambiguity in the definition of  $\Delta t$  the opposite direction would also be consistent with this data. Plot shows data over 3 sidereal days and is plotted against sidereal time. See Fig.20b for theoretical predictions for one sidereal day. The time of the year of the data is not identified. The fluctuations are evidence of turbulence associated with the gravitational in-flow towards the sun. Adapted from DeWitte [21].

Marais and the second cluster was due south in a building on Rue de la Paille. Digital phase comparators were used to measure changes in times between clocks within the same cluster and also in the propagation times of the RF signals. Time differences between clocks within the same cluster showed a linear phase drift caused by the clocks not having exactly the same frequency together with short term and long term noise. However the long term drift was very linear and reproducible, and that drift could be allowed for in analysing time differences in the propagation times between the clusters.

Changes in propagation times were observed and eventually observations over 178 days were recorded. A sample of the data, plotted against sidereal time for just three days, is shown in Fig.19. DeWitte recognised that the data was evidence of absolute motion but he was unaware of the Miller experiment and did not realise that the Right Ascension for maximum/minimum propagation time agreed almost exactly with Miller's direction  $(\alpha, \delta) = (5.2^h, -67^0)$ . In fact DeWitte expected that the direction of absolute motion should have been in the CMB direction, but that would have given the data a totally different sidereal time signature, namely the times for maximum/minimum would have been shifted by 6 hrs. The declination of the velocity observed in this DeWitte experiment cannot be determined from the data as only three days of data are available. However assuming exactly the same declination as Miller the speed observed by DeWitte appears to be also in excellent agreement with the Miller speed, which in turn is in agreement with that from the Michelson-Morley and Illingworth experiments, as shown in Fig.14.

Being 1st-order in  $v/c$  the Belgacom experiment is easily



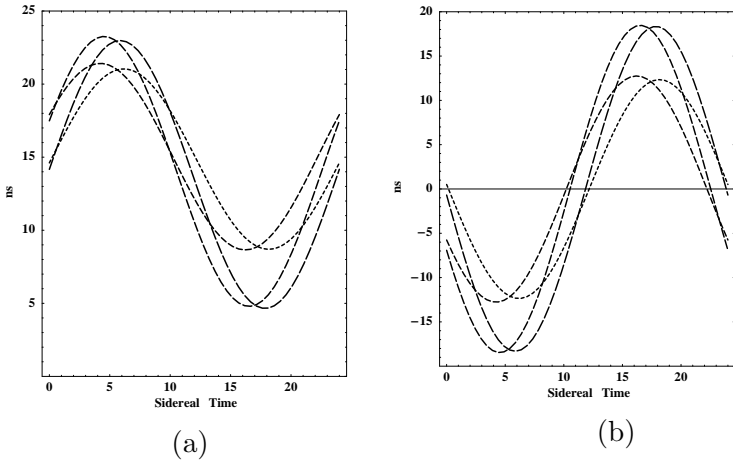


Figure 20: Theoretical predictions for the variations in travel time, in ns, for one sidereal day, in the DeWitte Brussels coaxial cable experiment for  $\mathbf{v}_{cosmic}$  in the direction  $(\alpha, \delta) = (5.2^h, -67^0)$  and with the Miller magnitude of 443 km/s, and including orbital and in-flow effects (but without turbulence). Shown are the results for four days: for the Vernal Equinox, March 21 (shortest dashes), and for 90, 180 and 270 days later (shown with increasing dash length). Figure (a) Shows change in one-way travel time  $t_0nv_P/c$  for signal travelling from N to S. Figure (b) shows  $\Delta t$ , as defined in (15), with an offset such that the average is zero so as to enable comparison with the data in Fig.19.  $\Delta t$  is twice the one-way travel time. For the direction opposite to  $(\alpha, \delta) = (5.2^h, -67^0)$  the same curves arise except that the identification of the months is different and the sign of  $\Delta t$  also changes. The sign of  $\Delta t$  determines which of the two directions is the actual direction of absolute motion. However the definition of the sign convention for  $\Delta t$  used by DeWitte is unclear.

analysed to sufficient accuracy by ignoring relativistic effects, which are 2nd-order in  $v/c$ . Let the projection of the absolute velocity vector  $\mathbf{v}$  onto the direction of the coaxial cable be  $v_P$  as before. Then the phase comparators reveal the difference between the propagation times in NS and SN directions. Consider the analysis with no Fresnel drag effect,

$$\begin{aligned}\Delta t &= \frac{L}{\frac{c}{n} - v_P} - \frac{L}{\frac{c}{n} + v_P}, \\ &= 2\frac{L}{c/n}n\frac{v_P}{c} + O\left(\frac{v_P^2}{c^2}\right) \approx 2t_0n\frac{v_P}{c}.\end{aligned}\quad (15)$$

Here  $L = 1.5$  km is the length of the coaxial cable,  $n = 1.5$  is the refractive index of the insulator within the coaxial cable, so that the speed of the RF signals is approximately  $c/n = 200,000$ km/s, and so  $t_0 = nL/c = 7.5 \times 10^{-6}$  sec is the one-way RF travel time when  $v_P = 0$ . Then, for example, a value of  $v_P = 400$ km/s would give  $\Delta t = 30$ ns. Because Brussels has a latitude of  $51^\circ$  N then for the Miller direction the projection effect is such that  $v_P$  almost varies from zero to a maximum value of  $|\mathbf{v}|$ . The DeWitte data in Fig.19 shows  $\Delta t$  plotted with a false zero, but shows a variation of some 28 ns. So the DeWitte data is in excellent agreement with the Miller's data<sup>8</sup>. The Miller experiment has thus been confirmed by a non-interferometer experiment.

The actual days of the data in Fig.19 are not revealed in Ref.[21] so a detailed analysis of the DeWitte data is not pos-

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<sup>8</sup>There is ambiguity in Ref.[21] as to whether the time variations in Fig.19 include the factor of 2 or not, as defined in (15). It is assumed here that a factor of 2 is included.

sible. Nevertheless theoretical predictions for various days in a year are shown in Fig.20 using the Miller speed of  $v_{cosmic} = 433$  km/s and where the diurnal effects of the earth's orbital velocity and the gravitational in-flow cause the range of variation of  $\Delta t$  and sidereal time of maximum effect to vary throughout the year. The predictions give  $\Delta t = 30 \pm 4$  ns over a year compared to the DeWitte value of 28 ns in Fig.19. If all of DeWitte's 178 days of data were available then a detailed analysis would be possible.

Ref.[21] does however reveal the sidereal time of the cross-over time, that is a 'zero' time in Fig.19, for all 178 days of data. This is plotted in Fig.21 and demonstrates that the time variations are correlated with sidereal time and not local solar time. A least squares best fit of a linear relation to that data gives that the cross-over time is retarded, on average, by 3.92 minutes per solar day. This is to be compared with the fact that a sidereal day is 3.93 minutes shorter than a solar day. So the effect is certainly cosmological and not associated with any daily thermal effects, which in any case would be very small as the cable is buried. Miller had also compared his data against sidereal time and established the same property, namely that up to small diurnal effects identifiable with the earth's orbital motion, features in the data tracked sidereal time and not solar time; see Ref.[13] for a detailed analysis.

The DeWitte data is also capable of resolving the question of the absolute direction of motion found by Miller. Is the direction  $(\alpha, \delta) = (5.2^h, -67^0)$  or the opposite direction? By doing a 2nd-order Michelson interferometer experiment Miller had to rely on the earth's diurnal effects in order to resolve this ambiguity, but his analysis of course did not take account of the gravitational

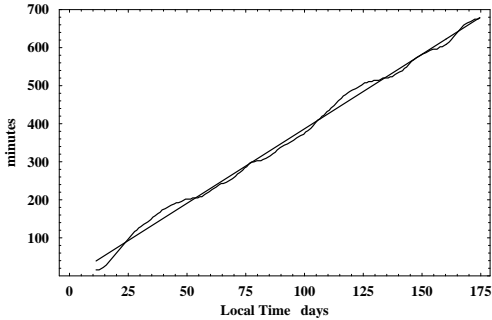


Figure 21: Plot of the negative of the drift of the cross-over time between minimum and maximum travel-time variation each day (at  $\sim 10^h \pm 1^h$  ST) versus local solar time for some 180 days. The straight line plot is the least squares fit to the experimental data, giving an average slope of 3.92 minutes/day. The time difference between a sidereal day and a solar day is 3.93 minutes/day. This demonstrates that the effect is related to sidereal time and not local solar time. The actual days of the year are not identified in Ref.[21]. Adapted from DeWitte [21].

in-flow effect, and so until a re-analysis of his data his preferred choice of direction must remain to be confirmed. The DeWitte experiment could easily resolve this ambiguity by simply noting the sign of  $\Delta t$ . Unfortunately it is unclear in Ref.[21] as to how the sign in Fig.19 is actually defined, and DeWitte does not report a direction expecting, as he did, that the direction should have been the same as the CMB direction.

The DeWitte observations were truly remarkable considering that initially they were serendipitous. They demonstrated yet again that the Einstein postulates were in contradiction with experiment. To my knowledge no physics journal has published

a report of the DeWitte experiment.

That the DeWitte experiment is not a gas-mode Michelson interferometer experiment is very significant. The value of the speed of absolute motion revealed by the DeWitte experiment of some 400 km/s is in agreement with the speeds revealed by the new analysis of various Michelson interferometer data, which use the recently discovered refractive index effect, see Fig.14. Not only was this effect confirmed by comparing results for different gases, but the re-scaling of the older  $v_M$  speeds to  $v = v_M/\sqrt{n^2 - 1}$  speeds resulting from this effect are now confirmed.

## 2.10 The Torr-Kolen Experiment: 1981

A coaxial cable experiment similar to but before the DeWitte experiment was performed at the University of Utah in 1981 by Torr and Kolen [20]. This involved two rubidium vapor clocks placed approximately 500m apart with a 5 MHz sinewave RF signal propagating between the clocks via a nitrogen filled coaxial cable maintained at a constant pressure of  $\sim 2$  psi. This means that the Fresnel drag effect is not important in this experiment. Unfortunately the cable was orientated in an East-West direction which is not a favourable orientation for observing absolute motion in the Miller direction, unlike the Brussels North-South cable orientation. There is no reference to Miller's result in the Torr and Kolen paper, otherwise they would presumably not have used this orientation. Nevertheless there is a projection of the absolute motion velocity onto the East-West cable and Torr and Kolen did observe an effect in that, while the round

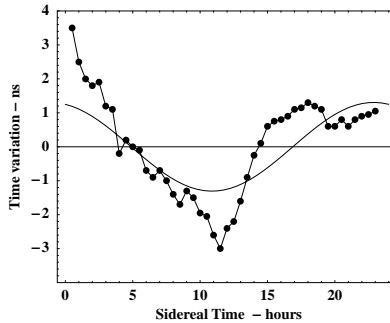


Figure 22: Data from the 1981 Torr-Kolen experiment at Logan, Utah [20]. The data shows variations in travel times (ns), for local times, of an RF signal travelling through 500m of coaxial cable orientated in an E-W direction. Actual days are not indicated but the experiment was done during February-June 1981. Results are for a typical day. For the 1st of February the local time of 12:00 corresponds to 13:00 sidereal time. The predictions are for February for a cosmic speed of 433 km/s in the direction  $(\alpha, \delta) = (5.2^h, -72^0)$ , and including orbital and in-flow velocities but without theoretical turbulence.

speed time remained constant within 0.0001%*c*, typical variations in the one-way travel time were observed, as shown in Fig.22 by the data points. The theoretical predictions for the Torr-Kolen experiment for a cosmic speed of 433 km/s in the direction  $(\alpha, \delta) = (5.2^h, -67^0)$ , and including orbital and in-flow velocities, are shown in Fig.22. As well the maximum effect occurred, typically, at the predicted times. So the results of this experiment are also in remarkable agreement with the Miller direction, and the speed of 433 km/s which of course only arises after re-scaling the Miller speeds for the effects of the gravitational in-flow. As well Torr and Kolen reported fluctuations in both the magnitude and time of the maximum variations in travel time just as DeWitte observed some 10 years later. Again we argue that these fluctuations are evidence of genuine turbulence in the in-flow as discussed in Sect.2.12. So the Torr-Kolen experiment again shows strong evidence for the new theory of gravity, and which is over and above its confirmation of the various observations of absolute motion.

## 2.11 Galactic In-flow and the CMB Frame

Absolute motion (AM) of the solar system has been observed in the direction  $(\alpha, \delta) = (5.2^h, -67^0)$ , up to an overall sign to be sorted out, with a speed of 433 km/s. This is the velocity after removing the contribution of the earth's orbital speed and the sun in-flow effect. It is significant that this velocity is different to that associated with the Cosmic Microwave Background (CMB) relative to which the solar system has a speed of 369 km/s in the direction  $(\alpha, \delta) = (11.20^h, -7.22^0)$ , see [18]. This CMB ve-

locity is obtained by finding the preferred frame in which this thermalised  $3^0\text{K}$  radiation is isotropic, that is by removing the dipole component. The CMB velocity is a measure of the motion of the solar system relative to the universe as a whole, or at least a shell of the universe some 15Gyrs away, and indeed the near uniformity of that radiation in all directions demonstrates that we may meaningfully refer to the spatial structure of the universe. The concept here is that at the time of decoupling of this radiation from matter that matter was on the whole, apart from small observable fluctuations, at rest with respect to the quantum-foam system that is space. So the CMB velocity is the motion of the solar system with respect to space *universally*, but not necessarily with respect to the *local* space. Contributions to this velocity would arise from the orbital motion of the solar system within the Milky Way galaxy, which has a speed of some 250 km/s, and contributions from the motion of the Milky Way within the local cluster, and so on to perhaps larger clusters.

On the other hand the AM velocity is a vector sum of this *universal* CMB velocity and the net velocity associated with the *local* gravitational in-flows into the Milky Way and the local cluster. If the CMB velocity had been identical to the AM velocity then the in-flow interpretation of gravity would have been proven wrong. We therefore have three pieces of experimental evidence for this interpretation (i) the refractive index anomaly discussed previously in connection with the Miller data, (ii) the turbulence seen in all detections of absolute motion, and now (iii) that the AM velocity is different in both magnitude and direction from that of the CMB velocity, and that this CMB velocity does not display the turbulence seen in the AM velocity.

That the AM and CMB velocities are different amounts to



the discovery of the resolution to the ‘dark matter’ conjecture. Rather than the galactic velocity anomalies being caused by such undiscovered ‘dark matter’ we see that the in-flow into non spherical galaxies, such as the spiral Milky Way, will be non Newtonian [2]. As well it will be interesting to determine, at least theoretically, the scale of turbulence expected in galactic systems, particularly as the magnitude of the turbulence seen in the AM velocity is somewhat larger than might be expected from the sun in-flow alone. Any theory for the turbulence effect will certainly be checkable within the solar system as the time scale of this is suitable for detailed observation.

It is also clear that the time of observers at rest with respect to the CMB frame is absolute or universal time. This interpretation of the CMB frame has of course always been rejected by supporters of the SR/GR formalism. As for space we note that it has a differential structure, in that different regions are in relative motion. This is caused by the gravitational in-flow effect locally, and as well by the growth of the universe.

## 2.12 In-Flow Turbulence and Gravitational Waves

The velocity flow-field equation, in [2], is expected to have solutions possessing turbulence, that is, fluctuations in both the magnitude and direction of the gravitational in-flow component of the velocity flow-field. Indeed all the Michelson interferometer experiments showed evidence of such turbulence. The first clear evidence was from the Miller experiment, as shown in Fig.9.

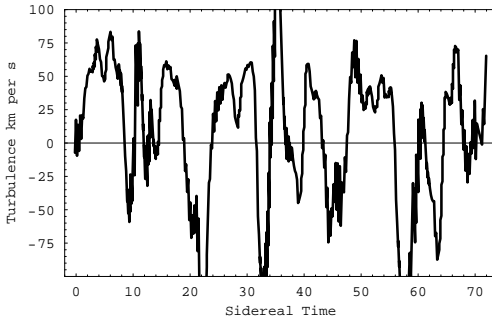


Figure 23: Speed fluctuations determined from Fig.19 by subtracting a least squares best fit of the forms shown in Fig.20b. A 1ns variation in travel time corresponds approximately to a speed variation of 27km/s. The larger speed fluctuations actually arise from a fluctuation in the cross-over time, that is, a fluctuation in the direction of the velocity. This plot implies that the velocity flow-field is turbulent. The scale of this turbulence is comparable to that evident in the Miller data, as shown in Fig.9 and Fig.24a.

Miller offered no explanation for these fluctuations but in his analysis of that data he did running time averages, as shown by the smoother curves in Fig.9. Miller may have in fact have simply interpreted these fluctuations as purely instrumental effects. While some of these fluctuations may be partially caused by weather related temperature and pressure variations, the bulk of the fluctuations appear to be larger than expected from that cause alone. Even the original Michelson-Morley data in Fig.5 shows variations in the velocity field and supports this interpretation. However it is significant that the non-interferometer DeWitte data also shows evidence of turbulence in both the magnitude and direction of the velocity flow field, as shown in

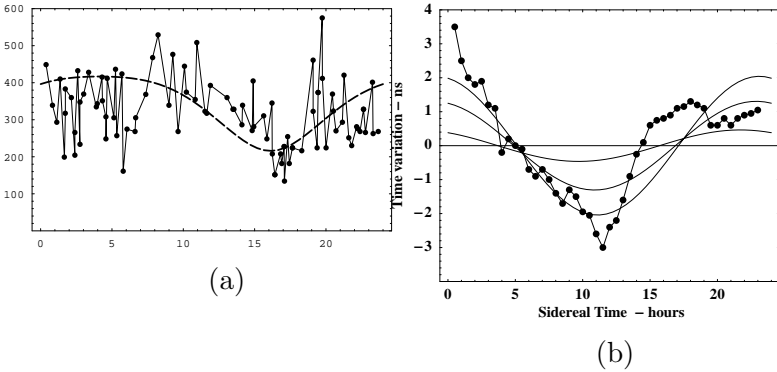


Figure 24: (a) The absolute projected speeds  $v_P$  in the Miller experiment plotted against sidereal time in hours for September 1925, showing the variations in speed caused by the gravitational wave turbulence. and (b) similar variations in travel times when the declination is varied by  $\pm 10^0$  about the direction  $\alpha = 5.2^h, \delta = -67^0$ , for a cosmic speed of 433 km/s in the Torr-Kolen experiment.

Fig.23. Just as the DeWitte data agrees with the Miller data for speeds and directions the magnitude fluctuations, shown in Fig.23, are very similar in absolute magnitude to, for example, the Miller speed turbulence shown in Fig.24a. As well the orientation of the Torr-Kolen coaxial cable is very sensitive to the directional changes associated with the turbulence. Being almost at  $90^0$  to the direction of absolute motion, any variation in that direction produces significant effects, as shown in Fig. 24b where the declination is varied by  $\pm 10^0$ . Indeed Torr and Kolen [20] reported significant fluctuations in the coaxial cable travel times from day to day, as expected.

It therefore becomes clear that there is strong evidence from these three experiments for these fluctuations being evidence of physical turbulence in the flow field. The magnitude of this turbulence appears to be somewhat larger than that which would be caused by the in-flow of quantum foam towards the sun, and indeed following on from Sect.2.11 some of this turbulence may be associated with galactic in-flow into the Milky Way. This in-flow turbulence is a form of gravitational wave and the ability of gas-mode Michelson interferometers to detect absolute motion means that experimental evidence of such a wave phenomena has been available for a considerable period of time, but suppressed along with the detection of absolute motion itself. Of course flow equations of the form in [2] do not exhibit those gravitational waves of the form that have been predicted to exist based on the Einstein equations, and which are supposed to propagate at the speed of light. All this means that gravitational wave phenomena is very easy to detect and amounts to new physics that can be studied in much detail.

## 2.13 Vacuum Michelson Interferometers

Over the years vacuum-mode Michelson interferometer experiments have become increasingly popular, although the motivation for such experiments appears to be increasingly unclear. The first vacuum interferometer experiment was planned by Joos [15] in 1930, but because of technical problems helium was actually used, as discussed in section 2.7. The first actual vacuum experiment was by Kennedy and Thorndike [23]. The result was actually unclear but was consistent with a null effect as predicted by both the quantum-foam physics and the Einstein physics. Only Newtonian physics is disproved by such experiments. These vacuum interferometer experiments do give null results, with increasing confidence level, as for example in Refs.[23, 24, 25, 26], but they only check that the Lorentz contraction effect completely cancels the geometrical path-length effect in vacuum experiments, and this is common to both theories. So they are unable to distinguish the new physics from the Einstein physics. Nevertheless recent works [25, 26] continue to claim that the experiment had been motivated by the desire to look for evidence of absolute motion, despite effects of absolute motion having been discovered as long ago as 1887. The ‘null results’ are always reported as proof of the Einstein formalism. Unfortunately the analysis of the data from such experiments is always by means of the Robertson [27] and Mansouri and Sexl formalism [28], which purports to be a generalisation of the Lorentz transformation if there is a preferred frame. However in [2] we have seen that absolute motion effects, that is the existence of a preferred frame, are consistent with the usual Lorentz transformation, based as it is on the restricted Einstein mea-

surement protocol. A preferred frame is revealed by gas-mode Michelson interferometer experiments, and then the refractive index of the gas plays a critical role in interpreting the data. The Robertson and Mansouri-Sexl formalism contains no contextual aspects such as a refractive index effect and is thus totally inappropriate to the analysis of so called ‘preferred frame’ experiments.

It is a curious feature of the history of Michelson interferometer experiments that it went unnoticed that the results fell into two distinct classes, namely vacuum and gas-mode, with recurring non-null results from gas-mode interferometers.

## 2.14 Solid-State Michelson Interferometers

The gas-mode Michelson interferometer has its sensitivity to absolute motion effects greatly reduced by the refractive index effect, namely the  $k^2 = n^2 - 1$  factor in (1), and for gases with  $n$  only slightly greater than one this factor has caused much confusion over the last 115 years. So it would be expected that passing the light beams through a transparent solid with  $n \approx 1.5$  rather than through a gas would greatly increase the sensitivity. Such an Michelson interferometer experiment was performed by Shamir and Fox [29] in Haifa in 1969. This interferometer used light from a He-Ne laser and used perspex rods with  $L = 0.26\text{m}$ . The experiment was interpreted in terms of the supposed Fresnel drag effect, which has a drag coefficient given by  $b = 1 - 1/n^2$ . The light passing through the solid was supposed to be ‘dragged’ along in the direction of motion of the solid with a velocity  $\Delta\mathbf{V} = b\mathbf{v}$  additional to the usual  $c/n$  speed. As well the Michel-

son geometrical path difference and the Lorentz contraction effects were incorporated into the analysis. The outcome was that no fringe shifts were seen on rotation of the interferometer, and Shamir and Fox concluded that this negative result “*enhances the experimental basis of special relativity*”.

The Shamir-Fox experiment was unknown to us<sup>9</sup> at Flinders university when in 2002 several meters of optical fibre were used in a Michelson interferometer experiment which also used a He-Ne laser light source. Again because of the  $n^2 - 1$  factor, and even ignoring the Fresnel drag effect, one would have expected large fringe shifts on rotation of the interferometer, but none were observed. As well in a repeat of the experiment single-mode optical fibres were also used and again with no rotation effect seen. So this experiment is consistent with the Shamir-Fox experiment. Re-doing the analysis by including the supposed Fresnel drag effect, as Shamir and Fox did, makes no material difference to the expected outcome. In combination with the non-null results from the gas-mode interferometer experiments along with the non-interferometer experiment of DeWitte it is clear that transparent solids behave differently to a gas when undergoing absolute motion through the quantum foam. Indeed this in itself is a discovery of a new phenomena.

The most likely explanation is that the physical Fitzgerald-Lorentz contraction effect has a anisotropic effect on the refractive index of the transparent solid, and this is such as to cause a cancellation of any differences in travel time between the two arms on rotation of the interferometer. In this sense a transpar-

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<sup>9</sup>This experiment was performed by Professor Warren Lawrance, an experimental physical chemist with considerable laser experience.

ent solid medium shares this outcome with the vacuum itself.

### 3 Conclusions

We have shown here that seven experiments, so far, have clearly revealed experimental evidence of absolute motion. As well these are all consistent with respect to the direction and speed of that motion. This clearly refutes the fundamental postulates of the Einstein reinterpretation of the relativistic effects that had been developed earlier by Lorentz and others. Indeed these experiments are consistent with the Lorentzian interpretation of the special relativistic effects in which reality displays both absolute motion and relativistic effects. It is absolute motion that actually causes these relativistic effects. Data from the five Michelson interferometer fringe-shift experiments had never been properly analysed until now. That analysis requires that the Fitzgerald-Lorentz contraction effect be taken into account, as well as the effect of the gas on the speed of light in the interferometer. Only then does the fringe-shift data from air and helium interferometer experiments become consistent, and then also consistent with the two RF coaxial cable travel-time experiments. The seasonal changes in the Miller fringe-shift data reveal the orbital motion of the earth about the sun, as well as an in-flow of space past the earth into the sun. These results support the new theory of gravity. As well the large cosmic velocity of the solar system is seen to be different to the velocity associated with the Cosmic Microwave Background, which implies another gravitational in-flow, this time into the Milky Way. The fringe-shift data has also indicated the presence of



turbulence in these gravitational in-flows, and this amounts to the detection of gravitational waves. These are waves predicted by the new theory of gravity, and not those associated with the Hilbert-Einstein theory of gravity. As noted in [2] the Newtonian theory of gravity is deeply flawed, as revealed by its inability to explain a growing number of gravitational anomalies, but which are explained by the new theory. In particular the borehole  $g$  anomaly and the rotation velocity curves of spiral galaxies, together with the absence of this effect in ordinary elliptical galaxies, have been explained. These flaws arose because the solar system was too special, because of its high spherical symmetry, to have revealed the full range of phenomena that is gravity. General Relativity ‘inherited’ these flaws, and so is itself flawed. As discussed in [2] the clear-cut checks of General Relativity were actually done in systems also with high spherical symmetry.

## 4 Acknowledgements

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