# Voltage recovery in intermediate band solar cells

Pablo G. Linares , Antonio Martí , Elisa Antolín , Corrie D. Farmer , Íñigo Ramiro , Colin R. Stanley , Antonio Luque

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## ABSTRACT

The intermediate band solar cell (IBSC) is based on a novel photovoltaic concept and has a limiting efficiency of 63.2%, which compares favorably with the 40.7% efficiency of a conventional, single junction solar cell. It is characterized by a material hosting a collection of energy levels within its bandgap, allowing the cell to exploit photons with sub-bandgap energies in a two-step absorption process, thus improving the utilization of the solar spectrum. However, these intermediate levels are often regarded as an inherent source of supplementary recombination, although this harmful effect can in theory be counteracted by the use of concentrated light. We present here a novel, low-temperature characterization technique using concentrated light that reveals how the initially enhanced recombination in the IBSC is reduced so that its open-circuit voltage is completely recovered and reaches that of a conventional solar cell.

### 1. Introduction

The concept of the intermediate band solar cell (IBSC), first proposed in 1997 [1], is based on the introduction of a collection of energy levels within the bandgap of a semiconductor, which are generally referred to as the intermediate band (IB). The IB allows the absorption of sub-bandgap photons that would otherwise be wasted in photovoltaic conversion and leads to an increase in the current generated by the IBSC. Conventionally, an increase of the photocurrent can be achieved by reducing the bandgap of a single gap solar cell, but this is accompanied by a reduction of the voltage. In contrast, the current increase in the IBSC can be achieved without degradation of the voltage (*V*) [2,3]. For fundamental thermodynamic reasons [4] this requires that two photons are used for the generation of an electron–hole pair [1].

The carrier recombination in an IBSC has often been regarded as being inherently larger than that in a single gap solar cell with the same bandgap because of the additional electronic de-excitation paths originating from the intermediate levels. Increased recombination in a solar cell is associated with a drop in output voltage, which can be counteracted with the use of concentrated light [5]. In this work, we demonstrate experimentally how the IBSC recombination can be reduced at low temperature and under high-irradiance operation, so that it does not exceed the recombination in a conventional solar cell, leading to the recovery of the initially degraded open-circuit voltage ( $V_{\rm OC}$ ) and verifying our theoretical

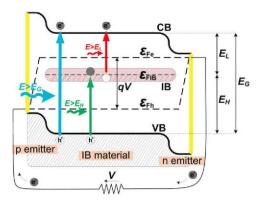
prediction. This is particularly promising for IBSCs to be used in concentrated photovoltaic (CPV) systems.

Fig. 1 presents an IBSC band diagram, where low-energy photon absorption is produced by means of the IB via two-step absorption processes involving transitions through the  $E_{\rm L}$  and  $E_{\rm H}$  sub-bandgaps. Proper operation of any IBSC requires the preservation of the output voltage, for which there must be three independent electronic populations corresponding to the valence band (VB), the IB and conduction band (CB), each represented by its own quasi Fermi level (QFL):  $\varepsilon_{\rm Fh}$ ,  $\varepsilon_{\rm FIB}$  and  $\varepsilon_{\rm Fe}$ . If the IB material is isolated from the contacts by two conventional semiconductors, the  $V_{\rm OC}$  will only be limited by the fundamental bandgap,  $E_{\rm G}$ , and not by any of the lower  $E_{\rm H}$  or  $E_{\rm L}$  sub-bandgaps. Further discussion of the design and operation of the IBSC can be found in Refs. [3,6–8].

The experimental samples used in this work were fabricated by embedding InAs quantum dot (QD) regions within a GaAs host, with the IB created from the confined energy levels of the QDs. The resulting QD-IBSC technology [9–12] exhibits a thermally connected IB–CB transition under room-temperature (RT) operation, which means that the electronic transfer through the  $E_{\rm L}$  subbandgap is characterized by a thermalization rate much higher than the photogeneration rate [13]. This difference makes a large IB–CB QFL split difficult to achieve at RT; therefore, low-temperature characterization is essential for observing a  $V_{\rm OC}$  that is not limited by the  $E_{\rm H}$  sub-bandgap.

# 2. Methods

Several designs for engineered IB materials have already been proposed, some of which have also been developed, allowing the



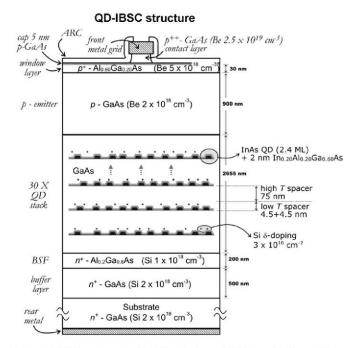
**Fig. 1.** Band diagram of an IBSC showing the VB, IB and CB along with their associated QFLs, labeled  $\epsilon_{\rm Fh},\,\epsilon_{\rm FIB}$  and  $\epsilon_{\rm Fe},\,$  respectively. Photons of energies above each of the sub-bandgaps ( $E_{\rm G},E_{\rm H}$  and  $E_{\rm L}$ ) pump electrons through each of the three possible transitions.

verification of several fundamental IBSC principles [14–16]. These IB materials can be divided into three main families: QD-based materials [10,17], bulk semiconductors [18] (which, in turn, are subdivided into thiospinels [19], highly mismatched semiconductor alloys [20,21] thin-film semiconductors [22] and others) and molecular-based materials [23].

In this work, we focus on QD-IBSCs, which to date have mainly been based on In(Ga)As/GaAs materials for the QD/barrier system. Although this material combination is not optimal in terms of efficiency, it was first identified as being a mature and ready-to-manufacture technology. In practice, the InAs QDs are characterized by a low height-base aspect ratio and sizes that are far from the calculated optimum (approximately 3 nm for spherical dots [24]). As a result, several confined levels appear in these nanostructures in addition to the fundamental level, thereby reducing the efficiency of devices due to increased recombination [25].

Both the excess QD-IBSC recombination and its influence on the  $V_{\rm OC}$  can be quantified. This quantification is performed experimentally by measuring photogenerated current density  $(J_L)-V_{OC}$  pairs at different concentrations, which, according to the superposition principle, describe recombination in the device without the influence of a series resistance. The fact that this measurement is acquired at zero amperes also contributes to the negligible influence of other spurious effects that may occur, e.g., non-ohmic effects at the contacts. For this purpose, we have developed a novel low-temperature, high-concentration characterization technique based on a pulsed light source (a xenon flash lamp), a fast computer-controlled data acquisition unit and a four quadrant source meter for the reverse voltage polarization of the cell. This setup can reach an irradiance equivalent to 10,000 suns when focused on an area of approximately  $5 \times 5 \text{ mm}^2$  with a concentrating lens and operated at RT. A closed-cycle He cryostat was used for the low-temperature concentration measurements. However, the irradiance can only reach 1,000 suns in this mode, because of the need for different optics to direct the light into the cryostat chamber.

Two different types of solar cells were fabricated by molecular beam epitaxy (MBE) for this work: a 30 layer InAs/GaAs QD-IBSC and a GaAs reference solar cell (the control sample). A simple IBSC requires only an IB material sandwiched between p- and n-emitters for the exploitation of the full  $V_{\rm OC}$  potential. A detailed sketch of the layer structure of our QD-IBSC is presented in Fig. 2, together with the specifications of the composition, doping and thickness of each layer. The bottom layer consists of a 625  $\mu$ m thick n-type GaAs wafer on top of which the device was grown. Next, a buffer layer was grown to smooth the surface. Then, an  $n^+$ -  $Al_{0.2}Ga_{0.8}As$  that acts as the n-emitter of the cell, was added.



**Fig. 2.** QD-IBSC layer structure including doping and thickness data for each layer. The IB region consists of 30 stacked QD layers, each of which is separated by a thick spacer, which incorporates the Si  $\delta$ -doping.

The IB material was grown next and consists of 30 layers of selfassembled InAs QDs built up by Stranski-Krastanov growth and separated by thick, intrinsic GaAs spacers, which, as an alternative to the strain compensation technique [11,12], are meant to dilute the accumulated vertical strain and thus prevent dislocations. A large spacer thickness also prevents carrier tunneling between contiguous QD layers [14]. A quaternary In<sub>0.2</sub>Al<sub>0.2</sub>Ga<sub>0.6</sub>As capping was added to preserve the QD height/base aspect ratio, to redshift the E<sub>H</sub> transition [26] and to eliminate the CB discontinuity with the GaAs barrier layer and thus, prevent the formation of an additional quantum well (QW) [27] in an effort to isolate the IB from the CB. Si δ-doping was included to half-fill the IB with electrons, so there are both empty states to enable the electrons to be pumped from the VB to the IB and filled states to assist in the IB-CB transition [8]. After the IB material region, a p-type GaAs emitter completes the electrically active part of the solar cell. Finally, an AlGaAs window, a p+-GaAs contact layer and an antireflective coating (ARC) complete the IBSC structure. The reference cell structure is identical except for the lack of the QD-IB region and the substitution of the  $n^+$ -Al<sub>0.2</sub>Ga<sub>0.8</sub>As emitter by a  $2 \times 10^{17}$  cm<sup>-3</sup> doped, 3.1  $\mu$ m n-type base layer on top of a 200 nm  $\,n^+$ -Al $_{0.2}{\rm Ga}_{0.8}{\rm As}$  BSF layer. We will later discuss the impact of this substitution in the interpretation of our results.

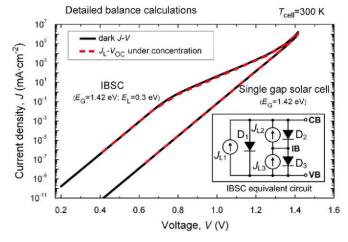
# 3. Results and discussion

From detailed balance calculations at RT, the electrical behavior of both the IBSC and the reference cell operating at the radiative limit can be inferred (the sun is treated as a black body at 6000 K). At 1 sun operation, the IBSC efficiency is marginally higher than the reference cell efficiency (34.5% compared to 30.7%), assuming a fundamental bandgap of  $E_G$ =1.42 eV at RT and an IB located 0.3 eV below the CB, which is a realistic assumption for InAs/GaAs QD-IBSCs. However, under high irradiance, this same IBSC structure has a significantly higher potential than the reference cell, e.g., when both are operated at 1,000 suns, the IBSC efficiency limit is 47.2% compared to 35.8%. This

larger efficiency difference is caused by the more rapid increase of  $V_{\rm OC}$  for the IBSC compared to the reference cell, i.e., the reference cell  $V_{\rm OC}$  is approximately 150 mV higher than the IBSC  $V_{\rm OC}$  under 1 sun operating conditions, whereas the  $V_{\rm OC}$  difference decreases to 8 mV for a concentrated illumination of 1,000 suns.

The dark curve of a solar cell provides an insight into the recombination taking place in the device. It is conventionally divided into two recombination regimes: the low-voltage regime, dominated by Shockley-Read-Hall recombination in the spacecharge region and characterized by a low exponential slope in the range of e/2kT (i.e., an ideality factor  $m_2=2$ ); and the high-voltage regime, dominated by recombination in the neutral zones and characterized by an exponential slope of e/kT ( $m_1=1$ ) at low injection. However, as shown in the detailed balance calculations of Fig. 3 (resulting in a purely m=1 dark current), the electronic recombination of an IBSC differs from the conventional solar cell case. In black, solid lines, the dark current density (J)-V characteristics for both an IBSC ( $E_G = 1.42 \text{ eV}$ ,  $E_H = 1.12 \text{ eV}$  and  $E_L$ =0.3 eV) and its reference cell are represented. Within the low J range, up to  $\sim$ 1 mA cm $^{-2}$ , both curves remain parallel and are separated by a significant gap, i.e., the IBSC recombination is much larger because, even with the IB assumed thermally separated from the CB (and radiatively linked), it is dominated in this range by only one of the electronic transitions to or from the IB ( $D_2$  and  $D_3$  in the inset of Fig. 3). Between J=1 mA cm<sup>-2</sup> and  $J=2.5\times10^4$  mA cm<sup>-2</sup>, an operating regime with a lower exponential slope (and influenced by the second of the electronic transitions to or from the IB) becomes dominant and leads to the recovery of  $V_{OC}$ . Above the current density of a typical GaAs solar cell operating under 1,000 suns illumination ( $2.5 \times 10^4$  mA cm<sup>-2</sup>), the two recombination characteristics nearly converge, meaning that the VB-CB transition (D<sub>1</sub> in the inset) is now dominating the IBSC recombination, leading to the recovery of the IBSC  $V_{OC}$ , as demonstrated experimentally in this work.

The  $J_{\rm L}$ – $V_{\rm OC}$  characteristics of the IBSC and the reference cell calculated with the detailed balance model under illumination are represented as red, dashed curves in Fig. 3 for a broad range of concentrated light levels. These  $J_{\rm L}$ – $V_{\rm OC}$  curves, which superimpose their respective dark characteristics (at least to a fairly accurate extent), are the ones that can be compared in practice for the analysis of the solar cell recombination, since the series resistance present in actual device dark curves would otherwise make the interpretation of the recombination difficult.



**Fig. 3.** Detailed balance J–V characteristics of an IBSC ( $E_G$ =1.42 eV and  $E_L$ =0.3 eV) and a conventional solar cell with the same bandgap. The dark J–V characteristics are represented by black solid lines, and the  $J_1$ – $V_{OC}$  pairs are represented by red, dashed lines. The inset shows a simplified version of the IBSC electric equivalent circuit. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

The experimental results were analyzed and compared to the previous calculations. For this purpose, the  $J_L$ – $V_{\rm OC}$  characteristics of both the QD-IBSC and the reference cell were obtained at different light concentration levels and at four different temperatures. Fig. 4(a) shows the RT experiment, in which the recombination associated with the QD-based technology still appears to be too large (non-radiative recombination is also present) and the voltage recovery effect is not observed, i.e., both  $I_L$ - $V_{OC}$  curves remain approximately parallel. This is consistent with the fact that, under RT conditions, a large IB-CB QFL split is difficult to achieve, and the QD-IBSC behaves as a low-bandgap conventional solar cell [13]. The dark J-V curves are presented in dark, solid lines in order to verify the fulfillment of the superposition principle in the range where they are not strongly affected by the series resistance. In Fig. 4(b), the experimental temperature was lowered to  $T=150 \,\mathrm{K}$ , and the first evidence of an IBSC recombination actually approaching the reference cell recombination is observed. The latter shows the trend of IBSC recombination moving with a slope higher than e/kT toward convergence with the reference cell recombination as predicted by theory. The maximum  $V_{OC}$  obtained in this measurement exceeds the photon energy absorption threshold of approximately 1.05 eV, (corresponding to the  $E_{\rm H}$  sub-bandgap energy), as experimentally observed in the device quantum efficiency (not shown in this work, although extensively reviewed in Ref. [13]).

Fig. 4(c) shows the results of an experiment performed at T=77 K, in which the IBSC voltage recovery is complete at a high concentration;  $V_{\rm OC}$  values close to 1.4 eV are reached by both cells, and the IBSC recombination matches that of the reference cell, thus verifying the proposed IBSC voltage recovery model. Under these conditions, an exponential slope of e/kT fits the low-voltage range of the QD-IBSC  $J_{\rm L}$ - $V_{\rm OC}$  signature, which may be interpreted as an operating regime close to the radiative limit. Nevertheless, the low-voltage range of the reference cell  $J_{\rm L}$ - $V_{\rm OC}$  does not behave in such radiative way, since it shows a much lower exponential slope.

Finally, Fig. 4(d) presents the same measurements at T=20 K, showing the full IBSC recombination characteristic lying below the reference cell recombination. This may be justified by the use of the  $n^+$ -Al<sub>0.2</sub>Ga<sub>0.8</sub>As rear n-type emitter for the QD-IBSC, which decreases the contribution to the total recombination provided by this emitter, in contrast to the use of a conventional n-GaAs emitter in the reference cell and supports that recombination is dominated by the emitters and not by the QD-region. The largest IBSC output voltage value measured at this temperature corresponds to  $V_{\rm OC}$ =1.506 V. It is exceptionally high, although consistently lower than the QD-IBSC  $E_{\rm G}/e$  (1.516 eV), which accounts for the bandgap of a properly operating IBSC at T=20 K.

#### 4. Conclusions

It was predicted that the increased recombination and the associated strong reduction of the device  $V_{\rm OC}$  observed in the IBSC operating at 1 sun can be counteracted by the use of concentrated illumination. Theory has been verified experimentally using low temperatures, which are necessary to slow down the fast RT IB–CB thermalization rate. The trend of the QD–IBSC  $J_{\rm L}$ – $V_{\rm OC}$  characteristics under concentrated light towards an exponential slope closer to e/kT is identified as the mechanism responsible for voltage recovery. It can be explained by a recombination initially dominated by IB–related diodes that evolves toward a recombination dominated by the fundamental transition, i.e., from CB to VB, making the IBSC recombination equivalent to that of a single gap solar cell. The latter has also been verified experimentally in this work when low measurement temperatures are employed, so an IB–CB QFL split

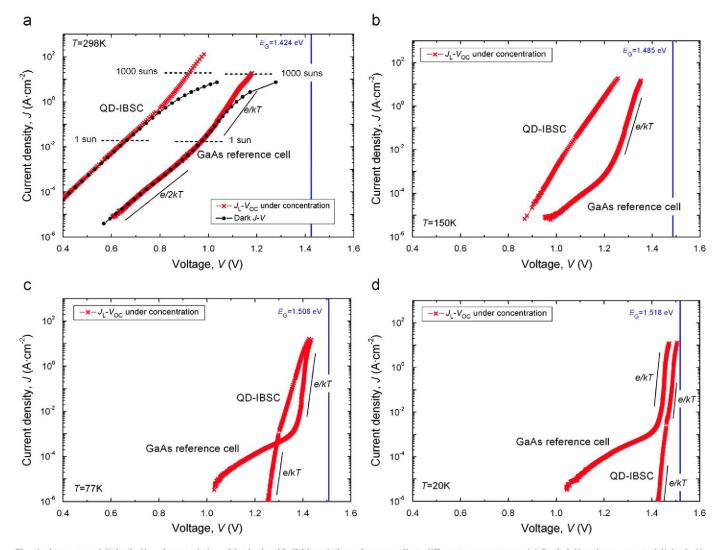


Fig. 4. Concentrated light  $J_L$ – $V_{OC}$  characteristics of both the QD-IBSC and the reference cell at different temperatures. (a) Dark J–V and concentrated light  $J_L$ – $V_{OC}$  characteristics measured at T=298 K. The concentration levels are indicated with horizontal gray dashed lines. (b), (c) and (d) show concentrated light  $J_L$ – $V_{OC}$  measurements at T=150 K, T=77 K and T=20 K, respectively. The solar cells fundamental bandgap is represented with a solid, blue line for each temperature and the e/kT and e/2 kT exponential slopes are also indicated when present. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

can be easily achieved and a properly operating IBSC can be obtained, even without concentration. Furthermore, an enhanced QD-IBSC electrical performance (with respect to its reference cell) is described under these conditions.

In addition, the maximum measured IBSC  $V_{\rm OC}$  (1.5 V) corresponds to a QFL separation that is very close to the material bandgap (1.52 eV at T= 20 K) divided by the electron charge. This exceeds the photon energy absorption threshold. An output voltage that is not limited by its absorption threshold is a proof of the voltage preservation operation principle.

We can conclude that the new low-temperature, concentrated light  $J_{\rm L}$ – $V_{\rm OC}$  characterization method presented here is a powerful tool for the analysis of solar cell recombination in IBSCs because series resistance, which usually dramatically disguises the low-temperature dark J–V characteristic, does not affect it.

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