

A Comparison of Some Numerical Methods for the Advection-Diffusion Equation

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This paper describes a comparison of some numerical methods for solving the advection-diffusion (AD) equation which may be used to describe transport of a pollutant. The one-dimensional advection-diffusion equation is solved by using cubic splines (the natural cubic spline and a "special" AD cubic spline) to estimate first and second derivatives, and also by solving the same problem using two standard finite difference schemes (the FTCS and Crank-Nicolson methods). Two examples are used for comparison; the numerical results are compared with analytical solutions. It is found that, for the examples studied, the finite difference methods give better point-wise solutions than the spline methods.

1 Introduction

Pepper et al. (1) and Okamoto et al. (2) solve the one-dimensional advection equation by using a spline interpolation technique that they call a quasi-Lagrangian cubic spline method. In section 4 of (3), Ahmad & Kothiyari solve the one dimensional advection-diffusion equation by using cubic spline interpolation for the advection component and the Crank-Nicolson scheme for the diffusion component. Sastry (6) uses a cubic spline technique to approximate the solution of the one-dimensional diffusion equation.

In this paper the one-dimensional advection-diffusion (advection-dispersion) equation is solved by using cubic spline interpolation to obtain estimates for the advection and diffusion terms. The same problem is investigated using finite difference schemes. Numerical results are shown and compared with analytical solutions.

2 Governing Equation and Numerical Methods

2.1 Governing Equation

The volumetric concentration of a pollutant in a moving, turbulent fluid may be described by the advection-diffusion equation

$$\frac{\partial C}{\partial t} + \nabla \cdot (\mathbf{V}C) = \nabla \cdot (\bar{K} \otimes \nabla C) \quad (1)$$

Here, $C(x, y, z, t)$ is the concentration (mass per unit volume) of pollutant at point (x, y, z) in Cartesian coordinates, at time t . The vector \mathbf{V} is the fluid velocity field and \bar{K} is the eddy-diffusivity or dispersion tensor.

In this study, we consider one-dimensional motion with constant speed u and dispersion D which gives the 1-D advection-diffusion equation:

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} = D \frac{\partial^2 C}{\partial x^2} \quad (2)$$

with appropriate initial and boundary conditions. Here, $C(x, t)$ is the concentration of the pollutant at point x ($0 \leq x \leq L$) and time t , u is the constant wind speed in the x direction and D is the diffusivity coefficient in the x direction. Several combinations of boundary conditions are possible. We distinguish three cases, all for a finite domain:

Case 1:

$$\begin{aligned} C(x, 0) &= f(x) \\ C(0, t) &= c_0 \\ C(L, t) &= c_L \end{aligned} \quad (3)$$

Case 2:

$$\begin{aligned} C(x, 0) &= f(x) \\ C(0, t) &= c_0 \\ \frac{\partial C}{\partial x}(0, t) &= 0 \end{aligned} \quad (4)$$

Case 3:

$$\begin{aligned} C(x, 0) &= f(x) \\ C(0, t) &= c_0 \\ q(0, t) = (Cu - D \frac{\partial C}{\partial x})(0, t) &= \text{constant} \end{aligned} \quad (5)$$

where c_0, c_L are constant concentration values, while the quantity $q = uC - D \frac{\partial C}{\partial x}$ is the mass flux of pollutant per unit cross-sectional area, and includes both of the advective and dispersive components. All cases correspond to a fixed constant concentration at the left-hand end, together with one of: a constant concentration at the right-hand end (Case 1), advective pollutant inflow only at the left-hand end (Case 2) or fixed constant pollutant influx there (Case 3).

For construction of interpolating cubic splines at any step in a numerical approximation procedure, 2 extra conditions are required apart from the C_i values at the mesh points x_i . The "natural" cubic spline requires the condition that $\frac{\partial^2 C}{\partial x^2} = 0$ on

the interpolants at $x = 0, L$. How does this requirement impact on the boundary conditions?

For all cases above, $C = c_0$ at $x = 0$ implies that

$$\frac{\partial C}{\partial t}(0, t) = 0,$$

and therefore

$$u \frac{\partial C}{\partial x} - D \frac{\partial^2 C}{\partial x^2} = 0$$

at $x = 0$, i.e. $\frac{\partial q}{\partial x}(0, t) = 0$. This implies that $q(0^-, t) = q(0^+, t)$. Assuming that the pollutant source is well-mixed, then $C = c_0$ upstream and $q = uc_0$ at $x = 0^-$. Because $C = c_0$ at $x = 0^+$, this means that $\frac{\partial C}{\partial x} = 0$ at $x = 0^+$ and so $\frac{\partial^2 C}{\partial x^2} = 0$ at $x = 0^+$. Cases 1 and 2 therefore are consistent with the boundary condition at $x = 0$ corresponding to the natural cubic spline, while Case 3 is not. Hence the motivation for developing a "special" cubic spline scheme for the advection-dispersion equation, as outlined in a later section of this paper.

2.2 Finite difference (FTCS) method

Finite difference schemes involve calculating approximate values of the unknown function at a finite number of (mesh- or grid-) points in the domain. Here we let $0 = x_1 \leq x_j \leq x_{N+1} = L$ be the grid points in x-domain. The time is divided into equal steps of size Δt , with time $t_n = n\Delta t$. Derivatives are approximated by truncated Taylor Series expansions. For our purposes, we use an explicit forward difference estimate for the time derivative (FT), and central difference approximations for the space derivatives (CS) that both have the same truncation error; hence the acronym FTCS. The approximate solution of the governing equation using the finite difference scheme satisfies:

$$\frac{C_j^{n+1} - C_j^n}{\Delta t} + u \frac{C_{j+1}^n - C_{j-1}^n}{2\Delta x} = D \frac{C_{j+1}^n - 2C_j^n + C_{j-1}^n}{(\Delta x)^2} + \mathcal{O}(\Delta t, (\Delta x)^2) \quad (6)$$

for $j = 2, 3, \dots, N$, while for $j = N + 1$, one-sided forms of the difference formulae are required:

$$\begin{aligned} & \frac{C_{N+1}^{n+1} - C_{N+1}^n}{\Delta t} + u \frac{3C_{N+1}^n - 4C_N^n + C_{N-1}^n}{2\Delta x} \\ & = D \frac{2C_{N+1}^n - 5C_N^n + 4C_{N-1}^n - C_{N-2}^n}{(\Delta x)^2} + \mathcal{O}(\Delta t, (\Delta x)^2) \end{aligned} \quad (7)$$

all for $n = 0, 1, 2, \dots$, where the initial values are $C_j^0 = C((j-1)\Delta x, 0)$.

Hindmarsh (7) and Sousa (8), show that the conditions that the finite difference scheme in Equation (6) is stable are

$$2 \frac{D\Delta t}{\Delta x^2} \leq 1,$$

and

$$\left(\frac{u\Delta t}{\Delta x}\right)^2 \leq 2\frac{D\Delta t}{\Delta x^2}.$$

Rearrangement and simplification of both of these conditions gives the restriction on the size of the time step in terms of the parameters and the grid-size:

$$\Delta t \leq \min\left\{\frac{\Delta x^2}{2D}, \frac{2D}{u^2}\right\}$$

2.3 The Crank-Nicolson method

The Crank-Nicolson scheme approximates the governing equation by using central differences in time; the spatial derivatives are estimated by the average of their values at time step n and step $n + 1$, in the form:

$$\begin{aligned} & \frac{C_j^{n+1} - C_j^n}{\Delta t} + u\left(\frac{1}{2}\left(\frac{C_{j+1}^{n+1} - C_{j-1}^{n+1}}{2\Delta x} + \frac{C_{j+1}^n - C_{j-1}^n}{2\Delta x}\right)\right) \\ &= D\left(\frac{1}{2}\left(\frac{C_{j+1}^{n+1} - 2C_j^{n+1} + C_{j-1}^{n+1}}{(\Delta x)^2} + \frac{C_{j+1}^n - 2C_j^n + C_{j-1}^n}{(\Delta x)^2}\right)\right) + \mathcal{O}(\Delta t, \Delta x^2), \end{aligned} \quad (8)$$

with a similar equation for the right-hand end-point. The Crank-Nicolson scheme is unconditionally stable (4).

2.4 Cubic spline method ("natural" cubic spline)

For this work the definition of the "natural" cubic spline includes:

- (i) The interpolating spline segments are cubic polynomial functions on each sub-interval $[x_j, x_{j+1}]$, $j = 1, 2, \dots, N$, and the segments agree with the function values at the grid-points;
- (ii) the first and second derivatives of the cubic spline segments are continuous at the internal points;
- (iii) the second derivatives of the cubic spline segments at the first and the last grid points are equal to zero.

The approximate solution of the governing equation using the cubic spline method satisfies:

$$\frac{C_j^{n+1} - C_j^n}{\Delta t} + uP_j^n = DQ_j^n. \quad (9)$$

for $j = 1, 2, \dots, N + 1$; $n = 0, 1, 2, \dots$ where P_j^n is the first derivative and Q_j^n the second derivative of the cubic spline function at the point x_j at time $n\Delta t$. Equation (9) can be written in the explicit form:

$$C_j^{n+1} = C_j^n + D\Delta t Q_j^n - u\Delta t P_j^n. \quad (10)$$

The values of the slopes P_j^n can be obtained by solving the following system of simultaneous equations (derived by manipulation of the equations which result from

the continuity conditions for the spline segments; see (5) for details of algebraic working):

$$\begin{bmatrix} 2 & 1 & 0 & 0 & 0 & \cdot & \cdot & \cdot & 0 & 0 & 0 \\ \alpha_2 & 2 & \mu_2 & 0 & 0 & \cdot & \cdot & \cdot & 0 & 0 & 0 \\ 0 & \alpha_3 & 2 & \mu_3 & 0 & \cdot & \cdot & \cdot & 0 & 0 & 0 \\ \cdot & & & & & \cdot & & & & & \cdot \\ \cdot & & & & & \cdot & & & & & \cdot \\ \cdot & & & & & \cdot & & & & & \cdot \\ 0 & 0 & 0 & 0 & 0 & \cdot & \cdot & \cdot & \alpha_{N-1} & 2 & \mu_{N-1} \\ 0 & 0 & 0 & 0 & 0 & \cdot & \cdot & \cdot & 0 & 1 & 2 \end{bmatrix} \begin{bmatrix} P_1^n \\ P_2^n \\ P_3^n \\ \cdot \\ \cdot \\ \cdot \\ P_N^n \\ P_{N+1}^n \end{bmatrix} = \begin{bmatrix} d_1^n \\ d_2^n \\ d_3^n \\ \cdot \\ \cdot \\ \cdot \\ d_N^n \\ d_{N+1}^n \end{bmatrix} \quad (11)$$

where

$$\begin{aligned} d_1^n &= 3\left(\frac{C_2^n - C_1^n}{x_2 - x_1}\right) \\ d_i^n &= 3\frac{\mu_j}{h_{j+1}}(C_{j+1}^n - C_j^n) + 3\frac{\alpha_j}{h_j}(C_j^n - C_{j-1}^n) \quad \text{for } j=2,3,\dots,N \\ d_{N+1}^n &= 3\left(\frac{C_{N+1}^n - C_N^n}{x_{N+1} - x_N}\right) \end{aligned}$$

and where $\alpha_j = \frac{h_{j+1}}{(h_j + h_{j+1})}$, $\mu_j = 1 - \alpha_j = \frac{h_j}{(h_j + h_{j+1})}$, $h_{j+1} = x_{j+1} - x_j$ and $h_j = x_j - x_{j-1}$.

The values of Q_j^n are the second derivatives of cubic spline at points x_j for $j = 2, 3, \dots, N$, at time $n\Delta t$. For the natural cubic spline it is assumed that $s_1''(x_1) = s_n''(x_{n+1}) = 0$ (i.e. $Q_1^n = Q_{N+1}^n = 0$). Then we have:

$$Q_j^n = 6\frac{C_{j+1}^n - C_j^n}{(x_{j+1} - x_j)^2} - 4\frac{P_j^n}{x_{j+1} - x_j} - 2\frac{P_{j+1}^n}{x_{j+1} - x_j} \quad (12)$$

for $j = 2, 3, \dots, N$.

2.5 Cubic spline method ("Special A-D" cubic spline)

The boundary condition $C = c_0$ at $x = 0$, implies that

$$\frac{\partial C}{\partial t}(0, t) = 0, \quad (13)$$

and so

$$\left(u\frac{\partial C}{\partial x} - D\frac{\partial^2 C}{\partial x^2}\right)(0, t) = 0. \quad (14)$$

In this section we present a cubic spline interpolation scheme that satisfies the condition (12).

The requirement is that $uP_1^n = DQ_1^n$ and $uP_{N+1}^n = DQ_{N+1}^n$. These imply that

$Q_1^n = \frac{u}{D}P_1^n$ and $Q_{N+1}^n = \frac{u}{D}P_{N+1}^n$ (compare with natural cubic spline where $Q_1^n = 0; Q_{N+1}^n = 0$). The values of P_j^n can then be calculated from the following system:

$$AP = d \quad (15)$$

where

$$A = \begin{bmatrix} 2 + \frac{u(x_2-x_1)}{2D} & 1 & 0 & 0 & 0 & \dots & 0 & 0 & 0 \\ \alpha_2 & 2 & \mu_2 & 0 & 0 & \dots & 0 & 0 & 0 \\ 0 & \alpha_3 & 2 & \mu_3 & 0 & \dots & 0 & 0 & 0 \\ \vdots & & & & & \ddots & & & \vdots \\ \vdots & & & & & & & & \vdots \\ \vdots & & & & & & & & \vdots \\ 0 & 0 & 0 & 0 & 0 & \dots & \alpha_{n-1} & 2 & \frac{\mu_{n-1}}{2D} \\ 0 & 0 & 0 & 0 & 0 & \dots & 0 & 1 & \frac{4D-u(x_{N+1}-x_N)}{2D} \end{bmatrix}$$

$$P = \begin{bmatrix} P_1^n \\ P_2^n \\ P_3^n \\ \vdots \\ \vdots \\ \vdots \\ P_N^n \\ P_{N+1}^n \end{bmatrix}$$

and

$$d = \begin{bmatrix} d_1^n \\ d_2^n \\ d_3^n \\ \vdots \\ \vdots \\ \vdots \\ d_N^n \\ d_{N+1}^n \end{bmatrix}$$

where

$$d_1^n = 3\left(\frac{C_2^n - C_1^n}{x_2 - x_1}\right)$$

$$d_j^n = 3\frac{\mu_j}{h_{j+1}}(C_{j+1}^n - C_j^n) + 3\frac{\alpha_j}{h_j}(C_j^n - C_{j-1}^n) \quad \text{for } j=2,3,\dots,N$$

$$d_{N+1}^n = 3\left(\frac{C_{N+1}^n - C_N^n}{x_{N+1} - x_N}\right)$$

and $\alpha_j = \frac{h_{j+1}}{(h_j + h_{j+1})}$ and $\mu_j = 1 - \alpha_j = \frac{h_j}{(h_j + h_{j+1})}$. For this case $Q_j^n, j = 2, 3, \dots, N$ can be obtained from

$$Q_j^n = 6\frac{C_{j+1}^n - C_j^n}{(x_{j+1} - x_j)^2} - 4\frac{P_j^n}{(x_{j+1} - x_j)} - 2\frac{P_{j+1}^n}{(x_{j+1} - x_j)} \quad (16)$$

while Q_1^n and Q_{N+1}^n are calculated directly using the formulae already given above.

3 Numerical Experiments

Two examples are solved by the various methods outlined above, and the calculated numerical approximations are compared with the analytical solutions and with each other. The idea is to try to find the method which gives the best estimates for solutions of the advection-dispersion equation. Example 1 is a boundary value problem on a finite domain, as in Case 1 defined in Section 2.1 above. Example 2 is for a semi-infinite domain, but a finite-domain solution is sought.

3.1 Example 1.

The one-dimensional advection-diffusion equation:

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} = D \frac{\partial^2 C}{\partial x^2} \quad (17)$$

is to be solved with the boundary and initial conditions:

$$\begin{aligned} C(0, t) &= 0 \\ C(L, t) &= 100 \\ C(x, 0) &= \frac{100x}{L} \quad 0 \leq x \leq L \end{aligned}$$

The analytical solution is (4):

$$C(x, t) = 100 \left[\frac{e^{\frac{Px}{L}} - 1}{e^P - 1} + \frac{4\pi e^{\frac{Px}{2L}} \sinh(P/2)}{e^P - 1} \sum_{m=1}^{\infty} A_m + 2\pi e^{\frac{Px}{2L}} \sum_{m=1}^{\infty} B_m \right] \quad (18)$$

where P is the Peclet number

$$P = \frac{uL}{D}$$

and the coefficients A_m, B_m are given by

$$\begin{aligned} A_m &= (-1)^m \frac{m}{\beta_m} \sin\left(\frac{m\pi x}{L}\right) e^{-\lambda_m t} \\ B_m &= \left[(-1)^{m+1} \frac{m}{\beta_m} \left(1 + \frac{P}{\beta_m}\right) e^{\frac{-P}{2}} + \frac{mP}{\beta_m^2} \right] \sin\left(\frac{m\pi x}{L}\right) e^{-\lambda_m t}, \end{aligned}$$

where

$$\beta_m = \left(\frac{P}{2}\right)^2 + (m\pi)^2$$

and

$$\lambda_m = \frac{u^2}{4D} + \frac{m^2\pi^2 D}{L^2} = \frac{D\beta_m}{L^2}$$

In this study, we assume that $L = 1.0$ m, $D = 0.01$ m²s⁻¹, $u = 0.1$ ms⁻¹, $\Delta x = 0.1$ m (i.e. $N = 10$) and $\Delta t = 0.01$ s. The numerical solutions and the analytical solution of this example are shown in Figure 1. The continuous line is the steady-state (long-time) solution.

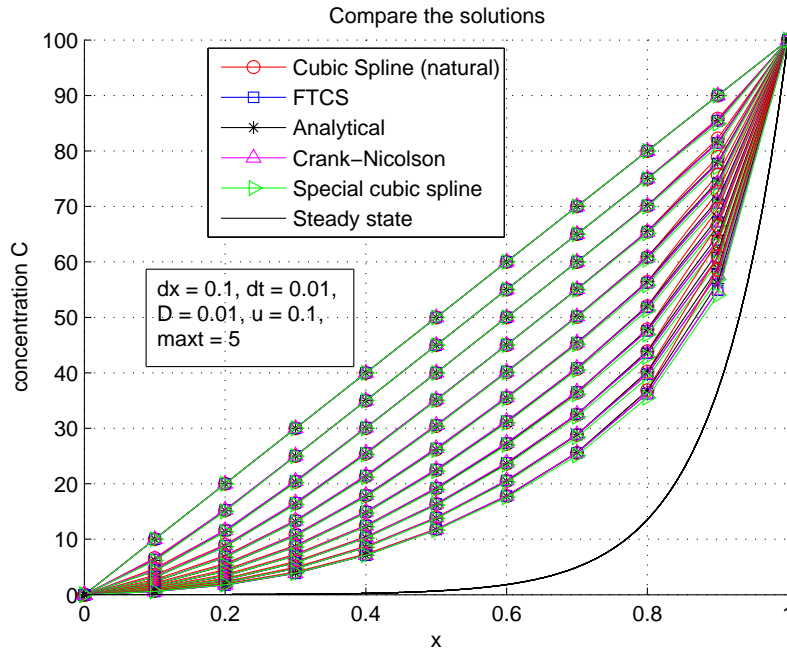


Figure 1: Comparisons of the solutions for Example 1 using various methods

It can be seen in Figure 1 that the numerical solutions given by all of the methods decrease with time from the initial (linear profile) condition. The calculated values are shown at times $t = 0(0.5)5s$. The time $T = 5s$ is still too early for the steady state to be approached. The errors in the numerical results are shown in Figure 2.

Figure 2 shows the error in the numerical solutions from each of the methods when compared with the analytical solution, for the three cases $N = 10$, $N = 20$ and $N = 40$, corresponding to $\Delta x = 0.1, 0.05$ and 0.025 respectively. Comparisons are made for the solutions at time $t = 5s$ and $t = 25s$. The latter time corresponds very closely to steady state. From Figure 2(a)(i) for $T = 5$ and $\Delta x = 0.1$, the errors are all small except near $x = 1$. The error for the natural cubic spline is positive and larger in magnitude than the errors for the other three methods, which are all negative. For smaller Δx [in (b)(i) $\Delta x = 0.05$ and (c)(i) $\Delta x = 0.025$] the errors all reduce in size and appear to be $\mathcal{O}(\Delta x^2)$.

For large T ($T = 25$ in Figure 2(a)(ii), (b)(ii) and (c)(ii), effectively steady-state) the error for the natural cubic spline scheme is again large and positive, while for the other three methods the errors are almost identical and negative, and this is the same for smaller Δx as well.

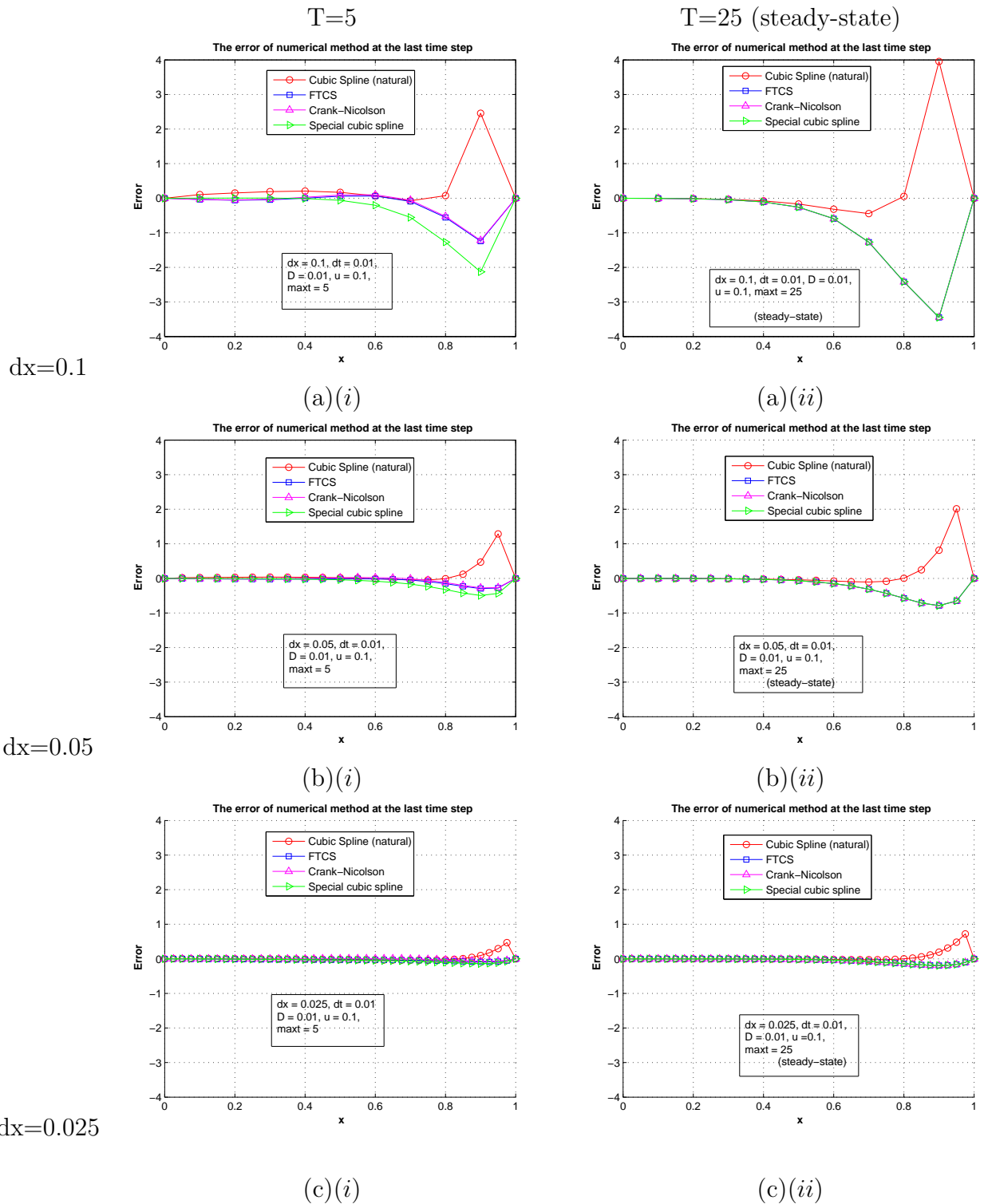


Figure 2: Comparison of the errors in the numerical solutions for Example 1 with the analytical solution. (a) $dx = 0.1$ [(i) $T = 5$, (ii) $T = 25$]. (b) $dx = 0.05$ [(i) $T = 5$, (ii) $T = 25$]. (c) $dx = 0.025$ [(i) $T = 5$, (ii) $T = 25$].

Because the solutions of the FTCS and Crank-Nicolson methods give better point-wise approximations to the analytical solutions than the "natural" cubic spline and "Special A-D" cubic spline schemes for Example 1, for the next example we will present only the solutions given by the FTCS method and the natural cubic spline schemes.

3.2 Example 2.

Solve the one-dimensional advection-diffusion equation:

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} = D \frac{\partial^2 C}{\partial x^2} \quad (19)$$

on a semi-infinite domain $x = [0, \infty)$, with the initial and boundary conditions:

$$\begin{aligned} C(x, 0) &= 0 \quad ; x \geq 0 \\ C(0, t) &= 1 \quad ; t > 0 \\ \frac{\partial C(\infty, t)}{\partial x} &= 0 \quad ; t > 0. \end{aligned}$$

The analytical solution to this problem is given by (9):

$$C(x, t) = \frac{1}{2} \left[\operatorname{erfc} \left(\frac{x - ut}{\sqrt{4Dt}} \right) + \exp \left(\frac{xu}{D} \right) \operatorname{erfc} \left(\frac{x + ut}{\sqrt{4Dt}} \right) \right]$$

where $\operatorname{erfc}(x)$ is the complementary error function defined by

$$\operatorname{erfc}(x) = 1 - \operatorname{erf}(x)$$

where

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-z^2} dz$$

In this example, the solution is required for a semi-infinite domain but we will solve this problem over a finite domain $x = [0, 2]$ by assuming that $L = 2 \text{ m}$, $D = 0.01 \text{ m}^2\text{s}^{-1}$, $u = 0.1 \text{ ms}^{-1}$, $\Delta x = 0.1 \text{ m}$ (i.e. $N = 20$) and $\Delta t = 0.01 \text{ s}$. We will then compare the solutions for $0 \leq x \leq 1$ only.

In this example the special cubic spline method is not presented. The conditions for a specialized A-D cubic spline method for this example are $P_1^n = 0$ and $Q_1^n = 0$, which are different from the conditions for the "Special A-D" cubic spline in Example 1.

The numerical solutions and the analytical solution for this example may be compared in Figure 3. Figure 4 shows the error of the numerical solutions when compared with the analytical solution, for the two cases $N = 10$ and $N = 40$, respectively. The values in Figure 4 are plotted every 50 time steps [i.e. at $t = 0(0.05)5\text{s}$].

In Figure 3, the numerical solutions are shown for $T = 5$; solutions given by

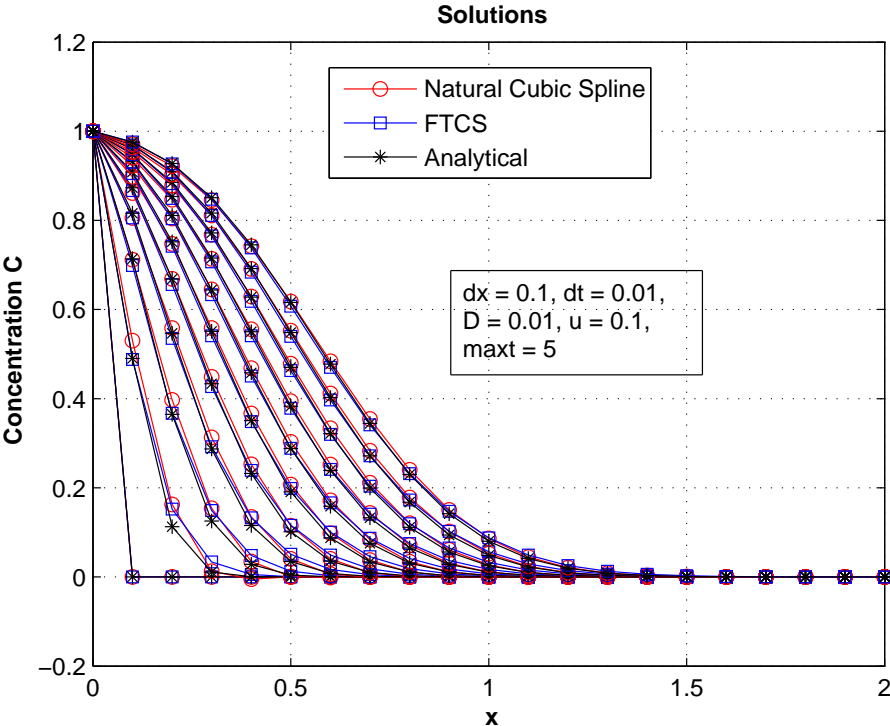


Figure 3: Comparisons of the solutions for Example 2 using various methods

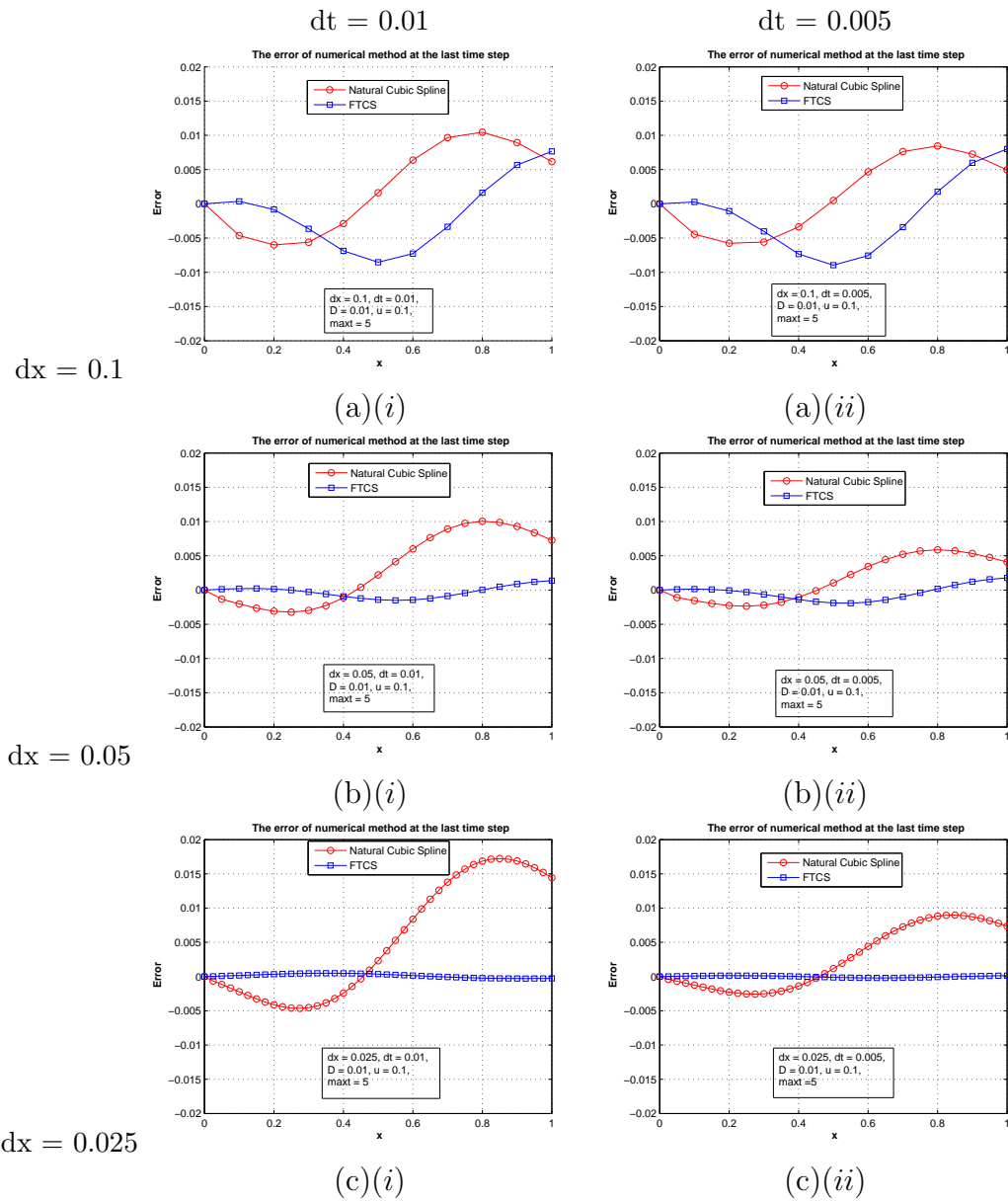


Figure 4: Comparison of the errors in the numerical solutions for Example 2 with the analytical solution at $T = 5$. (a) $dx = 0.1$ [(i) $dt = 0.01$, (ii) $dt = 0.005$]. (b) $dx = 0.05$ [(i) $dt = 0.01$, (ii) $dt = 0.005$]. (c) $dx = 0.025$ [(i) $dt = 0.01$, (ii) $dt = 0.005$].

all methods increase with time. The numerical solutions continue to increase to the steady state solution ($C(x, \infty) = 1$).

From Figure 4 (a)(i) for $\Delta x = 0.1$ and $dt = 0.01$, the errors for the natural cubic spline scheme are larger in magnitude than the errors for FTCS method. For smaller Δx [in (b)(i) $\Delta x = 0.05$ and (c)(i) $\Delta x = 0.025$] the errors all reduce in size and appear to be $\mathcal{O}(\Delta x^2)$.

For the small time step size $dt = 0.005$ in Figure 4 (a)(ii), (b)(ii) and (c)(ii), the error for the natural cubic spline is again larger than the FTCS method. However, as expected, the errors for the spline scheme decrease as the grid-size decreases.

4 Conclusion

Finite difference methods and cubic spline schemes for the one dimensional advection-diffusion equation have been presented. For the test examples studied, it has been found that the FTCS and the Crank-Nicolson finite-difference methods give better point-wise solutions than the "natural" cubic spline and "Special A-D" cubic spline schemes. However, the "Special A-D" cubic spline method gives better point-wise solutions than the "natural" cubic spline method.

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