



Research paper

## Prediction of rock block stability and scour depth in plunge pools

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### ABSTRACT

Block Stability in 3 Dimensions (BS3D) is an incremental-iterative algorithm and code to analyse general failure modes of rock blocks subjected to generic forces, including non-conservative forces such as water pressure. This algorithm is used to investigate the effect of high-velocity water jet impact on the stability of single rock blocks in plunge pools. After reviewing the literature, an approach is proposed to estimate the dynamic water pressure fluctuations at the pool bottom and in rock fractures. The approach is validated using an experimental study. In addition, by comparing the results of these analyses to additional experimental results, a scour threshold is introduced for plunge pools. It is found that if the maximum upward block displacement is larger than a quarter of the block height, the block is most probably removed from its mould. The proposed approach and threshold are validated using case histories that indicate excellent match between prediction and observed scour depth.

*Keywords:* Dam safety, fracture, rock, scour, spillway

### 1 Introduction

Standard dam risk assessment includes consideration of scour for the spillway and the energy dissipator. For concrete and masonry dams, it also includes the assessment of dam overtopping and the scour potential at the foundations. High-velocity plunging jets can result in scouring of the rock riverbed or the dam toe foundation. Assessment of the extent of scour is necessary to ensure the dam safety and to guarantee the stability of its abutments. Currently, there is no formulation for evaluating scour caused by general failure modes of rock blocks of general shape and subject to arbitrary loading (e.g. gravity, reinforcement, dam loads) and to the plunge pool water pressures. Available approaches to scour evaluation have limitations, which are explained below.

Bollaert's dynamic impulsion (DI) method (Bollaert 2002, Bollaert and Schleiss 2005) is limited to vertical translational failure (static failure mode) of parallelepiped rock blocks with one face at the plunge pool bottom. Indeed, roto-translational failures are common even for parallelepiped blocks subjected to pressure fluctuations (Fiorotto and Rinaldo 1992).

Annandale's (1995, 2006) erodibility index method is a classification method, as opposed to analytical or numerical methods based on mechanical principles, applicable to entire rock masses, but is not applicable to single rock blocks, which typically are critical for dam stability (Goodman and Powell 2003).

As for numerical methods of discontinuous rock masses, 3D distinct element code (3DEC) by Itasca (1999) has serious limitations in the dynamics of rigid bodies because it assumes that

Revision received 24 August 2011/Open for discussion until 30 June 2012.

the inertia tensor is always diagonal (Hart *et al.* 1988, Itasca 1999), i.e. rotations are assumed to be parallel to the applied resultant moment. In reality, rotations are parallel to the applied moment if and only if the applied resultant moment is parallel to a principal axis of inertia. 3D discontinuous deformation analysis (3D-DDA) by Jang and Lee (2002) as well as 3DEC use contact points at the vertices of the contact areas to calculate constraint forces: this does not allow for path-dependent behaviour of discontinuities. 3D-DDA assumes that all forces are conservative, and thus cannot handle follower forces, such as water pressure.

Herein, the stability of single rock blocks in plunge pools is investigated using Block Stability in 3D (BS3D), a Fortran code implemented by Asadollahi (2009) based on the algorithm introduced by Tonon (2007). BS3D follows the equilibrium path of a single rock block by using an incremental-iterative algorithm. It analyses general failure modes of rock blocks subject to generic forces, including non-conservative forces such as water forces. These are not conservative because water pressure is orthogonal to the surface onto which it is applied, and therefore water forces depend on the orientation of the surface itself (in conservative forces, the work done in moving a particle between two points is independent of the taken path). The algorithm uses consistent stiffness matrices that fully exploit the quadratic convergence of the adopted Newton–Raphson iterative scheme. The algorithm accounts for large block displacements and rotations, which, together with non-conservative forces render the stiffness matrix non-symmetric. The algorithm accounts for *in situ* stress, progressive failure, and fracture dilatancy, which introduce non-symmetric rank-one modifications to the stiffness matrix. Failure modes may originate from a limit point or from

dynamic instability (divergence or flutter). Equilibrium paths emanating from bifurcation points are followed by the algorithm, which identifies both static and dynamic failure modes; if a failure mode is detected, BS3D gives the initial acceleration of the rock block. BS3D has proven capable of detecting numerically challenging failure modes, such as rotations about only one point, and was validated using physical models and case histories (Tonon and Asadollahi 2008).

This work is organized as follows. Section 2 describes an approach to estimating pressure forces generated in plunge pools due to high-velocity jet impacts. Section 3 proposes a failure criterion for jointed rock masses. In Section 4, the scour model implemented in BS3D is calibrated and validated using the results of several experimental studies as well as case histories and prototypes, followed by conclusions in Section 5. A deterministic approach is used throughout.

## 2 Pressure forces generated in plunge pools

Studies on pressure fluctuations in plunge pools include those of Cola (1965), Hartung and Häusler (1973), Beltaos and Rajaratnam (1973), Mason and Arumugam (1985), Franzetti and Tanda (1987), Lopardo (1988), Armengou (1991), May and Willoughby (1991), Puertas and Dolz (1994), Ervine *et al.* (1997), Bollaert (2002), Bollaert and Schleiss (2003a, 2003b), or Ursino *et al.* (2003). The maximum pressure in a fracture is assumed to occur at the same time as the minimum pressure on the pool floor, and pressures are evaluated as stated in Table 1. The net upward pressure coefficients so calculated for different magnitudes of

Table 1 Equations used to calculate pressures at floor bottom and in rock features

Quantity and equation	Reference
$P_{\min, floor} = (C_{mean, floor} - C_{ext, floor}^-) \cdot \gamma_w \cdot \frac{V_j^2}{2g}$	Ervine <i>et al.</i> (1997)
$C_{mean, floor} = \begin{cases} \frac{38.4}{(1 + \beta)} \left(\frac{D_j}{Y}\right)^2 & \text{if } \frac{Y}{D_j} > 4 \\ 0.75 & \text{otherwise} \end{cases}$	Ervine <i>et al.</i> (1997)
$0.1 \leq C_{ext, floor}^- = \begin{cases} 1.8 \cdot e^{-0.20(Y/D_j)} & \text{if } \frac{Y}{D_j} > 4 \\ 0.1 \cdot e^{0.45(Y/D_j)} & \text{otherwise} \end{cases}$	Equations derived by authors based on data of Ervine <i>et al.</i> (1997) and Bollaert (2002)
$P_{\max, fracture} = (C_{mean, fracture} + C_{ext, fracture}^+) \cdot \gamma_w \cdot \frac{V_j^2}{2g}$	Bollaert (2002)
$C_{mean, fracture} = \begin{cases} \frac{38.4}{(1 + \beta)} \left(\frac{D_j}{Y}\right)^2 & \text{if } \frac{Y}{D_j} > 6 \\ 0.6 & \text{otherwise} \end{cases}$	Ervine <i>et al.</i> (1997) and Bollaert (2002)
$C_{ext, fracture}^+ = \begin{cases} 0.4 \cdot \ln\left(\frac{Y}{D_j}\right) & \text{if } \frac{Y}{D_j} < 10 \\ 187.3 \cdot \left(\frac{Y}{D_j}\right)^{2.4} + 0.15 & \text{otherwise} \end{cases}$	Equations derived by authors based on data of Ervine <i>et al.</i> (1997) and Bollaert (2002)

Note:  $V_j$  = jet velocity at pool surface,  $\gamma_w$  = unit weight of water,  $D_j$  = jet diameter at pool surface,  $Y$  = pool depth,  $\beta$  = free air content.

$Y/D_j$  were found to correspond well with those experimentally measured by Bollaert (2002) and to be in-line with the values obtained by Liu *et al.* (1998) by measuring the net uplift forces on scaled versions of rock blocks in plunge pools. Maximum pressure fluctuations, as opposed to root-mean-square values, are also typically used in the stability analysis of stilling basin slabs (Fiorotto and Rinaldo 1992, Bellin and Fiorotto 1995, Fiorotto and Salandin 2000).

The assumption that pressure fluctuations in each fracture are independent of each other is also based on considerations of the correlation function for pressure fluctuations because the integral scale of the correlation function for pressure fluctuations is small when compared with the typical dimension of a rock block (Melo *et al.* 2006); if the horizontal characteristic block dimension is at least twice the integral scale, then the pressure fluctuations at two fractures are independent (Fiorotto and Rinaldo 1992). This was confirmed in the facility developed by Bollaert (2002), which generated frequency spectra in good agreement with spectra of real high-velocity air–water jets. This assumption is a worst-case scenario because it entails that the minimum pressure at one fracture can occur simultaneously with the maximum pressure at another fracture, and vice versa. As a consequence, if a block has  $n$  faces, the stability analysis will be run  $2^n$  times to cover all possible combinations of pressure distributions on the block faces. Since the maximum and minimum dynamic pressure coefficients were obtained from 2 min records, their values must be doubled for a 24 h run time, which is considered representative of continuous operation of a plunge pool or a stilling basin (Ervine *et al.* 1997, Fiorotto and Salandin 2000).

### 3 Proposed failure criterion for jointed rock

#### 3.1 Impulsive nature of applied forces

The typical duration of maximum pressure fluctuations is of the order of 5 ms (Bollaert 2002), and there is 75% probability that pressure fluctuations persist less than 6 ms at the pool bottom (Manso 2006). For core jet conditions ( $Y/D_j < 4$ ), this value drops to 4 ms, whereas for transition and developed jet conditions ( $4 < Y/D_j < 6$  and  $Y/D_j > 6$ , respectively), this value is higher than 4 ms. Herein, the block is assumed to be subjected to its constant unbalanced force at failure for the maximum duration of an extreme pressure fluctuation, i.e. 5 ms. This assumption was successfully used for stilling basins (Fiorotto and Salandin 2000). In the absence of measurements or any other information regarding the maximum duration of an extreme pressure fluctuation, it is recommended to adopt a minimum of 5 ms and half of the fissure natural period as appropriate values.

BS3D follows the block in its static condition until equilibrium is possible between the active and reaction forces. As the block fails, the active force that cannot be in equilibrium with the constraints (nearby blocks) imparts an acceleration to the block (Tonon 2007). By applying rigorous rigid body dynamics, the formulation then computes the initial acceleration of the

block and its subsequent displacement by taking into account the inertia of the rock block. This displacement under maximum pressure fluctuations was experimentally observed by Yuditskii (1967, 1971), who presented the first conceptual model of the rock scouring process based on pressure fluctuations on rock blocks. His procedure is similar to that proposed here, in that it is based on an evaluation of the maximum instantaneous pressure separating a rock block from the matrix. For increasing pool depths, he compared the maximum pressure gradient amplitude originated by jet impact with a limit pressure value corresponding to the equilibrium condition. He also conducted more than 2000 tests focusing on the mechanisms of block ejection for varying scour depths, relative block size, block density, and joint thickness. One interesting observation of Yuditskii (1967, 1971) later made by Melo *et al.* (2006) is: “the block is ejected, neither by one pressure fluctuation of high amplitude nor by a succession of pressure fluctuations of high amplitude, but by one large average pressure that is established in the joint underneath the block following a small vertical displacement. The opening of the joint that allows this small vertical displacement is done by one pressure fluctuation of high amplitude”. What is interesting is that the flow and rock mass conditions that led to a small block displacement then led to block ejection from the rock mass. In other words, there is experimental evidence that a block is ejected if and only if it first fails according to one of the failure modes captured by the proposed stability analysis.

#### 3.2 Scour threshold

Below, a scour threshold is introduced for the cases in which the stability analysis predicts that a block of mass  $m$  fails in a translational mode, i.e. the dynamic failure mode obtained using BS3D gives an acceleration in one direction of magnitude  $a$ , without rotation. At failure, the block is assumed to be subjected by this constant acceleration for the maximum duration  $\Delta t$  of an extreme pressure fluctuation, i.e. 5 ms, causing a velocity  $V_{\Delta t} = ma$ .

After  $\Delta t$ , it is assumed that the block moves with the initial velocity  $V_{\Delta t}$  for a duration equal to  $t_{up} = T_c/2 - \Delta t$ , where  $T_c = 2L/c$  is the natural period of the characteristic length of the fractures that isolate the block (Annandale 2006),  $c$  is the pressure wave celerity, which may be estimated for air–water flow using the equation of Annandale (2006), and  $L$  is the characteristic fracture length ( $x_b + 2z_b$ ), where  $x_b$  and  $z_b$  are length and height of the block, respectively. The maximum upward displacement of the block  $h_{up}$  is thus

$$h_{up} = \frac{V_{\Delta t} \cdot (T_c - \Delta t) - g \cdot t_{up}^2}{2} \quad (1)$$

Bollaert (2002) proposed that a rock block is considered to be definitely removed from its matrix if  $h_{up}/z_b > 1$ . However, Bollaert and Schleiss (2005) found that the ultimate scour depth observed at Cabora Bassa dam corresponds to  $h_{up}/z_b = 0.20$ , which is inconsistent with the above criterion. Herein, BS3D was used to simulate the test data of Martins (1973). Based on

these, it is suggested that a block is most likely to be removed if  $h_{up}/z_b > 0.25$  (Section 4.2).

#### 4 Application, calibration, and validation of scour model

The experimental study of Federspiel *et al.* (2009) is used to validate the dynamic water pressures in Section 2, and BS3D to determine the stability of a rock block together with its displacements. Secondly, the failure criterion of Section 3 is calibrated by employing Martins' (1973) data. The entire scour model and BS3D algorithm are then validated using three case histories and prototypes.

##### 4.1 Experimental study

Federspiel *et al.* (2009) conducted a scaled experimental study to investigate the behaviour of a single rock block modelled as a hollow steel cube instrumented with a large series of pressure transducers, accelerometers, and displacement transducers. The sizes of the jet and the block were such that the whole block surface was subjected to characteristic pressure fluctuations and the block size was in the range of the prevailing turbulence structures. On prototypes, the jet dimensions are larger so that similar conditions apply to larger rock blocks when compared with the laboratory-instrumented block.

The scaled plunge pool consisted of a 3 m diameter, 1.4 m high cylindrical basin. The water was supplied with a 72 mm diameter cylindrical jet outlet. The maximum discharge was 120 l/s, corresponding to maximum jet outlet velocities of 30 m/s. The measurement box of length = width = 402 mm and height = 340 mm was composed of 20 mm thick steel plates. Inside the box, a large series of cavities housed pressure and displacement transducers. In its centre, a cavity of length = width = 202 mm and height = 201 mm housed the cubic highly-instrumented block of 200 mm side length. The width of the steel plates was optimized to obtain a block unit weight similar to real rock, i.e.  $\sim 25 \text{ kN/m}^3$ . On top of the block, holes were pre-perforated to fix the pressure transducers. Between the measurement box and the highly-instrumented block, a 3D fissure of 1 mm width was created. Both units were placed inside the basin that simulated the plunge pool.

The instrumentation consisted of a data acquisition system, 12 pressure transducers, two displacement transducers with a precision better than 0.01 mm, and an accelerometer transducer with a frequency range between 1 and 10 kHz. The pressure transducers fixed within the vertical plane allowed the pressure field around a block to be reconstructed. The displacement transducers and the accelerometer had fixed positions under the block in the "measurement box" and in the block, respectively.

Tests were performed with a flat bottom and a water jet centred in the middle of the highly-instrumented block. Two plunge pool water levels of  $Y = 0.1 \text{ m}$  and  $0.6 \text{ m}$ , and four jet outlet velocities of 19.6, 22.1, 24.6, and 27.0 m/s were tested. The  $Y = 0.1 \text{ m}$  water level in the plunge generated a core jet of

only  $Y/D_j = 1.39$ , while for  $Y = 0.6 \text{ m}$  resulted in a ratio of  $Y/D_j = 8.33$ . For each water level and jet velocity, three runs were performed. The data acquisition frequency was 1 kHz and the recording time 60 s (60,000 samples for each transducer).

The mean pressure coefficients  $C_p$  recorded directly under the jet axis for core jet impact were in good agreement with the theoretical curves of Ervine *et al.* (1997) and with previous pressure records of Bollaert (2002) and Manso (2006). The  $C_p$  values recorded away from the jet axis and inside the joints around the block were generally less than the mean pressures under the jet axis, as also observed by Bollaert (2002) and Manso (2006). For developed jet impact, however, the recorded values were higher than the theoretical curves and correspond to values for core jet impact. This is due to the jet deflecting the plunge pool water level and locally lowering it. Further, the measured pressure fluctuation coefficients  $C_p'$  were in good agreement with theory for core jets but differed slightly from developed jet theory. As such, the fluctuating component is lower than the theoretical curves, which again would correspond to a core jet rather than a developed jet.

Federspiel *et al.* (2009) recorded the maximum vertical block displacement  $h_{up,exp}$  and the maximum net uplift hydrodynamic pressure  $P_{net,exp}$  as the difference between the maximum dynamic water pressure on the bottom fracture and the minimum dynamic floor water pressure. The duration of the maximum pressure fluctuation was 1 ms. The accuracy of the equipment was  $\pm 1 \text{ ms}$ . The recorded displacements refer to a special case since the jet impact was perfectly aligned with the block centre. Therefore, the velocity change had a small effect and the pressure at the stagnation point kept the displacements to almost a constant value.

For each tested water level and jet velocity, BS3D analyses were performed to predict the maximum vertical block displacement  $h_{up}$  by using both the measured maximum net uplift hydrodynamic pressure from Federspiel's experiments and the hydrodynamic pressures determined as explained in Section 2. Because of the short test duration, the dynamic pressure magnitudes were estimated without considering a coefficient 2 accounting for long duration events in numerical simulations. In Federspiel's test apparatus, the jet velocity was adjusted by water discharge and the jet breakup length was very short, i.e.  $\sim 1 \text{ m}$ . Therefore,  $\beta = 1\text{--}2\%$  was used for the air content (Ervine *et al.* 1997). Consequently,  $c \cong 100 \text{ m/s}$  (Annandale 2006) and  $T_c = 0.012 \text{ s}$  for  $L = 0.6 \text{ m}$ .

In BS3D simulations, Young's modulus and Poisson's ratio of the steel plates forming the block and the mould were equal to 200 GPa and 0.3, respectively. Figure 1 compares the measured and BS3D-predicted maximum vertical block displacements. For both core and developed jets, BS3D predicts the maximum block displacement well if the hydrodynamic input pressures are equal to the net uplift dynamic water pressures measured by Federspiel *et al.* (2009). These results validate the BS3D algorithm excluding the estimation of the hydrodynamic pressure. The results of the analyses with dynamic water pressure estimated using the approach of Section 2 show that for core jets ( $Y/D_j < 6$ ) and developed jets ( $Y/D_j > 6$ ) the maximum vertical block

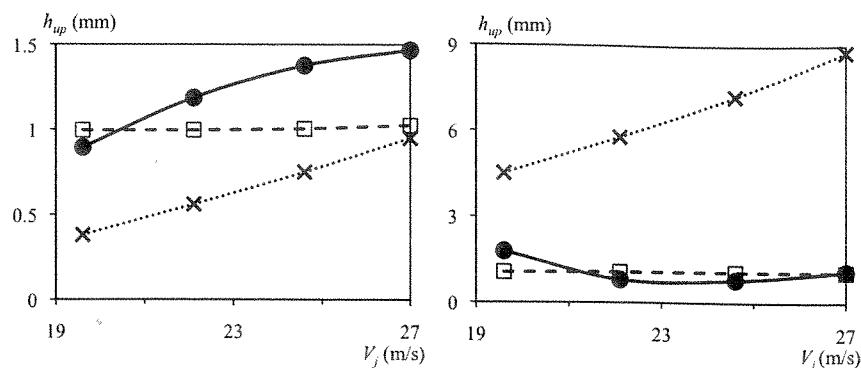


Figure 1 Comparison of maximum vertical block displacements ( $\square$ ) measured (Federspiel *et al.* 2009), predicted by BS3D using ( $\bullet$ ) measured dynamic pressures, and ( $\times$ ) dynamic pressures as in Section 2 for (a)  $Y = 0.1$  m (core jet  $Y/D_j = 1.39$ ), (b)  $Y = 0.6$  m (developed jet  $Y/D_j = 8.3$ )

displacement is underestimated and overestimated, respectively. These errors are caused by the discrepancy between the actual dynamic water pressures experienced by the block and the water pressures estimated using the approach of Section 2. However, this approach seems to be the best method currently found in the literature.

In Fig. 1, the dashed curve refers to predictions based on dynamic water pressures in Section 2. The dot-dashed curve refers to predictions based on measured water pressure. Loss of accuracy in BS3D, if any, should be quantified by comparing the dot-dashed curve with the measured displacements (continuous curve). In Fig. 1(b), the dot-dashed curve nearly collapses with the experimentally measured displacements, indicating that BS3D results are numerically accurate and, for a developed jet, the assumptions of Section 3 lead to excellent results. Since BS3D is free from numerical inaccuracies, the discrepancies in Fig. 1(a) are attributed to decreased validity of the assumptions in Section 3. In particular, the predictions are excellent for jet velocities of 20–22 m/s, and the relative error increases then to 50% as the velocity increases. In a core jet, this is attributable to fluctuations of the water velocity on top of the block during the time period  $\Delta t$ , whereas the model assumes constant pressure equal to the worst combination that causes uplift.

#### 4.2 Action of free jets on rocky river-beds

Martins (1973) describes a riverbed test facility made up of equal cubic blocks containing mortar of  $22 \text{ kN/m}^3$  unit weight, systematically arranged without cohesion. A total of 90 tests was conducted, resulting from three impact angles ( $40^\circ$ ,  $55^\circ$ , and  $70^\circ$ ), three jet diameters at the pool surface (5, 6, and 7 cm), two block widths ( $b = 3$  and 4.7 cm), and five depths of the water cushion. The free air content  $\beta$  ranged from 2% to 70% resulting in  $c \cong 25$  m/s for  $\beta \geq 30\%$  (Annandale 2006). Consequently,  $\beta = 30\%$  and  $c \cong 25$  m/s were adopted in BS3D simulations. For  $b = 3$  cm ( $L = 9$  cm),  $T_c = 8$  ms, and for  $b = 4.7$  cm ( $L = 14.1$  cm),  $T_c = 14$  ms. The dynamic pressures were estimated as described in Section 2 without considering the coefficient 2 accounting for long duration events in the numerical simulations.

The ratio of  $h_{up}/z_b$  was calculated for all 90 tests. For  $b = 3$  cm, the  $h_{up}/z_b$  ratio varied within  $[0.14, 0.55]$  with an average of 0.29, whereas for  $b = 4.7$  cm the  $h_{up}/z_b$  ratio was

within  $[0.13, 0.55]$  with an average of 0.25. Consequently, block is most likely removed from its mould if  $h_{up}/z_b > 0.25$ .

Scour depths were also predicted for all 90 tests using BS3D together with the method described in Section 2 to estimate hydrodynamic pressures caused by jet impact and the calibrated scour threshold ( $h_{up}/z_b > 0.25$ ). The ratio of predicted to measured scour depths was in the range of  $[0.44, 1.70]$  for  $b = 3.0$  cm with an average and standard deviation of 1.07 and 0.31, respectively. For  $b = 4.7$  cm, the range was  $[0.39, 1.52]$  with an average and standard deviation of 0.88 and 0.30, respectively. The results validate the ability of the suggested approach in predicting scour depth.

#### 4.3 Case study of Cabora Bassa dam

The Cabora Bassa dam is located on Zambezi River in Mozambique. Its tailwater level is at 225 m a.s.l. with a depth of near 50 m above the riverbed, whose elevation is at  $\sim 170$  m a.s.l. The rock is mainly fractured gneiss and gabbro and lamprophy dykes (Bollaert 2002, Bollaert and Schleiss 2005).

Hydraulic model tests at a 1/75 scale were conducted at Laboratório Nacional de Engenharia Civil (LNEC), Lisbon (Ramos 1982). A movable bed model was used, made of weak cemented gravel. An extensive survey of the scour hole showed that its deepest point was located at 158 m a.s.l., i.e. 22 m below the original riverbed. The outlet jet diameter was estimated as the equivalent hydraulic diameter of the  $6 \text{ m} \times 7.8 \text{ m}$  rectangular outlet, i.e. 7.7 m. The air concentration was high with  $\beta = 60\%$  resulting in  $c \cong 100$  m/s. The characteristic block length of 2 m (based on model tests performed at LNEC) yields  $T_c = 0.18$  s for  $L = 6$  m. The unit block weight was  $20 \text{ kN/m}^3$  (Bollaert 2002). Based on these values, two test series were performed using the DI method to estimate the ultimate scour depth using BS3D at either the Bollaert (2002) or the Bollaert and Schleiss's (2000) approach. The following results are compared with the measured value of 158 m a.s.l. and validate BS3D in predicting scour depth:

- Bollaert (2002) assumed  $V_j = 35$  m/s,  $D_j = 8$  m; the DI method gives an ultimate scour depth of 133 m a.s.l., whereas BS3D gives 161 m a.s.l.

- Bollaert and Schleiss (2005) assumed  $V_j = 42$  m/s,  $D_j = 7.2$  m; the DI method gives an ultimate scour depth of 152 m a.s.l., whereas BS3D's gives 155 m a.s.l.

#### 4.4 Two cases of scour in prototypes

Two published cases of scour in prototypes were analysed using BS3D and the method explained in Section 2 to estimate the hydrodynamic pressure. In each case, the maximum vertical displacements of the rock blocks at the given scour depth was calculated:

- (1) Picote prototype dam in Portugal (Cunha and Lencastre 1966): the side length of the cubic blocks was 1.05 m, the water cushion 36 m, the fall height 45 m with a discharge of  $7000$  m<sup>3</sup>/s,  $V_j = 29.71$  m/s and  $D_j = 17.32$  m. The measured scour depth was 19 m. Estimating  $\beta = 10\%$ , BS3D gave a maximum vertical block displacement of 0.311 m ( $h_{up}/z_b = 0.30$ ).
- (2) Kondopoga prototype dam, USSR (Kamenev 1965): the side length of the cubic blocks was 0.14 m, the water cushion 1.45 m, the fall height 11.4 m, with a discharge of  $70$  m<sup>3</sup>/s,  $V_j = 14.95$  m/s and  $D_j = 2.44$  m. The measured scour depth was 4.8 m. Estimating  $\beta = 2.5\%$ , BS3D gave a maximum vertical block displacement of 0.039 m ( $h_{up}/z_b = 0.28$ ).

Note that, at the reported ultimate scour depth, the ratios of  $h_{up}/z$  determined using BS3D agree with the failure criterion of Section 4.2 and the block is likely to be removed if  $h_{up}/z_b = 0.25$ .

## 5 Conclusions

The stability of single rock blocks in plunge pools was investigated. The introduced approach to estimate pressure forces generated in plunge pools due to high-velocity jet impact consists of assuming that the maximum/minimum pressures in a fracture occur at the same time as the minimum/maximum pressures on the pool floor. The block displacement was calculated by subjecting the block to its constant unbalanced force at failure applied for the maximum duration of an extreme pressure fluctuation, which was estimated using empirical equations. Experimental studies were used to validate this approach together with the ability of the BS3D method in determining the stability of rock blocks and their displacement. Based on these simulations, it was proposed to consider that a rock block is most likely to be removed if the ratio between maximum upward block displacement and the block height is larger than 1/4. The present approach was validated using three prototype case studies. The calculated scour depths match these observed, and can be obtained by using parameters determined from standard knowledge. In the cases examined (arrangements of cubic blocks), the present analyses indicate that scour and scour depth may be predicted if the failure mode is assumed to be a static translation.

## Notation

- $a$  = block acceleration caused by unbalance force (m/s<sup>2</sup>)  
 $c$  = pressure wave celerity (m/s)  
 $C$  = dynamic pressure fluctuation coefficient (dimensionless)  
 $D_j$  = jet diameter at pool surface (m)  
 $F$  = unbalanced force due to maximum dynamic pressure fluctuation (N)  
 $F_{\Delta t}$  = net impulse on block (N)  
 $g$  = gravity acceleration (m/s<sup>2</sup>)  
 $h_{up}$  = maximum upward block displacement (mm)  
 $L$  = characteristic fracture length (m)  
 $m$  = block mass (kg)  
 $P$  = dynamic pressure (N/m<sup>2</sup>)  
 $T_c$  = natural period of open-ended joint (s)  
 $V_j$  = jet velocity at pool surface (m/s)  
 $V_{\Delta t}$  = initial block velocity (m/s)  
 $x_b$  = block length (m)  
 $Y$  = pool depth (m)  
 $z_b$  = block height (m)

## Greek symbols

- $\beta$  = free air content (%)  
 $\gamma_w$  = unit weight of water (N/m<sup>3</sup>)  
 $\Delta t$  = maximum duration (s)  
 $\rho$  = density (kg/m<sup>3</sup>)

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