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Development of a global version of the gyrokinetic microturbulence code GENE

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Introduction

- The gyrokinetic code GENE [1], has been extended from its original flux-tube version to a global geometry.
- Includes radial variation of temperature and density profiles, as well as of magnetic geometry.
- Non-periodic boundary conditions allow for profile relaxation.
- Heat sources & sinks enable quasi-stationary microturbulence simulations.
- Interface with the MHD equilibrium code CHEASE [2,3].
- Various benchmarks, including comparisons with other global codes are presented.

Global GENE Model

- Field aligned coordinate system $\vec{X} = (x : \text{radial}, y: \text{binormal}, z: \text{parallel}) \Longrightarrow \vec{B}_0 = \mathcal{C}(x)\vec{\nabla}x \times \vec{\nabla}y$.
- Gyrokinetic equation with radial (x) variations of equilibrium quantities.
- Particle distribution function $f_j(\vec{X}, v_\parallel, \mu) = f_{0j} + f_{1j}$, with f_{0j} a local Maxwellian.
- Gyrokinetic equation is solved for the perturbed distribution function f_{1j}.
- Perturbed electrostatic and vector potentials $(\Phi_1, A_{1\parallel})$ are self-consistently computed through the quasineutrality (Q.N.) equation and parallel component of Ampère's law.
- Gyrokinetic ordering $|k_\parallel| \ll |k_\perp| \Longrightarrow$ Neglect $\partial/\partial z$ compared to $\partial/\partial x$ and $\partial/\partial y$.

The Gyrokinetic Equation

$$-\partial_t g_{1j} = \frac{1}{\mathcal{C}}\frac{B_0}{B^\star_{0\parallel}}\left[\frac{1}{L_{nj}} + \left(\frac{m_j v_\parallel^2}{2T_{0j}} + \frac{\mu B_0}{T_{0j}} - \frac{3}{2}\right)\frac{1}{L_{Tj}}\right] f_{0j}\,\partial_y\bar{\chi}_1 + \frac{1}{\mathcal{C}}\frac{B_0}{B^\star_{0\parallel}}\left(\partial_x\bar{\chi}_1\,\Gamma_{y,j} - \partial_y\bar{\chi}_1\,\Gamma_{x,j}\right)$$
$$+ \frac{B_0}{B^\star_{0\parallel}}\frac{\mu B_0 + m_j v_\parallel^2}{m_j\Omega_j}\left(\mathcal{K}_x\Gamma_{x,j} + \mathcal{K}_y\Gamma_{y,j}\right) - \frac{1}{\mathcal{C}}\frac{B_0}{B^\star_{0\parallel}}\frac{\mu_0 v_\parallel^2}{\Omega_j B_0}\frac{p_0}{L_p}\Gamma_{y,j} + \frac{\mathcal{C}v_\parallel}{B_0 J}\Gamma_{z,j} - \frac{\mathcal{C}\mu}{m_j B_0 J}\partial_z B_0\,\partial_{v_\parallel} f_{1j}\,,$$

- where $g_{1j} = f_{1j} + q_j v_\parallel \bar{A}_{1\parallel} f_{0j}/T_{0j}$, $\bar{\chi}_1 = \bar{\Phi}_1 - v_\parallel \bar{A}_{1\parallel}$, $\Gamma_{\alpha,j} = \partial_\alpha f_{1j} + q_j\partial_\alpha\bar{\Phi}_1 f_{0j}/T_{0j}$ for $\alpha = (x, y, z)$.
- The overbar denotes gyroaveraged quantities.
- Background density, temperature and pressure profiles: $n_{0j}(x)$, $T_{0j}(x)$, $p_0(x)$. Corresponding inverse logarithmic gradients: $L_A(x) = -(d\ln A/dx)^{-1}$ for $A = [n_j, T_j, p]$.
- $\mathcal{K}_x(x, z)$ and $\mathcal{K}_y(x, z)$ are related to curvature and gradients of $\vec{B}_0$. $J(x, z) = [(\vec{\nabla}x \times \vec{\nabla}y)\cdot\vec{\nabla}z]^{-1}$ is the Jacobian.
- $\Omega_j(x, z) = q_j B_0/m_j$, and $B^\star_{0\parallel}(x, z, v_\parallel) = B_0 + (m_j/q_j)v_\parallel(\vec{\nabla}\times\vec{b}_0)\cdot\vec{b}_0$, with $\vec{b}_0 = \vec{B}_0/B_0$.

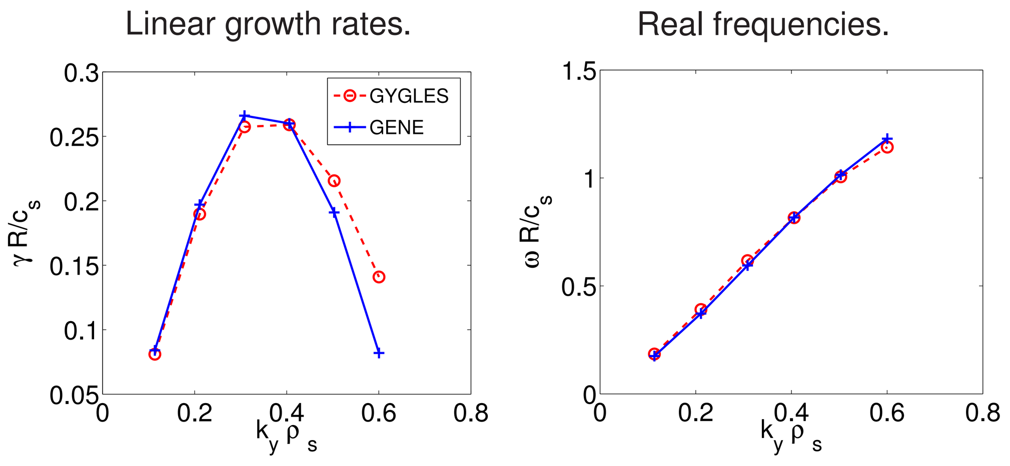
Benchmarking and Code Comparisons

Codes Used for Comparisons

- Comparison with linear and non-linear global PIC codes GYGLES [4] and ORB5 [5] based on δf scheme.
- Global GENE :
  - Solving in direct space except y-direction for which Fourier representation is used.
  - Derivatives in real space computed with finite differences.
  - Dirichlet radial boundary conditions.
  - Direct space anti-aliasing scheme in radial direction.
  - Direct space integral gyroaveraging operator in radial direction.

Linear ITG Spectra for CYCLONE Base Case [6] with Adiabatic Electrons

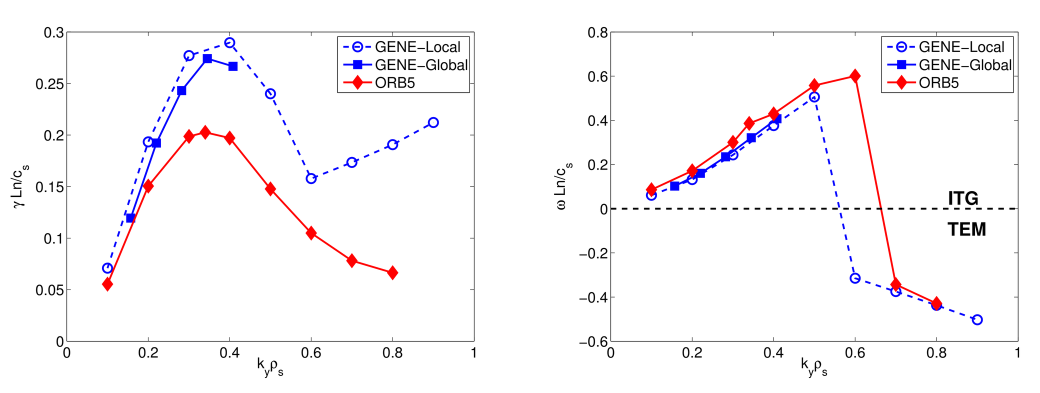
CYCLONE parameters with adiabatic electrons : $a/R = 0.36$, $\rho^\star = \rho_s/a = 1/180$, $q = 0.85 + 2.4(x/a)^2$, $T_i/T_e = 1$, peaked T and n profiles with $R/L_{Ti}(x_0) = 6.96$, $R/L_n(x_0) = 2.2$, and $x_0 = 0.5a$.



- Good agreement on growth rates and real frequencies.
- Remaining discrepancies at high k_y can be assigned to differences in the field solvers (2nd order expansion in $k_\perp\rho_s$ in GYGLES, all orders kept in GENE).

Linear ITG-TEM Spectra for CYCLONE Base Case with Kinetic Electrons

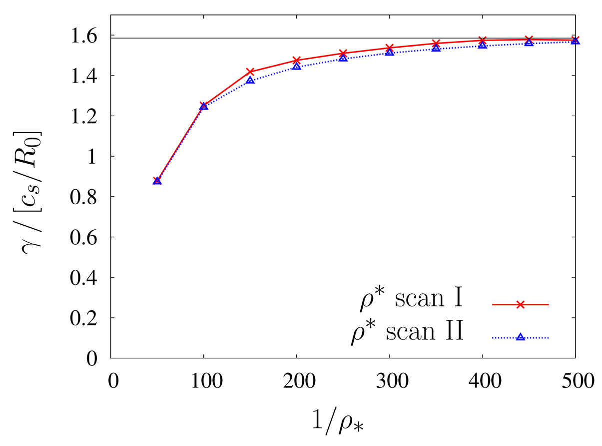
- CYCLONE parameters with kinetic electrons and mass ratio $m_i/m_e = 400$.



- Transition from ITG to TEM at higher $k_y\rho_i$.
- Differences between global GENE and ORB5 results may be related to ORB5 treating only trapped electrons kinetically (adiabatic response for passing), while GENE treats electrons fully kinetically.
- Resolution for global GENE simulations: $(320 \times 64 \times 64 \times 32)$ in the $(x, z, v_\parallel, \mu)$ directions $\Longrightarrow$ High resolutions in $(x, v_\parallel, \mu)$ required for resolving non-adiabatic response of passing electrons at mode rational surfaces.

Linear $\rho^\star$ scan with Kinetic Electrons and EM Effects

- CYCLONE-like parameters with finite $\beta = 2.5\%$.
- Simulations are carried out considering kinetic electrons (proton-electron mass ratio) and both potentials $(\Phi_1, A_{1\parallel})$.

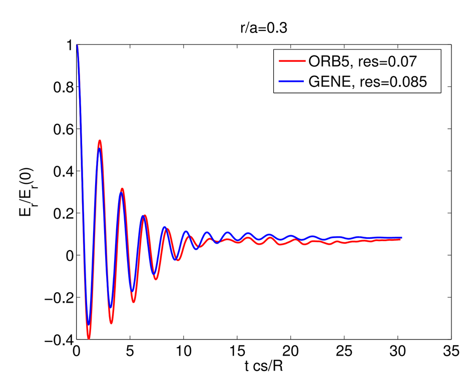


Growth rate γ for $\rho^\star$ scan at $k_y\rho_s = 0.28$. Radial width of the simulation annulus is kept fixed with respect to (I) the Larmor radius, and (II) the minor radius. Local flux-tube result in black.

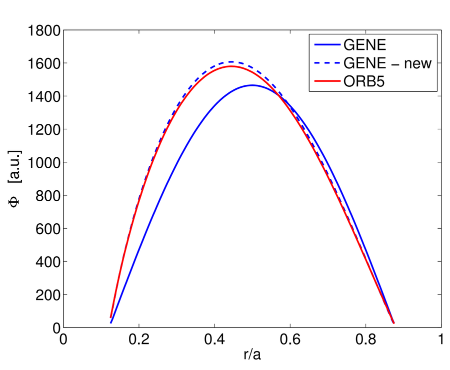
- $\beta = 2.5\% \Longrightarrow$ Kinetic ballooning modes dominate.
- Local flux-tube limit recovered by global code in limit $\rho^\star \to 0$.

Rosenbluth-Hinton Test

Parameters : $a/R = 0.1$, $\rho^\star = 1/180$, $q = 1 + 0.75(x/a)^2$, $T_i/T_e = 1$, $R/L_T = R/L_n = 0$, $f_1(t = 0) = cos(\pi x/lx)$. Adiabatic electrons.



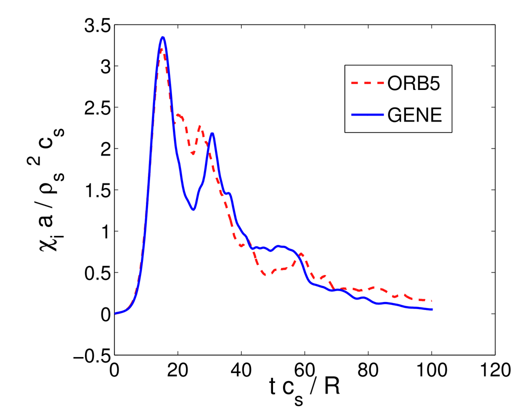
Evolution of the flux-surface-averaged, radial electric field.



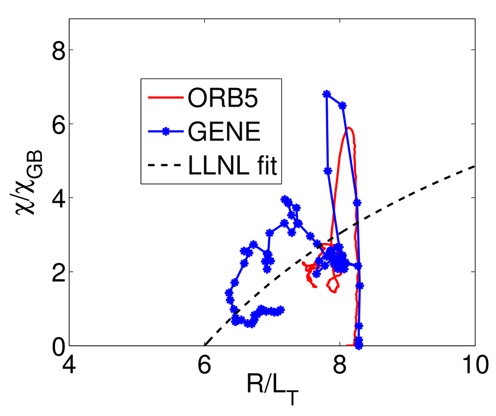
Initial electrostatic potential, solution to the Q.N. equation.

- Good agreement obtained for GAM frequency and damping rate, as well as for residual.
- Remaining discrepancies related to $\rho^\star$ approximations in GENE, in particular in the gyroaveraging appearing in Q.N. equation.
- After correcting these $\rho^\star$ approximations on gyroaveraging:
  - Very good agreement is reached on the Q.N. solution.
  - However, zonal modes become unstable! (under investigation).
- Current simulation results are thus still obtained using the uncorrected gyroaveraging operator.

Non-Linear ITG Simulations without Sources $\Longrightarrow$ Relaxation



Evolution of ion heat diffusivity χ_i for CYCLONE parameters with peaked grad. profiles.



$(R/L_T, \chi_i)$ trace for CYCLONE parameters with flat grad. profiles [7].

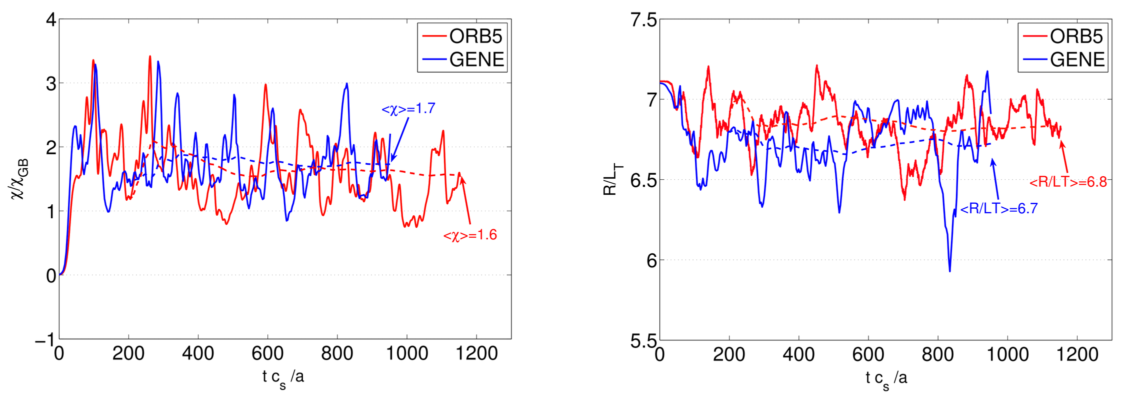
- Same initial conditions $\Longrightarrow$ Remarkable agreement: Time traces of the first burst are essentially identical.
- Global GENE recovers well the non-linear relaxation traces in the $(R/L_T, \chi_i)$ plane published in [7].

Non-Linear ITG Simulations with Sources $\Longrightarrow$ Quasi-Stationary Microturbulence

- Radially dependent heat source/sink over whole system, conserving surface-averaged density and parallel momentum:

$$\frac{df_1}{dt} = -\gamma_h\left[\langle f_1(\vec{X}, |v_\parallel|, \mu)\rangle - \langle f_0(\vec{X}, |v_\parallel|, \mu)\rangle\frac{\langle\int d\vec{v}\langle f_1(\vec{X}, |v_\parallel|, \mu)\rangle\rangle}{\langle\int d\vec{v}\langle f_0(\vec{X}, |v_\parallel|, \mu)\rangle\rangle}\right].$$

- Relaxation coefficient $\gamma_h \sim 10^{-1}\gamma_{ITG}$ $\Longrightarrow$ Background temperature profile is approximately maintained, while avoiding direct effect on microturbulence.
- CYCLONE parameters with flat grad. profiles.
- Numerical resolution: $(120 \times 48 \times 16 \times 48 \times 16)$ in the $(x, y, z, v_\parallel, \mu)$ directions.



Time evolution of (a) heat diffusivity χ_i, and (b) temperature gradient R/L_{T_i} for CYCLONE parameters with heat sources/sinks.

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