The Hilbert matrix as a singular integral operator*)

By JAMES ROVNYAK in Charlottesville (Virginia, U.S.A.)

In this note we show that the theory of the Hilbert matrix $\{1/(j+k+1)\}_{j,k=0}^{\infty}$ can be derived from the theory of selfadjoint singular integral operators, i.e. operators in $L^2(-\infty, \infty)$ of the form

(1)
$$S: F(x) \to A(x)F(x) + (\pi i)^{-1} P.V. \int_{-\infty}^{+\infty} (x-t)^{-1} \overline{K}(x)K(t)F(t) dt$$

where A(x), $K(x) \in L^{\infty}$ $(-\infty, \infty)$ and A(x) is real valued. If $f(x) \in L^2(0, \infty)$ its Mellin transform is

$$F(x) = \lim_{a \to 0} \int_{a}^{1/a} t^{-1/2 + ix} f(t) dt$$

where the limit is taken in the metric of $L^2(-\infty,\infty)$. We have then

$$2\pi f(x) = \lim_{A \to \infty} \int_{-A}^{+A} x^{-1/2 - it} F(t) dt$$

where the limit is taken in the metric of $L^2(0, \infty)$, and

$$\|F\|^2 = 2\pi \|f\|_+^2.$$

The subscript "+" indicates that the underlying Hilbert space is $L^2(0, \infty)$ rather than $L^2(-\infty, \infty)$. We also say that f(x), F(x) is a Mellin transform pair.

Theorem. Let $p(x), q(x) \in L^2(0, \infty)$ and let P(x), Q(x) be the corresponding Mellin transforms in $L^2(-\infty, \infty)$. Assume $P(x), Q(x) \in L^{\infty}(-\infty, \infty)$ and P(x) is real valued (equivalently $\overline{p}(x) = x^{-1}p(x^{-1})$). For x > 0, t > 0 set

$$L(x, t) = t^{-1}p(xt^{-1}) + 2\int_{1}^{\infty} q(t\lambda)\bar{q}(x\lambda) d\lambda.$$

(A) If $f(x) \in L^2(0, \infty)$ then

$$\int_{0}^{\infty} L(x,t)f(t) \, dt = \lim_{a \to 0} \int_{a}^{1/a} L(x,t)f(t) \, dt$$

*) Research supported by NSF Grant GP 8981.

11 A.

J. Rovnyak

exists in the metric of $L^2(0, \infty)$, and the linear operator defined in $L^2(0, \infty)$ by

$$T: f(x) \to \int_0^\infty L(x, t) f(t) dt$$

is bounded and selfadjoint.

(B) The operator T is unitarily equivalent to the singular integral operator S in $L^2(-\infty, \infty)$ defined by (1) with coefficients

$$A(x) = P(x) + |Q(x)|^2, \quad K(x) = Q(x).$$

Proof. The Hilbert transformation in $L^2(-\infty, \infty)$ is defined by

H: *F*(*x*) → (*πi*)⁻¹ P.V.
$$\int_{-\infty}^{+\infty} (x-t)^{-1} F(t) dt$$
.

Let J denote multiplication by the function

$$\sigma(x) = \begin{cases} 1, \ 0 < x < 1 \\ -1, \ 1 < x < \infty \end{cases}$$

in $L^2(0, \infty)$. For any Mellin transform pairs f(x), F(x) and g(x), G(x) we have then

$$\langle HF, G \rangle = 2\pi \langle Jf, g \rangle_+$$
.

See TITCHMARSH [8, Chapt. V].

If f(x), F(x) and g(x), G(x) are Mellin transform pairs and F(x) is essentially bounded, then

$$h(x) = \int_{0}^{\infty} t^{-1} f(xt^{-1})g(t) dt, \quad H(x) = F(x)G(x)$$

is a Mellin transform pair. See e.g. [8, p. 90].

Let f(x), F(x) and g(x), G(x) be Mellin transform pairs. It is easy to see that

$$B(f,g) = \int_{0}^{\infty} \left(\int_{0}^{\infty} t^{-1} p(xt^{-1}) f(t) dt \right) \bar{g}(x) dx + 2 \int_{1}^{\infty} \left(\int_{0}^{\infty} q(t\lambda) f(t) dt \right) \left(\int_{0}^{\infty} \bar{q}(x\lambda) \bar{g}(x) dx \right) d\lambda$$

defines a bounded symmetric bilinear form on $L^2(0, \infty)$. Therefore $B(f, g) = = \langle Tf, g \rangle_+$ for some bounded selfadjoint operator T. If f(x), g(x) vanish outside a compact subinterval of $(0, \infty)$, then

$$\langle Tf, g \rangle_+ = \int_0^\infty \left(\int_0^\infty L(x, t) f(t) \, dt \right) \bar{g}(x) \, dx,$$

the rearrangement of integrals being justified by absolute convergence. From this we deduce (A).

348

In the same notation we have

$$\begin{split} \langle Tf,g\rangle_{+} &= \int_{0}^{\infty} \int_{0}^{\infty} t^{-1} p(xt^{-1}) f(t) \bar{g}(x) \, dt \, dx + \\ &+ 2 \int_{1}^{\infty} \left(\int_{0}^{\infty} q(t\lambda) f(t) \, dt \right) \left(\int_{0}^{\infty} \bar{q}(x\lambda) \bar{g}(x) \, dx \right) d\lambda - \\ &- \int_{0}^{\infty} \left(\int_{0}^{\infty} q(t\lambda) f(t) \, dt \right) \left(\int_{0}^{\infty} \bar{q}(x\lambda) \bar{g}(x) \, dx \right) d\lambda + \\ &+ \int_{0}^{\infty} \left(\int_{0}^{\infty} q(t\lambda) f(t) \, dt \right) \left(\int_{0}^{\infty} \bar{q}(x\lambda) \bar{g}(x) \, dx \right) d\lambda = \int_{0}^{\infty} \int_{0}^{\infty} t^{-1} p(xt^{-1}) f(t) \bar{g}(x) \, dt \, dx + \\ &+ \int_{0}^{\infty} \left(\int_{0}^{\infty} q(t\lambda) f(t) \, dt \right) \left(\int_{0}^{\infty} \bar{q}(x\lambda) \bar{g}(x) \, dx \right) d\lambda - \\ &- \int_{0}^{\infty} \sigma(\lambda) \left(\int_{0}^{\infty} q(t\lambda) f(t) \, dt \right) \left(\int_{0}^{\infty} \bar{q}(x\lambda) \bar{g}(x) \, dx \right) d\lambda = \\ &= (2\pi)^{-1} \{ \langle PF, G \rangle + \langle QF, QG \rangle + \langle HQF, QG \rangle \} = (2\pi)^{-1} \langle SF, G \rangle, \end{split}$$

where S is the operator (1) with coefficients as in (B). The theorem follows.

The Hilbert matrix. If p(x) = 0, $q(x) = \frac{1}{2}e^{-x/2}$, then

 $L(x, t) = (x+t)^{-1}e^{-(x+t)/2}.$

In this case the matrix of T with respect to the orthonormal basis in $L^2(0, \infty)$ consisting of weighted Laguerre polynomials $\varphi_j(x) = L_j(x)e^{-x/2}$ (j=0, 1, 2, ...) is the Hilbert matrix $\{1/(j+k+1)\}_{j,k=0}^{\infty}$. See ROSENBLUM [4, 7]. By our theorem T is unitarily equivalent to the operator S defined by (1) with coefficients

$$K(x) = 2^{-1/2 + ix} \Gamma(\frac{1}{2} + ix), \quad A(x) = |K(x)|^2 = \frac{1}{2}\pi/\cosh \pi x.$$

From the theory of singular integral operators (PINCUS [2, 3], ROSENBLUM [6] we see at once that the spectrum of S is purely absolutely continuous, it consists of the interval $[0, \pi]$, and has uniform multiplicity 1 on the interval. In particular this gives a new proof of Hilbert's inequality

$$\left|\sum_{j,k=0}^{\infty} x_j \bar{x}_k / (j+k+1)\right| \leq \pi \sum_{j=0}^{\infty} |x_j|^2$$

and the fact that π is the best constant (see HARDY, LITTLEWOOD, and PÓLYA [1, Chapt. IX]). ROSENBLUM [5] diagonalizes the Hilbert matrix by a different method. His method has the advantage that it can be adapted to the generalized Hilbert matrix $\{1/(j+k+1-\nu)\}_{j,k=0}^{\infty}$ where ν is a real parameter which is not a positive integer.

11*

J. Rovnyak: Hilbert matrix

Miscellaneous examples. (1) If p(x) = 0, $q(x) = 2^{-\frac{1}{2}}(1+x)^{-1}$ then

$$L(x, t) = (x-t)^{-1} \log \left[(1+t^{-1})/(1+x^{-1}) \right].$$

In the theorem

$$K(x) = 2^{-1/2} \pi / \cosh \pi x$$
, $A(x) = K(x)^2 = \frac{1}{2} \pi^2 / \cosh^2 \pi x$.

So T has absolutely continuous spectrum with uniform multiplicity 1 on $[0, \pi^2]$. (2) If p(x) = 0, $q(x) = 2^{-1/2}e^{-x^2}$ then

$$L(x, t) = (x^{2} + t^{2})^{-1/2} \operatorname{erfc} \left((x^{2} + t^{2})^{1/2} \right)$$

where

$$\operatorname{erfc} x = \int_{x}^{\infty} e^{-t^2} dt, \quad x > 0.$$

We have

$$K(x) = 2^{-3/2} \Gamma(\frac{1}{4} + \frac{1}{2}ix), \quad A(x) = 2^{-3} \left| \Gamma(\frac{1}{4} + \frac{1}{2}ix) \right|^{2}$$

The operator T has absolutely continuous spectrum with multiplicity 1 on [0, c] where

$$c = \Gamma(1/4)^2/4 = 3.28626...$$

(3) If p(x) = 0, $q(x) = 2^{-\frac{1}{2}}(1+x^2)^{-1}$ then

 $L(x, t) = [x \operatorname{arc} \operatorname{cot} x - t \operatorname{arc} \operatorname{cot} t]/(x^2 - t^2),$

 $K(x) = \frac{1}{2}\pi/[\cosh{(\frac{1}{2}\pi x)} + i\sinh{(\frac{1}{2}\pi x)}], A(x) = \frac{1}{4}\pi^2/\cosh{(\pi x)}.$

The operator T has absolutely continuous spectrum with multiplicity 1 on $[0, \frac{1}{2}\pi^2]$.

References

- [1] G. H. HARDY, J. E. LITTLEWOOD, and G. PÓLYA, *Inequalities*, second edition (Cambridge, 1964).
- [2] J. D. PINCUS, Commutators, generalized eigenfunction expansions and singular integral operators, *Trans. Amer. Math. Soc.*, **121** (1966), 358–377.
- [3] A singular Riemann—Hilbert problem, Summer Institute on Spectral Theory and Statistical Mechanics, 57–92 (Brookhaven National Laboratory, 1965).
- [4] M. ROSENBLUM, On the Hilbert matrix. I, Proc. Amer. Math. Soc., 9 (1958, 137-140.
- [5] _____ On the Hilbert matrix. II, Proc. Amer. Math. Soc., 9 (1958), 581-585.
- [6] _____ A spectral theory for selfadjoint singular integral operators, Amer. J. Math., 88 (1966), 314-328.
- [7] _____ Selfadjoint Toeplitz operators, Summer Institute on Spectral Theory and Statistical Mechanics, 135–157 (Brookhaven National Laboratory, 1965).
- [8] E. C. TITCHMARSH, Theory of Fourier Integrals, second edition (London, 1948).

UNIVERSITY OF VIRGINIA

(Received September 20, 1969)