

1 Mechanisms of strain accommodation in zircon in a shear zone

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12

### 13 **Abstract**

14 Zircon grains with the evidence of crystal-plastic deformation are often found in deformed  
15 terrestrial and lunar rocks. In this study we present microstructural data of plastically-deformed  
16 zircon crystals that were analyzed in a kinematic context of the respective shear zone. The aims  
17 are to describe how orientation of zircon grains in a simple shear affects the zircon deformation  
18 mechanisms.

19 Careful microstructural analyses of zircon crystals in-situ with scanning electron backscatter  
20 diffraction (EBSD) mapping show strong geometric relationships between orientations of: (i) the

21 <c> of plastically-deformed zircon crystals, (ii) the misorientation axes of plastically-deformed  
22 zircon crystals and (iii) the XYZ directions of the kinematic framework of a sample. All  
23 plastically-deformed zircon crystals have <c> parallel to the XY (foliation plane); crystals with  
24 <c> oriented at a higher angle than ca. 15° to the XY plane are undeformed or fractured.  
25 Furthermore, zircons that have <c> aligned parallel or normal to X direction (stretching lineation)  
26 within the XY plane develop misorientation and rotation axes parallel to [001] and form the  
27 <100>{010} slip system and twist boundaries parallel to {001}. Zircons with the <c> aligned at  
28 45° with respect to X within XY plane can develop two misorientation axes, in ideal case parallel  
29 to [001] and [100], and form a combination of tilt and twist low-angle boundaries. These  
30 relationships describe the strong geometric control of the macroscopic kinematic rotation axis on  
31 the slip systems in zircon, regardless of the grain size and shape. It suggest a tool for constraining  
32 an orientation of the plastically deformed zircon with respect to a bulk kinematic rotation axis of  
33 the deformation event.

34 Relationships between zircon deformation mechanisms and macroscopic kinematic frame  
35 have important implications for zircon geochronology: if the deformation events result in zircon  
36 distortion and rejuvenation of zircon isotopic system, they may be dated. Our microstructural  
37 study suggests a criteria for the reliable sample selection for such isotopic dating.

38

## 39 **Introduction**

40 Crystal-plastic deformation in zircon has been documented for different geological settings:  
41 syn-magmatic deformation (Reddy et al., 2009; Timms and Reddy, 2009), deformation related to  
42 ductile shear zones formation (Timms et al., 2006, 2011; Reddy et al., 2006, 2007; Timms and

43 Reddy, 2009; Flowers et al., 2010; Kaczmarek et al., 2011; Piazzolo et al., 2012; MacDonald et al.,  
44 2013; Kovaleva et al., 2014), deformation related to the seismic activity (Austrheim and Corfu,  
45 2009; Kovaleva et al., 2015), impact-related lattice distortion (Leroux et al., 1999; Moser et al.,  
46 2009, 2011; Nemchin et al., 2009; Timms et al., 2012; Grange et al., 2013). Most often crystal-  
47 plastic deformation of zircon is documented in the ductile shear zones and is thus believed to be  
48 genetically connected with the deformation event (e.g. Kaczmarek et al., 2011; Piazzolo et al.,  
49 2012; Kovaleva et al., 2014).

50         Microstructural analyses of naturally-deformed crystals show that misorientation axes  
51 associated with low-angle boundaries have a certain geometric relationship with the macroscopic  
52 kinematic rotation axis. Geometric control of the kinematic framework of the shear zone on the  
53 activity of crystal slip systems has been described for calcite (Bestmann and Prior, 2002; Reddy  
54 and Buchan, 2005), quartz (Menegon et al., 2011) and sillimanite (Piazzolo and Jaconelli, 2013).  
55 For the simple shear deformation, it has been demonstrated that the rotation axes in plastically-  
56 deformed crystals of calcite and sillimanite are parallel to the kinematic rotation axis of the sample  
57 (Reddy and Buchan, 2005; Piazzolo and Jaconelli, 2013). The assumption that zircon should  
58 behave similar to that was inherited by the model of Kaczmarek et al. (2011) describing plastic  
59 deformation of zircon in a simple shear framework. However, the misorientation axes of zircons in  
60 their sample lie in different orientations and are independent of the deformation framework, which  
61 is explained by a post-deformation rigid body rotation of zircon crystals. Thus it was not possible  
62 to reconstruct the initial orientation of grains with respect to the kinematic rotation axes, neither to  
63 find the corresponding relationships. Kovaleva et al. (2014) have demonstrated crystal-plastically  
64 deformed zircon grains with  $\langle c \rangle$  parallel to the stretching lineation and with  $\langle 100 \rangle \{010\}$  active

65 slip systems, and suggested that this crystal-plastic deformation mechanism was induced by the  
66 specific orientation of the zircons crystals in the shear zone.

67 In our study we present plastically-deformed zircons from the high-strained rock and  
68 confirm that the slip systems geometry is strongly controlled by the orientation of zircon crystals  
69 in the simple shear. We describe the zircon deformation evolution in a shear zone, explain its  
70 mechanisms, and emphasize the importance of studying crystals in situ, taking into account their  
71 orientation to the XYZ framework. Such studies allow to constrain a genetic link of the  
72 deformation microstructures in zircon with the deformation event; and have the implications for  
73 mineral isotopic geochronology. Plastically-deformed zircon grains could be available for direct  
74 dating of the corresponding deformation event.

75

#### 76 **Sampling locality and sample description**

77 Sampling took place in the Western Tauern Window, Eastern Alps (Zillertal, Tyrol,  
78 Austria). In the Tauern Window, continental and oceanic rocks of the Penninic and sub-Penninic  
79 nappe sequences are exposed, which represent the footwall of the Austroalpine nappe stack.  
80 Nappe stacking and predominant final metamorphism are related to the closure of the Alpine  
81 Neotethys and subsequent continental collision in late Cretaceous-Tertiary (Miller et al., 2007).  
82 Samples were collected from the “Zillertaler Kern” lobe of the “Zentralgneis” formation (see  
83 Selverstone et al., 1991 and reference therein).

84 The magmatic protholiths of the "Zentralgneis" formation are uppermost Devonian to lower  
85 Permian in age. Three magmatic “pulses” of potassium-rich and calc-alkaline granites, felsic and

86 intermediate volcanites and tonalitic/granodioritic plutonites can be distinguished (Veselá et al.,  
87 2011). The granitoids intruded into pre-Carboniferous, partly poly-metamorphic basement rocks  
88 consisting of various schists, para- and orthogneisses, amphibolites and meta-ophiolites. In the  
89 Zillertal section Variscan amphibolite facies regional metamorphism has been overprinted at  
90 greenschist- to amphibolite-facies metamorphic conditions of 0.5-0.7 GPa and 550-600 °C at ca.  
91 30 Ma and re-equilibrated the rocks (Selverstone, 1991; Pennacchioni and Mancktelow, 2007).  
92 Metamorphic (re)crystallization was accompanied by the formation of ductile extensional shear  
93 zones that represent a pure strike slip systems (Pennacchioni and Mancktelow, 2007). The  
94 sampled shear zone (Fig. 1A) formed during this latter tectono-metamorphic event and the rocks  
95 were deformed under greenschist- to amphibolite-facies metamorphism under a simple shear with  
96 a minimum of pure shear component.

97         We have sampled the ~1 m thick ductile shear zone (47°01'18.129"N/11°50'26.709"E),  
98 which is exposed on the NE slope of the Zemmbach river side valley and represents strongly  
99 foliated quartz-biotite orthogneiss that hosts two adjacent deformed dykes, a leucocratic aplitic  
100 dyke and a highly-deformed melanocratic, presumably meta-lamprophyric dyke in the core of the  
101 shear zone (Fig. 1A), which was sampled. Analyzed rock contains indicators of the well-  
102 developed kinematic framework including foliation, stretching lineation, sigma-clasts and folding  
103 of compositional layering (Fig. 1B); the folds axis coincides with the stretching lineation.  
104 Foliation plane is used here as the kinematic reference plane (Bestmann and Prior, 2002). The  
105 sample is chlorite-rich (>50% modal content) due to intense retrogression of biotite; it is  
106 composed of mafic layers consisting of chlorite, biotite and titanite, alternating with the fine-  
107 grained plagioclase-quartz layers with an average thickness of 1-6 mm; the compositional layering

108 is isoclinally folded (Fig. 1B). The accessory minerals are calcite, pyrite, zircon, rutile and titanite.  
109 Zircon is present as small (10-30  $\mu\text{m}$ ) euhedral crystals (Figs. 2, 3) with degraded CL-zonation,  
110 with the 1:1 to 1:3 aspect ratio, and are mostly hosted by biotite or chlorite.

111

## 112 **Methodology and data representation**

### 113 *Sample preparation*

114 Thin sections were cut at an acute angle of about  $45^\circ$  to the stretching lineation, in order to  
115 fully document all relevant deformation microstructures in plastically-deformed zircons. I.e.  
116 sample cut parallel to the lineation may not reveal microstructures and subgrain walls that are  
117 normal to the kinematic rotation axis of the sample, and thus it would become impossible to  
118 document the corresponding geometric relationships. More specific, if we assume that the  
119 suggestion given in Kaczmarek et al. (2011) that the rotation axis in the deformed zircon grain is  
120 parallel to the kinematic rotation axis of the sample, there should be twist subgrain walls that are  
121 by definition normal to the rotation axis. These walls will be not revealed by a classical sample cut  
122 normal to the kinematic rotation axis. Therefore, the samples were cut normal to the XY plane but  
123 at  $45^\circ$  to the XZ plane.

124 Zircons were examined in-situ using polished thin sections, mechanically prepared with 0.25  
125  $\mu\text{m}$  diamond paste and subsequently chemically polished with alkaline colloidal silica solution  
126 (Köstrosol 3530; pH 9.2-10) on an active rotary head polishing machine for 4 hours.

127

### 128 *Electron Backscatter Diffraction (EBSD) analysis and Forescattered electron (FSE) imaging*

129 Zircon crystals were examined for potential crystal-plastic deformation structures using  
130 orientation contrast images (Trimby and Prior, 1999; Prior et al., 1999). These were taken using a  
131 forescatter-electron detector (FSD) mounted on the EBSD-tube of an FEI Quanta 3D FEG  
132 instrument (Center of Earth Sciences, University of Vienna, Austria), which is equipped with a  
133 Schottky field emission electron source. Electron beam conditions were 15 kV accelerating  
134 voltage, 2.5-4 nA probe current using the analytic mode. Stage settings were at 70° tilt and 14-16  
135 mm working distance. After identification of the potentially deformed crystals, EBSD orientation  
136 mapping was applied to selected zircon crystals. The FEI Quanta 3D FEG instrument is equipped  
137 with an EDAX Pegasus Apex 4 system consisting of a Digiview IV EBSD camera and an Apollo  
138 XV silicon drift detector for EDX analysis. EDX intensities and EBSD data were collected  
139 contemporaneously using the OIM data collection software v6.21. An EBSD camera binning of  
140 4x4 was used at exposure times of 50 - 130 milliseconds. As zircon yielded a significantly higher  
141 EBSD signal intensity than the matrix phases (used for the background calibration), the EBSD  
142 camera exposure time was significantly reduced after background collection in order to avoid  
143 signal oversaturation during zircon analysis. Therefore, the matrix phases collected at these  
144 settings yielded a very weak pattern contrast. Hough parameters were set to a binned pattern size  
145 of 9x9 pixels, a Theta step size of 1° and a Rho-fraction of 74-86%. After applying a 9x9  
146 convolution mask 3 - 15 bands with a minimum pattern contrast of 200 and at a minimum peak  
147 distance of 3-10 pixels in Hough space were used for indexing. At the given settings indexing  
148 rates were between 6 and 24 points per second. Orientation maps were obtained from the beam  
149 scanning in hexagonal grid mode at step sizes of 0.1 – 0.16 micrometer.

150 The raw indexing for zircon phase shows a very good quality of more than 99.99%. In some  
151 cases, after EBSD data collection, the maps were recalculated based on chemical composition of  
152 phases with the OIM v6.21 software.

153

#### 154 *EBSD data representation*

155 The EBSD data are represented in the sample reference frame as false color-coded  
156 cumulative misorientation maps, with colors showing the relative angular misorientation of each  
157 data point with respect to a user-selected single reference point within the crystal (indicated by a  
158 white star marker; Fig. 2). Another mode is so-called “local misorientation” EBSD maps (Fig.  
159 3A), where each pixel is color-coded according to the mean misorientation of the respective data  
160 point relative to its neighboring points. The orientations of the crystallographic axes are plotted as  
161 lower hemisphere equal area projections and are color-coded according to the corresponding  
162 EBSD map (Figs. 2, 3B). Cumulative EBSD maps and pole figures were produced using the  
163 EDAX OIM v6.2.1 Analysis software, whereas the local misorientation EBSD maps together with  
164 the visualizing of misorientation axis orientation and density contours in the inverse pole figures  
165 were generated with the MTEX toolbox for MATLAB (Bachmann et al., 2010, 2011; Mainprice et  
166 al., 2011). Presented misorientation axes were calculated with a threshold starting from 1° of  
167 misorientation, however for the small misorientations of 1-2° large error is possible (Prior et al.,  
168 1999; Reddy and Buchan, 2005; Reddy et al., 2007). By the term “misorientation axis” we imply  
169 the crystallographic direction around which two subgrains of the same grain are rotated with  
170 respect to each other. By the term “rotation axis” we mean the least dispersed crystallographic  
171 axis of the deformed grain in the pole figure, around which all the other grain’s axes are rotating.



172

173 *Weighed Burgers Vector (WBV) calculations*

174 We also gained insight into geometrically necessary dislocation densities using Weighted  
175 Burgers Vector (WBV) calculations (Wheeler et al., 2009; 2012). WBV quantifies the total  
176 Burgers vector for all the dislocations passing through the user-selected rectangular region in the  
177 EBSD map (the “integral form”, according to Wheeler et al., 2009). This can be expressed in  
178 terms of lattice vectors and then divided by the sample region area to measure dislocation density  
179 including Burgers vector direction. The three numbers, listed for each selected subarea, are the *a*,  
180 *b* and *c* components of WBV, measured in  $(\mu\text{m})^{-2}$ . Rectangular areas with three WBV components  
181 were calculated over the EBSD maps with the MATLAB toolbox CrystalScape 1.3 based on the  
182 method described in Wheeler et al. (2009, 2012). For this goal the maps were transformed to a  
183 rectangular grid and the Euler angles were recalculated accordingly with the Channel software  
184 (method described in Kovaleva et al., 2015). However, for the better visual representation, the  
185 rectangular areas were superimposed on top of hexagonal “local misorientation” EBSD map.

186 A non-homogeneous distribution of WBV values in analyzed zircon grain (Fig. 3A)  
187 indicates that the plastic deformation in zircon crystal was a post-growth process, and it was  
188 caused by directed external differential stress (e.g. MacDonald et al., 2013).

189

190 **Microstructural and crystallographic zircon data**

191 Zircons in the sample are mostly decoupled from the host matrix, which is indicated by  
192 normal and reverse drag of the surrounding biotite and by the open voids parallel to the grains

193 faces (see Kovaleva et al., 2014). Decoupling allows the inhomogeneous distribution of stress  
194 within the grains hosted by a rheologically softer material (e.g. see numerical modelling by  
195 Schmid and Podladchikov, 2005) and may lead to their crystal-plastic deformation (Kenkmann,  
196 2000; Kovaleva et al., 2014). Microstructural data for the zircons from a sampled meta-  
197 lamprophyre demonstrate that some of the grains are crystal-plastically deformed (Figs. 2, 3; see  
198 also Kovaleva et al., 2014, Figs. 4, 5 and 9 there); the amount of plastically deformed grains varies  
199 from 27 to 38% of all zircon grains in the sample.

200         Finite deformation pattern of the deformed zircon grains is characterized by the presence of  
201 the strain-free subgrains separated by low-angle boundaries and rotated with respect to each other;  
202 total misorientation of the subgrains with respect to each other ranges from 3 to 10° (Fig. 2, 3A).  
203 Low-angle boundaries traces represent a continuous network of step- or zigzag-shaped lines;  
204 subgrains represent angular, irregular-shaped domains from 1 to 10 µm in size and with ragged  
205 boundaries (Figs. 2, 3A).

206         Zircons in the analyzed sample generally do not indicate crystallographic preferred  
207 orientation and oriented rather randomly (Fig. 4). However, some of the grains that are elongate  
208 and sit in the flanks of the compositional folds (e.g. grains 03 and 04, locations marked in Fig. 1B)  
209 demonstrate crystallographic preferred orientation. C-axes of those grains are parallel to the fold  
210 axes and to the stretching lineation, which may indicate the syn-deformation rigid body rotation of  
211 zircon grains in the areas of the highest differential stress (e.g. Jeffery, 1922; Mancktelow et al.,  
212 2002). Zircon grains that are not crystal-plastically deformed have  $\langle c \rangle$  at an angle  $\geq 15^\circ$  to the  
213 foliation plane (Fig. 4, open circles). All plastically-deformed crystals from this sample have  $\langle c \rangle$   
214 roughly aligned in the foliation plane (Figs. 2, 3B, 4, solid circles; foliation is subhorizontal).

215 Crystallographic orientations of  $\langle c \rangle$  of the plastically-deformed zircons 03 and 04 coincide with  
216 the orientation of the stretching lineation (Figs. 2A, 3B); grain 15 is oriented with  $\langle c \rangle$  at  
217 approximately  $45^\circ$  to the lineation (Fig. 2B); grain 24 has  $\langle c \rangle$  oriented at approximately  $90^\circ$  to the  
218 lineation (Fig. 2C).

219 Presented crystal-plastically deformed zircons have misorientation axes clustering around  
220 [001] crystallographic direction (Figs. 2, 3A inset). Grain 15 has additional strong cluster of  
221 misorientation axes around [100] (Fig. 2B). Misorientation axes of the grain 24, besides strong  
222 clustering around [001], are dispersed in the pole figure forming a few smaller maxima. Grains 03,  
223 24 and 04 reveal rotation of the crystallographic axes around the [001] direction in the pole figures  
224 (Figs. 2A-B, 3B). Crystallographic axes in grain 15 are dispersed in a star shape (so-called  
225 “asterism” of crystallographic axes, Moser et al., 2009).

226

## 227 **Discussion**

### 228 *Analyses of the active slip systems*

229 Presence of the low-angle boundaries gives an opportunity to determine active slip system(s)  
230 of dislocations that may be associated with the certain low angle boundary. Low-angle boundaries  
231 are the result of crystal recovery, when geometrically necessary dislocations that are formed to  
232 accommodate lattice strain assemble together by dislocation creep and form dislocation walls. The  
233 geometry of low-angle boundaries, therefore, reflect the slip systems geometry that can be  
234 reconstructed (e.g. Reddy et al., 2007; Kaczmarek et al., 2011; Timms et al., 2012). However, we  
235 should emphasize that it is not always possible to unequivocally determine what slip system is

236 related to a subgrain boundary, providing only two-dimensional analysis. Therefore, we suggest  
237 slip systems that are consistent with a specific boundary.

238 Grain 03 accommodates a lot of strain by developing a dense network of low-angle  
239 boundaries that cause the total misorientation of subgrains of about  $10^\circ$  with respect to each other  
240 (Fig. 2A). The traces of low-angle boundaries form zigzag lines with a complicated geometry. We  
241 suggest that there are several slip systems operating and changing each other along the low-angle  
242 boundaries. All these slip systems have rotation axes around [001] (Fig. 2A, pole figure), which is  
243 the most frequent and energetically-preferable rotation axis in zircon (e.g. Leroux et al., 1999;  
244 Reddy et al., 2007; Kovaleva et al., 2014). According to their complicated geometry, subgrain  
245 boundaries in this grain probably represent both tilt and twist walls (see, for example, analyses of  
246 subgrain boundary traces in Reddy et al., 2007).

247 In grain 15 low-angle boundaries are stretching in two general directions and likely occupy  
248 planes (100) and (010) (Fig. 2B, “LAB-1” and “LAB-2” accordingly), and misorientation axes are  
249 directed parallel to [001] and  $\langle 100 \rangle$  crystallographic directions. This imply several possibilities  
250 for active slip systems:

- 251 a. If the rotation axis for low-angle boundary 1 is [001], then it is a tilt  
252 boundary and the slip system should be [100](010).
- 253 b. If the rotation axis for low-angle boundary 1 is [100], it’s a twist boundary.
- 254 c. If the rotation axis for low-angle boundary 1 is [010], it’s a tilt boundary  
255 with the slip system [100](001).
- 256 d. If the rotation axis for the low-angle boundary 2 is [001], it’s a tilt boundary  
257 with the slip system [010](100).

258 e. If the rotation axis for the low-angle boundary 2 is [010], it's a twist  
259 boundary.

260 f. If the rotation axis for the low-angle boundary 2 is [100], it's a tilt boundary  
261 with the slip system [010](001).

262 Taking into account that low-angle boundary traces in grain 15 do not follow straight lines  
263 and rather represent zigzags, the low-angle boundary network most likely represents a system of  
264 interconnected tilt and twist walls with the slip geometries  $\langle 100 \rangle \{010\}$  and  $\langle 100 \rangle \{001\}$ .  
265 Asterism of the crystallographic axes also evidence that there are several slip systems operating at  
266 the same time (Moser et al., 2009).

267 In grain 24 (Fig. 2C) one of the low-angle boundaries ("LAB-1") is parallel to (100)  
268 crystallographic plane. Misorientation axis parallel to [001] implies slip system [100]{010}  
269 operating in this crystal. The other low-angle boundary, tracing in the lower part of the grain from  
270 left to right, has a complicated geometry and most likely represents a result of combination of  
271 several slip systems. Minor clusters of misorientation axes around high Miller-indices directions  
272 (Fig. 2C) are either a result of analytical error due to low misorientation angles (e.g. Prior, 1999;  
273 Reddy and Buchan, 2005); or evidence of minor slips activating to accommodate strain and to  
274 connect main tilt and twist walls that have rotation around [001]. The latter is also supported by a  
275 minor asterism of crystallographic axes in the pole figure (Fig. 2C).

276 In grain 04, traces of low-angle boundary 1 (Fig. 3A, "LAB-1") are parallel to (100) plane  
277 (Fig. 3B). Misorientation axes in this grain are coinciding with rotation axis and are parallel to  
278 [001] (Fig. 3A, inset). Such geometry implies a tilt boundary correlated with the [100]{010} slip  
279 system. Low-angle boundary 2 (Fig. 3A, "LAB-2") seems to be parallel to the (110) plane, which

280 implies a tilt boundary with the slip system  $\langle -110 \rangle \{110\}$ . The conclusion is supported by the  
281 Weighted Burgers Vector (WBV) calculations (Wheeler et al., 2009, 2012).

282 WBV is measured in  $[\mu\text{m}^{-2}]$  and expressed as three numbers,  $a$ ,  $b$  and  $c$ , which are the  
283 components of summary Burgers vector across the user-selected selected rectangular area. In the  
284 Figure 3A several selected areas with corresponding WBVs values are presented. In the lower left  
285 portion of the crystal the WBV values are low (Fig. 3A, gray rectangles), that means that the  
286 dislocation density is low and the analyzed domain is strain-free. In the lower right portion WBV  
287 is dominated by  $a$  component (dotted rectangle) which implies dislocations with  $[100]$  slip  
288 direction (slip system  $\langle 100 \rangle \{010\}$ ). WBV analysis across low-angle boundary 1 (“LAB-1”)  
289 shows domination of  $b$  component (dashed rectangles) that implies dislocations with slip along  
290  $[010]$  direction (slip system  $\langle 010 \rangle \{100\}$ ). WBV analyses across low-angle boundary 2 (“LAB-  
291 2”) shows high values for  $a$  and  $b$  components, implying slip system  $\langle -110 \rangle \{110\}$  with slip along  
292  $[-110]$  plane (Fig. 3).

293 In naturally-deformed crystals it is difficult to find pure edge or screw dislocations. Almost  
294 all dislocations have both screw and edge components to them and thus are called mixed  
295 dislocations (Poirier, 1985). Mixed dislocations make up most of the dislocations encountered in  
296 natural samples and they build networks of low-angle boundaries. To summarize our data,  
297 subgrain walls in analyzed zircons appear as zigzag lines that reflects an interplay between tilt  
298 dislocations with the slip systems  $\langle 100 \rangle \{010\}$  and  $\langle -110 \rangle \{110\}$  and rotation axis  $[001]$  and the  
299 twist dislocations with the rotation axis  $[001]$ . In one case we also observed the possible presence  
300 of tilt dislocations with the slip system  $\langle 100 \rangle \{001\}$  and rotation axis  $[010]$ . In order to  
301 accommodate external strain generated in a shear zone, zircon grains develop low-angle boundary

302 network with switching of the glide directions. Rotation axis [001] and slip systems  $\langle 100 \rangle \{010\}$   
303 and  $\langle 100 \rangle \{001\}$  are the most frequently observed in plastically-deformed zircons (e.g. Leroux,  
304 1999; Reddy et al., 2007; Kovaleva et al., 2014, 2015), and  $\langle -110 \rangle \{110\}$  can be regarded as a  
305 newly described slip system (Fig. 3B, “LAB-2”).

306 Slip along  $\langle 100 \rangle \{010\}$  with rotation around [001] is the most energetically preferable  
307 geometry of a zircon and is easily activated when zircon has a specific orientation: normal or  
308 parallel to the kinematic rotation axis of a shear zone (Figs. 4; 5, cases i and iii). Along the  $\langle c \rangle$   
309 zircon atomic structure consists of chains of alternating edge-sharing  $\text{SiO}_4$  tetrahedra and  $\text{ZrO}_8$   
310 dodecahedra that are joined laterally by edge-sharing dodecahedra (Robinson et al., 1971; Finch  
311 and Hanchar, 2003). In order to develop a slip in the zircon crystal lattice it is easier to break the  
312 bonds between  $\text{SiO}_4$  tetrahedra and  $\text{ZrO}_8$  dodecahedra along  $\langle 100 \rangle$ , than between the strongly  
313 bonded  $\text{ZrO}_8$  dodecahedra along [001]. However, if long axis of zircon has a complicated  
314 orientation with respect to the kinematic rotation axis, like in case of grain 15 (Fig. 2B) which  $\langle c \rangle$   
315 is oriented at  $45^\circ$  to it, crystal lattice develops two orthogonal misorientation axes in order to  
316 accommodate the strain (Figs. 4; 5, case ii).

317 Selective plastic deformation of zircon grains that are aligned within XY plane (Fig. 4)  
318 could be explained with the critical resolved shear stress (CRSS) that is more easily reached along  
319 the specific planes, if their orientation is favorable (Hobbs, 1985). The preferable slip along [100]  
320 or [001] planes in zircon crystallographic structure is facilitated by the specific crystallographic  
321 orientation with respect to a local stress field.

322

323 *Geometric regularities derived from the observations of natural samples*

324 In the conditions of the simple-shear, kinematic rotation axis is assumed to be orthogonal to  
325 X-Z plane of the sample, parallel to foliation plane and normal to lineation direction (Reddy and  
326 Buchan, 2005). Schematic sketch revealing the character of zircon deformation in a shear zone is  
327 presented in Fig. 5.

328 The following regularities can be derived from the presented zircons deformed in a simple  
329 shear, regardless of their crystal shape and aspect ratio:

330 **a.** Zircon crystals are only plastically-deformed if their  $\langle c \rangle$  aligned in the XY  
331 (foliation) plane (Fig. 4, solid circles). Zircon crystals with c-axis oblique to the XY (foliation)  
332 plane ( $>15^\circ$ ) either not deformed or fractured (Fig. 4, open circles).

333 **b.** Zircon crystals aligned to XY with  $\langle c \rangle$  parallel or normal to X (stretching  
334 lineation), i.e. normal or parallel to the kinematic rotation axis, mostly develop misorientation and  
335 rotation axes [001] (Fig. 4, grains 03, 04, 24, and closely oriented).

336 **c.** Zircon crystals with  $\langle c \rangle$  aligned in XY at  $45^\circ$  to X develop two misorientation axes  
337 (Fig. 4), which in case of grain 15 are [100] and [001].

338 **d.** Strain in plastically-deformed zircons is accommodated by formation of a continuous  
339 network of low-angle boundaries, which are the result of combination of tilt and twist dislocations  
340 with  $\langle 100 \rangle \{010\}$ ,  $\langle 100 \rangle \{001\}$ ,  $\langle -110 \rangle \{110\}$ , etc. slip geometry.

341 The reconstruction of the kinematic rotation axis of a deformation event that formed a shear  
342 zone is important for understanding the tectonic evolution of geological units in the Earth's crust.  
343 Based on natural data we have demonstrated a strong kinematic control on the geometry of the  
344 dominant slip systems in deformed zircon minerals.



345 Using the derived regularities (a)-(c) it is possible to make the reconstruction of a  
346 macroscopic tectonic frame that was causing zircon crystal-plastic deformation, and vice versa,  
347 the crystallographic orientation of zircon grain within. This gives an opportunity to attribute  
348 crystal-plastic deformation of zircon to a specific metamorphic/deformation event and thus,  
349 possibly, to derive the timing and  $P$ - $T$  conditions of the latter, even for the detrital and inherited  
350 grains. Because different lattice distortion patterns in zircon are usually restricted to a specific  
351 stress-strain conditions (e.g. Kovaleva et al., 2014), therefore, they may represent a certain  
352 snapshot of the potentially complex deformation history. For example, lattice distortion pattern in  
353 analyzed grains (interconnected network of low-angle boundaries that separate strain-free  
354 subgrains) is attributed to an upper-greenschist to amphibolite facies of metamorphism. Thus,  
355 using these patterns, we can judge about the peak metamorphic conditions of the respective shear  
356 zone. The fact that some of zircon crystals indicate crystal preferred orientation points to the rigid  
357 body rotation *prior* crystal-plastic deformation, which reflects some earlier stages of rock-  
358 forming/deformation process.

359

#### 360 *Implications for zircon geochronology and further research trends*

361 Our observations have important implications for zircon geochronology. It has been shown  
362 by a number of authors that crystal-plastic deformation can dramatically affect the content of trace  
363 elements in the domains of subjected zircon lattice and, therefore, can cause isotopic system  
364 resetting (e.g. Flowers et al., 2010; MacDonald et al., 2013; Moser et al., 2009, 2011; Piazzolo et  
365 al., 2012; Reddy et al., 2006, 2007, 2009; Reddy and Timms, 2010; Timms and Reddy, 2009;  
366 Timms et al., 2011, 2012). With the help of plastically-deformed zircon it is possible to resolve the

367 ages of such deformation events as shearing, seismic activity and meteorite impacts in terrestrial  
368 and even lunar rocks (e.g. Austrheim and Corfu, 2009; Moser et al., 2009, 2011; Nemchin et al.,  
369 2009). Based on our study, we suggest the following criteria of sample selection for the zircon  
370 mineral dating. In accordance with our regularity (a), crystals, aligned with the foliation plane of a  
371 shear zone are likely to be plastically-deformed, and thus might be (partially)rejuvenated; crystals  
372 that are at a high angle to foliation plane are unlikely to be deformed, and thus are not rejuvenated  
373 and preserve older ages. These crystal-orientation systematics potentially allow determining the  
374 age of the deformation event using in situ dating of zircons that are aligned in the foliation plane.  
375 On the other hand, if the goal is the mineral isotopic dating of the undistorted age (e.g. age of  
376 zircon crystallization), one should strictly avoid zircon grains that are aligned in the foliation  
377 plane, in order to exclude any possible age disturbances. The detailed and careful isotopic study in  
378 order to test this speculative suggestion is still waiting for its time.

379

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384

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503

#### 504 **Figure captions**

505 Figure 1. A. Field photograph of the sampled outcrop: sample of the meta-lamprophyric  
506 dyke hosted by granitic gneiss and adjacent to the deformed aplitic dyke, highly deformed and  
507 strained. B. Thin section photographs, foliation, compositional layering and folds are clearly  
508 visible. Numbers indicate positions of the in-situ studied zircon grains, presented in this paper.



509 Figure 2. Deformed zircon crystals: microstructural data. Left column: cumulative EBSD  
510 maps, misorientation of every pixel is shown with respect to the user-selected reference point,  
511 marked as a white star; middle column: misorientation axes distribution plots, numbers are  
512 crystallographic directions; right column: pole figures with lower hemisphere projections of zircon  
513 crystallographic directions, color coded as corresponding EBSD map, labels indicate main  
514 crystallographic axes, black lines are reconstruction of the subgrain boundary planes. Foliation  
515 plane in figures is subhorizontal, stretching lineation is oriented at about  $45^\circ$  to the image plane  
516 (see description in the text). Direction of the lineation is highlighted by a red circle.

517 Figure 3. A. Local misorientation map of the deformed grain 04, with the superimposed  
518 Weighed Burgers Vector (WBV) components for the highlighted rectangular subareas. Gray  
519 rectangles show the areas with WBV that is comparatively low. Dotted and dashed rectangles  
520 show areas with WBV dominated by *a* or *b* components accordingly, black rectangles show areas  
521 with WBV with mixed components. In the lower right inset distribution of the misorientation axes  
522 for grain 04 is shown. B. Reconstruction of the low-angle boundaries and slip systems of the grain  
523 04. Thick lines outside the circle indicate the direction of the low angle boundary traces; solid  
524 lines – reconstruction of low angle boundary planes; dashed lines – reconstruction of slip plane for  
525 the low-angle boundaries pointed in A. In black are the elements that correspond to the low angle  
526 boundary 1 (“LAB-1”), in gray – to low angle boundary 2 (“LAB-2”). Small circle highlights  
527 rotation and misorientation axis.

528 Figure 4. The pole figure with  $\langle c \rangle$  (c-axis) positions of analyzed zircon grains from the  
529 studied sample. Labels in angle brackets indicate misorientation axes for the corresponding grains.  
530 Gray dashed line shows direction of foliation, gray dashed circle shows direction of lineation.  
531 Solid black circles correspond to  $\langle c \rangle$  directions of grains that are plastically-deformed, empty

532 circles – to undeformed or fractured zircon grains. From the pole figure is clear that the  $\langle c \rangle$  of all  
533 plastically deformed grains are roughly aligned with the foliation plane.

534 Figure 5. Schematic sketch showing zircon deformation evolution that strongly depends on  
535 its orientation in the macroscopic kinematic frame. Cases (i), (ii) and (iii) indicate plastic  
536 deformation; (i)  $\langle c \rangle$  is parallel to the macroscopic kinematic rotation axis, misorientation axis is  
537 parallel to [001]. (ii): grain with  $\langle c \rangle$  at an angle  $45^\circ$  to the kinematic rotation axis develops two  
538 rotation axes, ideally [001] and [100]. (iii)  $\langle c \rangle$  is normal to the kinematic rotation axis, rotation  
539 axis is parallel to [001]. (iv) – grains with c-axes at a high angle to the foliation develop fractures  
540 or are not deformed.









