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MEAN ANISOTROPY OF HOMOGENEOUS GAUSSIAN RANDOM FIELDS AND ANISOTROPIC NORMS OF LINEAR TRANSLATION-INVARIANT OPERATORS ON MULTIDIMENSIONAL INTEGER LATTICES¹

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Sensitivity of output of a linear operator to its input can be quantified in various ways. In Control Theory, the input is usually interpreted as disturbance and the output is to be minimized in some sense. In stochastic worst-case design settings, the disturbance is considered random with imprecisely known probability distribution. The prior set of probability measures can be chosen so as to quantify how far the disturbance deviates from the white-noise hypothesis of Linear Quadratic Gaussian control. Such deviation can be measured by the minimal Kullback-Leibler informational divergence from the Gaussian distributions with zero mean and *scalar* covariance matrices. The resulting *anisotropy* functional is defined for *finite power* random vectors. Originally, anisotropy was introduced for *directionally generic* random vectors as the relative entropy of the normalized vector with respect to the uniform distribution on the unit sphere. The associated *a-anisotropic* norm of a matrix is then its maximum *root mean square* or *average energy* gain with respect to finite power or directionally generic inputs whose anisotropy is bounded above by $a \geq 0$. We give a systematic comparison of the anisotropy functionals and the associated norms. These are considered for unboundedly growing fragments of homogeneous Gaussian random fields on multidimensional integer lattice to yield *mean anisotropy*. Correspondingly, the anisotropic norms of finite matrices are extended to bounded linear translation invariant operators over such fields.

Keywords: Gaussian Random Field, Kullback-Leibler Informational Divergence, Mean Anisotropy, Anisotropic Norm.

AMS (MOS) subject classification. 60G60, 94A17, 60B12, 47B35.

1 Introduction

The sensitivity of the output of a given linear operator to its input can be quantified in many different ways. This issue is important in the situations, normally studied by Control Theory, where the input plays the role of a disturbance and it is desirable to minimize the output in some sense. In turn, this last is associated with a certain performance criterion and depends on assumptions made on the input.

For deterministic disturbances, the largest singular value of the operator can be used. In application to dynamic systems, this approach is employed by H_∞ control theory, e.g. [26, 8, 7, 16] to mention a few. Alternatively, if the disturbance is a random vector with homoscedastic uncorrelated entries, then an appropriate measure of the sensitivity is the trace norm of the operator. This “white noise” hypothesis is the principal supposition in Wiener-Hopf-Kalman filtering and Linear Quadratic Gaussian (LQG) control theories [11, 25, 1, 5, 13].

In more realistic situations, one is confronted by statistical uncertainty where the disturbance can be considered random, but with imprecisely known probability distribution. The associated set of probability measures constitutes the prior information on the disturbance. This leads to stochastic worst-case design problems which nowadays form a wide area of research, see e.g. [21, 14] and references therein.

Among various settings which are possible within the paradigm, we choose the one where the prior set of probability distributions serves to quantify how far the disturbance is expected to deviate from the white-noise hypothesis of LQG control. As a measure of such deviation we use the minimal Kullback-Leibler informational divergence [9, Chapter 5] of the probability distribution of a random vector from the Gaussian distributions with zero mean and *scalar* covariance matrices.

The resulting functional, called *anisotropy*, is well defined for absolutely continuously

distributed square integrable (or briefly, *finite power*) random vectors. The so-defined anisotropy functional was studied in [23] and is not dissimilar to the power-entropy construct considered in [2] for scalar random variables. The sensitivity of a linear operator can then be described by its *a-anisotropic norm* defined as the maximum *root mean square gain* of the operator with respect to random vectors whose anisotropy is bounded above by a given nonnegative parameter a . The corresponding worst-case input turns out to be Gaussian distributed with zero mean. In [23] this approach was used to develop a robust performance analysis of control systems evolving on a finite discrete time horizon.

The anisotropy-based approach to quantitative description of the statistical uncertainty in entropy theoretic terms for the purposes of robust control was proposed in [19] and [22], where anisotropy of a random vector was defined in a different way, as the relative entropy of the normalized vector with respect to the uniform distribution on the unit sphere. The associated *a-anisotropic norm* is the maximum *average energy gain* with respect to *directionally generic* disturbances for which the normalized vector is well-defined and absolutely continuously distributed on the sphere. In [22], the anisotropy functional was also extended to stationary Gaussian sequences by computing it for increasingly long fragments of the sequence and taking an appropriate limit to obtain *mean anisotropy* per unit time.

The present paper is aimed at a more systematic comparison of the anisotropy functionals and anisotropic norms and at generalization of the aforementioned constructs to bounded linear translation-invariant operators over vector-valued homogeneous Gaussian random fields on multi-dimensional integer lattices. These results can find applications in robust recovery of multivariate functions by noise corrupted data, e.g. in image processing, and in robust control of flexible structures.

The paper is organized as follows. Sections 2 and 3 provide definitions and basic properties of the anisotropy functionals for the classes of directionally generic and finite power random vectors. Complementing the results of [23, Section 2.2], the functionals are compared in Section 4 where a class of *quasigaussian* random vectors is described for which the anisotropies share the same value. In Section 5, the anisotropies are computed for zero mean Gaussian random vectors. Section 6 gives definitions and basic properties of the anisotropic norms of matrices induced by the aforementioned anisotropy functionals. In Sections 7 and 8, the anisotropies are considered for fragments of a homogeneous Gaussian random field on a multidimensional integer lattice obtained by restricting the field to finite subsets of the lattice. In Section 8, it is shown that as the subsets tend to infinity in the sense of van Hove, widely used in Statistical Mechanics of lattice systems [15, 18], the properly normalized anisotropies have a common limit, the *mean anisotropy* of the field. In Sections 9 and 10, the anisotropic norm is defined for bounded linear translation invariant operators over homogeneous Gaussian random fields, and formulas are given for computing the norm. In Section 11, an asymptotic connection of the norm is established with those of finite dimensional projections of the operator associated with finite subsets of the lattice. In Sections 12 and 13, proofs of the main theorems are given along with subsidiary lemmas.

2 Directionally Generic Random Vectors

Recall that for probability measures M and N on a common measurable space (X, \mathcal{E}) , the Kullback-Leibler informational divergence [9, p. 89] of M with respect to (wrt) N is defined as

$$\mathbf{D}(M\|N) = \begin{cases} \mathbf{E} \ln \frac{dM}{dN} & \text{if } M \ll N \\ +\infty & \text{otherwise} \end{cases}.$$

Here, $\mathbf{E}(\cdot)$ denotes expectation in the sense of M , and $dM/dN : X \rightarrow \mathbb{R}_+$ is the Radon-Nikodym derivative in the case of absolute continuity of M wrt N written as $M \ll N$. By the Divergence Inequality [9, Lemma 5.2.1 on p. 90], the quantity $\mathbf{D}(M\|N)$ is always nonnegative and is only zero if $M = N$.

If M and N are probability distributions of random elements ξ and η or are given by their probability density functions (pdf) f and g wrt a common dominating measure, we shall, slightly abusing notations, occasionally replace the symbols M or N in $\mathbf{D}(M\|N)$ with ξ , f or η , g , respectively.

Definition 2.1: Say that a \mathbb{R}^m -valued random vector W , defined on an underlying probability space $(\Omega, \mathcal{F}, \mathbf{P})$, is *directionally generic* if $\mathbf{P}(W = 0) = 0$ and the probability distribution of $W/|W|$ is absolutely continuous wrt to the uniform distribution U_m on the unit sphere $\mathbb{S}_m = \{s \in \mathbb{R}^m : |s| = 1\}$.

Denote by \mathbb{D}_m the class of m -dimensional directionally generic random vectors. *Anisotropy* of $W \in \mathbb{D}_m$ was defined in [22] as the quantity

$$\mathbf{A}_o(W) = \mathbf{D}(Q\|U_m) = \int_{\mathbb{S}_m} g(s) \ln g(s) U_m(ds). \quad (2.1)$$

Here, $g = dQ/dU_m$ is the pdf of $V = W/|W|$ wrt U_m , and Q is the probability distribution of V expressed in terms of the distribution P of W as

$$Q(B) = P(\mathbb{R}_+ B), \quad B \in \mathcal{S}_m, \quad (2.2)$$

where $\mathbb{R}_+ B = \{rs : r \in \mathbb{R}_+, s \in B\}$ is a cone in \mathbb{R}^m , and \mathcal{S}_m denotes the σ -algebra of Borel subsets of \mathbb{S}_m .

By the Divergence Inequality, the anisotropy $\mathbf{A}_o(W)$ is always nonnegative and is only zero if $Q = U_m$. Clearly, $\mathbf{A}_o(W)$ is invariant under transformations $W \mapsto \varphi RW$, where φ is a positive scalar random variable and $R \in \text{so}(m)$ is a nonrandom orthogonal $(m \times m)$ -matrix. In particular, $\mathbf{A}_o(W)$ is invariant wrt nonrandom permutations of the entries of W . Therefore, $\mathbf{A}_o(W)$ can also be interpreted as an information theoretic measure of *directional nonuniformity* of P , i.e. noninvariance of Q under the group of rotations.

For example, any random vector W , distributed absolutely continuously wrt the m -dimensional Lebesgue measure mes_m , is directionally generic. In this case, the pdf g of $W/|W|$ is expressed in terms of the pdf f of W as

$$g(s) = \int_0^{+\infty} f(rs) R_m(dr), \quad s \in \mathbb{S}_m. \quad (2.3)$$

Here, R_m is an absolutely continuous measure on \mathbb{R}_+ defined by

$$R_m(dr) = S_m r^{m-1} dr, \quad (2.4)$$

where

$$S_m = \text{mes}_{m-1} \mathbb{S}_m = \frac{2\pi^{m/2}}{\Gamma(m/2)}, \quad (2.5)$$

and $\Gamma(\lambda) = \int_0^{+\infty} u^{\lambda-1} \exp(-u) du$ denotes the Euler gamma function.

3 Finite Power Random Vectors

Denote by \mathbb{L}_2^m the class of square integrable \mathbb{R}^m -valued random vectors distributed absolutely continuously wrt mes_m . Elements of the class will be called briefly *finite power* random vectors. Clearly, any $W \in \mathbb{L}_2^m$ is directionally generic in the sense of Definition 2.1, i.e. $\mathbb{L}_2^m \subset \mathbb{D}_m$. Although the last inclusion is strict, for any $W \in \mathbb{D}_m \setminus \mathbb{L}_2^m$ there exists a positive scalar random variable φ such that $\varphi W \in \mathbb{L}_2^m$.

Based on a power-entropy construct considered in [2] for scalar random variables, a definition of anisotropy $\mathbf{A}(W)$ of $W \in \mathbb{L}_2^m$, alternative to (2.1), was proposed in [23] as

$$\mathbf{A}(W) = \min_{\lambda > 0} \mathbf{D}(W \| p_{m,\lambda}) = \frac{m}{2} \ln \left(\frac{2\pi e}{m} \mathbf{E}|W|^2 \right) - \mathbf{h}(W), \quad (3.1)$$

where

$$\mathbf{h}(W) = - \int_{\mathbb{R}^m} f(w) \ln f(w) dw \quad (3.2)$$

is the differential entropy [4, p. 229] of W , and f is its pdf wrt mes_m . In (3.1), $p_{m,\lambda}$ denotes the Gaussian pdf on \mathbb{R}^m with zero mean and *scalar* covariance matrix λI_m ,

$$p_{m,\lambda}(w) = (2\pi\lambda)^{-m/2} \exp \left(-\frac{|w|^2}{2\lambda} \right).$$

In general, denote by $\mathbb{G}^m(C)$ the class of \mathbb{R}^m -valued Gaussian distributed random vectors with zero mean and covariance matrix C . In the case $\det C \neq 0$, the corresponding pdf is

$$f(w) = (2\pi)^{-m/2} (\det C)^{-1/2} \exp \left(-\frac{1}{2} \|w\|_{C^{-1}}^2 \right), \quad (3.3)$$

where $\|x\|_Q = \sqrt{x^T Q x}$ denotes the (semi-) norm of a vector x induced by a positive (semi-) definite symmetric matrix Q . The lemma below shows that the anisotropy functional (3.1) is qualitatively similar to (2.1).

Lemma 3.1: [23, Lemma 1]

- (a) The anisotropy $\mathbf{A}(W)$ defined by (3.1) is invariant under rotation and central dilatation of W , i.e. $\mathbf{A}(\lambda U W) = \mathbf{A}(W)$ for any $\lambda \in \mathbb{R} \setminus \{0\}$ and any $U \in \text{so}(m)$;
- (b) For any positive definite symmetric $C \in \mathbb{R}^{m \times m}$,

$$\min \{ \mathbf{A}(W) : W \in \mathbb{L}_2^m, \mathbf{E}(W W^T) = C \} = -\frac{1}{2} \ln \det \frac{mC}{\text{Tr} C}, \quad (3.4)$$

where the minimum is attained only at $W \in \mathbb{G}^m(C)$;

- (c) For any $W \in \mathbb{L}_2^m$, $\mathbf{A}(W) \geq 0$. Moreover, $\mathbf{A}(W) = 0$ iff $W \in \mathbb{G}^m(\lambda I_m)$ for some $\lambda > 0$.

By Lemma 3.1(c) which essentially replicates the definition (3.1), the anisotropy $\mathbf{A}(W)$ is an information theoretic distance of the probability distribution of W from the Gaussian distributions with zero mean and scalar covariance matrices. At the same time, $\mathbf{A}(W)$ quantifies noninvariance of the distribution under the group of rotations.

4 Quasi-Gaussian Random Vectors

Denote by \mathbb{L}_2^+ the class of square integrable \mathbb{R}_+ -valued random variables, distributed absolutely continuously wrt mes_1 . For any $\xi \in \mathbb{L}_2^+$ with pdf α wrt the measure R_m given by (2.4)–(2.5), define the quantity

$$\mathbf{a}_m(\xi) = \frac{m}{2} \ln \left(\frac{2\pi e}{m} \mathbf{E}\xi^2 \right) - \mathbf{b}_m(\xi), \quad (4.1)$$

where

$$\mathbf{b}_m(\xi) = - \int_0^{+\infty} \alpha(r) \ln \alpha(r) R_m(dr) \quad (4.2)$$

is the differential entropy of ξ wrt R_m . A variational meaning of (4.1) is clarified immediately below.

Lemma 4.1: [23, Lemma 2] *For any $\xi \in \mathbb{L}_2^+$, the functional \mathbf{a}_m , defined by (4.1)–(4.2), is representable as*

$$\mathbf{a}_m(\xi) = \min_{\lambda > 0} \mathbf{D}(\xi \| \sqrt{\lambda\eta}). \quad (4.3)$$

Here, η is a χ_m^2 -distributed random variable, with χ_m^2 denoting the χ^2 -law with m degrees of freedom [24, pp. 183–184], and the minimum is attained at $\lambda = \mathbf{E}\xi^2/m$.

By the variational representation (4.3) and by the Divergence Inequality, the quantity (4.1) is always nonnegative, with $\mathbf{a}_m(\xi) = 0$ iff $m\xi^2/\mathbf{E}\xi^2$ is χ_m^2 -distributed as is $|W|^2$ for $W \in \mathbb{G}^m(I_m)$. The lemma below links together the two definitions of anisotropy given in the previous sections.

Lemma 4.2: [23, Theorem 1] *For any $W \in \mathbb{L}_2^m$, the anisotropies (2.1) and (3.1) are related as*

$$\mathbf{A}(W) = \mathbf{A}_o(W) + \mathbf{I}(\rho; \sigma) + \mathbf{a}_m(\rho) \quad (4.4)$$

where $\mathbf{a}_m(\rho)$ is the functional (4.1) applied to $\rho = |W|$, and $\mathbf{I}(\rho; \sigma)$ is the mutual information [4, p. 231] between ρ and $\sigma = W/|W|$.

The representations (4.3) and (4.4) imply that

$$\mathbf{A}_o(W) \leq \mathbf{A}(W) \quad \text{for all } W \in \mathbb{L}_2^m. \quad (4.5)$$

Definition 4.1: A random vector $W \in \mathbb{L}_2^m$ is called *quasigaussian* if $|W|$ and $W/|W|$ are independent, and $m|W|^2/\mathbf{E}|W|^2$ is χ_m^2 -distributed.

Denote the class of the quasigaussian random vectors by \mathbb{Q}^m . Clearly, $\mathbb{G}^m(\lambda I_m) \subset \mathbb{Q}^m$ for any $\lambda > 0$. By Lemmas 4.1 and 4.2,

$$\mathbb{Q}^m = \{W \in \mathbb{L}_2^m : \mathbf{A}_o(W) = \mathbf{A}(W)\}. \quad (4.6)$$

Also note that for any $W \in \mathbb{D}_m$, $\mathbf{A}_o(W) = \inf \mathbf{A}(\varphi W)$ where the infimum is taken over all the positive scalar random variables φ such that $\varphi W \in \mathbb{L}_2^m$.

5 Anisotropy of Gaussian Random Vectors

Lemma 5.1: For $W \in \mathbb{G}^m(C)$ with $\det C \neq 0$, the anisotropies (2.1) and (3.1) satisfy the relations

$$\mathbf{A}_\circ(W) = -\frac{1}{2} \ln \det \frac{C}{\exp(2\mathbf{E} \ln \|\zeta\|_C)} \leq -\frac{1}{2} \ln \det \frac{mC}{\text{Tr } C} = \mathbf{A}(W), \quad (5.1)$$

where ζ is a random vector distributed uniformly on the unit sphere \mathbb{S}_m .

Proof: Plugging the Gaussian pdf (3.3) in (2.3) and using (2.5), obtain that the pdf of $V = W/|W|$ wrt U_m takes the form

$$\begin{aligned} g(s) &= \frac{2^{1-m/2}}{\Gamma(m/2)} (\det C)^{-1/2} \int_0^{+\infty} r^{m-1} \exp\left(-\frac{1}{2}(r\|s\|_{C^{-1}})^2\right) dr \\ &= (\det C)^{-1/2} \|s\|_{C^{-1}}^{-m}. \end{aligned}$$

Hence, (2.1) reads

$$\mathbf{A}_\circ(W) = -\frac{1}{2} \ln \det C - m\mathbf{E} \ln \|V\|_{C^{-1}}. \quad (5.2)$$

Introducing the random vector $Z = C^{-1/2}W \in \mathbb{G}^m(I_m)$, where $C^{1/2}$ is a matrix square root of C , obtain that

$$\|V\|_{C^{-1}} = \frac{\|W\|_{C^{-1}}}{|W|} = \frac{|Z|}{\|Z\|_C} = \left\| \frac{Z}{|Z|} \right\|_C^{-1}. \quad (5.3)$$

Since $\zeta = Z/|Z|$ is uniformly distributed on the unit sphere \mathbb{S}_m , from (5.3) it follows that

$$\mathbf{E} \ln \|V\|_{C^{-1}} = -\mathbf{E} \ln \|\zeta\|_C.$$

This last equality and (5.2) imply the left-most equality in (5.1). The equality on the right of (5.1) is a corollary of Lemma 3.1(b), while the inequality follows from the general relationship (4.5), concluding the proof.

6 Anisotropic Norms of Matrices

Let $F \in \mathbb{R}^{p \times m}$ be interpreted as a linear operator with \mathbb{R}^m -valued random input W and \mathbb{R}^p -valued output $Z = FW$. For any $a \in \mathbb{R}_+$, consider the *a-anisotropic* norms of F associated with the anisotropy functionals (2.1) and (3.1),

$$\|F\|_{a,\circ} = \sup \{ \mathbf{N}_\circ(F, W) : W \in \mathbb{D}_m, \mathbf{A}_\circ(W) \leq a \}, \quad (6.1)$$

$$\|F\|_a = \sup \{ \mathbf{N}(F, W) : W \in \mathbb{L}_2^m, \mathbf{A}(W) \leq a \}. \quad (6.2)$$

Here,

$$\mathbf{N}_\circ(F, W) = \sqrt{\mathbf{E}(|FW|/|W|)^2} = \mathbf{N}_\circ(F, W/|W|) \quad (6.3)$$

characterizes the *average energy gain* of F wrt W and is well-defined as soon as $\mathbf{P}(W = 0) = 0$, while

$$\mathbf{N}(F, W) = \sqrt{\mathbf{E}|FW|^2/\mathbf{E}|W|^2} \quad (6.4)$$

measures the *root mean square gain* of F wrt a square integrable input W . Clearly, the norms (6.1) and (6.2) are nondecreasing in $a \in \mathbb{R}_+$ and satisfy

$$\begin{aligned} \|F\|_{0,\circ} &= \|F\|_0 = \|F\|_2/\sqrt{m}, \\ \lim_{a \rightarrow +\infty} \|F\|_{a,\circ} &= \lim_{a \rightarrow +\infty} \|F\|_a = \|F\|_\infty, \end{aligned}$$

where $\|F\|_2 = \sqrt{\text{Tr}(F^T F)}$ and $\|F\|_\infty$ are respectively the Frobenius norm and the largest singular value of F .

Lemma 6.1: *For any $a \in \mathbb{R}_+$ and $F \in \mathbb{R}^{p \times m}$, the a -anisotropic norms (6.1) and (6.2) satisfy*

$$\|F\|_{a,\circ} \leq \|F\|_a.$$

Proof: Let W be a square integrable random vector satisfying $\mathbf{P}(W = 0) = 0$ and such that $|W|$ and $V = W/|W|$ are independent. Then $\mathbf{E}|FW|^2 = \mathbf{E}(|FV|^2 |W|^2) = \mathbf{E}|FV|^2 \mathbf{E}|W|^2$, and consequently, by (6.3) and (6.4),

$$\mathbf{N}(F, W) = \sqrt{\mathbf{E}|FV|^2} = \mathbf{N}_\circ(F, W). \quad (6.5)$$

In particular, (6.5) holds for any quasigaussian $W \in \mathbb{Q}^m$ (see Definition 4.1). Combining this last property with (4.6), obtain that

$$\begin{aligned} \|F\|_a &\geq \sup\{\mathbf{N}(F, W) : W \in \mathbb{Q}^m, \mathbf{A}(W) \leq a\} \\ &= \sup\{\mathbf{N}_\circ(F, W) : W \in \mathbb{Q}^m, \mathbf{A}_\circ(W) \leq a\} \\ &= \sup\{\mathbf{N}_\circ(F, W) : W \in \mathbb{D}^m, \mathbf{A}_\circ(W) \leq a\} = \|F\|_{a,\circ}, \end{aligned}$$

thereby concluding the proof.

7 Fragments of Random Fields

Denote by $\mathbb{G}F^{m,n}(S)$ the class of \mathbb{R}^m -valued homogeneous Gaussian random fields $W = (w_x)_{x \in \mathbb{Z}^n}$ on the n -dimensional integer lattice \mathbb{Z}^n with zero mean and spectral density function (sdf) $S : \Omega_n \rightarrow \mathbb{C}^{m \times m}$, where $\Omega_n = [-\pi, \pi]^n$. Since S can be extended to \mathbb{R}^n (2π)-periodically in each of its n variables, Ω_n is identified with n -dimensional torus. For any $\omega \in \Omega_n$, the matrix $S(\omega)$ is Hermitian and positive semi-definite, and satisfies $S(-\omega) = (S(\omega))^T$. The corresponding covariance function $\mathbb{Z}^n \ni x \mapsto c_x \in \mathbb{R}^{m \times m}$ is

$$c_x = \mathbf{E}(w_x w_0^T) = (2\pi)^{-n} \int_{\Omega_n} \exp(i\omega^T x) S(\omega) d\omega. \quad (7.1)$$

Definition 7.1: Say that $W \in \mathbb{G}F^{m,n}(S)$ is strictly regular if $\text{ess inf}_{\omega \in \Omega_n} \lambda_{\min}(S(\omega)) > 0$, where $\lambda_{\min}(\cdot)$ denotes the smallest eigenvalue of a Hermitian matrix.

Clearly, the strict regularity is a stronger property than standard regularity [17, pp. 27–29 and Theorem 3.2.2 on p. 30]. For simplicity, we shall assume throughout that the covariance function is absolutely summable, i.e.

$$\sum_{x \in \mathbb{Z}^n} \|c_x\|_\infty < +\infty. \quad (7.2)$$

Under this condition, the sdf S is continuous on the torus Ω_n and so also are the functions $\Omega_n \ni \omega \mapsto \lambda_{\min}(S(\omega))$, $\lambda_{\max}(S(\omega))$, with $\lambda_{\max}(\cdot)$ denoting the largest eigenvalue of a

Hermitian matrix. In this case, the strict regularity of W is equivalent to nonsingularity of $S(\omega)$ for all $\omega \in \Omega_n$.

Denote by $\mathcal{Z}_n = \{X \subset \mathbb{Z}^n : 0 < \#X < +\infty\}$ the class of nonempty finite subsets of \mathbb{Z}^n , where $\#(\cdot)$ stands for the counting measure. For any $X \in \mathcal{Z}_n$, the restriction of W to X is identified with the $\mathbb{R}^{m\#X}$ -valued Gaussian random vector

$$W_X = (w_x)_{x \in X}. \tag{7.3}$$

The order in which the random vectors w_x are “stacked” one underneath the other in (7.3) is not essential for what follows. However, to avoid ambiguity, the set X will be assumed lexicographically ordered. The spectrum of the covariance matrix

$$C_X = \mathbf{E}(W_X W_X^T) = \text{block}_{x,y \in X}(c_{x-y}) \tag{7.4}$$

is invariant under translations of $X \in \mathcal{Z}_n$ since for any $z \in \mathbb{Z}^n$ there exists a permutation matrix Π of order $m\#X$ such that $C_{X+z} = \Pi C_X \Pi^T$. If the random field W is strictly regular, then $\det C_X > 0$ for any $X \in \mathcal{Z}_n$. This implication follows from the spectral bounds

$$\text{ess inf}_{\omega \in \Omega_n} \lambda_{\min}(S(\omega)) \leq \lambda_{\min}(C_X) \leq \lambda_{\max}(C_X) \leq \text{ess sup}_{\omega \in \Omega_n} \lambda_{\max}(S(\omega)), \tag{7.5}$$

where, under the assumption (7.2), ess inf and ess sup can be replaced with \min and \max . Applying Lemma 5.1 to (7.3) and using the identity $\text{Tr } C_X = \text{Tr } c_0 \#X$, obtain that

$$\mathbf{A}_\circ(W_X) = -\frac{1}{2} \ln \det \frac{C_X}{\exp(2 \mathbf{E} \ln \|\zeta_X\|_{C_X})} \leq -\frac{1}{2} \ln \det \frac{mC_X}{\text{Tr } c_0} = \mathbf{A}(W_X), \tag{7.6}$$

where ζ_X is a random vector, distributed uniformly on the unit sphere $\mathbb{S}_{m\#X}$. It turns out that, when divided by $\#X$, both anisotropies in (7.6) have a common limit as the set X tends to infinity in a sense specified below.

8 Definition of Mean Anisotropy

With every $X \in \mathcal{Z}_n$, associate the function $D_X : \mathbb{Z}^n \rightarrow [0, 1]$ by

$$D_X(z) = \frac{\#((X+z) \cap X)}{\#X}. \tag{8.1}$$

It is worth noting that $\#X D_X$ is the *geometric covariogram* [12, p. 22] of the set X wrt the counting measure $\#$. Clearly, $\text{supp } D_X = \{x - y : x, y \in X\}$. A probabilistic interpretation of D_X is as follows. Let ξ_X and η_X be independent random vectors each distributed uniformly on X . Then the probability mass function (pmf) of $\theta_X = \xi_X - \eta_X$ is expressed in terms of (8.1) as

$$\mathbf{P}(\theta_X = z) = (\#X)^{-2} \#\{(x, y) \in X^2 : x - y = z\} = \frac{D_X(z)}{\#X}. \tag{8.2}$$

Recall that a sequence of sets $X_k \in \mathcal{Z}_n$, labeled by positive integers $k \in \mathbb{N}$, is said to tend to infinity in the sense of van Hove [18, p. 45] if $\lim_{k \rightarrow +\infty} D_{X_k}(z) = 1$ for every

$z \in \mathbb{Z}^n$. The property induces a topological filter on the class \mathcal{Z}_n and is denoted by $\nearrow \infty$.

By the identity $\#X = \sum_{z \in \mathbb{Z}^n} D_X(z)$ which follows from (8.2), a necessary condition for $X \nearrow \infty$ is $\#X \rightarrow +\infty$. A simple example of a sequence which tends to infinity in the sense of van Hove is provided by the discrete hypercubes $X_k = ([0, k] \cap \mathbb{Z})^n$ since for such sets, $D_{X_k}(z) = \prod_{j=1}^n \max(0, 1 - |z_j|/k) \rightarrow 1$ as $k \rightarrow +\infty$ for any $z = (z_j)_{1 \leq j \leq n} \in \mathbb{Z}^n$.

Theorem 8.1: *Let $W \in \mathbb{G}F^{m,n}(S)$ be strictly regular and let its covariance function be absolutely summable. Then the asymptotic behaviour of the anisotropies (2.1) and (3.1) of the random vectors W_X in (7.3) is described by*

$$\lim_{X \nearrow \infty} \frac{\mathbf{A}_c(W_X)}{\#X} = \lim_{X \nearrow \infty} \frac{\mathbf{A}(W_X)}{\#X} = -\frac{1}{2(2\pi)^n} \int_{\Omega_n} \ln \det \frac{mS(\omega)}{\text{Tr } c_0} d\omega. \quad (8.3)$$

The proof of the theorem is given in Section 12. The common limit on the right of (8.3) will be referred to as *mean anisotropy* of the field W and denoted by $\mathbf{A}(W)$.

Example 8.1: Compute the mean anisotropy of $W \in \mathbb{G}F^{m,n}(S)$ with covariance function

$$c_z = c_0 \exp\left(-\sum_{k=1}^n \frac{|z_k|}{\rho_k}\right), \quad z = (z_k)_{1 \leq k \leq n} \in \mathbb{Z}^n,$$

where $c_0 \in \mathbb{R}^{m \times m}$ is a positive definite symmetric matrix, and ρ_1, \dots, ρ_n are positive reals, with ρ_k interpreted as a *correlation radius* of W along the k -th coordinate axis in \mathbb{R}^n . The corresponding sdf S is given by

$$S(\omega) = c_0 \prod_{k=1}^n \sigma_k(\omega_k), \quad \omega = (\omega_k)_{1 \leq k \leq n} \in \Omega_n.$$

Here, for every $1 \leq k \leq n$, the function $\sigma_k : \Omega_1 \rightarrow \mathbb{R}_+$, defined by

$$\sigma_k(u) = \frac{1 - \alpha_k^2}{1 + \alpha_k^2 - 2\alpha_k \cos u}, \quad \alpha_k = \exp(-1/\rho_k),$$

is sdf of a stationary scalar Gaussian sequence $(\xi_t)_{t \in \mathbb{Z}}$ with zero mean and covariance function $\mathbf{E}(\xi_t \xi_0) = \alpha_k^{|t|}$. Applying the Szego-Kolmogorov formula and using the Markov property of the sequence together with the Normal Correlation lemma, obtain

$$\exp\left(\frac{1}{2\pi} \int_{-\pi}^{\pi} \ln \sigma_k(\omega) d\omega\right) = \mathbf{Var}(\xi_0 \mid (\xi_t)_{t < 0}) = \mathbf{Var}(\xi_0 \mid \xi_{-1}) = 1 - \alpha_k^2, \quad (8.4)$$

where $\mathbf{Var}(\cdot \mid \cdot)$ denotes conditional variance. Clearly, the random field W satisfies the assumptions of Theorem 8.1 and, by (8.4), its mean anisotropy defined in (8.3) reads

$$\begin{aligned} \mathbf{A}(W) &= -\frac{1}{2} \ln \det \frac{m c_0}{\text{Tr } c_0} - \frac{m}{4\pi} \sum_{k=1}^n \int_{-\pi}^{\pi} \ln \sigma_k(\omega) d\omega \\ &= -\frac{1}{2} \ln \det \frac{m c_0}{\text{Tr } c_0} - \frac{m}{2} \sum_{k=1}^n \ln(1 - \alpha_k^2). \end{aligned}$$

Here, the right-most sum behaves asymptotically like $-\sum_{k=1}^n \ln \rho_k$ if the correlation radii are all large.

9 Anisotropic Norm of LTI Operators

Denote by $\ell_2^{r,n} = \{V = (v_x)_{x \in \mathbb{Z}^n} \in (\mathbb{R}^r)^{\mathbb{Z}^n} : \|V\|_2 = \sqrt{\sum_{x \in \mathbb{Z}^n} |v_x|^2} < +\infty\}$ the Hilbert space of square summable maps of \mathbb{Z}^n to \mathbb{R}^r . Let $\mathcal{L}_\infty^{p \times m, n}$ stand for the Banach space of bounded linear translation-invariant operators $F : \ell_2^{m,n} \rightarrow \ell_2^{p,n}$ equipped with the norm

$$\|F\|_\infty = \sup_{W \in \ell_2^{m,n}} \frac{\|FW\|_2}{\|W\|_2}. \tag{9.1}$$

Here, the output $Z = (z_x)_{x \in \mathbb{Z}^n} = FW$ of $F \in \mathcal{L}_\infty^{p \times m, n}$ relates to the input $W = (w_x)_{x \in \mathbb{Z}^n}$ by

$$z_x = \sum_{y \in \mathbb{Z}^n} f_{x-y} w_y, \quad x \in \mathbb{Z}^n, \tag{9.2}$$

where $\mathbb{Z}^n \ni x \mapsto f_x \in \mathbb{R}^{p \times m}$ is the impulse response function. The operator is identified with the transfer function $F : \Omega_n \rightarrow \mathbb{C}^{p \times m}$ defined by

$$F(\omega) = \sum_{x \in \mathbb{Z}^n} f_x \exp(-ix^T \omega). \tag{9.3}$$

The \mathcal{L}_∞ -norm of this last, $\text{ess sup}_{\omega \in \Omega_n} \|F(\omega)\|_\infty$, coincides with (9.1) and, upon rescaling, is an upper bound for the L_2 -norm,

$$\|F\|_2 = \sqrt{(2\pi)^{-n} \int_{\Omega_n} \text{Tr } H(\omega) d\omega} = \sqrt{\sum_{x \in \mathbb{Z}^n} \|f_x\|_2^2} \leq \sqrt{m} \|F\|_\infty, \tag{9.4}$$

where the map $H : \Omega_n \rightarrow \mathbb{C}^{m \times m}$ is defined by

$$H(\omega) = (F(\omega))^* F(\omega). \tag{9.5}$$

The inequality on the right of (9.4) becomes an equality iff there exists $\lambda \in \mathbb{R}_+$ such that $H(\omega) = \lambda I_m$ for mes_n -almost all $\omega \in \Omega_n$.

Definition 9.1: An operator $F \in \mathcal{L}_\infty^{p \times m, n}$ is called nonround if $\|F\|_2 < \sqrt{m} \|F\|_\infty$.

If $F \in \mathcal{L}_\infty^{p \times m, n}$ and its input $W \in \mathbb{G}F^{m,n}(S)$, then the convergence of the series (9.2) is understood in mean square sense and the output satisfies $Z \in \mathbb{G}F^{p,n}(FSF^*)$. In particular,

$$\mathbf{E}|z_0|^2 = (2\pi)^{-n} \int_{\Omega_n} \text{Tr}(H(\omega)S(\omega)) d\omega.$$

Recalling the relations (7.1) and $\mathbf{E}|w_0|^2 = \text{Tr } c_0$, quantify the root mean square gain of F wrt W by

$$\mathbf{N}(F, W) = \sqrt{\frac{\mathbf{E}|z_0|^2}{\mathbf{E}|w_0|^2}} = \sqrt{\frac{\int_{\Omega_n} \text{Tr}(H(\omega)S(\omega)) d\omega}{\int_{\Omega_n} \text{Tr } S(\omega) d\omega}}.$$

For every $a \geq 0$, define the a -anisotropic norm of the operator F as

$$\|F\|_a = \sup \{\mathbf{N}(F, W) : \mathbf{A}(W) \leq a\}. \tag{9.6}$$

Here, the supremum is taken over all the strictly regular homogeneous Gaussian random fields W whose mean anisotropy (8.3) is bounded above by a . Denote by $\mathbb{W}_a(F)$ the corresponding set of the *worst-case* inputs W at which the supremum in (9.6) is attained.

10 Computing Anisotropic Norm

Assuming the operator $F \in \mathcal{L}_\infty^{p \times m, n}$ fixed, for notational convenience let

$$Q = [0, \|F\|_\infty^{-2}]. \quad (10.1)$$

Recalling (9.5), define the functions $\mathcal{A}, \mathcal{N}, \Phi, \Psi$ on Q by

$$\mathcal{A}(q) = \frac{m}{2}(\ln \Phi(q) - \Psi(q)), \quad (10.2)$$

$$\mathcal{N}(q) = \left(\frac{1}{q} \left(1 - \frac{1}{\Phi(q)} \right) \right)^{1/2}, \quad (10.3)$$

$$\Phi(q) = \frac{1}{m(2\pi)^n} \int_{\Omega_n} \text{Tr}(I_m - qH(\omega))^{-1} d\omega, \quad (10.4)$$

$$\Psi(q) = -\frac{1}{m(2\pi)^n} \int_{\Omega_n} \ln \det(I_m - qH(\omega)) d\omega. \quad (10.5)$$

Here, \mathcal{N} is extended to 0 by continuity as $\mathcal{N}(0) = \lim_{q \rightarrow 0+} \mathcal{N}(q) = \|F\|_2 / \sqrt{m}$. The functions (10.2)–(10.5) are all analytic, nondecreasing in $q \in Q$ and take values in \mathbb{R}_+ , $(\|F\|_2 / \sqrt{m}, \|F\|_\infty)$, $[1, +\infty)$ and \mathbb{R}_+ , respectively.

Following the technique of a randomized singular value used in the proof of [6, Lemma 4], one verifies that if F is nonround in the sense of Definition 9.1, then \mathcal{A} is strictly increasing and convex on the interval Q , with $\mathcal{A}(0) = 0$ and $\mathcal{A}(q) \rightarrow +\infty$ as $q \rightarrow \|F\|_\infty^{-2}$. These properties imply that the function \mathcal{A} is invertible and its inverse $\mathcal{A}^{-1} : \mathbb{R}_+ \rightarrow Q$ is strictly increasing and concave.

Theorem 10.1: *Let $F \in \mathcal{L}_\infty^{p \times m, n}$ be nonround. Then for any $a \geq 0$, the a -anisotropic norm (9.6) is expressed in terms of the functions (10.2) and (10.3) as*

$$\|F\|_a = \mathcal{N}(\mathcal{A}^{-1}(a)). \quad (10.6)$$

The corresponding set of worst-case inputs is

$$\mathbb{W}_a(F) = \bigcup_{\lambda > 0} \mathbb{G}F^{m, n} (\lambda(I_m - \mathcal{A}^{-1}(a)H)^{-1}). \quad (10.7)$$

The proof of the theorem is similar to that of [6, Theorem 3] and therefore omitted. Using the remark made after the proof of [23, Theorem 2], one verifies that the norm $\|F\|_a$ is concave in $a \in \mathbb{R}_+$.

Example 10.1: For a given $Y \in \mathcal{Z}_n$, consider a Y -averaging operator $F \in \mathcal{L}_\infty^{1 \times 1, n}$ with impulse response

$$f_x = \frac{\mathcal{I}_Y(x)}{\#Y}, \quad x \in \mathbb{Z}^n,$$

where $\mathcal{I}_Y : \mathbb{Z}^n \rightarrow \{0, 1\}$ is the indicator function of the set Y . The corresponding transfer function is

$$F(\omega) = \frac{1}{\#Y} \sum_{y \in Y} \exp(-iy^T \omega), \quad \omega \in \Omega_n.$$

Clearly, $\|F\|_2 = 1/\sqrt{\#Y}$ and $\|F\|_\infty = 1$. The complex conjugate \overline{F} is the characteristic function (cf) of the uniform distribution on Y . Hence, $H = |F|^2 : \Omega_n \rightarrow [0, 1]$ is cf of

$\theta = \xi - \eta$, where ξ and η are independent random vectors each distributed uniformly on Y . The corresponding pmf is given by $D_Y/\#Y$ as in (8.2). Now let $\Theta = (\theta_k)_{k \in \mathbb{N}}$ be a sequence of independent $(Y - Y)$ -valued random vectors, each with the cf H . Associate with Θ a random walk $\Sigma = (\sigma_k)_{k \in \mathbb{Z}_+}$ on \mathbb{Z}^n defined by

$$\sigma_k = \sum_{j=1}^k \theta_j.$$

For every $k \in \mathbb{Z}_+$, H^k is cf of σ_k , and hence, $(2\pi)^{-n} \int_{\Omega_n} (H(\omega))^k d\omega = \mathbf{P}(\sigma_k = 0)$. Therefore, the function (10.4) takes the form

$$\Phi(q) = (2\pi)^{-n} \int_{\Omega_n} \frac{d\omega}{1 - qH(\omega)} = \sum_{k \in \mathbb{Z}_+} q^k \mathbf{P}(\sigma_k = 0). \tag{10.8}$$

Denote by $\tau = \min \{k \in \mathbb{N} : \sigma_k = 0\}$ the first recurrence time for the random walk Σ . By the well-known identity for Markov chains, (10.8) is expressed in terms of the moment generating function of τ as

$$\Phi(q) = \frac{1}{1 - \mathbf{E}q^\tau}.$$

11 Connection with Anisotropic Norms of Matrices

A connection of the anisotropic norm of the operator F with those of finite matrices (see Section 6) is established below. To formulate the statement, for every $X \in \mathcal{Z}_n$, introduce a matrix $F_X \in \mathbb{R}^{\#X \times m\#X}$ by appropriately restricting the impulse response function of F ,

$$F_X = \text{block}(f_{x-y})_{x,y \in X}. \tag{11.1}$$

If P_X and M_X are the orthogonal projectors in $\ell_2^{p,n}$ and $\ell_2^{m,n}$ to the subspaces of signals whose support is contained in X , then $P_X F M_X = F_X M_X$.

Theorem 11.1: *Let $F \in \mathcal{L}_\infty^{p \times m, n}$ be nonround and let its impulse response function be absolutely summable, i.e. $\sum_{x \in \mathbb{Z}^n} \|f_x\|_\infty < +\infty$. Then for every $a \geq 0$, the $(a\#X)$ -anisotropic norms (6.1) and (6.2) of the matrices (11.1) have the a -anisotropic norm of F in (9.6) as their common limit,*

$$\lim_{X \nearrow \infty} \|F_X\|_{a\#X, \circ} = \lim_{X \nearrow \infty} \|F_X\|_{a\#X} = \|F\|_a. \tag{11.2}$$

The theorem is proved in Section 13.

12 Proof of Theorem 8.1

For any $X \in \mathcal{Z}_n$ and $r \in \mathbb{N}$, introduce the function $E_{X,r} : \mathbb{Z}^{rn} \rightarrow [0, 1]$ which maps a vector $y = (y_k)_{1 \leq k \leq rn}$, formed by $y_1, \dots, y_r \in \mathbb{Z}^n$, to

$$E_{X,r}(y) = \frac{\# \left(X \cap \bigcap_{j=1}^r \left(X + \sum_{k=1}^j y_k \right) \right)}{\#X}. \tag{12.1}$$

Comparison with (8.1) shows that $E_{X,1}(z) = E_{X,2}(z, -z) = D_X(z)$ for any $z \in \mathbb{Z}^n$. By (12.1),

$$\begin{aligned} 1 - E_{X,r}(y_1, \dots, y_r) &= \frac{\#\bigcup_{j=1}^r \left(X \setminus \left(X + \sum_{k=1}^j y_k \right) \right)}{\#X} \\ &\leq \sum_{j=1}^r \frac{\#(X \setminus (X + \sum_{k=1}^j y_k))}{\#X} = r - \sum_{j=1}^r D_X \left(\sum_{k=1}^j y_k \right). \end{aligned}$$

Therefore, the definition of convergence in the sense of van Hove (see Section 8) yields

$$\lim_{X \nearrow \infty} E_{X,r}(y) = 1 \quad \text{for all } r \in \mathbb{N}, y \in \mathbb{Z}^{rn}. \quad (12.2)$$

For notational convenience in the sequel, introduce the set

$$\mathbb{O}_r = \left\{ (z_k)_{1 \leq k \leq r} \in \mathbb{Z}^{rn} : z_1, \dots, z_r \in \mathbb{Z}^n, \sum_{k=1}^r z_k = 0 \right\}. \quad (12.3)$$

Lemma 12.1: *Let the covariance function of $W \in \mathbb{G}F^{m,n}(S)$ in (7.1) be absolutely summable. Then for any $r \in \mathbb{N}$, the matrices (7.4) satisfy*

$$\lim_{X \nearrow \infty} \frac{\text{Tr } C_X^r}{\#X} = (2\pi)^{-n} \int_{\Omega_n} \text{Tr} (S(\omega))^r d\omega. \quad (12.4)$$

Proof: Define the function $\varphi : \Omega_{rn} \rightarrow \mathbb{R}$ which maps a vector $\omega = (\omega_k)_{1 \leq k \leq r}$, formed by $\omega_1, \dots, \omega_r \in \Omega_n$, to

$$\varphi(\omega) = \text{Tr} (S(\omega_1) \times \dots \times S(\omega_r)) = \sum_{y \in \mathbb{Z}^{rn}} \psi_y \exp(-iy^T \omega). \quad (12.5)$$

Here, for any $y = (y_k)_{1 \leq k \leq r}$ formed by $y_1, \dots, y_r \in \mathbb{Z}^n$, the Fourier coefficient ψ_y is given by

$$\psi_y = (2\pi)^{-rn} \int_{\Omega_{rn}} \varphi(\omega) \exp(iy^T \omega) d\omega = \text{Tr} (c_{y_1} \times \dots \times c_{y_r}). \quad (12.6)$$

In these notations, (7.1) and (7.4) imply that for any $X \in \mathcal{Z}_n$,

$$\begin{aligned} \text{Tr } C_X^r &= \sum_{x_1, \dots, x_r \in X} \text{Tr} (c_{x_1 - x_2} c_{x_2 - x_3} \times \dots \times c_{x_{r-1} - x_r} c_{x_r - x_1}) \\ &= \sum_{x_1, \dots, x_r \in X} \psi_{(x_1 - x_2, x_2 - x_3, \dots, x_{r-1} - x_r, x_r - x_1)}. \end{aligned}$$

Hence, recalling (12.1) and (12.3), obtain

$$\frac{\text{Tr } C_X^r}{\#X} = \sum_{y \in \mathbb{O}_r} \psi_y E_{X,r}(y). \quad (12.7)$$

By the inequality $|\text{Tr } A| \leq m \|A\|_\infty$ which holds for any $A \in \mathbb{R}^{m \times m}$ and by submultiplicativity of $\|\cdot\|_\infty$, (12.6) implies that

$$\sum_{y \in \mathbb{Z}^{rn}} |\psi(y)| \leq m \sum_{y_1, \dots, y_r \in \mathbb{Z}^n} \prod_{k=1}^r \|c_{y_k}\|_\infty \leq m \left(\sum_{x \in \mathbb{Z}^n} \|c_x\|_\infty \right)^r.$$

Consequently, the assumption of the lemma assures that $\sum_{x \in \mathbb{O}_r} |\psi_y| < +\infty$, thus legitimating the passage to the limit under the sum in (12.7) on a basis of (12.2),

$$\lim_{X \nearrow \infty} \frac{\text{Tr } C_X^r}{\#X} = \sum_{y \in \mathbb{O}_r} \psi_y \lim_{X \nearrow \infty} E_{X,r}(y) = \sum_{y \in \mathbb{O}_r} \psi_y. \tag{12.8}$$

It now remains to note that by (12.3), (12.5) and (12.6),

$$\sum_{y \in \mathbb{O}_r} \psi_y = (2\pi)^{-n} \int_{\Omega_n} \underbrace{\varphi(\omega, \dots, \omega)}_{r \text{ times}} d\omega = (2\pi)^{-n} \int_{\Omega_n} \text{Tr}(S(\omega))^r d\omega$$

which, in combination with (12.8), immediately yields (12.4), thereby completing the proof.

Note that the assertion of Lemma 12.1 for the particular case $r = 2$ can be established in a much simpler way. Indeed, by (7.4), (8.1) and by Parseval's equality,

$$\begin{aligned} \frac{\text{Tr } C_X^2}{\#X} &= \frac{1}{\#X} \sum_{x,y \in X} \|c_{x-y}\|_2^2 = \sum_{z \in \mathbb{Z}^n} D_X(z) \|c_z\|_2^2 \\ &\rightarrow \sum_{z \in \mathbb{Z}^n} \|c_z\|_2^2 = (2\pi)^{-n} \int_{\Omega_n} \text{Tr}(S(\omega))^2 d\omega \quad \text{as } X \nearrow \infty. \end{aligned} \tag{12.9}$$

Lemma 12.2: *Let $W \in \mathbb{G}F^{m,n}(S)$ be strictly regular and let (7.2) hold. Then the matrices (7.4) satisfy*

$$\lim_{X \nearrow \infty} \frac{\ln \det C_X}{\#X} = (2\pi)^{-n} \int_{\Omega_n} \ln \det S(\omega) d\omega \tag{12.10}$$

Proof:By (7.5), under the assumptions of the lemma, for any $X \in \mathbb{Z}_n$, the spectrum of C_X is entirely contained in the interval

$$\Delta = \left[\min_{\omega \in \Omega_n} \lambda_{\min}(S(\omega)), \max_{\omega \in \Omega_n} \lambda_{\max}(S(\omega)) \right] \tag{12.11}$$

which is separated from zero and bounded. Since the logarithm function is expandable on the interval to a uniformly convergent power series, application of Lemma 12.1 to the series yields (12.10).

Note that the assertion of Lemma 12.2 under weaker assumptions is well-known in the case $n = 1$ for Toeplitz forms [10], and is closely related to Szego-Kolmogorov formula for Shannon entropy rate in stationary Gaussian sequences. For the multivariate case $n > 1$, it is worth pointing out the links to the mean entropy results for Gibbs-Markov random fields [18, pp. 44–47].

Lemma 12.3: *Let ζ be uniformly distributed on the unit sphere \mathbb{S}_r , and let $C \in \mathbb{R}^{r \times r}$ be a positive semi-definite symmetric matrix. Then*

$$\mathbf{E} \|\zeta\|_C^2 = \frac{\text{Tr } C}{r}, \tag{12.12}$$

$$\mathbf{Var} \|\zeta\|_C^2 = \frac{2}{r+2} \left(\frac{\text{Tr } C^2}{r} - \left(\frac{\text{Tr } C}{r} \right)^2 \right), \tag{12.13}$$

where $\mathbf{Var}(\cdot)$ denotes the variance of a random variable. Moreover, if $\det C \neq 0$ and $r \geq 3$, then

$$\mathbf{E}\|\zeta\|_C^{-2} \leq \frac{r-2}{2r} \left(\frac{\Gamma(\frac{r-2}{2r})}{\Gamma(1/2)} \right)^r (\det C)^{-1/r}. \tag{12.14}$$

Proof: Since the uniform distribution on \mathbb{S}_r is invariant under the group of rotations, $\|\zeta\|_C^2$ has the same distribution as $\sum_{k=1}^r \lambda_k \zeta_k^2$, where λ_k are the eigenvalues of C , and ζ_k are the entries of ζ . Denote $\tau_k = \zeta_k^2$, $1 \leq k \leq r$. By definition, $\sum_{k=1}^r \tau_k = 1$. The random variables $\tau_1, \dots, \tau_{r-1}$ have the $(r-1)$ -variate Dirichlet distribution $D(\underbrace{1/2, \dots, 1/2}_{r \text{ times}}; 1/2)$ [24, p. 177] with pdf

$$\frac{\Gamma(r/2)}{(\Gamma(1/2))^r} \left(1 - \sum_{k=1}^{r-1} t_k \right)^{-1/2} \prod_{k=1}^{r-1} t_k^{-1/2}$$

wrt mes_{r-1} on the simplex $\{(t_1, \dots, t_{r-1}) \in \mathbb{R}_+^{r-1} : \sum_{k=1}^{r-1} t_k \leq 1\}$. By a straightforward calculation (also see [24, p. 179]), for all $1 \leq j \neq k \leq r$,

$$\mathbf{E}\tau_j = \frac{\Gamma(r/2)}{(\Gamma(1/2))^r} \frac{\Gamma(3/2)(\Gamma(1/2))^{r-1}}{\Gamma(r/2+1)} = \frac{1}{r}, \tag{12.15}$$

$$\mathbf{E}\tau_j^2 = \frac{\Gamma(r/2)}{(\Gamma(1/2))^r} \frac{\Gamma(5/2)(\Gamma(1/2))^{r-1}}{\Gamma(r/2+2)} = \frac{3}{r(r+2)}, \tag{12.16}$$

$$\mathbf{E}(\tau_j \tau_k) = \frac{\Gamma(r/2)}{(\Gamma(1/2))^r} \frac{(\Gamma(3/2))^2 (\Gamma(1/2))^{r-2}}{\Gamma(r/2+2)} = \frac{1}{r(r+2)}. \tag{12.17}$$

From (12.15) it follows that

$$\mathbf{E}\|\zeta\|_C^2 = \frac{1}{r} \sum_{k=1}^r \lambda_k = \frac{\text{Tr } C}{r} \tag{12.18}$$

which coincides with (12.12). Furthermore, taking (12.16) and (12.17) into account, obtain that

$$\begin{aligned} \mathbf{E}\|\zeta\|_C^4 &= \sum_{k=1}^r \lambda_k^2 \mathbf{E}\tau_k^2 + \sum_{1 \leq j \neq k \leq r} \lambda_j \lambda_k \mathbf{E}(\tau_j \tau_k) \\ &= \frac{3}{r(r+2)} \sum_{k=1}^r \lambda_k^2 + \frac{1}{r(r+2)} \left(\left(\sum_{k=1}^r \lambda_k \right)^2 - \sum_{k=1}^r \lambda_k^2 \right) \\ &= \frac{1}{r(r+2)} (2\text{Tr } C^2 + (\text{Tr } C)^2). \end{aligned} \tag{12.19}$$

The equalities (12.18) and (12.19) immediately imply that

$$\mathbf{Var} \|\zeta\|_C^2 = \mathbf{E}\|\zeta\|_C^4 - (\mathbf{E}\|\zeta\|_C^2)^2 = \frac{2}{r^2(r+2)} (r\text{Tr } C^2 - (\text{Tr } C)^2)$$

which coincides with (12.13). To prove (12.14), note that, by the geometric-arithmetic mean inequality,

$$\sum_{k=1}^r \lambda_k \tau_k \geq r \left(\prod_{k=1}^r \lambda_k \tau_k \right)^{1/r} = r(\det C)^{1/r} \prod_{k=1}^r \tau_k^{1/r}.$$

Therefore, for any $r \geq 3$,

$$\begin{aligned} \mathbf{E}\|\zeta\|_{\mathcal{C}}^{-2} &= \mathbf{E}\left(\sum_{k=1}^r \lambda_k \tau_k\right)^{-1} \leq \frac{1}{r(\det C)^{1/r}} \mathbf{E} \prod_{k=1}^r \tau_k^{-1/r} \\ &= \frac{1}{r(\det C)^{1/r}} \frac{\Gamma(r/2)}{(\Gamma(1/2))^r} \frac{(\Gamma(1/2 - 1/r))^r}{\Gamma(r/2 - 1)} \\ &= (\det C)^{-1/r} \left(\frac{1}{2} - \frac{1}{r}\right) \left(\frac{\Gamma(1/2 - 1/r)}{\Gamma(1/2)}\right)^r \end{aligned}$$

which yields (12.14), completing the proof.

Remark 12.1: Note that the \mathcal{C} -independent multiplier on the right of (12.14) is convergent,

$$\frac{r-2}{2r} \left(\frac{\Gamma(\frac{r-2}{2r})}{\Gamma(1/2)}\right)^r \rightarrow \frac{1}{2} \exp(-(\ln \Gamma)'(1/2)) \quad \text{as } r \rightarrow +\infty, \tag{12.20}$$

where $(\ln \Gamma)'(\lambda) = \Gamma'(\lambda)/\Gamma(\lambda)$ is the logarithmic derivative of the Euler gamma function.

Proof of Theorem 8.1: For any $X \in \mathcal{Z}_n$, let ζ_X be uniformly distributed on the unit sphere $\mathbb{S}_{m\#X}$. Then (7.6) reads

$$\mathbf{A}_o(W_X) = -\frac{1}{2} \ln \det C_X + \frac{m\#X}{2} \mathbf{E} \ln \|\zeta_X\|_{\mathcal{C}_X}^2. \tag{12.21}$$

Let us show that $\|\zeta_X\|_{\mathcal{C}_X}^2$ is mean square convergent as $X \nearrow \infty$. Applying (12.12) of Lemma 12.3 and recalling (7.4), obtain

$$\mathbf{E}\|\zeta_X\|_{\mathcal{C}_X}^2 = \frac{\text{Tr } C_X}{m\#X} = \frac{\text{Tr } c_0}{m}. \tag{12.22}$$

Combining (12.13) with (12.9) yields

$$\begin{aligned} \mathbf{Var} \|\zeta_X\|_{\mathcal{C}_X}^2 &= \frac{2}{m\#X + 2} \left(\frac{\text{Tr } C_X^2}{m\#X} - \left(\frac{\text{Tr } C_X}{m\#X}\right)^2 \right) \\ &\leq \frac{2 \text{Tr } C_X^2}{(m\#X)^2} \sim \frac{2}{m^2\#X} \sum_{z \in \mathbb{Z}^n} \|c_z\|_2^2 \rightarrow 0 \quad \text{as } X \nearrow \infty. \end{aligned}$$

This relation and (12.22) imply the convergence of $\|\zeta_X\|_{\mathcal{C}_X}^2$ to $\text{Tr } c_0/m$ in mean square sense and consequently, in probability,

$$\|\zeta_X\|_{\mathcal{C}_X}^2 \xrightarrow{p} \frac{\text{Tr } c_0}{m} \quad \text{as } X \nearrow \infty. \tag{12.23}$$

Recalling that under assumptions of the theorem, the spectra of the matrices C_X are all contained in (12.11), obtain that $\|\zeta_X\|_{\mathcal{C}_X}^2 \in \Delta$ for any $X \in \mathcal{Z}_n$, and hence, the random variables $|\ln \|\zeta_X\|_{\mathcal{C}_X}|$ are uniformly bounded. Therefore, by the Dominated Convergence Theorem, (12.23) implies

$$\lim_{X \nearrow \infty} \mathbf{E} \ln \|\zeta_X\|_{\mathcal{C}_X}^2 = \ln \frac{\text{Tr } c_0}{m}. \tag{12.24}$$

Assembling (12.10) of Lemma 12.2 and (12.24), from (12.21) obtain that

$$\frac{\mathbf{A}_o(W_X)}{\#X} \rightarrow -\frac{1}{2(2\pi)^n} \int_{\Omega_n} \ln \det S(\omega) d\omega + \frac{m}{2} \ln \frac{\text{Tr } c_0}{m} \quad \text{as } X \nearrow \infty$$

which gives (8.3), completing the proof.

Remark 12.2: If the random field W is not strictly regular, then the dominated convergence argument that was used to derive (12.24) from (12.23) fails. In this case, however, the condition

$$\liminf_{X \nearrow \infty} \frac{\ln \det C_X}{\#X} > -\infty, \quad (12.25)$$

if it holds, implies the uniform integrability of $\ln \|\zeta_X\|_{C_X}^2$ as $X \nearrow \infty$. Indeed, by the identity $\exp(|\ln u|) = \max(u, 1/u)$ for any $u > 0$,

$$\mathbf{E} \exp(|\ln \|\zeta_X\|_{C_X}^2|) \leq \mathbf{E} \|\zeta_X\|_{C_X}^2 + \mathbf{E} \|\zeta_X\|_{C_X}^{-2}.$$

Hence, applying (12.22), (12.14) and (12.20), obtain that

$$\begin{aligned} \limsup_{X \nearrow \infty} \mathbf{E} \exp(|\ln \|\zeta_X\|_{C_X}^2|) &\leq \frac{\text{Tr } c_0}{m} \\ &+ \frac{1}{2} \exp\left(-(\ln \Gamma)'(1/2) - \frac{1}{m} \liminf_{X \nearrow \infty} \frac{\ln \det C_X}{\#X}\right) \\ &< +\infty. \end{aligned}$$

By [20, Lemma 3 on p. 190], this last relationship implies the uniform integrability of $\ln \|\zeta_X\|_{C_X}^2$. Therefore, by [3, Theorem 5.4 on p. 32], fulfillment of (12.25) is enough to derive (12.24) from (12.23).

13 Proof of Theorem 11.1

Note that $\|F_X\|_\infty \leq \|F_Y\|_\infty$ for any $X, Y \in \mathcal{Z}_n$ satisfying $X \subseteq Y + z$ for some $z \in \mathbb{Z}^n$, and

$$\sup_{X \in \mathcal{Z}_n} \|F_X\|_\infty = \|F\|_\infty. \quad (13.1)$$

Furthermore,

$$\lim_{X \nearrow \infty} \frac{\|F_X\|_2}{\sqrt{\#X}} = \|F\|_2. \quad (13.2)$$

Indeed, using (8.1) and (9.4), obtain that

$$\begin{aligned} \frac{\|F_X\|_2^2}{\#X} &= \frac{1}{\#X} \sum_{x, y \in X} \|f_{x-y}\|_2^2 = \sum_{z \in \mathbb{Z}^n} D_X(z) \|f_z\|_2^2 \\ &\rightarrow \sum_{z \in \mathbb{Z}^n} \|f_z\|_2^2 = \|F\|_2^2 \quad \text{as } X \nearrow \infty. \end{aligned}$$

For every $X \in \mathcal{Z}_n$, associate with (11.1) a positive semi-definite symmetric matrix $H_X \in \mathbb{R}^{m\#X \times m\#X}$ as

$$H_X = F_X^T F_X = \text{block}_{x, y \in X} \left(\sum_{z \in X} f_{z-x}^T f_{z-y} \right). \quad (13.3)$$

Lemma 13.1: *Let the impulse response function of $F \in \mathcal{L}_\infty^{p \times m, n}$ be absolutely summable, i.e. $\sum_{x \in \mathbb{Z}^n} \|f_x\|_\infty < +\infty$. Then for any $r \in \mathbb{N}$, the matrices (13.3) satisfy*

$$\lim_{X \nearrow \infty} \frac{\text{Tr } H_X^r}{\#X} = (2\pi)^{-n} \int_{\Omega_n} \text{Tr} (H(\omega))^r d\omega, \tag{13.4}$$

where the map H is defined by (9.5).

Proof: Introduce the function $\varphi : \Omega_{2rn} \rightarrow \mathbb{C}$ which maps a vector $\omega = (\omega_k)_{1 \leq k \leq 2r}$, formed by $\omega_1, \dots, \omega_{2r} \in \Omega_n$, to

$$\begin{aligned} \varphi(\omega) &= \text{Tr} ((F(\omega_1))^* F(\omega_2) \times \dots \times (F(\omega_{2r-1}))^* F(\omega_{2r})) \\ &= \sum_{z \in \mathbb{Z}^{2rn}} \psi_z \exp(-i\omega^T z). \end{aligned} \tag{13.5}$$

Here, F is the transfer function defined in (9.3), and for any $z = (z_k)_{1 \leq k \leq 2r}$ formed by $z_1, \dots, z_{2r} \in \mathbb{Z}^n$, the Fourier coefficient ψ_z is given by

$$\psi_z = (2\pi)^{-2rn} \int_{\Omega_{2rn}} \exp(i\omega^T z) \varphi(\omega) d\omega = \text{Tr} \left(f_{-z_1}^T f_{z_2} \times \dots \times f_{-z_{2r-1}}^T f_{z_{2r}} \right). \tag{13.6}$$

From (13.3) and (13.6), it follows that

$$\begin{aligned} \text{Tr } H_X^r &= \sum_{x_1, \dots, x_r, y_1, \dots, y_r \in X} \text{Tr} \left(f_{y_1-x_1}^T f_{y_1-x_2} f_{y_2-x_2}^T f_{y_2-x_3} \right. \\ &\quad \left. \times \dots \times f_{y_{r-1}-x_{r-1}}^T f_{y_{r-1}-x_r} f_{y_r-x_r}^T f_{y_r-x_1} \right) \\ &= \sum_{x_1, \dots, x_r, y_1, \dots, y_r \in X} \psi_{(x_1-y_1, y_1-x_2, \dots, x_r-y_r, y_r-x_1)} \end{aligned} \tag{13.7}$$

Hence, recalling (12.1) and (12.3), obtain

$$\frac{\text{Tr } H_X^r}{\#X} = \sum_{z \in \mathbb{O}_{2r}} \psi_z E_{X, 2r}(z). \tag{13.8}$$

By the same reasoning as in the proof of Lemma 12.1, the absolute summability of the impulse response function implies a similar property for the Fourier coefficients (13.6),

$$\sum_{z \in \mathbb{Z}^{2rn}} |\psi_z| \leq m \left(\sum_{x \in \mathbb{Z}^n} \|f_x\|_\infty \right)^{2r}.$$

In particular, $\sum_{z \in \mathbb{O}_{2r}} |\psi_z| < +\infty$, thus validating the passage to the limit under the sum in (13.8) on a basis of (12.2),

$$\lim_{X \nearrow \infty} \frac{\text{Tr } H_X^r}{\#X} = \sum_{z \in \mathbb{O}_{2r}} \psi_z \lim_{X \nearrow \infty} E_{X, 2r}(z) = \sum_{z \in \mathbb{O}_{2r}} \psi_z. \tag{13.9}$$

It now remains to note that by (9.5), (12.3), (13.5) and (13.6),

$$\sum_{z \in \mathbb{O}_{2r}} \psi_z = (2\pi)^{-n} \int_{\Omega_n} \underbrace{\varphi(\omega, \dots, \omega)}_{2r \text{ times}} d\omega = (2\pi)^{-n} \int_{\Omega_n} \text{Tr} (H(\omega))^r d\omega,$$

which, in combination with (13.9), immediately implies (13.4), thereby concluding the proof.

Under the assumptions of Lemma 13.1, for any $r \in \mathbb{N}$,

$$\lim_{X \nearrow \infty} \left(\frac{\text{Tr } H_X^r}{m\#X} \right)^{1/r} = \left(\frac{1}{m(2\pi)^n} \int_{\Omega_n} \text{Tr}(H(\omega))^r d\omega \right)^{1/r}. \quad (13.10)$$

By the well-known properties of Hölder norms, for any positive semi-definite Hermitian $A \in \mathbb{C}^{d \times d}$, the quantity $(\text{Tr } A^r / d)^{1/r}$ is nondecreasing in $r \in \mathbb{N}$ and converges to $\lambda_{\max}(A)$ as $r \rightarrow +\infty$. Applying the monotone convergence argument to both sides of (13.10), obtain the following refinement of (13.1),

$$\lim_{X \nearrow \infty} \|F_X\|_\infty = \|F\|_\infty. \quad (13.11)$$

Proof of Theorem 11.1: Recalling (13.3), for any $X \in \mathcal{Z}_n$ introduce the functions

$$\mathcal{A}_X(q) = \frac{m\#X}{2} (\ln \Phi_X(q) - \Psi_X(q)), \quad (13.12)$$

$$\mathcal{N}_X(q) = \left(\frac{1}{q} \left(1 - \frac{1}{\Phi_X(q)} \right) \right)^{1/2}, \quad (13.13)$$

$$\Phi_X(q) = \frac{\text{Tr}(I_{m\#X} - qH_X)^{-1}}{m\#X}, \quad (13.14)$$

$$\Psi_X(q) = -\frac{\ln \det(I_{m\#X} - qH_X)}{m\#X}. \quad (13.15)$$

These are similar to (10.2)–(10.5) and well-defined on the interval

$$Q_X = [0, \|F_X\|_\infty^{-2})$$

which, by (13.1), contains (10.1). Comparison of (13.2) with (13.11) shows that the nonroundness of the operator F implies

$$\|F_X\|_2 < \sqrt{m\#X} \|F_X\|_\infty \quad (13.16)$$

for all $X \in \mathcal{Z}_n$ large enough in the sense of van Hove convergence to infinity. For any such X , the function $\mathcal{A}_X : Q_X \rightarrow \mathbb{R}_+$ is strictly increasing and has a well-defined inverse $\mathcal{A}_X^{-1} : \mathbb{R}_+ \rightarrow Q_X$. Application of [23, Theorem 4] to the matrix (11.1) gives that for any $a \in \mathbb{R}_+$,

$$\|F_X\|_{a\#X} = \mathcal{N}_X(\mathcal{A}_X^{-1}(a\#X)). \quad (13.17)$$

Expanding (13.14) and (13.15) to the Taylor series

$$\begin{aligned} \Phi_X(q) &= \frac{1}{m\#X} \sum_{r \in \mathbb{Z}_+} \text{Tr}(qH_X)^r, \\ \Psi_X(q) &= \frac{1}{m\#X} \sum_{r \in \mathbb{N}} \frac{\text{Tr}(qH_X)^r}{r}, \end{aligned}$$

whose convergence is uniform in $X \in \mathcal{Z}_n$ for any $q \in Q$, and using Lemma 13.1, obtain that $\Phi_X(q)$ and $\Psi_X(q)$ (and so also $\mathcal{A}_X(q)/\#X$ and $\mathcal{N}_X(q)$) converge to $\Phi(q)$ and $\Psi(q)$

(respectively, to $\mathcal{A}(q)$ and $\mathcal{N}(q)$) as $X \nearrow \infty$. Moreover, a slightly refined reasoning shows that the convergence holds for the derivatives of the functions as well and is uniform in q over any compact subset of the interval (10.1). Therefore,

$$\mathcal{A}_X^{-1}(a\#X) = \left(\frac{\mathcal{A}_X}{\#X}\right)^{-1}(a) \rightarrow \mathcal{A}^{-1}(a) \quad \text{as } X \nearrow \infty. \quad (13.18)$$

Combining the last property with (13.17) and recalling (10.6) of Theorem 10.1, obtain that

$$\lim_{X \nearrow \infty} \|F_X\|_{a\#X} = \mathcal{N}(\mathcal{A}^{-1}(a)) = \|F\|_a. \quad (13.19)$$

This proves the right-most equality in (11.2). By Lemma 6.1, the left-most equality in (11.2) will be proved if we show that

$$\liminf_{X \nearrow \infty} \|F_X\|_{a\#X, \circ} \geq \|F\|_a. \quad (13.20)$$

For this purpose, consider a random vector $W_{X,a} \in \mathbb{G}^{m\#X}(\Sigma_{X,a})$ where

$$\Sigma_{X,a} = (I_{m\#X} - q_{X,a}H_X)^{-1}, \quad q_{X,a} = \mathcal{A}_X^{-1}(a\#X). \quad (13.21)$$

By [23, Theorem 4], $W_{X,a}$ is a worst-case input for F_X in the sense of the $(a\#X)$ -anisotropic norm (6.2) associated with the root mean square gain (6.4). That is,

$$\|F_X\|_{a\#X} = \mathbf{N}(F_X, W_{X,a}), \quad \mathbf{A}(W_{X,a}) = a\#X.$$

Applying to $W_{X,a}$ the inequality (4.5), obtain that $\mathbf{A}_\circ(W_{X,a}) \leq a\#X$. Consequently, the average energy gain of F_X wrt $W_{X,a}$ provides a lower bound for the $(a\#X)$ -anisotropic norm of F_X defined by (6.1), i.e.

$$\|F_X\|_{a\#X, \circ} \geq \mathbf{N}_\circ(F_X, W_{X,a}). \quad (13.22)$$

To compute the gain, write $W_{X,a} = \Sigma_{X,a}^{1/2}V_X$ where $V_X \in \mathbb{G}^{m\#X}(I_{m\#X})$. By (13.3) and (13.21), the corresponding output $Z_{X,a} = F_XW_{X,a}$ satisfies

$$|Z_{X,a}|^2 = \|W_{X,a}\|_{H_X}^2 = \|W_{X,a}\|_{(I_{m\#X} - \Sigma_{X,a}^{-1})/q_{X,a}}^2 = \frac{|W_{X,a}|^2 - |V_X|^2}{q_{X,a}}.$$

Hence,

$$\begin{aligned} \mathbf{N}_\circ(F_X, W_{X,a}) &= \sqrt{\mathbf{E}(|Z_{X,a}|/|W_{X,a}|)^2} \\ &= \sqrt{\frac{1 - \mathbf{E}(|V_X|/|W_{X,a}|)^2}{q_{X,a}}} \\ &= \sqrt{\frac{1 - \mathbf{E}\|\zeta_X\|_{\Sigma_{X,a}}^{-2}}{q_{X,a}}} \end{aligned} \quad (13.23)$$

where $\zeta_X = V_X/|V_X|$ is uniformly distributed on the unit sphere $\mathbb{S}_{m\#X}$. Using the convergence

$$\lim_{X \nearrow \infty} q_{X,a} = \mathcal{A}^{-1}(a) \quad (13.24)$$

which by (13.21) replicates (13.18), one verifies, similarly to the proof of Theorem 8.1, that

$$\begin{aligned} \|\zeta_X\|_{\Sigma_{X,a}}^2 &\xrightarrow{p} \frac{1}{m(2\pi)^n} \int_{\Omega_n} \text{Tr} (I_m - \mathcal{A}^{-1}(a)H(\omega))^{-1} d\omega \\ &= \Phi(\mathcal{A}^{-1}(a)) \quad \text{as } X \nearrow \infty. \end{aligned} \quad (13.25)$$

Since $\Sigma_{X,a} - I_{m\#X}$ is positive semi-definite, $\|\zeta_X\|_{\Sigma_{X,a}} \geq 1$. By the Dominated Convergence Theorem, (13.25) then implies that

$$\lim_{X \nearrow \infty} \mathbf{E} \|\zeta_X\|_{\Sigma_{X,a}}^{-2} = 1/\Phi(\mathcal{A}^{-1}(a)).$$

Plugging the last convergence and (13.24) in (13.23), obtain

$$\lim_{X \nearrow \infty} \mathbf{N}_o(F_X, W_{X,a}) = \sqrt{\frac{1}{\mathcal{A}^{-1}(a)} \left(1 - \frac{1}{\Phi(\mathcal{A}^{-1}(a))}\right)} = \mathcal{N}(\mathcal{A}^{-1}(a)) = \|F\|_a.$$

In combination with (13.22), this implies (13.20), thereby concluding the proof.

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