

## at 30, 40, 80 and 158 A·GeV

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A non-monotonic energy dependence of the  $K^+/\pi^+$  ratio with a sharp maximum close to 30 A·GeV is observed in central Pb+Pb collisions [1]. Within a statistical model of the early stage [2], this is interpreted as a sign of the phase transition to a QGP, which causes a sharp change in the energy dependence of the strangeness to entropy ratio. This observation naturally motivates us to study the production of multistrange hyperons ( $\Xi$ ,  $\Omega$ ) as a function of the beam energy.

Furthermore it was suggested that the kinematic freeze-out of  $\Omega$  takes place directly at QGP hadronization. If this is indeed the case, the transverse momentum spectra of the  $\Omega$  directly reflect the transverse expansion velocity of a hadronizing QGP [3, 4].

In this report we show preliminary NA49 results on  $\Omega^-$  and  $\bar{\Omega}^+$  production in central Pb+Pb collisions at 40 and 158 A·GeV and compare them to measurements of  $\Xi^-$  and  $\bar{\Xi}^+$  production in central Pb+Pb collisions at 30, 40, 80 and 158 A·GeV.

The NA49 detector [5] is a large acceptance hadron spectrometer at the CERN SPS, consisting of four TPCs. Two of them, the Vertex TPCs (VTPC), are inside a magnetic field for the determination of particle momenta and charge. The ionisation energy loss (dE/dx) measurements in the two Main TPCs (MTPC), which are outside the magnetic field, are used for mass determination. Central collisions were selected by a trigger using information from a downstream calorimeter (VCAL), which measures the energy of the projectile spectator nucleons.

In Fig.1 the NA49  $\bar{\Omega}^+/\Omega^-$  and  $\bar{\Xi}^+/\Xi^-$  ratios as a function of the collision energy ( $\sqrt{s_{NN}}$ ) are shown and compared to results of NA57 [6, 7] and STAR [8, 9]. The NA49 and NA57 results measured at the same energies are consistent. The data show a clear increase of the  $\bar{\Omega}^+/\Omega^-$  ratio from a value of about 0.4 at SPS energies to about 1 at RHIC energies. The  $\bar{\Xi}^+/\Xi^-$  ratio also increases from SPS energies to about

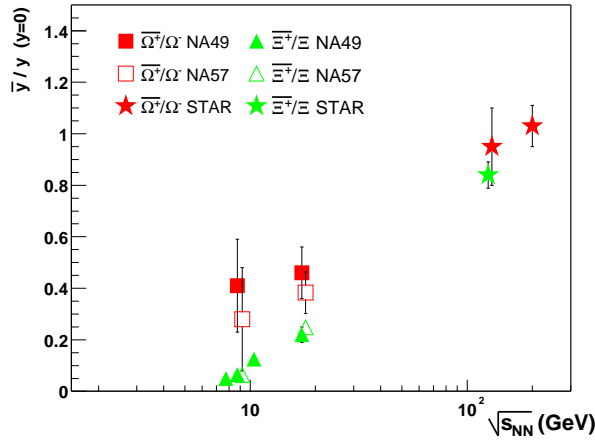


Figure 1: The antihyperon/hyperon ( $\bar{Y}/Y$ ) ratio at midrapidity in the SPS-RHIC energy range.

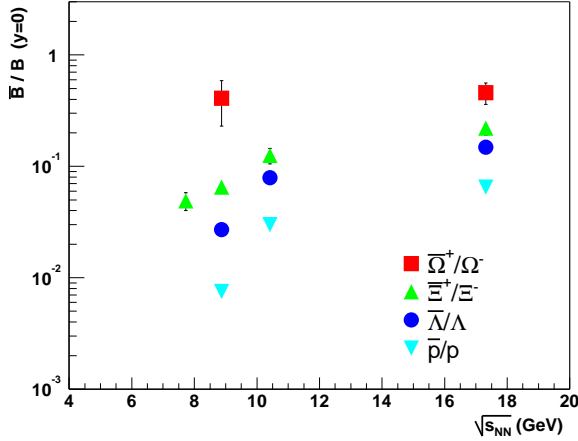


Figure 2: The antibaryon/baryon ratio ( $\bar{B}/B$ ) at midrapidity in the SPS energy range measured by NA49.

The midrapidity  $\bar{\Omega}^+/\Omega^-$  ratio is estimated to be  $0.46 \pm 0.1$  and  $0.41 \pm 0.18$  for central Pb+Pb collisions at 158 and 40 A·GeV, respectively. The values for the midrapidity  $\bar{\Xi}^+/\Xi^-$  ratio are estimated to be  $0.13 \pm 0.02$ ,  $0.065 \pm 0.05$  and  $0.0049 \pm 0.009$  for central Pb+Pb collisions at 80, 40 and 30 A·GeV, respectively. In Fig.2 the antibaryon/baryon ratios are shown as a function of the beam energy in the SPS energy domain. In addition to  $\bar{\Xi}/\Xi$  and  $\bar{\Omega}/\Omega$  ratios the results on  $\bar{\Lambda}/\Lambda$  [10] and  $\bar{p}/p$  [11] are shown. The energy dependence of  $\bar{B}/B$  ratios gets weaker with increasing strangeness content.

Fig. 3 shows the energy dependence of the midrapidity

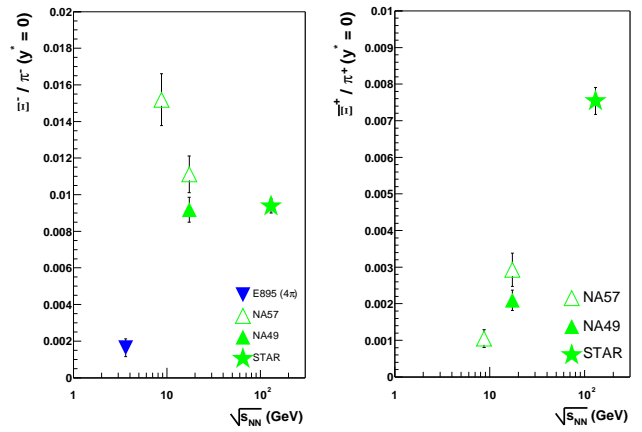


Figure 3: Energy dependence of the midrapidity  $\Xi^-/\pi^-$  (left) and  $\bar{\Xi}^+/\pi^+$  (right) ratio in central Pb+Pb and Au+Au collisions.

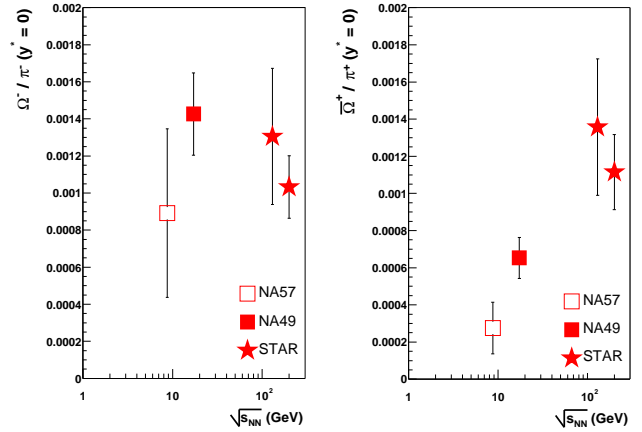


Figure 4: Energy dependence of the midrapidity  $\Omega^-/\pi^-$  (left) and  $\bar{\Omega}^+/\pi^+$  (right) ratio in central Pb+Pb and Au+Au collisions.

$\Xi^-/\pi^-$  (left) and  $\bar{\Xi}^+/\pi^+$  (right) ratio in central Pb+Pb and Au+Au collisions. The  $\Xi^-/\pi^-$  ratio suggests that there is a non-monotonic energy dependence at SPS energies. The  $\bar{\Xi}^+/\pi^+$  ratio increases with energy from SPS to RHIC energies. The  $\Omega^-/\pi^-$  ratio shown in Fig. 4 seems to be energy independent, but the  $\bar{\Omega}^+/\pi^+$  ratio shows again an increase from SPS to RHIC energies.

At 158 A·GeV, a high statistics data sample of central Pb+Pb collisions is available, which allows us to obtain fully corrected spectra of  $\Omega^-$  and  $\bar{\Omega}^+$ . The transverse mass spectra are fitted by an exponential function :

$$\frac{1}{m_t} \frac{d^2 N}{dm_t dy} = C \cdot e^{-m_t/T}, \quad (1)$$

where the fit parameters are a normalization factor  $C$  and the inverse slope parameter  $T$ . The slope parameter is similar for  $\Omega^-$  and  $\bar{\Omega}^+$  :  $T(\Omega^-) = 276 \pm 23$  MeV and  $T(\bar{\Omega}^+) = 285 \pm 39$  MeV [11]. Our values

agree with those measured by the NA57 collaboration ( $T(\Omega^-) = 280 \pm 16$  MeV and  $T(\bar{\Omega}^+) = 324 \pm 29$  MeV) [12]. The large acceptance of the NA49 experiment allows us to measure the  $\Omega^-$  ( $\bar{\Omega}^+$ ) spectra in a large rapidity interval. The rapidity distributions for  $\Omega^-$  and  $\bar{\Omega}^+$  are obtained by extrapolating  $p_t$  spectra using the exponential. Both  $y$ -spectra were fitted by a Gaussian. The width of the  $\Omega^-$  distribution ( $\sigma(\Omega^-) = 1.0 \pm 0.2$ ) seems to be larger than the one of the  $\bar{\Omega}^+$  ( $\sigma(\bar{\Omega}^+) = 0.7 \pm 0.1$ ). Mean multiplicities in full phase-space were estimated as integrals over measured points corrected for the missing rapidity coverage using the Gaussian parametrisations. The resulting yields are  $\langle\Omega^-\rangle = 0.47 \pm 0.07$  and  $\langle\bar{\Omega}^+\rangle = 0.15 \pm 0.02$ , where the errors are statistical only.

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