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# Power flow in the interior and exterior of cylindrical coated nano-particles

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## Abstract

Electrically small configurations consisting of resonant active coated nano-particles of a circularly cylindrical shape are considered in the case of a magnetic line source excitation at optical frequencies. The active particles are comprised of a canonical gain-impregnated silica nano-core covered concentrically with a silver, gold, or copper nano-shell. Particular attention is devoted to both the direction and the magnitude of the power flow density inside and outside of such particles. The results for the active coated nano-particles are related to those of the corresponding passive designs at optical frequencies, and are further contrasted to the results for the electrically small structures based on epsilon-negative materials at radio frequencies.

## 1. Introduction

Numerous potential applications exploiting metamaterials (MTMs), as well as plasmonic materials, have been proposed [1, 2], with particular attention devoted to sub-wavelength MTM-based [3]-[5] and plasmonic-based [6]-[10] particles. The gain inclusion in the active designs was found to lead to the so-called super-resonant states, which for a plane wave incidence resulted in scattering cross-sections vastly surpassing values predicted by the geometrical size of the CNPs [6], [10]. Similarly, enormously large radiated powers were obtained in the case of localized excitation sources [8], [9]. These are all desirable features which make the plasmonic-based CNPs attractive candidates for nano-sensing, -lasing and -amplifying applications; see e.g., a recent review [11] and the works references therein for an excellent overview of applications and potentials of plasmonic-based devices.

To better understand the interaction between the localized sources and plasmonic-based particles, and their overall radiation process, the present work investigates the power flow associated with various types of circularly cylindrical active coated nano-particles (CNPs) excited by a magnetic line source (MLS) at optical frequencies. Specifically, the active particles are comprised of a gain-impregnated silica nano-core covered concentrically with a silver, gold, or copper nano-shell where the size-dependency of the permittivity of this nano-shell is taken into account. The gain model for the active CNPs is a canonical, constant

frequency, gain model. The previously reported super-resonant states of the proposed active CNPs in [9], and the associated large levels of the radiated power, are correlated with the correspondingly large localizations of the power density near specific parts of the shell-interfaces and its corresponding large magnitude inside the gain region. The resulting interactions with the gain material lead to a strong and distinct dipolar power radiation in all of these electrically small, active super-resonant CNP designs. These active CNP power flow results are compared to those obtained for the corresponding passive structures at optical frequencies, and are further contrasted to epsilon-negative (ENG) sub-wavelength resonator designs at radio frequencies. The latter sheds further light on important differences between the optical and radio frequency cases. We note that interesting power flow results were reported in [5] for MTM-based cylindrical configurations for a finite radius line source excitation, as well as in [10] for active, and [11], [12] for passive nano-particles of spherical shape, against which the present findings for the cylindrical active CNPs will be displayed. Throughout the work, the time factor  $\exp(j\omega t)$ , with  $\omega$  being the angular frequency, and  $t$  being the time, is assumed and suppressed.

## 2. Configuration and methods of analysis

We consider the CNP configuration depicted in Figure 1; it consists of an infinite circular cylindrical nano-core (Region 1) covered concentrically with an infinite cylindrical nano-shell (Region 2). The inner and outer radii of the nano-shell are denoted by  $\rho_1$  and  $\rho_2$ , respectively, and the CNP is imbedded in free-space (region 3) with the permittivity,  $\epsilon_0$ , permeability,  $\mu_0$ , and wave number  $k_0 = \omega\sqrt{\epsilon_0\mu_0} = 2\pi/\lambda$ , with  $\lambda$  being the free-space wavelength. The CNP is illuminated by the field provided by an arbitrarily located and infinitely long magnetic line source (MLS) possessing a constant magnetic current  $I_m$  [V/m]. The nano-core and nano-shell regions of the CNP may, in general, contain an arbitrary simple, lossy and dispersive material. However, for the present purposes, specific materials are chosen for these two regions, as explained later in Section 3 of this manuscript.

A circular cylindrical coordinate system  $(\rho, \phi, z)$  and an associated rectangular coordinate system  $(x, y, z)$  are introduced such that their origins coincide with the cross-sectional center of the CNP, and such that the entire CNP-MLS configuration is infinite in the  $z$ -direction, with the MLS being parallel to the axis of the cylinders. The coordinates of the observation point are  $(\rho, \phi)$ , while those of the MLS are  $(\rho_s, \phi_s)$ .

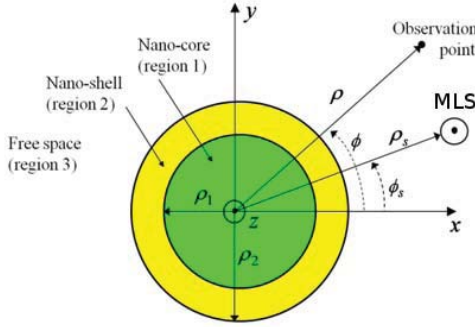


Figure 1: The MLS-excited CNP.

The analytical solution for the problem in Figure 1 has been derived in [4], and only its main features will be outlined in here. For the MLS excitation case, the magnetic field in all regions (incident and scattered fields in the region containing the MLS, and total fields in the remaining regions) contains only a  $z$ -component, i.e.,  $\mathbf{H} = \hat{z}H_z$ , while the electric field in all regions has the form of  $\mathbf{E} = \hat{\rho}E_\rho + \hat{\phi}E_\phi$ , where  $E_\rho$  and  $E_\phi$  are its  $\rho$ - and  $\phi$ -components. The unknown field components are expanded in cylindrical wave functions, and thus contain a number of unknown expansion coefficients. The latter follow in a straightforward manner from the application of the electromagnetic boundary conditions at the two interfaces of the CNP. These field solutions were used in [4] to perform an extensive study of electrically small passive MTM-based particles of cylindrical shape. To account for the near- and far-field properties of a number of cylindrical active plasmonic-based CNP configurations, the necessary gain levels for the so-called super-resonant state were determined recently [9]. To further address the intriguing electromagnetic properties of these particles, the field solutions from [4] are presently used to explore the flow of electromagnetic power inside and outside the CNP. This is done in terms of the Poynting vector  $\mathbf{S}$ , which is given by

$$\begin{aligned} \mathbf{S} &= (1/2) \text{Re}\{\mathbf{E} \times \mathbf{H}^*\} = \hat{x}S_x + \hat{y}S_y \\ &= \hat{x}[S_\rho \cos \phi - S_\phi \sin \phi] + \hat{y}[S_\rho \sin \phi + S_\phi \cos \phi], \end{aligned} \quad (1)$$

where  $S_x$  and  $S_y$  are the  $x$ - and  $y$ -components of  $\mathbf{S}$ , while  $S_\rho = \text{Re}\{E_\phi H_z^*\}/2$  and  $S_\phi = -\text{Re}\{E_\rho H_z^*\}/2$  are its  $\rho$ - and  $\phi$ -components. The asterisk  $*$  designates complex

conjugation. In Section 3, attention will be devoted to both the amplitude and the direction of the Poynting vector determined with (1).

### 3. Results and discussions

As reported in [6], [8]-[10], highly resonant sub-wavelength CNPs can be designed by including gain in their interiors, e.g., their nano-cores. As to the cylindrical CNPs considered in [9], the nano-core was in all cases made of silica, had a radius of 24 nm, and was coated with a 6 nm thick plasmonic (silver, gold, or copper) nano-shell. The permittivity of the silica nano-core was expressed as  $\epsilon_1 = (n^2 - \kappa^2 - 2jn\kappa)\epsilon_0$  and was comprised of two contributions: a contribution from its refractive index in the considered frequency region ( $n = \sqrt{2.05}$ ) and a contribution from the canonical gain model. The parameter  $\kappa$  (termed the optical loss or gain constant) determines the nature of the nano-core: lossless (and passive) for  $\kappa = 0$ , lossy and passive for  $\kappa > 0$  (optical loss constant) and active for  $\kappa < 0$  (optical gain constant). As to the plasmonic nano-shell permittivity, its empirical values, which account for its size-dependence and were used successfully in [6], were employed. In particular, for the silver-based CNP configuration excited by a MLS, the so-called super-resonant state, with a large total radiated power, was identified for  $\kappa = -0.175$  and the free-space excitation wavelength  $\lambda = 577.70$  nm [9]. This total radiated power enhancement is shown in Figure 2 in terms of the normalized radiation resistance (NRR) as a function of the wavelength when the MLS is located at 12 nm along the positive  $x$ -axis and has the current amplitude:  $I_m = 1$  [V/m].

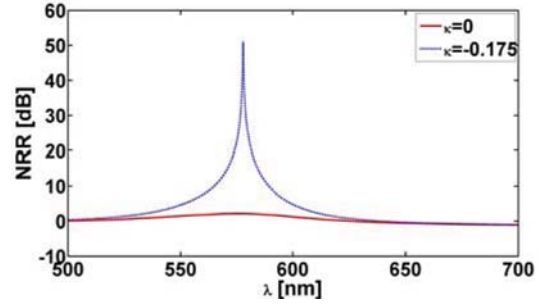


Figure 1: The NRR of a silver-based passive and super-resonant active CNP.

Note that the NRR equals the total power radiated by the MLS in the presence of the CNP normalized with the power radiated by the MLS alone in free-space. As can be seen in the figure, a peak value of the NRR  $\approx 51$  dB is found for the super-resonant active CNP at 577.70 nm. This should be compared to a NRR of around 2 dB for the corresponding lossless and passive CNP ( $\kappa = 0$ ), which also is included in the figure. As explained in [9] in terms of the near-field distributions of the magnetic field, the super-resonance behavior reported in Figure 2 for the active silver-based CNP is due to a strong excitation of the resonant dipole

mode inside the CNP. The existence of the resonant dipole mode inside an electrically small CNP (diameter size of approximately  $\lambda/10$ ) is a result of juxtaposing of two materials, one with positive permittivity (silica nano-core) and the other with a negative permittivity (silver nano-shell) [1]. The strong excitation of this dipole mode is heavily assisted by the gain-impregnation of the silica-nano core; its presence overcomes the plasmonic material losses [6].

The super-resonant state of the active Ag-based CNP is next illustrated in Figure 3a). The magnitude (color) and direction (arrows) of the Poynting vector (1) is shown when the MLS is located at 12 nm along the positive  $x$ -axis.

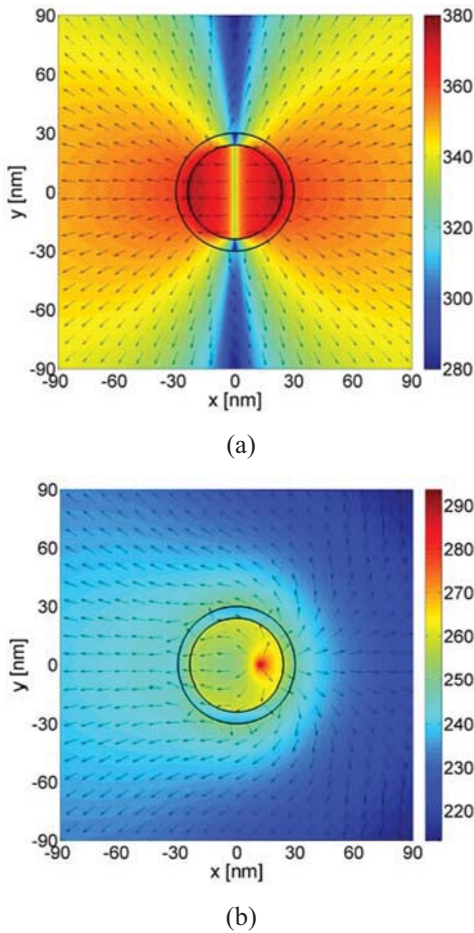


Figure 3: The magnitude (color) and direction (arrows) of the power flow density for the super-resonant active Ag-based CNP (a) and for the corresponding passive Ag-based CNP (b). In both cases, the radius of the silica nano-core is 24 nm, the thickness of the plasmonic nano-shell is 6 nm, and the MLS is located at 12 nm along the positive  $x$ -axis. Note that the dynamic range in (a) is larger than that in (b). The curves representing the cylindrical surfaces of the CNP are included in the figure.

A very clear and strong dipolar pattern, as though it is being produced by a dipolar source of radiation centered at the origin of the CNP, is observed in this figure. The power density is rather strong and homogeneous in strength and

direction inside the CNP. It is mainly directed along the negative and positive  $x$ -direction in the left and right halves, respectively, of the nano-core. This is particularly true near the inner surface of the nano-shell, where one witnesses a profound localization of the power density within the CNP. The interaction of the large fields in the electrically small resonant CNP with the canonical gain inside the nano-core leads to a clear and strong dipolar radiated power pattern. This radiated power behavior is illustrated with by an outward-propagating Poynting vector from the surface of the CNP in Figure 3a). While the magnitude of the Poynting vector suggests the pattern is symmetric, its direction reveals a slight asymmetry which is attributed to the off-set MLS location in this case. In contrast to the super-resonant Ag-based CNP results, Figure 3b) shows the magnitude and the direction of the corresponding passive Ag-based CNP. No excitation of a resonant dipole mode is in evidence; this is consistent with the very low radiated power levels observed, e.g., in Figure 2.

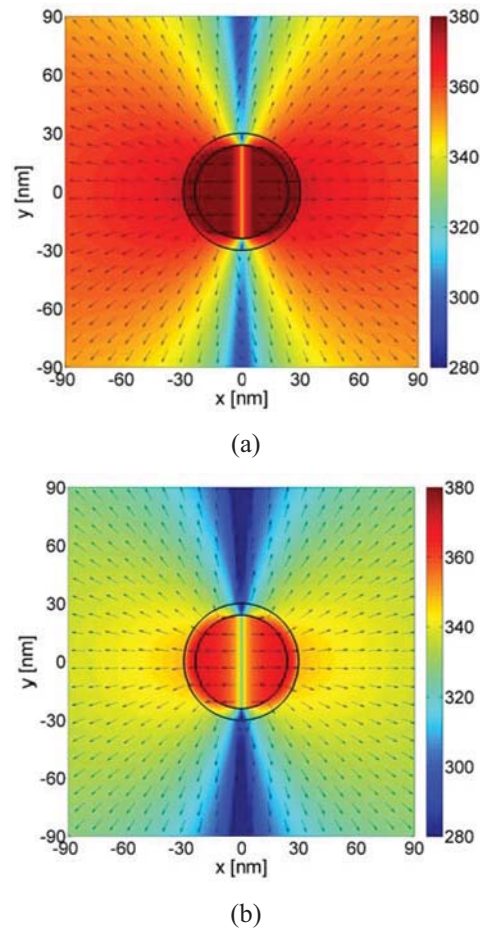


Figure 4: The magnitude (color) and direction (arrows) of the power flow density for the super-resonant active Au-based CNP (a) and Cu-based CNP (b). In both cases, the radius of the silica nano-core is 24 nm, the thickness of the plasmonic nano-shell is 6 nm, and the MLS is located at 12 nm along the positive  $x$ -axis. The curves representing the cylindrical surfaces of the CNP are included in the figure.

Besides the Ag-based CNP configuration considered above,

the corresponding super-resonant CNPs with their nano-shell composed of either gold or copper have also been studied [9]. For the Au-based CNP, the super-resonance was identified for  $\kappa = -0.262$  and  $\lambda = 669.39$  nm, with a resulting NRR  $\approx 56$ dB. For the Cu-based CNP, it was identified for  $\kappa = -0.310$  and  $\lambda = 662.19$  nm, with a resulting NRR  $\approx 46$ dB. The behavior of the latter two CNPs with the MLS located at 12 nm along the positive  $x$ -axis is illustrated further with the power flow density results reported for the Au-based CNP in Figure 4a) and for the Cu-based CNP in Figure 4b). Again, the power flow clearly has a dipolar form. However, its magnitude for the Au-based CNP surpasses the values for the other two CNPs given in Figures 3a) and 4b). Because it has the highest losses, the corresponding value for the super-resonant Cu-based CNP is the lowest of all of the super-resonant cases. The large levels of the power flow density for the Au-based CNP correspond to a strong localization of the power density within the CNP, which in turn leads to the largest NRR values of the three super-resonant CNP cases. The relatively low levels for the Cu-based CNP are perfectly in line with the correspondingly lowest NRR values.

We note that for all super-resonant CNPs cases considered above, the power flow has also been observed in an extended range up to 200 nm in both the  $x$ - and  $y$ -directions. Throughout such a range, a clear dipolar pattern, suggesting an outwardly propagating power flow density away from the CNP, was observed. Moreover, power flow results were obtained for super-resonant CNPs ( $\kappa = -0.175$ ) for wavelengths shorter ( $\lambda = 567.70$ ) and longer ( $\lambda = 587.70$ ) than the super-resonant state wavelength ( $\lambda = 577.70$ ). Although they are not included here, these results still show dipolar patterns with an outwardly propagating power flow density similar to those found in Figures 3a) and 4, but with a lower strength. Thus, contrary to power flows for active [10] and passive [13] nano-particles of spherical shape, which were found to exhibit specific vortices near their plasmon resonances, this is not the case for the (infinite and, hence, effectively two-dimensional) active *super-resonant* CNPs of cylindrical shape investigated here.

The above results for the electrically small super-resonant cylindrical CNPs, which employ the naturally occurring ENG-based metals at optical frequencies, are next contrasted to the analogous electrically small structures at radio frequencies which must utilize ENG metamaterials. To this end, Figure 5a) shows the NRR as a function of frequency,  $f$ , for a resonant sub-wavelength ENG shell with an inner radius of 6 mm and an outer radius of 10.03 mm. The ENG shell is modeled by a lossless Drude dispersion model [1] adjusted to select the relative permittivity of the shell to be  $-3.0$  at the design frequency of 300 MHz (giving the diameter of the particle to be  $\lambda/50$ ). The so-selected geometrical and electrical parameters of the shell are in compliance with the well-known resonance conditions in electrically small structures of cylindrical shape [3], [4]. Figure 5a), which was obtained when the MLS was located along the  $x$ -axis at 3mm, shows a resonance phenomenon with the NRR  $\approx 17$

dB at 300 MHz. The associated power flow density (magnitude and direction) is depicted in Figure 5b). It is found to differ significantly from any of the super-resonant active CNP cases. In this case, the power flow lines are seen to exhibit an interesting behavior, including vortices, at the three “deep-blue” spots – one to the left of the particle and two inside the shell-region. In fact, the power flow *outside* of the resonant ENG-based particle resembles the one for the lossless and passive CNP in Figure 3b), rather than that of the super-resonant active CNP. In the latter, the significantly smaller physical size of the particle, as well as the gain inclusion, are the driving factors behind the clear and strong outward power flow and the dipolar shape of the response.

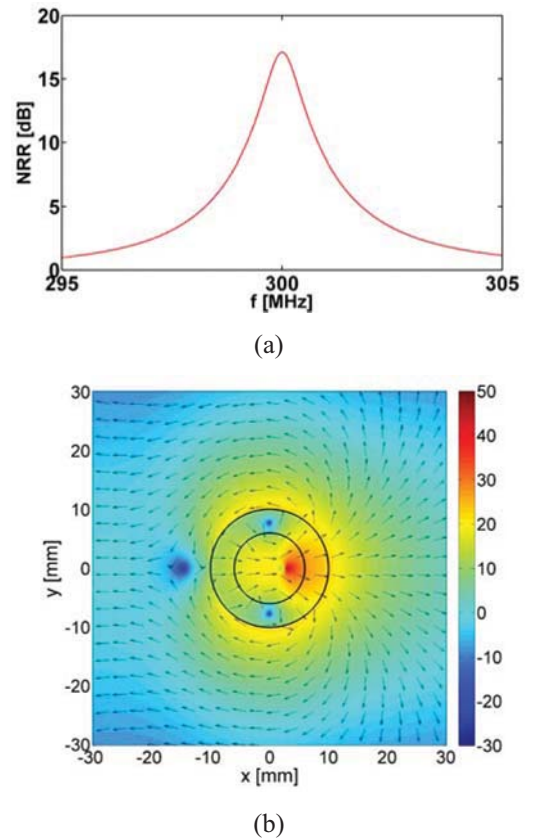


Figure 5: (a) The NRR as a function of frequency,  $f$ , for an ENG shell. (b) The magnitude (color) and direction (arrows) of the power flow density for the resonant ENG-shell at 300 MHz. In (a) and (b), the MLS is at 3 mm along the positive  $x$ -axis. Note that the curves representing the cylindrical surfaces of the ENG shell are also included in (b).

#### 4. Summary and conclusions

The power flow density (magnitude and direction) in the interior and exterior of electrically small configurations consisting of resonant active coated nano-particles (CNPs) of a circularly cylindrical shape was examined. The source of excitation was taken to be an infinitely long magnetic line source (MLS) at optical frequencies. The active particles were comprised of a canonical gain-impregnated silica nano-

core covered concentrically with a silver, gold, or copper nano-shell. The previously reported super-resonant states, excited in silver-, gold-, and copper-based CNPs and leading to large powers extracted from the MLS for a given value of the current along it, were confirmed with the power flow density results. For all cases, the power density was found to be rather strong and homogeneous in strength and direction inside the CNP; this being particularly so near the inner surface of the nano-shell, where a profound localization of the power density within the CNP was observed. The interaction of such large fields with the canonical gain inside the nano-core was found to lead to a clear and strong dipolar power pattern, with an outwardly propagating power flow density. These effects were found to be most notable for the gold-based CNP, as it is the least lossy case of all at the resonance wavelength. Moreover, no peculiar behavior in terms of certain vortices was found for the cylindrical *active* and *super-resonant* active CNPs at specific wavelengths that are shorter or longer than the resonant wavelength, this appearing in contrast to the observations for spherical particles [10], [13]. The results for the active CNPs at optical frequencies were furthermore contrasted to those obtained at radio frequencies for the corresponding epsilon-negative based particles. The latter resulted in power flow behavior significantly different from any of the super-resonant active optical CNP cases. In fact, the power flow lines in the radio frequency case were found to exhibit an interesting behavior including vortices at specific locations in the interior and exterior of the particle.

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