

Monitoring moisture movements in building materials using x-ray attenuation

Influence of beam-hardening of polychromatic x-ray photon beams

Pease, Bradley Justin; Scheffler, Gregor A.; Janssen, Hans

Published in:
Construction and Building Materials

Publication date:
2012

Document Version
Early version, also known as pre-print

[Link back to DTU Orbit](#)

Citation (APA):
Pease, B. J., Scheffler, G. A., & Janssen, H. (2012). Monitoring moisture movements in building materials using x-ray attenuation: Influence of beam-hardening of polychromatic x-ray photon beams. Construction and Building Materials.

DTU Library

Technical Information Center of Denmark

General rights

Copyright and moral rights for the publications made accessible in the public portal are retained by the authors and/or other copyright owners and it is a condition of accessing publications that users recognise and abide by the legal requirements associated with these rights.

- Users may download and print one copy of any publication from the public portal for the purpose of private study or research.
- You may not further distribute the material or use it for any profit-making activity or commercial gain
- You may freely distribute the URL identifying the publication in the public portal

If you believe that this document breaches copyright please contact us providing details, and we will remove access to the work immediately and investigate your claim.

Monitoring moisture movements in building materials using x-ray attenuation: Influence of beam-hardening of polychromatic x-ray photon beams

Brad J. Pease^{a,*}, Gregor A. Scheffler^b, Hans Janssen^c

^aTechnical University of Denmark (DTU), Department of Civil Engineering, Brovej - Building 118, DK-2800 Kgs. Lyngby, Denmark

^bXella Technologie- und Forschungsgesellschaft, Hohes Steinfeld 1, 14797 Kloster Lehmin, Germany

^cBuilding Physics Section, Department of Civil Engineering, K.U.Leuven, Kasteelpark Arenberg 40, B-3001 Heverlee, Belgium

Abstract

X-ray attenuation measurements are commonly used as a non-destructive method to monitor internal concentration changes of moisture (i.e., moisture content) and other chemical compounds in porous building materials. The technique provides direct measurements of moisture content changes through analysis with a composite model consisting of a dry porous material and a thickness of water equivalent to the moisture content of the material. The current formulation of this composite model relies on certain assumptions, including a monochromatic x-ray photon beam source (i.e., x-ray photons of a single, consistent energy) and that interactions between the x-ray photons and the materials (water and porous material) are independent. However, x-ray sources typically used by researchers in this field of study produce x-ray photon beams over a spectrum of energy levels, or polychromatic x-ray photons. Implications of this inconsistency are introduced and discussed. This paper presents both an overview of fundamental descriptions of the x-ray attenuation measurement technique and results from a parametric experimental study of various porous construction materials, including calcium silicate board, aerated autoclaved concrete, clay brick, cementitious materials, and wood. Results from the parametric investigation indicate the attenuation coefficient of water is dependent on the type and thickness of the porous material.

Keywords: Moisture content, X-ray, Attenuation coefficient, Porous media

1. Introduction

Moisture, and more specifically varying moisture conditions, often significantly affects structural performance and durability of construction materials. For example, moisture movements may affect warping of timber studs [1], workability [2] and plastic cracking [3] of fresh cementitious materials, maturation of bond strength of lime mortars [4], aging of bitumen binders [5], freeze-thaw damage of concretes [6], etc. Furthermore, moisture influences various chemical, biological and thermal characteristics of materials. In general, moisture and moisture movements are key determining factors in the durability of structures constructed from a multitude of materials. Therefore, experimental approaches for accurate non-destructive monitoring of moisture conditions are essential. Various non-destructive methods are available [7]; however, x-ray

attenuation measurements are emerging as a popular technique for monitoring moisture movements in various porous materials, particularly in cracked materials [8–21]. Measurement speed and resolution, ease of two-dimensional (or three-dimensional [22]) imaging, availability of facilities, and direct measurement of changes in moisture content are key advantages.

X-ray technologies are commonplace in the medical and security fields; however, application of x-ray attenuation measurements for monitoring of moisture movements in porous media is relatively complex. Changes in moisture content necessitate comparison of multiple x-ray measurements captured at different states (i.e., times). Currently, descriptions of basic and fundamental controlling mechanisms and influential factors are somewhat lacking.

Quantification of moisture content in porous materials from x-ray attenuation measurements is achieved by decomposing a water-containing (saturated or partially saturated) porous material into a dry porous material of the same thickness and a water layer with a thickness

*Corresponding author

Email address: bjp@byg.dtu.dk (Brad J. Pease)

equivalent to the moisture content [14, 16–18, 20]. Using the composite system, equations have been derived allowing for calculation of moisture content changes from x-ray measurements taken at different states. The derived expressions however depend upon certain assumptions. Firstly, calculations assume the incoming x-ray photons are of a single, constant energy. Secondly, the quantification of moisture content in porous material with x-ray attenuation is based on a non-interacting composite concept where the attenuation of x-ray photons provided by the porous material and the water are assumed independent. This paper demonstrates these assumptions are not generally valid, and that the interaction between material and water needs to be accounted for to accurately quantify moisture contents in porous materials via x-ray attenuation. Further, the common use of polychromatic x-ray source, which produce x-ray photons over a spectrum of energies, needs to be accounted for to accurately quantify moisture contents in porous materials via x-ray attenuation.

The first section of the paper summarizes the physics of x-ray attenuation, including both monochromatic and polychromatic x-ray photon beams. Key terminology and common misconceptions within the civil/materials engineering communities, particularly the various definitions of attenuation coefficient, are introduced. In the remaining section the composite concept typically utilized to relate changes in x-ray photon counts to changes in moisture content is challenged. Factors affecting the composite concept, including primarily beam hardening and the resulting dependence of the attenuation coefficient of water on the 'parent' material, are discussed. While specific results presented are dependent on the type of materials investigated (calcium silicate board, aerated autoclaved concrete, clay brick, cementitious materials, and wood here) and also likely on the x-ray system used, vital conclusions on and potential pitfalls of the general measurement technique are discussed in this paper.

2. Fundamentals of the x-ray attenuation measurement technique

Fig. 1(a) illustrates the required equipment and the behaviour of an x-ray photon beam during x-ray attenuation measurements. For such measurements a specimen is placed between an x-ray source, which produces a conical beam of x-ray photons, and an x-ray detector (camera used here). Additional details on the interactions between x-ray photons and the specimen are described in a following section. Fig. 1(b) illustrates the commonly accepted composite concept, consisting of a

dry or partially saturated specimen and a thickness of water, to describe the effect moisture content changes have on x-ray attenuation measurements [8, 14, 16–18, 20, 23]. Using this composite system along with several assumptions, a relationship can be derived to convert changes in x-ray photon counts to changing moisture content of porous materials. One key assumption in the derivation this relationship is that the porous material does not impact the x-ray attenuation provided by water, and vice versa; therefore, the model illustrated in Fig. 1(b) is referred to as a 'non-interacting' composite concept. The following sections describe the x-ray attenuation measurement technique in detail, provide the derivation the photon count-moisture content relationship, and introduce other common assumptions.

2.1. Interactions between x-ray photons and building materials

As illustrated in Fig. 1(a) an x-ray photon beam with an initial intensity, I_0 is produced by an x-ray source and passes through a specimen. A portion of the photons interact with the specimen, reducing the initial intensity to the transmitted intensity, I . The reduction in intensity for a monochromatic x-ray photon beam (i.e., photons of single and constant energy) is described by Eq. 2:

$$\frac{dI(x)}{dx} = -\mu \cdot I(x) \quad (1)$$

or in the more common form, known as the Beer-Lambert law [24]:

$$I = I_0 \cdot e^{-\mu \cdot t} \quad (2)$$

where μ is the linear attenuation coefficient [m^{-1}], $I(x)$ is the number of x-ray photons present at depth x , and t is the thickness of the specimen [m]. The units for intensity are counts or hits, simply indicating the number of x-ray photons.

The linear attenuation coefficient, μ is equivalent to the total probability of x-ray photon-material interaction per unit length of a material. An x-ray photon passing through a building material may interact on the atomic level in one of three ways – photoelectric effect, Compton scatter, or pair production [25]. Eq. 2 defines μ therefore as:

$$\mu(E, Z_{eff}, \rho) = \tau(E, Z_{eff}, \rho) + \sigma(\rho) + \kappa(E \geq 1022 \text{ keV}, \dots) \quad (3)$$

where τ , σ , and κ are the photoelectric, Compton, and pair production cross-sections [m^{-1}], respectively, all of which are characteristics of the material through which x-ray photons are passing.

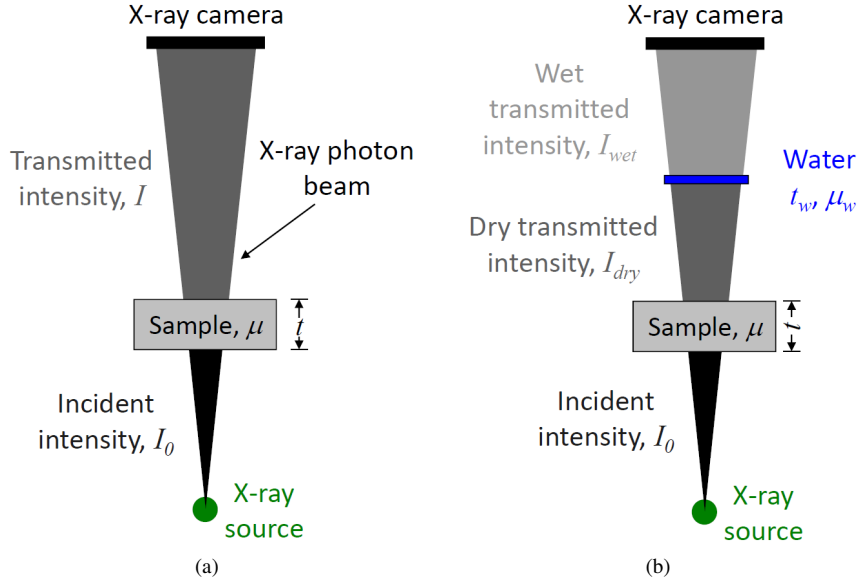


Figure 1: Schematic representation of a typical x-ray attenuation experimental setup. (a) illustrates the Beer-Lambert Law while (b) illustrates a composite experimental system consisting of a dry specimen and an amount of liquid water to represent movement of moisture.

133 The photoelectric effect describes one possible inter- 159
 134 action between a photon and a bound atomic electron 160
 135 in the specimen, wherein the photon is absorbed and a 161
 136 photoelectron is produced. The probability of this inter- 162
 137 action (τ) increases with the effective atomic number, 163
 138 Z_{eff} and density, ρ [kg/m³] of the specimen and 164
 139 decreases significantly with increased x-ray photon energy, 165
 140 E [keV]. Compton scatter occurs when an x-ray photon 166
 141 collides with a free electron, causing a change in 167
 142 direction and reduction in energy of the photon. The 168
 143 probability of this interaction (σ) increases with ρ and 169
 144 is slightly influenced by E at energies far higher than 170
 145 typically used for the type of investigations discussed 171
 146 here [25]. Pair production, wherein the photon and an 172
 147 atomic nucleus interact to create an electron-positron, 173
 148 only occurs for $E > 1022$ keV (i.e., minimum kinetic 174
 149 energy required to initiate pair production) [25]. Such 175
 150 energies far exceed typically used x-ray source energy 176
 151 settings [8–23]. Therefore, for $E < 200$ keV the energy-, 177
 152 effective atomic number-, and density-dependent linear 178
 153 attenuation coefficient, μ can be described according 179
 154 to Eq. 4 [26–28]:

$$155 \mu(E, Z_{eff}, \rho) = \left(a \frac{Z_{eff}^{3.8}}{E^b} + c \right) \cdot \rho \quad (4)$$

156 where a and c are coefficients describing the photo- 183
 157 electric and Compton cross-sections, respectively and 184
 158 b describes the shape of the photoelectric cross-section 185

curve. The first term of Eq. 4 is equivalent to the photoelectric cross-section, τ in Eq. 3, while the second term is equivalent to Compton cross-section, σ . The Compton cross-section is assumed to be photon energy-independent below 200 keV [26–28].

Fig. 2 shows fitting results of Eq. 4 against tabulated data on linear attenuation coefficients for liquid water, μ_w , in [29]. Due to the photoelectric effect low energy x-ray photons are attenuated more efficiently, a trend inherent in all elements and compounds (see [e.g., 29]). Compton cross-section controls μ_w -values for photon energies greater than approximately 26 keV.

171 2.2. Impact of polychromatic x-ray sources

172 The x-ray source utilized here (described in [8]) and 173
 174 sources utilized by others [7, 9–22] produce x-ray photons 175
 176 over a spectrum of energies (i.e., polychromatic) limited 177
 178 to the user-defined x-ray energy setting, E_m as illustrated 179
 179 in Fig. 3(a). Furthermore, the x-ray camera used here, 180
 181 described below, and similar devices used by others [7, 14, 17, 18] 182
 182 detect only the presence of individual x-ray photons and not the 183
 183 energy of individual photons. The use of these non-energy discretizing 184
 184 x-ray cameras with a polychromatic x-ray source requires the 185
 185 introduction of an average attenuation coefficient, $\bar{\mu}$ as discussed 186
 186 below. Further, the polychromatic x-ray source leads to so-called 187
 187 beam hardening [30] illustrated in Fig 3(b), which must be 188
 188 considered.

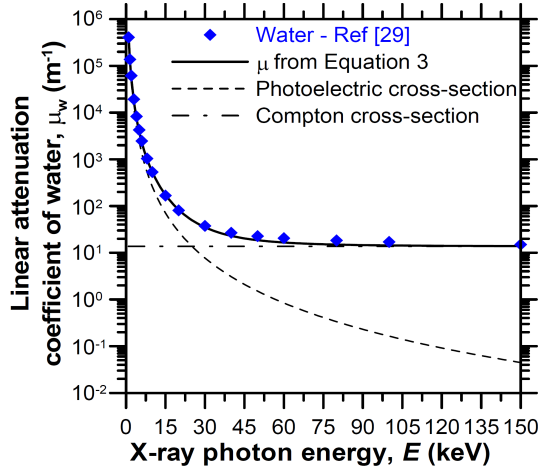


Figure 2: Fitting provided by Eq. 4 to tabulated data on μ_w from [29] including contributions from photoelectric and Compton cross-sections R^2 was 0.999 for μ between 0-200 keV.

186 Fig. 3(a) illustrates a possible x-ray spectrum pro-
 187 duced by an x-ray source, assuming an x-ray source
 188 energy (i.e., E_m) of 85 keV. The shape of the initial
 189 intensity spectrum presented in Fig. 3(a) comes from
 190 [31] where a photon energy detector measured the x-
 191 ray spectrum produced by the exact source utilized here.
 192 The shape of the photon energy spectrum, described by
 193 $\Psi_E(x)$, changes due to attenuation as it passes through
 194 increasing thicknesses, x of material. To illustrate this
 195 effect, the transmitted spectrum was calculated by ap-
 196 plying Eq. 2 at each x-ray energy level and using the fit-
 197 ted linear attenuation coefficients of water for a 15 mm
 198 thickness of water as shown in Fig. 3(b). Attenuation
 199 provided by the water progressively increases the average
 200 x-ray photon energy, meaning the remaining x-ray
 201 photons are ‘harder,’ hence the term beam hardening.

202 The x-ray camera used measures only the number of
 203 x-ray photons (i.e., intensity) as the sum of the area un-
 204 der the $\Psi_E(x)$ -curve, as shown in Fig. 3(a) for an initial
 205 intensity, I_0 . The x-ray camera does not distinguish be-
 206 tween photons of varying energies. Therefore, relating
 207 the initial and transmitted intensities, I_0 and I respec-
 208 tively, shown in Fig. 3(b) is not directly possible using
 209 Eq. 1 and 2. Eq. 1 must be modified, removing the lin-
 210 ear attenuation coefficient, μ and introducing an aver-
 211 age attenuation coefficient, $\bar{\mu}$ over the spectrum, $\Psi_E(x)$
 212 [32, 33] defined by:

$$213 \bar{\mu} = \frac{\int_0^{E_m} \mu(E, Z_{eff}, \rho) \cdot \Psi_E(x) dE}{\int_0^{E_m} \Psi_E(x) dE} \quad (5)$$

214 The x-ray spectrum, $\Psi_E(x)$, which can vary depend-
 215 ing on x-ray source (see spectra provided in references
 216 [8, 12, 14, 18]), clearly influences the average atten-
 217 uation coefficient. Fig. 3(c) shows calculated values of
 218 $\bar{\mu}$ as a function of water thickness, where the $\Psi_E(x)$ -
 219 curves were calculated as previously described. The
 220 broken line in Fig. 3(c) indicates the value of the lin-
 221 ear attenuation coefficient of water (μ_w) at the x-ray
 222 source energy setting of 85 keV. The blue diamonds show
 223 the average attenuation coefficient for varying thick-
 224 ness of water, as calculated using Eq. 5. The calcula-
 225 tions indicate average attenuation coefficients are poten-
 226 tially several orders of magnitude higher than linear at-
 227 tenuation coefficients. The average attenuation coeffi-
 228 cient is asymptotic to the linear attenuation coefficient
 229 at the maximum x-ray photon energy as a sufficient
 230 thickness would attenuate all but the most energetic
 231 photons.

232 Due to beam hardening, $\bar{\mu}$ is no longer a constant,
 233 but varies with the location in the material. Fig. 4 il-
 234 lustrates the thickness-dependency of the average at-
 235 tenuation coefficients from measurements on concrete.
 236 The specimen thickness-dependency of the average at-
 237 tenuation coefficient can be described by Eq. 6 [32, 33]:

$$237 \bar{\mu}(t) = \frac{\mu_0}{(1 + \lambda \cdot t)^2} = \frac{\bar{\mu}(\Psi_E(0), Z_{eff}, \rho)}{(1 + \lambda \cdot t)^2} \quad (6)$$

238 where μ_0 is the initial average attenuation coefficient
 239 [m^{-1}] and λ is the beam hardening coefficient of the
 240 material [m^{-1}]. The initial average attenuation co-
 241 efficient, μ_0 , may be considered the average atten-
 242 uation coefficient of an infinitely thin section of mate-
 243 rial, indicating μ_0 -values must depend upon the initial
 244 x-ray spectrum, $\Psi_E(0)$. Further, based on Eq. 4, μ_0 -
 245 values also likely depend on the density (ρ) and effec-
 246 tive atomic number (Z_{eff}) of the material. Therefore,
 247 the second definition for $\bar{\mu}$ is included in Eq. 6. Eqs.
 248 4 and 5 provide a means to derive an expression for
 249 $\bar{\mu}(\Psi_E(0), Z_{eff}, \rho)$; however, solutions are complicated
 250 due to a singularity when the x-ray energy is 0 keV.
 251 Therefore, Eq. 7 is suggested as a possible means to re-
 252 late μ_0 -values to the x-ray source energy setting, E_m
 253 and material properties, based upon Eq. 4:

$$254 \mu_0 = \bar{\mu}(\Psi_E(0), Z_{eff}, \rho) = \left(a^* \frac{Z_{eff}^{3.8}}{E_m^{b^*}} + c^* \right) \cdot \rho \quad (7)$$

255 where the terms a^* , b^* , and c^* are analogous to a , b , and
 256 c described in Eq. 4 for the initial average attenuation
 257 coefficient. The term $a^* \cdot Z_{eff}^{3.8}$ can be replaced with a
 258 single variable, α to simplify fitting of experimental re-
 259 sults presented later due to the complex stoichiometry of
 260 the materials tested. Values from Eq. 6 yield tangential

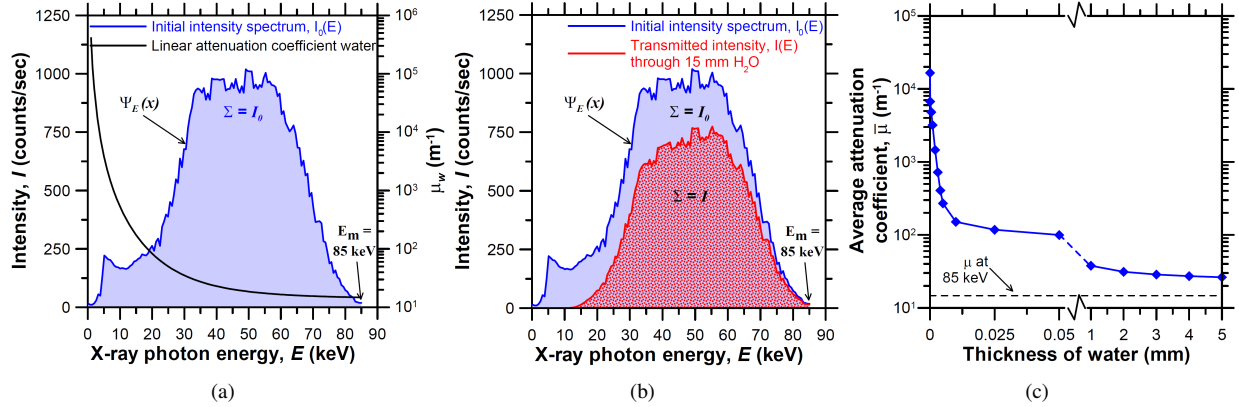


Figure 3: Illustration of influence of x-ray spectrum on x-ray attenuation measurements including: (a) possible x-ray spectrum produced by x-ray source with an x-ray source energy setting of 85 keV and the photon energy dependent linear attenuation coefficients, (b) calculated transmitted spectrum through a 15 mm thickness of water, and (c) comparison of average attenuation coefficients at varying thicknesses of water and the linear attenuation coefficient at 85 keV.

261 slopes of the solid red plot in Fig. 4; therefore, to sim-
 262 plify analysis the effective attenuation coefficient, μ_{eff} ,
 263 is utilized.

264 The effective attenuation coefficient, μ_{eff} [m⁻¹], is
 265 needed for use in the integrated Beer-Lambert law (Eq.
 266 1) to describe the x-ray attenuation provided by a par-
 267 ticular thickness of a material, as defined by Eq. 8:

$$\mu_{eff} = \frac{\int_0^t \bar{\mu} dx}{t} = \frac{\int_0^t \frac{\bar{\mu}(\Psi_E(0), Z_{eff}, \rho)}{(1+\lambda \cdot t)^2} dx}{t} = \frac{\bar{\mu}(\Psi_E(0), Z_{eff}, \rho)}{1 + \lambda \cdot t} \quad (8)$$

268 Combining Eqs. 8 and 7 provides an expression to de-
 269 scribe the effect of x-ray source energy and beam hard-
 270 ening on the effective attenuation coefficients of mate-
 271 rials which is measured using a polychromatic x-ray
 272 source and an x-ray camera:

$$\mu_{eff} = \frac{\left(\frac{\alpha}{E_m^{b^*}} + c^* \right) \cdot \rho}{1 + \lambda \cdot t} \quad (9)$$

274 The various attenuation coefficients illustrated in Fig.
 275 4 are analogous to another material behavior more fam-
 276 iliar to the civil engineering community, namely the
 277 various measurements of elastic modulus. In this anal-
 278 ogy the initial average attenuation coefficient coincides
 279 with the initial tangent elastic modulus, the average at-
 280 tenuation coefficient with the tangent elastic modulus,
 281 and the effective attenuation coefficient with the secant
 282 elastic modulus.

283 Figs. 3(c) and 4 indicate the potential for error in x-
 284 ray attenuation measurements if an incorrect attenuation

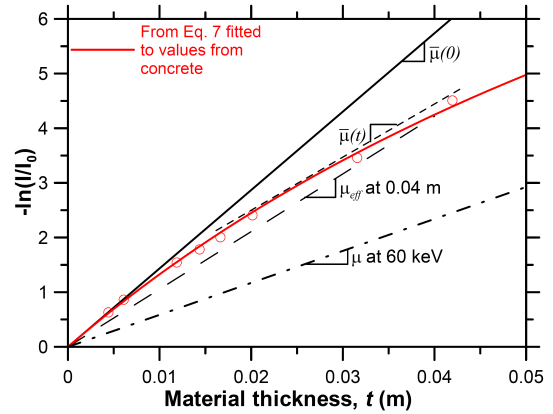


Figure 4: Descriptions of linear, average, and effective attenuation coefficients, μ , $\bar{\mu}$, and μ_{eff} , respectively. Values for $\bar{\mu}$ are the tangential slopes of the red line, which is fitted to the shown measured values for concrete (circle symbols) at a maximum x-ray photon energy, E_m of 60 keV. The linear attenuation coefficient, μ for concrete at 60 keV found in [29].

285 coefficient of water is used. Results from the study pre-
 286 sented in the second part of this paper further indicate
 287 the attenuation coefficient of water is affected by the
 288 type and thickness of the porous specimen.

2.3. Composite concept for assessment of concentration changes

289 Commonly, changes in moisture content (or changes
 290 in concentration of other materials [e.g., corrosion prod-
 291 ucts 23]) in a porous building material are described as
 292 the composite system shown in Fig. 1(b). Here an x-
 293 ray photon beam with initial intensity, I reduces to a
 294
 295

transmitted intensity through a wet material, I_{wet} in two drops: 1) some photons interact with the dry (or partially saturated) porous material resulting in a dry transmitted intensity, I_{dry} , and 2) additional photons interact with a water layer with thickness t_w , equivalent to the material's water content resulting in a wet transmitted intensity, I_{wet} . The composite system simulates water uptake by increasing the equivalent water thickness, while drying corresponds to a reduction in water thickness. Using Eq. 2 and assuming interactions between x-ray photons and the dry porous material only influence the quantity of photons (and not photon energy), I_{wet} may be expressed as:

$$I_{wet} = I_{dry} \cdot e^{-\mu_w \cdot t_w} = I_0 \cdot e^{-\mu \cdot t} \cdot e^{-\mu_w \cdot t_w} = I_0 \cdot e^{-\mu \cdot t - \mu_w \cdot t_w} \quad (10)$$

for a monochromatic beam, and as derived in [8, 14] moisture content changes can be directly calculated by measurements of I_{dry} and I_{wet} according to Eq. 11:

$$\Delta w = -\frac{\rho_w}{\mu_w \cdot t} \ln \left(\frac{I_{wet}}{I_{dry}} \right) \quad (11)$$

where Δw is a change in moisture content [kg/m^3], ρ_w the density of water [kg/m^3], and μ_w is the linear attenuation coefficient of water. Decreases in moisture content (i.e., drying) can also be directly calculated by Eq. 11 by inverting $\frac{I_{wet}}{I_{dry}}$ to $\frac{I_{dry}}{I_{wet}}$.

Eq. 11 also provides a measure of the ability to resolve changes in concentration using the x-ray attenuation measurement technique, which is defined as the change in concentration required to attenuate a single x-ray photon. Eq. 12, from [14], describes concentration resolution mathematically.

$$R_{\Delta w} = -\frac{\rho_w}{\mu_w \cdot t} \ln \left(\frac{I_{dry} - 1}{I_{dry}} \right) \quad (12)$$

Eqs. 11 and 12 are central to the quantification of moisture contents in porous materials by x-ray attenuation, and both make use of the attenuation coefficient of water. Typically the non-interactive composite concept is accepted and the effective attenuation coefficient of water is either determined in isolation without interference from the porous material or, in many cases, experimental descriptions do not adequately describe which measure of attenuation coefficient of water is used for calculations of changes in moisture content [12, 14, 15, 21, 22]. This paper will show the former to be incorrect and indicates the latter requires more attention in future implementations of the x-ray attenuation measurement technique.

3. Experimental program

The goal of the experiments presented in this paper was to investigate whether an interaction exists between the attenuation provided by the porous material and the stored water. The effect of type and thickness of the parent material was investigated. Details on the materials tested, the x-ray measurement system, and measurements conducted are provided in the following sections.

3.1. Materials

Materials tested included various porous building materials, acrylic, steel, and aluminium. Porous materials included aerated autoclaved concrete, calcium silicate, clay brick, cement paste, concrete, and wood. Specimens with varying thicknesses were created either by stacking individual thin plates of the material (e.g., aluminium, acrylic) or by cutting (e.g., cement paste, clay brick). Aalborg white portland cement was used in cement paste and concrete specimens. The cement paste (0.40 water-to-cement ratio) was cured submerged in water for 28 days. Additional details on the concrete used is available in the literature [16, 20].

3.2. X-ray facility

A GNI x-ray system located at the Technical University of Denmark [34], consisting of an x-ray source and camera affixed to three programmable axes, was used for all measurements described here.

The x-ray camera consisted of a 25x25 mm² cadmium telluride (CdTe) semiconductor detector bump bonded to 252x256 16-bit complementary metal-oxide semiconductor (CMOS) pixel circuits [35]. CdTe, a direct conversion x-ray detector, converts x-ray photons to electric signals and is capable of doing so at room temperatures. The 64,512 (252x256) CMOS pixel circuits record the location and number of photons detected over a set period of time, called the integration time (analogous to shutter time in visible-light photography). The number of photons detected during the integration time is the transmitted intensity, I (Eq. 2). CMOS pixel depth (maximum recordable I) was 65,536 (i.e., 2^{16}); therefore, integration times must be selected to avoid 'overflowing' pixels. The x-ray camera used is described in detail in [35].

Beyond transmitted x-ray photons, the x-ray camera detects background radiation sources, largely in the form of dark current. Other background sources (naturally occurring radioisotopes, cosmic rays, etc.) are minimal due to the materials tested and shielding chamber. Trial measurements of dark current indicated time and temperature variations, therefore, dark current was

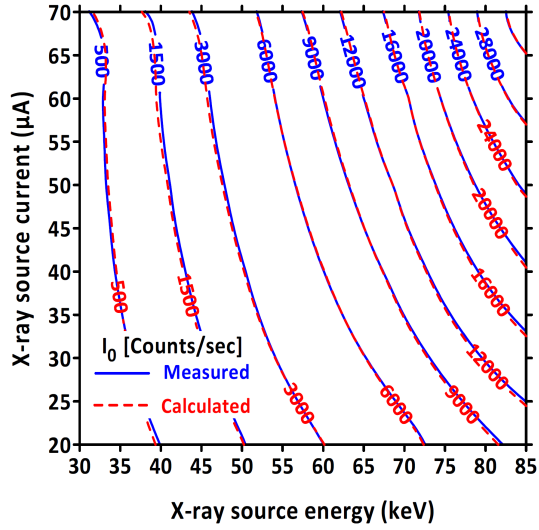


Figure 5: Contour plot of the initial intensity I_0 as a function of the input x-ray source energy and current. The blue contour lines indicate measured values, while red contours indicate calculated values from Eq. 13.

recorded prior to and subtracted from each x-ray attenuation measurement.

Additional details on the x-ray source [8] and other components of the x-ray equipment in its current configuration are available in the literature [20, 23].

3.3. Determination of attenuation coefficients of materials

Initial intensity, I_0 (transmitted intensity through air) was measured with varying x-ray source energies (30, 40, 50, 60, 70, 85 keV) and currents (20, 30, 40, 50, 60, 70 μA). Transmitted intensities, I were measured through varying thicknesses of the porous materials and isolated water using the same x-ray source settings to determine attenuation coefficients. The attenuation coefficient of water was also measured as illustrated in Fig. 1(b) by placing a 20 mm thickness of water (contained in acrylic containers) in series with specimens of varying material and thickness. All measurements conducted as part of this study were made with a warm-up time of 200 seconds, a stabilization time of 600 seconds, and an integration time of 20 seconds.

4. Results

4.1. Initial intensity of x-ray photon beam

Fig. 5 shows measured values of average initial x-ray photon intensity per second as measured by the x-ray camera (average of intensities measured by the 64,512

Table 1: Values of coefficients C_{1-6} in Eq. 13.

Fitting coefficient	Value
C_1	-0.0187
C_2	1.502×10^{-4}
C_3	-0.0800
C_4	2.0283
C_5	1.0607
C_6	1.00236

pixels) for varying x-ray source current and energies. Eq. 13, which relates initial intensity to x-ray source current, i and energy, E_m , was found to accurately fit measured initial intensity:

$$I_0 = (C_1 \cdot i) \cdot (C_2 \cdot E_m^4 + C_3 \cdot E_m^3 + C_4 \cdot E_m^2 + C_5 \cdot E_m + C_6) \quad (13)$$

where C_{1-6} are fitting parameters shown in Table 1. Eq. 13 is a polynomial fitting of measured results and therefore is only applicable for the investigated x-ray source inputs (i.e., $30 \leq E_m \leq 85$ keV and $20 \leq i \leq 70$ μA). This expression and the values for the coefficients may not be directly transferrable to other x-ray sources; however values for I_0 are needed to derive I_{dry} for determination of measurement resolution (i.e., Eq. 12).

4.2. Effective attenuation coefficients of materials

Fig. 6 illustrates the typical effects of x-ray source energy and materials thickness on effective attenuation coefficient, μ_{eff} . Measured values in Fig. 6 are indicated by symbols, while fits provided by Eq. 9 are shown as the solid lines. Table 2 provides values for α , b^* , c^* , λ along with coefficients of determination (R^2) of the resulting fits. Fig. 6(a) indicates μ_{eff} -values decrease as x-ray source energy, E_m increases, and that the x-ray source energy-effect can be accurately estimated by Eq. 7 proposed in Sect. 2.2. Fig. 6(b) shows μ_{eff} -values may be affected by material thickness, depending on the type of material. For example, the thickness of wood has a minimal impact on μ_{eff} -values, while increased clay brick thickness causes significant reductions in μ_{eff} -values.

Figs. 7(a)-(d), which illustrate both the material thickness, t and x-ray source energy, E_m effects on μ_{eff} -values, show measured and fitted (using Eq. 9) μ_{eff} -values for water, calcium silicate, clay brick, and aluminium, respectively. Measurement points are indicated by blue symbols with the measured μ_{eff} -values indicated by the adjacent number. The red isolines show the fitted μ_{eff} -values provided by Eq. 9. Fitting parameters are provided for all materials in Table 2, including water. The impact of hardening coefficient, λ is

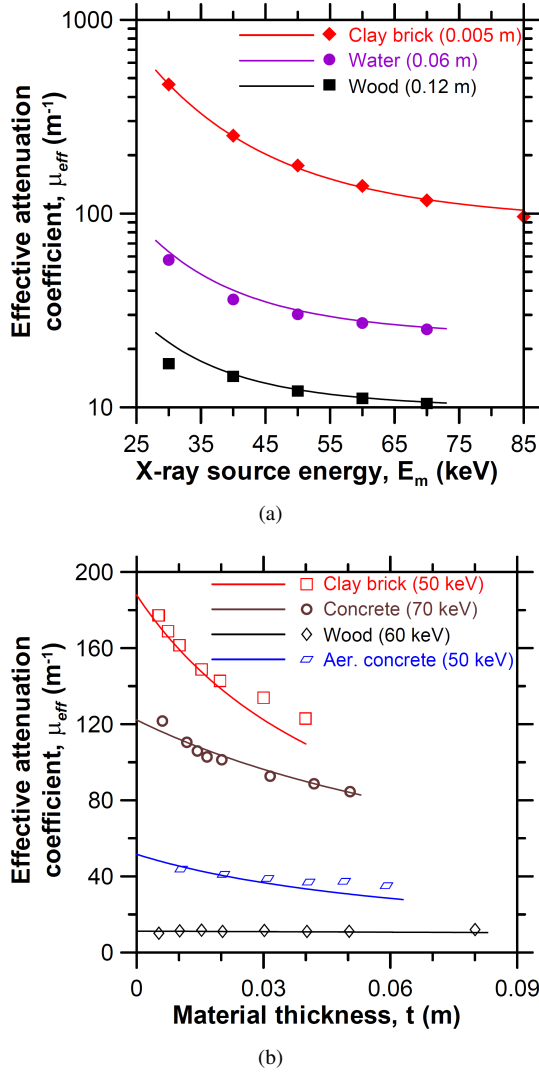


Figure 6: Influence of (a) x-ray source energy setting, E_m and (b) material thickness, t on effective attenuation coefficient. Measured values indicated by symbols, while lines indicate fits provided by Eqs. 7 and 8.

indicated by comparison of Figs. 7(a)-(d) where λ increases from 0.888 m⁻¹ in (a) to 8.94 m⁻¹ in (b), 17.8 m⁻¹ in (c) and 57.7 m⁻¹ in (d). As λ increases, the effect of specimen thickness on the μ_{eff} -isolines becomes increasingly more significant – with only minimal changes in μ_{eff} -values as the thickness of water increases (Fig. 7(a)). For aluminium (Fig. 7(d)), thickness increase drastically reduces μ_{eff} -values. Fig. 6(b) also illustrates the impact of material type and thickness on μ_{eff} -values.

4.3. Coupled effect of parent material on effective attenuation coefficient of water

Fig. 8(a) indicates the influence of water and parent material thickness on measured values of the effective attenuation coefficient of water, $\mu_{eff,w}$. Calcium silicate in various thicknesses (0, 0.02, 0.04, and 0.08 m) was used as the parent material in this case while water thickness was also varied (approximately 0.01, 0.02, 0.04, and 0.06 m). While the effect of the parent material's thickness is evident, no clear trend was identified concerning the impact of water thickness. Increasing the water thickness from 0.01 to 0.02 appears to result in an increased $\mu_{eff,w}$ -value, followed by a subsequent decrease with increased water thickness. Measurement variability may explain the apparent differences with water thickness. However, Fig. 8(a) clearly shows a coupling effect of the thickness of the calcium silicate on the effective attenuation coefficient of water. Therefore, the term coupled effective attenuation coefficient of water, $\mu_{eff,w}(t)$; which is a function of the parent material thickness, t ; is introduced here.

Fig. 8(b) shows this coupling effect of the parent material on the effective attenuation coefficient of water occurs for all porous materials tested. The symbols in Fig. 8(b) indicate measured effective attenuation coefficients of water when the x-ray photon beam passed through varying thickness of the various parent materials prior to interacting with the water. As shown in the figure, the measured effective attenuation coefficient of water decreases from an average of 24.6 m⁻¹ for a 0 mm thickness of parent material to as low as 17.8 m⁻¹ for a 23.6 mm thickness of cement paste. It should be noted that values of $\mu_{eff,w}(t)$ presented in Fig. 8(b) were measured using a constant 0.02 m thickness of water as Fig. 8(a) indicates water thickness is of less importance. The exponential decay model described in Eq. 14 was found to accurately fit measurements:

$$\mu_{eff,w}(t) = \mu_{eff,w}(\infty) + [\mu_{eff,w}(0) - \mu_{eff,w}(\infty)] \cdot e^{-\eta t} \quad (14)$$

where, $\mu_{eff,w}(\infty)$ is the coupled effective attenuation coefficient of water through an infinite thickness of the parent material [m⁻¹], $\mu_{eff,w}(0)$ is the coupled effective attenuation coefficient through a zero thickness of the parent material [m⁻¹], and η is the parent material's coupling effect [m⁻¹]. The lines in Fig. 8(b) indicate the fits provided by Eq. 14 for the various parent materials. Table 3 lists the fitted values for coefficients and the coefficients of determination for the fits provided by Eq. 14.

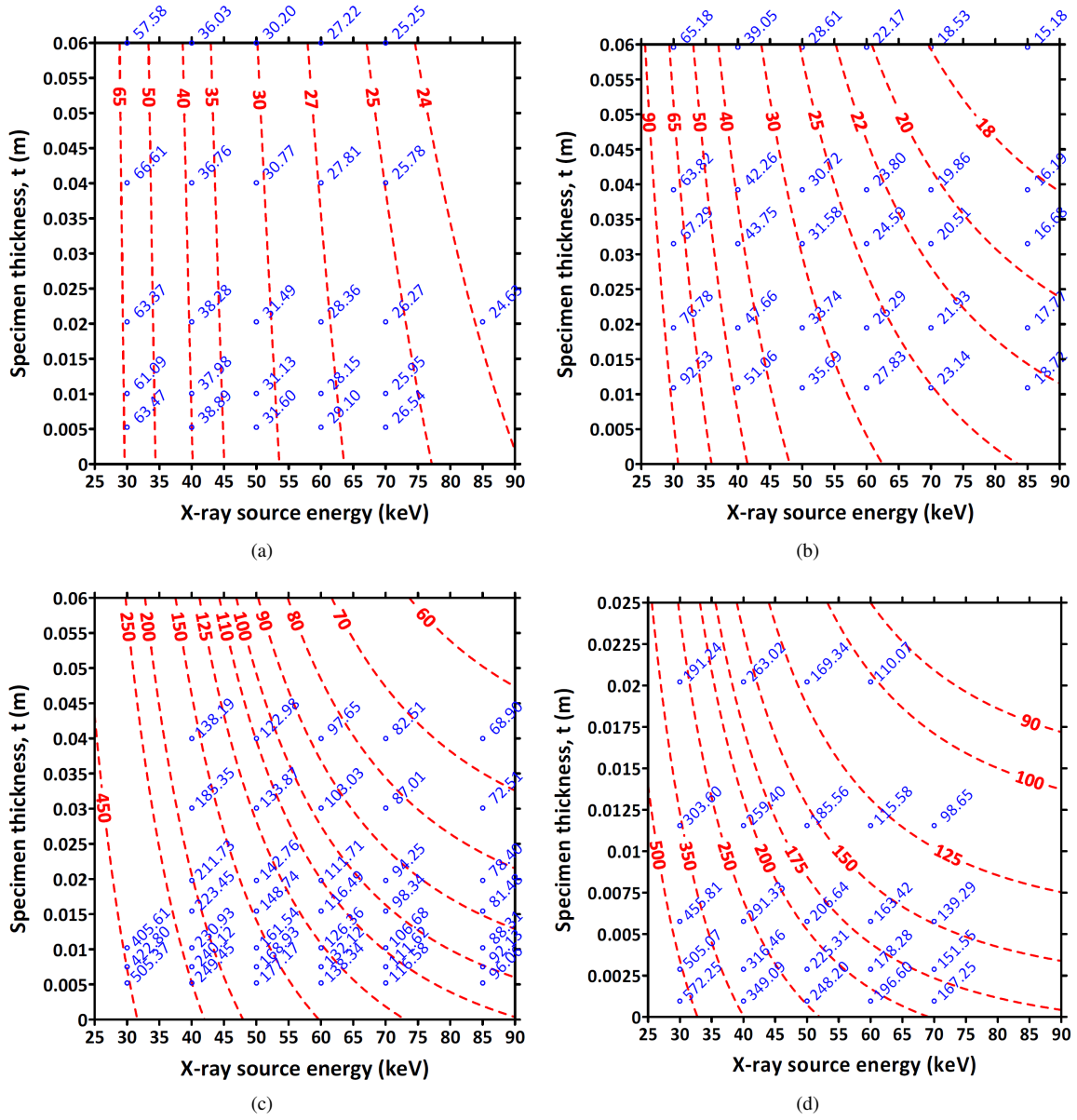


Figure 7: Comparison of measured and estimated (Eq. 9) effective attenuation coefficients for (a) water, (b) calcium silicate, (c) clay brick, and (d) aluminium as function of specimen thickness and x-ray source energy. Blue data points indicate measured points, while the red broken isolines indicate fitted values from Eq. 9.

Table 2: Values for parameters in Eq. 9 for various materials including the coefficient of determination (R^2) of the fits.

Parameter	Materials								
	Acrylic	Aerated concrete	Aluminium	Clay brick	Calcium silicate	Cement paste	Concrete	Wood	Water
ρ [kg/m ³]	1186	422	2700	1795	271	1440	2200	480	1000
α	484.4	3999.8	3112.8	4442.7	5228.4	6750.4	3840.0	481.2	790.3
b^*					2.9				
c^*	0.022	0.075	0.060	0.052	0.078	0.064	0.106	0.020	0.022
λ [m ⁻¹]	3.97	13.6	57.7	17.8	8.94	12.4	9.00	0.863	0.888
R^2	0.861	0.913	0.948	0.990	0.976	0.991	0.953	0.999	0.987

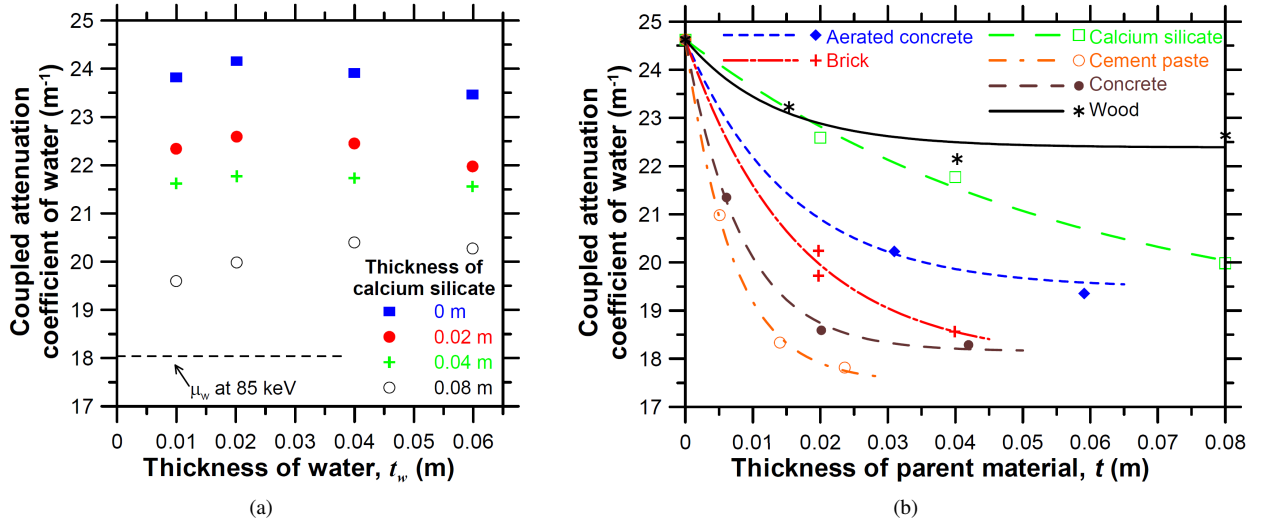


Figure 8: (a) Influence of water thickness on measured coupled attenuation coefficient of water with an x-ray source energy of 85 keV measured through calcium silicate parent specimen with varying thickness and (b) coupled attenuation coefficient of water with an x-ray source energy of 85 keV with varying types and thicknesses of parent materials with symbols indicating measured values and lines indicating fits provided by Eq. 14.

Table 3: Values for parameters in Eq. 14 for various materials including the coefficient of determination (R^2) of the fits.

Parameter	Materials							
	Acrylic	Aerated concrete	Aluminium	Clay brick	Calcium silicate	Cement paste	Concrete	Wood
$\mu_{eff,w}(\infty)$	21.6	19.5	18.5	18.0	18.6	17.5	118.2	22.4
$\mu_{eff,w}(0)$				24.63				
η	62.9	64.0	185.2	60.7	17.7	145.3	120.1	75.1
R^2	0.989	0.975	0.995	0.995	0.991	0.999	0.999	0.943

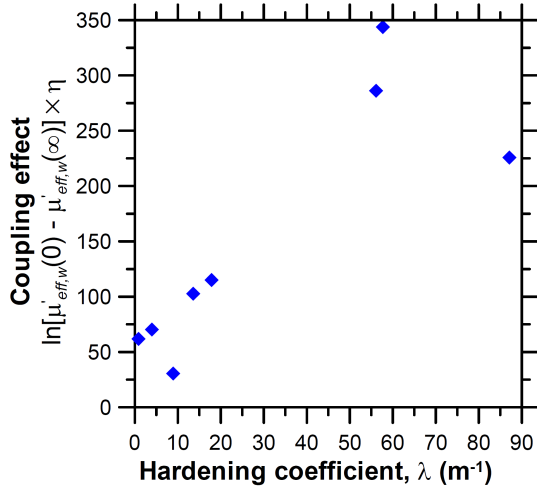


Figure 9: Hardening coefficient, λ of the parent materials versus the coupling effect of the parent material on the coupled effective attenuation coefficient of water.

5. Discussion

Comparison of results in Fig. 8(b) showing the coupling effect of the various porous materials and values for the hardening coefficient, λ of the porous materials presented in Table 2 indicates a potential relationship. Fig. 9 assesses this relationship, where the coupling effect is defined in accordance with Eq. 14 as $\ln[\mu_{eff,w}(0) - \mu_{eff,w}(\infty)] \cdot \eta$. Fig. 9 indicates the coupling effect tends to increase with the hardening coefficient of the porous material. Further investigations are required to verify and quantify this relationship. However, Fig. 9 and a comparison of results in Fig. 8(b) and Table 2 indicate a strong correlation between beam hardening and the reduction in the coupled effective attenuation coefficient of water.

Results presented in Fig. 8(b) necessitated the introduction a fourth attenuation coefficient of water, coupled effective ($\mu_{eff,w}(t)$), to describe the attenuation behaviour of water in composite systems. The lack of sufficient descriptions of fundamental mechanisms and the apparently previously undiscovered coupled effect has lead in many cases to unintentionally vague descriptions of experimental procedures and potential for mistakes when performing x-ray attenuation measurements with polychromatic x-ray sources. Specifically, a review of the literature yields little information on which attenuation coefficient of water was used for determination of moisture contents in many published results, and in some cases incorrect values being used. In [14, 15, 21, 22], the experimental procedure descriptions do not specify how or if an attenuation coefficient

was measured or if tabulated (i.e., linear) values were simply used. Others [8, 12, 18] have correctly indicated the effective attenuation coefficient of water should be measured and used, but do not indicate if they accounted for the coupled effect of the parent material. Eq. 15, a modification of Eq. 11, therefore describes an improved calculation for concentration changes using a polychromatic x-ray photon source:

$$\Delta w = -\frac{\rho_w}{\mu_{eff,w}(t) \cdot t} \ln\left(\frac{I_{wet}}{I_{dry}}\right) \quad (15)$$

where $\mu_{eff,w}(t)$ is the parent material thickness-dependent (coupled) effective attenuation coefficient of water. The coupled effective attenuation coefficient of the water therefore must be directly measured for individual investigations, and tabulated values for μ_w must be avoided. The following procedure describes one method for determining values of $\mu_{eff,w}(t)$ when using a polychromatic x-ray photon source:

1. Measure the transmitted intensity through a conditioned (partially saturated or dried) specimen with given material type and thickness and an empty thin-walled acrylic (or similar) container with known inner dimensions at user-defined x-ray source setting(s), or I_e .
2. Fill the acrylic container with water or other material investigated (e.g., corrosion products [23]) and again measure the transmitted intensity through the conditioned specimen, acrylic, and water at the same x-ray source setting(s), I_f .
3. Calculate the coupled effective attenuation coefficient of water as the quotient of $-\ln\left(\frac{I_f}{I_e}\right)$ over the water thickness, t_w .

Alternatively, values for I_e and I_f determined by measuring the transmitted intensity through a dry and saturated specimen, respectively, would likely yield the same value for $\mu_{eff,w}(t)$.

Fig. 10 illustrates the impact of using various attenuation coefficients when computing changes in moisture contents. In this example, a 0.40 water-to-cement ratio cement paste specimen (28 days curing submerged in 20°C water) was dried through solvent exchange and subsequently exposed to liquid water from the specimen's top surface. Transmitted intensities were recorded from the dried specimen and repeatedly after exposure to water. The moisture profiles shown were recorded at the same time (after 8 days of exposure to water), but moisture contents were computed using different attenuation coefficients for water. Attenuation coefficients for water included a tabulated (in [29]) value

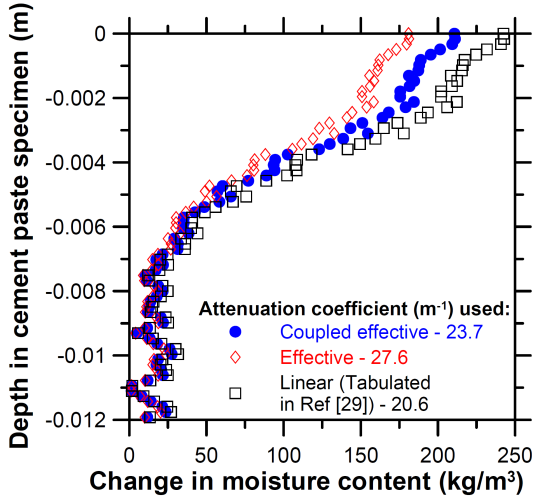


Figure 10: Example of impact of using different attenuation coefficients of water on measured change in moisture content. Data presented was collected after 8 days of exposure of liquid water to the surface of a dried (through solvent exchange) cement paste specimen using an x-ray source energy of 60 keV.

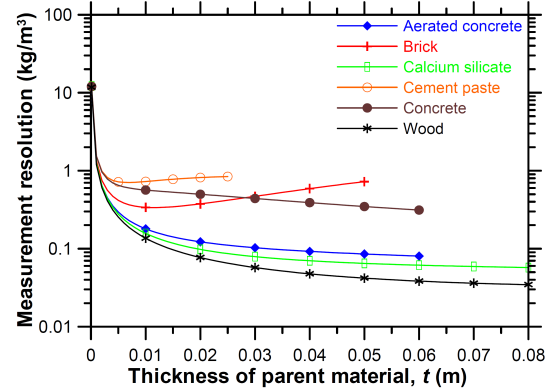


Figure 11: Resolution of moisture content changes for various materials and thicknesses with an x-ray source energy of 85 keV and current of $70 \mu\text{A}$ using fitted effective attenuation coefficients of the various building materials using Eq. 9 and coupled effective attenuation coefficient of water using Eq. 15.

589 for the linear attenuation coefficient of water at 60 keV, 20.6 m^{-1} ; the effective attenuation coefficient of water
 590 at 60 keV, 27.6 m^{-1} (according to Eq. 9 and the
 591 fitting parameters presented in Table 2); and the
 592 coupled effective attenuation coefficient, measured to be
 593 23.7 m^{-1} . Differences in moisture content (from correct
 594 values calculated using the coupled effective atten-
 595 uation coefficient) as large as 31.8 kg/m^3 were found in
 596 this case.
 597

598 Eqs. 11 and 15 indicate difference in the calcu-
 599 lated moisture content are inversely proportional to the
 600 various attenuation coefficients. Table 4 provides an
 601 overview of the percentage difference that would result
 602 from using incorrect attenuation coefficients of water
 603 assuming the x-ray source energy is 85 keV, the par-
 604 ent specimen thickness is 0.02 m, and a water thickness
 605 of 0.01 m was used to determine the effective atten-
 606 uation coefficient of water. For the materials investigated
 607 and with the stated assumptions, incorrect use of either
 608 the linear or effective attenuation coefficient for water
 609 would result in calculation errors of up to 26.9%. Table
 610 4 clearly indicates that even the use of the effective
 611 attenuation coefficient of water, $\mu_{eff,w}$, measured with-
 612 out the parent material results in potentially excessive
 613 errors.

614 Eq. 16, a modification of Eq. 12, describes the con-
 615 centration resolution when using the coupled effective

616 attenuation coefficient of water.

$$617 R_{\Delta w} = -\frac{\rho_w}{\mu_{eff,w}(t) \cdot t} \ln \left(\frac{I_{dry} - 1}{I_{dry}} \right) \quad (16)$$

618 As shown in Fig. 8, the thickness of the water, t_w used
 619 for determination of $\mu_{eff,w}(t)$ -values appears to have a
 620 minimal impact. However, the water thickness should
 621 likely closely resemble the expected maximum change
 622 in moisture content of the parent material. In this way
 623 the determination of $\mu_{eff,w}(t)$ will also act as a check
 624 to ensure the x-ray source settings and specimen thick-
 625 ness selected will be able to detect anticipated changes
 626 in moisture content.

627 Using the results presented above, the measurement
 628 resolution can be calculated using Eq. 16. Values for
 629 I_{dry} are computed using Eq. 1, where I_0 is determined
 630 from Eq. 13 and values for μ_{eff} of the various porous
 631 materials determined using Eq. 9 and the regression co-
 632 efficients in Table 2. Fig. 11 shows the concentration
 633 resolutions possible for the various parent materials us-
 634 ing an x-ray source energy of 85 keV and current of
 635 $70 \mu\text{A}$. The curves consider the beam hardening caused
 636 by the parent materials and the coupling effect on the
 637 effective attenuation coefficients of water, which occurs
 638 when polychromatic x-ray sources are used. For cement
 639 paste and clay brick optimal thicknesses are discovered,
 640 whereas for the other materials increasing the specimen
 641 thickness continues to improve measurement resolution
 642 for the studied conditions.

Table 4: Values for various attenuation coefficients of water in series with various parent materials and the potential error (%) in calculations of moisture content from x-ray attenuation measurements if the linear, μ_w or effective, $\mu_{eff,w}$ attenuation coefficient was used rather than the coupled effective attenuation coefficient, $\mu_{eff,w}(t)$.

Parent Material	Aerated concrete	Clay brick	Calcium silicate	Cement paste	Concrete	Wood
μ_w [m^{-1}]				18.04		
$\mu_{eff,w}$ [m^{-1}]				24.13		
$\mu_{eff,w}(t)$ [m^{-1}]	20.9	20.0	22.8	17.9	18.7	22.9
Error using μ_w (%)	15.8	10.6	26.5	0.7	3.9	26.9
Error using $\mu_{eff,w}$ (%)	13.5	17.4	5.5	25.8	22.4	5.2

6. Summary and conclusions

This paper presents a detailed description of the controlling mechanism behind x-ray attenuation measurement using both monochromatic and polychromatic x-ray sources and non-energy discretizing x-ray cameras. Differences in measurements taken with monochromatic and polychromatic sources are described. Various attenuation coefficients, including linear, average, effective, and coupled effective (for water only), were defined and described in order to shed light on possible mistakes made in previous investigations. It is concluded that using tabulated values of the linear attenuation coefficient, μ_w of water should not be used for calculating moisture content changes.

Additionally, a parametric investigation on material type, thickness and x-ray source energy settings on the ability to resolve changes in moisture content was presented. Results obtained through the parametric studies performed on aluminium, acrylic, clay brick, aerated autoclaved concrete, cement paste, concrete, calcium silicate, and wood indicated that:

- The effective attenuation coefficient of water is affected by the type and thickness of the parent specimen, necessitating the introduction of the coupled effective attenuation coefficient of water, $\mu_{eff,w}(t)$. Therefore, when utilizing the x-ray attenuation measurement technique to calculate moisture content (or other material concentration) changes, the attenuation coefficient of water must be directly measured while the parent specimen is in place. A procedure for determining $\mu_{eff,w}(t)$ is presented.
- It was found that calculations of change in moisture content (i.e., Eq. 15) may vary by up to 26.9% if the incorrect attenuation coefficient of water is used.
- Material-dependent optimization functions – in terms of ability to resolve water – have been de-

veloped for clay brick, aerated autoclaved concrete, cement paste, concrete, calcium silicate, and wood. These functions consider specimen thickness, x-ray source energy and current, beam hardening caused by the parent material, and the coupled effective attenuation coefficient of water. This extensive study provides a simple means for determination of proper specimen size and x-ray energy settings for a variety of materials.

Acknowledgements

The financial support of the Danish Expert Centre for Infrastructure Constructions, the Hans-Christian-Ørsted fund, and the funding by the Danish research council (project number 274-07-0393) is gratefully acknowledged.

References

- [1] M. Bäckström, M. Al-Emrani, R. Kliger, [Timber-framed partition walls and their restraining effect on warp in built-in wall studs - models for twist](#), *Construction and Building Materials* 23 (12) (2009) 3556 – 3563.
- [2] R. Hendrickx, S. Roels, K. V. Balen, [Measuring the water capacity and transfer properties of fresh mortar](#), *Cement and Concrete Research* 40 (12) (2010) 1650 – 1655.
- [3] P. Lura, B. Pease, G. Mazzotta, F. Rajabipour, J. Weiss, Influence of shrinkage-reducing admixtures on the development of plastic shrinkage cracks, *ACI Materials Journal* 104 (2) (2007) 187–194.
- [4] A. El-Turki, R. J. Ball, S. Holmes, W. J. Allen, G. C. Allen, [Environmental cycling and laboratory testing to evaluate the significance of moisture control for lime mortars](#), *Construction and Building Materials* 24 (8) (2010) 1392 – 1397.
- [5] L. Poulidakos, M. Partl, [Evaluation of moisture susceptibility of porous asphalt concrete using water submersion fatigue tests](#), *Construction and Building Materials* 23 (12) (2009) 3475 – 3484.
- [6] C. Hazaree, H. Ceylan, K. Wang, [Influences of mixture composition on properties and freeze-thaw resistance of rcc](#), *Construction and Building Materials* 25 (1) (2011) 313 – 319.
- [7] S. Roels, J. Carmeliet, H. Hens, O. Adan, H. Brocken, R. Cerny, Z. Pavlik, A. T. Ellis, C. Hall, K. Kumaran, L. Pel, R. Plagge,

- 720 A comparison of different techniques to quantify moisture content profiles in porous building materials, *Journal of Building Physics* 27 (4) (2004) 261.
- 721
- 722
- 723 [8] K. Hansen, S. Jensen, L. Gerward, K. Singh, Dual-energy x-ray absorptiometry for the simultaneous determination of density and moisture content in porous structural materials, in: *Proceedings of the 5th Symposium on Building Physics in the Nordic Countries*, Vol. 1, Gothenburg, Sweden, 1999, pp. 281–288.
- 724
- 725
- 726
- 727
- 728 [9] D. Bentz, M. Geiker, K. Hansen, Shrinkage-reducing admixtures and early-age desiccation in cement pastes and mortars, *Cement and Concrete Research* 31 (7) (2001) 1075–1085.
- 729
- 730
- 731 [10] J. Hu, P. Stroeven, X-ray absorption study of drying cement paste and mortar, *Cement and Concrete Research* 33 (3) (2003) 397–403.
- 732
- 733
- 734 [11] A. Bayer, H.-J. Vogel, R. K., Direct measurement of the soil water retention curve using x-ray absorption, *Hydrology and Earth System Sciences* 8 (1) (2004) 2–7.
- 735
- 736
- 737 [12] R. Baetig, R. Rémond, P. Perré, Measuring moisture content profiles in a board during drying: a polychromatic x-ray system interfaced with a vacuum/pressure laboratory kiln, *Wood Science and Technology* 40 (4) (2006) 261–274.
- 738
- 739
- 740
- 741 [13] P. Lura, D. P. Bentz, D. A. Lange, K. Kovler, A. Bentur, K. Van Breugel, Measurement of water transport from saturated pumice aggregates to hardening cement paste, *Materials and Structures/Materiaux et Constructions* 39 (293) (2006) 861–868.
- 742
- 743
- 744
- 745
- 746 [14] S. Roels, J. Carmeliet, Analysis of moisture flow in porous materials using microfocus x-ray radiography, *International Journal of Heat and Mass Transfer* 49 (25) (2006) 4762–4772.
- 747
- 748
- 749 [15] P. H. Baker, D. Bailly, M. Campbell, G. H. Galbraith, R. C. McLean, N. Poffa, C. H. Sanders, The application of x-ray absorption to building moisture transport studies, *Measurement* 40 (9–10) (2007) 951–959.
- 750
- 751
- 752
- 753 [16] B. Pease, J. Couch, M. Geiker, H. Stang, J. Weiss, Assessing the portion of the crack length contributing to water sorption using x-ray absorption measurements on concrete wedge splitting specimens, in: *ConcreteLife'09: Second International RILEM Workshop on Concrete Durability and Service Life Planning*, Haifa, Israel, 2009.
- 754
- 755
- 756
- 757
- 758
- 759 [17] M. Pour-Ghaz, F. Rajabipour, J. Couch, J. Weiss, Modeling fluid transport in cementitious systems with crack-like (notch) geometries, in: *ConcreteLife'09: Second International RILEM Workshop on Concrete Durability and Service Life Planning*, 2009.
- 760
- 761
- 762
- 763
- 764 [18] G. Sant, J. Weiss, Using x-ray absorption to assess moisture movement in cement-based materials, *Journal of ASTM International* 6 (9) (2009) 1–15.
- 765
- 766
- 767 [19] Z. Ye, M. Tirovic, M. Davies, P. Baker, M. Phillipson, C. Sanders, G. Galbraith, R. McLean, The testing of two methods for the moisture measurement of building fabrics via comparisons with data from an x-ray system, *Building and Environment* 44 (7) (2009) 1409–1417.
- 768
- 769
- 770
- 771
- 772 [20] B. Pease, Influence of concrete cracking on ingress and reinforcement corrosion, Ph.D. thesis, Technical University of Denmark, Kgs. Lyngby, Denmark (December 2010).
- 773
- 774
- 775 [21] S. Roels, K. Vandersteen, J. Carmeliet, Measuring and simulating moisture uptake in a fractured porous medium, *Advances in Water Resources* 26 (3) (2003) 237–246.
- 776
- 777
- 778 [22] J. Carmeliet, J.-F. Delerue, K. Vandersteen, S. Roels, Three-dimensional liquid transport in concrete cracks, *International Journal for Numerical and Analytical Methods in Geomechanics* 28 (7) (2004) 671–688.
- 779
- 780
- 781
- 782 [23] A. Michel, B. J. Pease, M. R. Geiker, H. Stang, J. F. Olesen, [Monitoring reinforcement corrosion and corrosion-induced cracking using non-destructive x-ray attenuation measurements](#), *Cement and Concrete Research* In Press, Corrected Proof (2011).
- 783
- 784
- 785
- 786
- 787 [24] G. Knoll, *Radiation Detection and Measurement*, J. Wiley, 1989.
- 788
- 789 [25] N. Tsoufanidis, *Measurement and detection of radiation*, Hemisphere Publishing Corp., 1983.
- 790
- 791 [26] M. Van Geet, *Optimisation of microfocus x-ray computer tomography for geological reserach with special emphasis on coal components (macerals) and fractures (cleats) characterisation*, Ph.D. thesis, KU Leuven, Heverlee (April 2001).
- 792
- 793
- 794 [27] N. Dyson, *X-rays in atomic and nuclear physics*, Longman Scientific & Technical, London, 1973.
- 795
- 796 [28] T. Curry, J. Dowdey, R. Murry, *Christensens Physics of Diagnostic Radiology*, Lea and Febiger, London, 1990.
- 797
- 798 [29] J. Hubbell, S. Seltzer, [Tables of x-ray mass attenuation coefficients and mass energy-absorption coefficients](#), Tech. Rep. NISTIR 5632, National Institute of Standards and Technology, Maryland, USA (2004).
- 799
- 800
- 801 [30] R. A. Brooks, G. D. Chiro, Beam hardening in x-ray reconstructive tomography, *Physics in Medicine and Biology* 21 (3) (1976) 390–398.
- 802
- 803
- 804 [31] J. Weiss, M. Geiker, K. Hansen, Using x-ray absorption to detect fluid ingress in cracked concrete, Tech. rep., Purdue University School of Civil Engineering, West Lafayette, IN, USA, 15 pp. (2004).
- 805
- 806
- 807 [32] C. Kleinschmidt, Analytical considerations of beam hardening in medical accelerator photon spectra, *Medical Physics* 26 (9) (1999) 1995–1999.
- 808
- 809
- 810 [33] J. Alles, R. F. Mudde, Beam hardening: analytical considerations of the effective attenuation coefficient of x-ray tomography, *Medical Physics* 34 (7) (2007) 2882–2889.
- 811
- 812
- 813 [34] [www.gni.dk](#), GNI X-ray System, Technical University of Denmark (2010).
- 814
- 815 [35] K. Spartiotis, J. Havulinna, A. Leppänen, T. Pansar, K. Puhakka, J. Pyyhtiä, T. Schulman, A cdte real time x-ray imaging sensor and system, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 527 (3) (2004) 478–486.
- 816
- 817
- 818
- 819
- 820
- 821

Nomenclature

α Fitting parameter equivalent to $a^* \cdot Z_{eff}^{3.8}$ [-]

Δw Change in moisture content [kg/m³]

η Parent material coupling effect [m⁻¹]

κ Pair production cross-section [m⁻¹]

λ Beam hardening coefficient [m⁻¹]

μ Linear attenuation coefficient [m⁻¹]

μ_0 Initial average attenuation coefficient [m⁻¹]

μ_w Linear attenuation coefficient of water [m⁻¹]

$\mu_{eff,w}$ Effective attenuation coefficient of water [m⁻¹]

$\mu_{eff,w}(0)$ Coupled effective attenuation coefficient of water with 0 m thickness of parent material, fitted [m⁻¹]

835	$\mu_{eff,w}(\infty)$ Coupled effective attenuation coefficient of	872	I_{dry} Transmitted intensity through a conditioned (dry or
836	water with infinite thickness of parent material, fitted	873	partially saturated) porous material [counts]
837	$[\text{m}^{-1}]$		
838	$\mu_{eff,w}(t)$ Coupled effective attenuation coefficient of	874	I_{wet} Transmitted intensity through a wet porous mate-
839	water $[\text{m}^{-1}]$	875	rial [counts]
840	μ_{eff} Effective attenuation coefficient $[\text{m}^{-1}]$	876	$R_{\Delta w}$ Resolution of change in moisture content $[\text{kg}/\text{m}^3]$
841	$\bar{\mu}$ Average attenuation coefficient $[\text{m}^{-1}]$	877	t Parent material (specimen) thickness [m]
842	Ψ_E Shape of photon energy spectrum [-]	878	t_w Thickness of water [m]
843	ρ Density $[\text{kg}/\text{m}^3]$	879	x Incremental thickness [m]
844	ρ_w Density of water $[\text{kg}/\text{m}^3]$	880	Z_{eff} Effective atomic number [-], $Z_{eff} = \sqrt[3.8]{\sum f_i Z_i^{3.8}}$
845	σ Compton cross-section $[\text{m}^{-1}]$	881	Z_i Atomic number of element i [-]
846	τ Photoelectric cross-section $[\text{m}^{-1}]$		
847	a Fitting parameter describing photoelectric cross-		
848	section [-]		
849	a^* Fitting parameter describing effective photoelectric		
850	cross-section [-]		
851	b Fitting parameter describing photoelectric cross-		
852	section[-]		
853	b^* Fitting parameter describing effective photoelectric		
854	cross-section[-]		
855	c Fitting parameter describing Compton cross-section		
856	[-]		
857	c^* Fitting parameter describing effective Compton		
858	cross-section [-]		
859	C_{1-6} Fitting parameters [-]		
860	E X-ray photon energy [keV]		
861	E_m X-ray source energy (maximum photon energy pro-		
862	duced by polychromatic x-ray source) [keV]		
863	f_i Fraction of element i [-]		
864	I Transmitted intensity [counts]		
865	i X-ray source current $[\mu\text{A}]$		
866	I_0 Initial intensity [counts]		
867	I_e Transmitted intensity through parent material and		
868	empty container [counts]		
869	I_f Transmitted intensity through parent material and		
870	container filled with water or other substance being in-		
871	vestigated [counts]		