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Plastic analysis of shear in beams, deep beams and corbels

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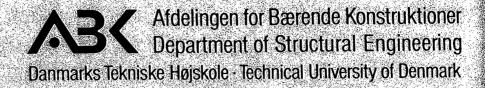
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Plastic Analysis of Shear in Beams Deep Beams and Corbels

Chen Ganwei

Serie R No 237 1988

PLASTIC ANALYSIS OF SHEAR IN BEAMS
DEEP SEAMS AND CORBELS

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Chen Ganwei

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Web Reinforcement

NOTATIONS

The symbols are defined when they first occur in the text. The frequently used notations are listed alphabetically below.

- a: Clear shear span
- A: Internal work
- A_c: Area of cross-section
- a_{CC}: Shear span, measured from concentrated load to the reaction for beams or measured from concentrated load to the face of column for corbels
- $\mathbf{A_S}$: Cross-sectional area of main tensile reinforcement
- ${
 m A}_{
 m CS}$: Cross-sectional area of compressive reinforcement
- $\mathbf{A}_{\mathrm{Swh}}$: Cross-sectional area of horizontal web reinforcement per unit depth
- $\mathbf{A}_{\mathbf{SWV}}$: Cross-sectional area of vertical web reinforcement per unit length
- b: Web width of beam
- bl: Width of compressive flange
- c: Internal cohesion of material
- C: Resultsant compressive force of concrete.

Parameter,
$$C = \begin{cases} 1-2\phi^* & \phi^* \le \frac{1}{2} \\ 0 & \phi^* > \frac{1}{2} \end{cases}$$

c₁, c₂: Constants

Parameter,
$$C_1 = \begin{cases} 1-2(n^*+\phi^*) & n^*+\phi^* < \frac{1}{2} \\ 0 & n^*-\phi^* \le \frac{1}{2} \le n^*+\phi^* \\ 1-2(n^*-\phi^*) & \frac{1}{2} < n^*-\phi^* \end{cases}$$

v: Coefficient of variation

Effective depth of beam. Length of line

D: Parameter,
$$D = \begin{cases} \frac{\mu - 2\phi^*}{\lambda} & \phi^* \le \frac{\mu}{2} \\ 0 & \phi > \frac{\mu}{2} \end{cases}$$

$$D_o: Parameter, D_o = \begin{cases} \frac{\mu_o - 2T/bhf_c^*}{\lambda_o} & T/bhf_c^* \le \frac{\mu_o}{2} \\ 0 & T/bhf_c^* \le \frac{\mu_o}{2} \end{cases}$$

$$D_1: Parameter, D_1 = \begin{cases} \frac{\mu - 2(n^* + \phi^*)}{\lambda_o} & n^* + \phi^* \le \frac{\mu}{2} \le n^* + \phi^* \\ 0 & n^* - \phi^* \le \frac{\mu}{2} \le n^* + \phi^* \end{cases}$$

 $\mathbf{d}_{\mathbf{C}}$: Distance from extreme compression fibre to centroid of compression reinforcement

- dg: Stirrup depth
- $f_{
 m A}$: Separation resistance
- : Uniaxial compressive strength of concrete
- f_{C}^{*} : Plastic compressive strength of concrete, defined as $f_{C}^{*} = \nu f_{C}$

 $F_{\mathtt{Se}}$: Effective prestress acting on the section

ft: Tensile strength of concrete

 f_t^* : plastic tensile strength of concrete, defined as $f_t^* = v_t f_t$

 $\mathbf{f}_{\mathbf{y}}$: Yield strength of tensile reinforcement

 f_{yc} : Yield strength of compressive reinforcement

 f_{yw} : Yield strength of web reinforcement

 $\mathbf{f}_{\mathbf{ywh}}$: Yield strength of horizontal web reinforcement

h: Depth of cross-section

hl: Depth of compressive flange

h*: Effective shear depth

k: Material constant, $k = (1+\sin\varphi)/(1-\sin\varphi)$. Length of loading platen

k₁,k₂: Constants

1: Material constant, $l = 1-(k-1)f_t/f_c$

Lo: Span of beam

m: Material constant, $m = 1 - (k+1) f_t / f_c$

Mp: Bending yield moment

mp: Effective dimensionless bending yield moment, defined as

 $m_p^* = M_p/bd^2f_c^*$

n: Case Madison, steamed to dringer of tests

N: Normal force

 n^* : Effective normal force degree, defined as $n^* = N/bhf_C^*$

P: External load

Function of curved yield line: Polar coordinate

H

ro, r1: Parameters in polar coordinate system

 r_{m} : Length of line

 s_y : Equivalent yield stress of web reinforcement, defined as $s_y = \lambda_{swv} f_{yw}/b$ or $s_y = \lambda_{swh} f_{yw}/b$

t: Tangential direction. Distance

T: Resultant force of tensile reinforcement

Relative displacement in yield line

!

V: Ultimate load. Volume

 v_{cal} : Calculated ultimate load

 $v_{\text{t'}} v_{\text{test}}$: Observed ultimate load in test

W: Internal plastic work per unit volume (dissipation)

WE: External work

W_I: Internal work

 w_1 : Dissipation per unit length

x: Depth of rectangular stress block. Cartesian coordinate.
Distance

- xo: Distance
- x: Mean
- y: Distance measured from bottom of cross section to centroid of tensile reinforcement. Cartesian coordinate
- y_{o} : Distance from top of cross section to neutral axis
- :: Internal moment lever arm
- : Angle. Ratio, $\alpha = d_c/d$
- α_1 : Ratio, $\alpha_1 = \phi_{hw}^*/\phi_b^*$
- $\alpha_{\rm m}, \beta, \tau$: Angles
- nt: Shear strain in n,t-coordinate system
- $\delta, \delta_{m}, \delta_{s}$: Relative displacement rates.
- A; Width of deforming zone idealized as yield line
- 61,62,63: Principal strains
- en'et: Strains in n- and t-directions, respectively
- η : Rotation rate. Ratio, $\eta = \frac{\lambda}{\mu}$
- θ: Angles. Polar coordinate
- θ_0 , θ_1 : Angles

- λ : Proportionality factor in flow law. Material constant, $\lambda = 1 (k-1)\rho^*$
- λ_1 : Proportionality factor in flow law
- λ_0 : Material constant, $\lambda_0 = 1 k \frac{\sigma_1}{t} + \frac{\sigma_1}{t}$
- μ : Coefficient of friction, μ = tan φ . Material constant μ = 1-(k+1) ρ *
- o: Material constant, $\mu_0 = 1-k \frac{\sigma_1}{t_c} \frac{\sigma_1}{t_c}$
- $\mathfrak v$: Effectiveness factor for the compressive strength of concrete
- v_{b} : Effectiveness for bending
- $\mathfrak{v}_{\mathsf{t}}$: Effectiveness factor for the tensile strength of concrete

Parameter,
$$\xi = \begin{cases} \frac{k-1}{4k}, \\ n^* - \phi^*, \end{cases}$$

'n

$$n^{*}+\phi^{*} < \frac{k-1}{4k}$$

$$n^{*}-\phi^{*} \le \frac{k-1}{4k} \le n^{*}+\phi^{*}$$

$$\frac{k-1}{4k} < n^{*}+\phi^{*}$$

- ρ^* : Effectiveness factor, $\rho^* = f_t^*/f_c^*$
- σ : Normal stress. Standard deviation
- $\sigma_{1}, \sigma_{2}, \sigma_{3}$: Principal stresses
- $\sigma_{ extsf{I}}, \sigma_{ extsf{II}}$: Principal stresses
- σ_b,σ_V: Stresses
- $\sigma_{\mathbf{x}'}\sigma_{\mathbf{y}}$: Normal stresses in x,y-coordinate system

Shear stress in x,y-coordinate system

6 as $\varphi = A_{S}/bh$ Angle of friction for concrete. Reinforcement ratio, defined

ช์. Reinforcement ratio for bending, defined as $\phi_{\rm b}={
m A_S/bd}$

wh: φ_{WV} : Vertical web reinforcement ratio, defined as $\varphi_{\mathrm{WV}} = \mathrm{A}_{\mathrm{SWV}}/\mathrm{b}$

Horizontal web reinforcement ratio, defined as $\phi_{
m wh} = {
m A_{Swh}/b}$

Effective reinforcement degree, defined as $\phi^* = A_s f_y/bh f_c^*$

* ..*

٠<u>.</u>* Effective reinforcement degree for bending, $\phi_{b}^{*} = A_{s}f_{y}/bdf_{c}^{*}$

ບໍ່.* Effective compressive reinforcement degree, $\phi_{_{_{\mathbf{C}}}}^{*}=\mathrm{A_{_{\mathbf{SC}}}}_{_{\mathbf{YC}}}/\mathrm{bdf}_{_{\mathbf{C}}}^{*}$

₽.* Effective horizontal reinforcement degree, $\phi_h^* = A_s f_y/bhf_c^* +$ Aswhfyw/bfc

ф_{ћw}: Effective horizontal web reinforcement degree, ϕ_{hw}^* = Aswhfyw/bfc*

∢.* Effective vertical web reinforcement degree, ϕ_V^* Aswvfyw /bfc*

 ϕ_{VO}^* : Critical web reinforcement degree

×. Parameter, x = a ≤ tanφ

 $\frac{a}{h} \geq \tan \varphi$

CHAPTER I. INTRODUCTION

[77.2], [78.3] and [79.1]. Denmark [67.1], [74.3], [75.1], [75.4], [76.3], [76.6], [77.1], concrete in shear, which had been developed since the 1960's in for concrete inspired by the achievements of plastic theory of reinforcement in 1985, I became interested in limit state analysis During shear tests on deep concrete beams without web

conventional slender beams and relatively low reinforcement ratio has been widely used in such structures as tanks, bins, silos and Deep beam is the beam having a depth comparable to its span. It factor in deep beams without web reinforcement. because of the bigger concrete web. Shear is very often a dominant relatively in buildings. Because of the geometry, deep beams may have low shear span-to-depth ratio compared

accumulated experience may be sufficiently safe, but they need not more than empirical formulae deduced from a large number of tests. the results obtained. Just as most building codes contain little report specific test series and supply empirical formulae to fit on the subject is almost overwhelming. Most papers, however on shear behavior and shear capacity of deep beams. The literature be particularly economical. The traditional design approach for deep concrete beams based on extensively for decades. There have been conducted a lot of work corbels. Shear in reinforced concrete deep beams has been studied independent kind of structural members different from beams and For a long time deep beams were classified and treated as an

concerned, the only difference between beams, deep analysis in beams with that of deep beams and corbels. According advantages. Firstly, as a general theory, it unifies the shear rational mathematical theory of plasticity has considerable to the limit state analysis, so far as the shear problem is Compared to the hitherto dominating empirical methods, beams and

ε

Angle

corbels is the shear span-to-depth ratio. Normal beams have relatively large shear span-to-depth ratios, while corbels have very small shear span-to-depth ratios and deep beams lie between them. Secondly, as a rational mathematical model describing the mechanism of the shear members, it leads to a thorough understanding of the failure mechanism in concrete. Therefore, by means of this theory one can calculate the ultimate shear capacity effectively and simply, and also dimension and detail the shear members economically.

modified simple solutions for conventional beams are then presencomplete solutions are compared with the simple solutions and the non-zero tension cut off have been derived in Chapter IV. The as a perfectly plastic modified Coulomb-material with a small, but and corbels, because they neglect the tensile strength of concrete difficult and uncertain when dealing with empirical methods. An such an extension is more or less impossible or at least extremely within related areas not directly covered by cases studied, while in the web. Therefore, the exact solutions regarding the concrete corresponding v formula in [79.4] are not suitable for deep beams solutions [78.2] with the so-called square yield locus and solved by this theory. It has been found that the simple plastic of deep beams with and without web reinforcement can also be beams, it is natural to investigate whether the ultimate strength remarkably well the shear failure load of reinforced concrete plastic approach is simple, consistent and able to concrete can be found in [87.3]. Since many examples [73.7], example of extending the plastic solution to fibre reinforced Furthermore, it is easy to extend the theory into applications [76.6], [77.7], [78.5] and [80.2] have fully proved that the predict

Results derived by plastic theory have to be modified in an empirical manner because of the well known fact that concrete is not perfectly plastic and the whole failure surface will not be fully active. To take this anto account in a simple way the uniaxial compressive strength $f_{\rm C}$ is reduced by the effectiveness factor $\nu_{\rm c}$ which must be correlated with the experimental evidence.

Therefore, in this paper a big effort has been done to develop a more general and rational approach for evaluating an approximate $\nu\text{-value}$ in some shear cases, especially for shear members without web reinforcement.

The theoretical calculations are compared with a wide range of available test results of beams, deep beams and corbels found in literature. A satisfactorily good agreement has been found.

4

CHAPTER II. PLASTIC THEORY FOR CONCRETE

2.1 Failure and yield criteria

Failure and yield criteria are the criteria that must be satisfied for the stress field at yielding. The purpose of formulating failure and yield criteria for a material is to enable us to determine their behavior by means of simple tests from which we can evaluate the risk of failure or yielding for complicated stress fields.

Since our knowledge of the structure and composition of materials does not yet enable us to develop the failure criteria based on known natural laws, most failure criteria appear as hypotheses whose application to various materials will have to be evaluated from tests.

2.1.1 Coulomb's failure hypothesis and Coulomb material

Coulomb's failure hypothesis was presented in 1773 by C.A. Coulomb, who had remarked that failure in stone prisms subjected to unlaxial compression took place along certain surfaces. These surfaces are called sliding surfaces, and this type of failure is known as sliding failure.

Sliding failure is assumed to occur in a section if the shear stress $|\tau|$ in this section exceeds the sliding resistance, which can be determined by two contributions. One contribution is the internal cohesion, denoted c, which is a material constant. The other contribution comes from a kind of internal friction and equals a certain fraction μ of the normal stress σ in the sliding plane. The quantity μ is called the coefficient of friction. If σ is a compressive stress, it gives a possitive contribution; if σ is a tensile stress, it gives a negative contribution. The condition for sliding failure therefore is

where c and μ are positive constant and σ is counted positive as a tensile stress. A material complying with failure condition (2.1.1) is called a Coulomb material. Coulomb's failure hypothesis is governed entirely by two parameters, e.g. c and μ .

Fig. 2.1 shows Coulomb's failure hypothesis depicted in a $(\sigma,~\tau)$ coordinate system. The figure includes Mohr's circles for the stresses at a point at which the failure hypothesis is satisfied.

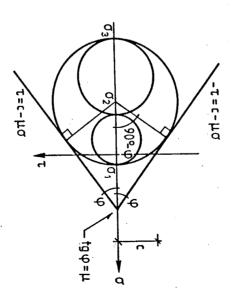


Figure 2.1. Coulomb's failure hypothesis with Mohr's circles.

The angle φ given by tg $\varphi = \mu$ is called the angle of friction.

A more general hypothesis was formulated by O. Mohr in 1882. Mohr assumed that failure occurs when the stresses in a section satisfy the condition

 $f(\sigma, \tau) = 0$

where $f(\sigma, \ au)$ is a function characteristic of the material, and

where σ and τ are normal stress and the shear stress, respectively, in the section. Coulomb's failure hypothesis is thus a special case of Mohr's theory.

2.1.2 Coulomb's modified failure hypothesis and modified Coulomb material

For a large group of materials it appears that reasonable failure conditions are obtained by combining Coulomb's sliding hypothesis with a bound for the maximum tensile stress. The resulting failure criterion makes it natural to distinguish between two failure modes, sliding failure and separation failure. In bothe cases the name of the failure refers to what we imagine the relative motion between particles on each side of the failure surface to be. At sliding failure there is motion parallel to the failure surface, while at the separation failure motion is perpendicular to the failure surface.

Separation failure occurs when the biggest tensile stress σ_1 in a section exceeds the separation resistance f_A , i.e.

Q1 = fA

This failure hypothesis, which results from a combination of [2.1.1] and (2.1.3) is called Coulomb's modified failure hypothesis.

A material complying with (2.1.1) and (2.1.3) is called a modified coulomb material. As is clear, Coulomb's modified failure hypothesis requires three parameters to be determined, e.g., c, μ and f_A . Fig. 2.2 shows the modified Coulomb's failure hypothesis illustrated in a (σ, τ) -coordinate system.

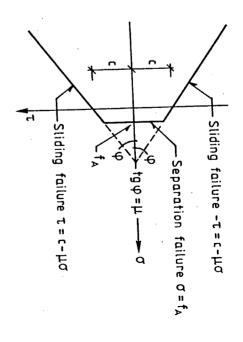


Figure 2.2. Failure criterion for a modified Coulomb material.

The straight lines, which represent the failure criterion, divide the plane into two regions. When the stresses in a section the plane into two regions. When the stresses in a section correspond to a Mohr's circle lying within the boundary lines, no correspond to a Mohr's circle stresses corresponding to circles failure will occur, while stresses combinations which in fact touching the lines represent stress combinations which in fact involve failure. The failure mode depends on whether the contact point lies on the lines $|\tau|=c-\mu\sigma$, which involve sliding failure, or on the line $\sigma=f_A$, which involves separation failure.

By means of drawing Mohr's circle it is easy to investigate whether the stress field in a point, given by the principal whether the stress field in a point, given by the principal stresses σ_1 , σ_2 and σ_3 , where $\sigma_1 > \sigma_2 > \sigma_3$, will cause failure. Because on Mohr's circles the points closest to the boundary lines lie on the circle with $(\sigma_1 - \sigma_3)$ as diameter, we have only to focus on the points on this circle. If the circle with diameter $(\sigma_1 - \sigma_3)$ lies within the boundary lines, failure will not occur

Figure 2.3. Mohr's circles of principal stresses and the failure criterion.

If the coircle with diameter $(\sigma_1-\sigma_3)$ touches the boundary lines corresponding to sliding failure, a sliding failure may occur in two sections for reasons of symmetry, (see Fig. 2.4).

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Figure 2.4. Mohr's circles at sliding failure.

If the circle touches the line corresponding to separation failure, a separation failure will then take place, (see Fig. 2.5).

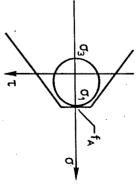


Figure 2.5. Mohr's circles at separation failure.

By means of the Fig. 2.4, we find that equation (2.1.1) for sliding failure can be written in principal stress form in the following way:

$$\frac{1}{2}(\sigma_1 - \sigma_3) = c \cos \varphi - \frac{1}{2}(\sigma_1 + \sigma_3) \sin \varphi$$

(2.1.4)

which is valid for $\sigma_1 > \sigma_2 > \sigma_3$. If the mutual magnitudes of the principal stresses are altered, equation (2.1.4) must be altered correspondingly.

Introducing

$$k = (\frac{\cos\varphi}{1 - \sin\varphi})^2 = \frac{1 + \sin\varphi}{1 - \sin\varphi} = \tan^2(\frac{\pi}{4} + \frac{\varphi}{2})$$

(2.1.5)

The condition for sliding failure can be written in the simple form

$$k\sigma_1 - \sigma_3 = 2c\sqrt{k} \tag{2.1.6}$$

while the condition for separation failure is

$$\sigma_1 = f_A \tag{2.1.7}$$

The uniaxial compression strength f_c of a material is determined by a test, where the stress field at failure is given by $(\sigma_1, \ \sigma_2, \ \sigma_3) = (0, \ 0, \ -f_C)$. Since the uniaxial compression test always involves sliding failure (see Fig. 2.6), we find from equation (2.1.6)

$$-\sigma_3 = f_C = 2c\sqrt{k}$$
 (2.1.8)

whereupon equation (2.1.6) can be written as

$$\lambda \sigma_1 - \sigma_3 = f_C \tag{2.1.9}$$

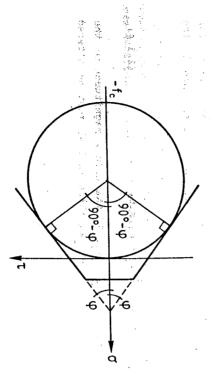
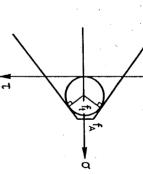
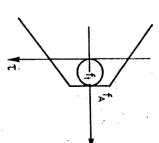


Figure 2.6. Mohr's circle at pure compression.

In a similar way the uniaxial tensile strength of a material f_t can be introduced as the stress field at failure that is given by $(\sigma_1, \ \sigma_2, \ \sigma_3) = (f_t, \ 0, \ 0)$.

As seen clearly from Fig. 2.7, the tensile test holds the possibility of sliding failure as well as separation failure.





sliding failure.

Separation failure.

Figure 2.7. Mohr's circle at pure tension.

In the case of sliding failure we have, applying equation (2.1.6)

$$f_{t} = \frac{1}{k} f_{c}$$
 (2.1.10)

In the case of separation failure we get, applying equation (2.1.7)

$$f_{\mathbf{A}} = f_{\mathbf{A}} \tag{2.1.11}$$

This means that the tensile failure is a sliding failure when

$$\frac{1}{K} f_C < f_A$$
 (2.1.12)

and a separation failure when

(2.1.13)

In the case of $rac{1}{k}$ $f_{_{f C}}$ < $f_{_{f A}}$, the equation (2.1.6) can then be written as

$$\frac{1}{1} - \frac{\sigma_3}{1} = 1$$
 (2.1.14)

2.1.3 Failure criterion for concrete

A lot of tests have shown that in many respects concrete can be considered as a modified Coulomb material, the parameter k being about 4 [84.1]. The agreement has been found to be reasonable. The modified Coulomb's failure hypothesis can thus be used as yield condition for concrete.

With k = 4, we find the constant $\mu = \tan \varphi$

$$\mu = 0.75 \tag{2.1.15}$$

corresponding to an angle of friction

$$\psi = 37^{\circ} \tag{2.1.16}$$

From equation (2.1.8), we get

$$c = \frac{f_C}{2\sqrt{k}} = \frac{1}{4}f_C \tag{2.1.17}$$

In the tests we find that both the tension failure and the shear failure are separation failure. If the tensile strength is $\mathbf{f}_{\xi'}$ we therefore have

$$f_{\mathbf{A}} = f_{\mathbf{t}} \tag{2.1.18}$$

The failure condition (2.1.6) and (2.1.7) then get the final form

$$4\sigma_1 - \sigma_3 = f_C$$
 (2.1.19)

$$\sigma_1 = f_A = f_t \tag{2.1.20}$$

The failure condition for concrete in the (σ, τ) -coordinate system is shown in Fig. 2.8.

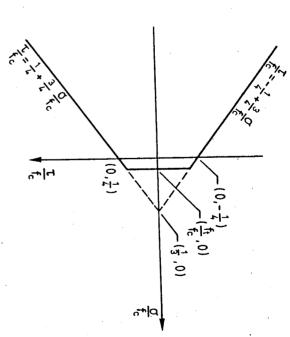


Figure 2.8. Failure condition for concrete in a $(\sigma, \ \tau)$ -coordinate system.

It's worth to notice that Coulomb's modified failure hypothesis does not completely describe the failure criterion for concrete. According to Coulomb's failure hypothesis the intermediate principal stress σ_2 has no effect on the carrying capacity. However, tests show that it has an effect in disagreement with the hypothesis. Some biaxial compression tests show that $\sigma_3/f_C = -1.2$ to -1.4 for $\sigma_1/f_C = -0.6$, whereas Coulomb's failure hypothesis gives $\sigma_3/f_C = -1.0$. Biaxial tensile tests and tensile/compressive tests also show that the Coulomb's modified failure hypothesis is not always convex and is a little on the unsafe side when tensile stresses occur.

Even then, the hypothesis must anyway be described as usable because the deviations are often moderate, and because there is a great dependence of the strength properties on the test circumstances [77.1]. Furthermore, in plastic theory we do not need a very accurate failure condition because of the large empirical modifications necessary for getting agreement between tests and theory (the v-value).

In the following chapters we will throughout assume that the concrete is a rigid, perfectly plastic material with the modified Coulomb's failure criterion as yield condition and the associated flow rule (normality condition) as constitutive equations.

2.2 Theory of plasticity

A rigid-plastic material is defined as a material in which no deformations occur for stresses up to a certain limit, the yield point. The yield condition describes the states of stress which may result in failure of the material. For stresses at the yield point, arbitrarily large deformations are possible without any change in the stresses.

A rigid-plastic material does not exist in reality. However, it is possible to use the model if the plastic strains are much larger than the elastic strains.

For an isotropic material we may express the failure condition in terms of the principal stresses as

$$f(\sigma_1, \ \sigma_2, \ \sigma_3) = 0$$
 (2.2.1)

This means that stresses rendering f < 0 correspond to stresses that can be sustained by the material and thus these stress combinations give no strain, while f > 0 corresponds to stress fields which the material cannot sustain and f = 0 means that plastic deformation may take place.

The relationship between the strains is then determined by the yield law (also called associated flow rule)

$$\epsilon_{\perp} = \lambda \frac{\partial f}{\partial \sigma_{\perp}}$$
 (i = 1, 2, 3) (2.2.2)

In a 3 dimensional representation, with σ_1 , σ_2 and σ_3 as coordinate axes, (2.2.1) produces a so-called yield surface. If we

regard (2.2.2) as a vector in the σ_1 , σ_2 , σ_3 -coordinate system, it is a normal vector to the yield surface at the point $(\sigma_1, \ \sigma_2, \ \sigma_3)$. Equation (2.2.2) is therefore called the normality condition.

As the plastic work $\sigma_1 \cdot \epsilon_1$ is assumed to be non-negative, it is seen that

Thus the vector (e₁, e₂, e₃) becomes an outward-directed normal to the yield surface. If the yield surface is a differentiable surface, the normality criterion uniquely determines the direction of the strain vector for a given stress field on the yield surface. If the yield surface consists of piece wise differentiable surfaces, the strain vector for stress fields lying on the curve of intersection of two surfaces must be located in the angle between the normals to the adjacent surfaces.

The load carrying capacity of a body consisting of a rigid-plastic material is the load at which plastic deformations become possible. The load carrying capacity is also called the yield load or the failure load. For determination of the yield load of a rigid-plastic body the following extremum principles are very useful.

First we define some important concepts which will be used in formulating the extremum principles.

A safe and statically admissible stress field is a stress field corresponding to stresses within or on the yield surface and satisfying the equilibrium conditions and the statical boundary conditions.

A geometrically admissible failure mechanism is a deformation field that satisfies the compatibility conditions and the geometrical boundary conditions.

The extremum principles can be described as follows:

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The Upper-Bound Theorem:

geometrically admissible failure mechanism equal to the yield load of the body. load found from the work equation is greater than or for an arbitrary,

The Lower-Bound Theorem:

Any load corresponding to a safe and statically admissible stress field is smaller than or equal to the yield load of the body.

The Uniqueness Theorem:

lower bound being the yield load of the body. then an exact solution has been found, the coinciding upper and If the lowest upper bound and the highest lower bound coincide,

2.3 Internal work for plain concrete

separating two rigid parts. One part is moving relatively to the material. A yield line in concrete is a kinematical discontinuity Here, concrete is identified as a rigid-plastic modified Coulomb yield line, (Fig. 2.9). other with the displacement rate ν inclined at the angle α to the

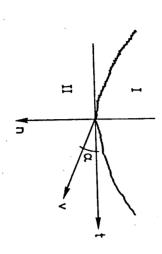


Figure 2.9. Yield line in concrete.

deformation zone, (Fig. 2:10). The discontinuity is a mathematical idealization of a narrow

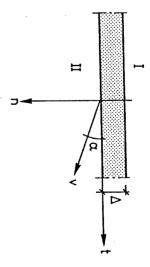


Figure 2.10. Mathematical idealization of a narrow deforming zone.

2.11). concrete having the width \mathtt{A} , the length one and breadth \mathtt{b} , (Fig. Now let us consider a narrow deforming zone of yield line in plain

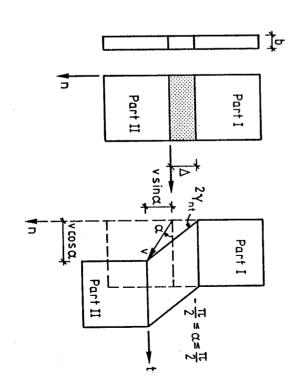


Figure 2.11. Deforming zone of yield line in plain concrete.

The strains in this narrow deforming zone are

$$\epsilon_{\rm n} = \frac{1}{\Delta}$$

$$\epsilon_{\rm t} = 0$$

$$2 \gamma_{\rm nt} = \frac{v \cos \alpha}{\Delta}$$

The principal strains are found from (2.3.1)

$$\epsilon_1 = \frac{V}{2\Lambda}(\sin\alpha + 1)$$
 (2.3.2)

 $\epsilon_2 = \frac{\mathrm{v}}{2\mathrm{h}}(\sin\alpha - 1)$

From equations (2.3.2) and (2.3.3) we find that

$$\epsilon_1 - \epsilon_2 = \frac{V}{\Lambda}$$
 (2.3.4)
 $\frac{\epsilon_1}{\epsilon_2} = -\frac{1 + \sin\alpha}{1 - \sin\alpha}$ (2.3.5)

corresponding to the strains $(\epsilon_1,\ \epsilon_2,\ \epsilon_3)$ is The internal work (dissipation) for a material deformed

$$A = \int_{V} (\sigma_1 \epsilon_1 + \sigma_2 \epsilon_2 + \sigma_3 \epsilon_3) dv = \int_{V} W dv$$
 (2.3.6)

W is the internal work per unit volume being

$$W = \sum_{i=1}^{n} \sigma_{i} \epsilon_{i}$$
 (i = 1, 2, 3)

(2.3.7)

Using Mohr's strain circle, we find

$$tan 2\beta = cot\alpha$$

first principal axis, (Fig. 2.12). Here, $oldsymbol{eta}$ is the angle between the displacement rate vector and the

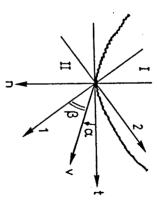


Figure 2.12. The relation between the relative displacement rate vector and the principal axes.

solving equation (2.3.8) yields

(2.3.3)

$$\beta = \frac{1}{2}(\frac{\pi}{2} - \alpha) \tag{2}$$

relative displacement rate vector and the yield line normal. It means the first principal axis bisects the angle between the

2.3.1 The dissipation of concrete in plane strain fields

simplicity we assume ϵ_3 = 0 and σ_1 \geq σ_3 \geq σ_2 . In this case the yield criterion is composed of the following two criteria strains in directions perpendicular to a plane are zero, A plane strain field is defined as a strain field in which the

$$k\sigma_1 - \sigma_2 - f_c = 0$$
 (2.3.10)

$$f_{+} = 0$$
 (2.3.11)

equation (2.3.11) corresponds to separation failure. The equation (2.3.10) corresponds to sliding failure while

(2.3.8)

For the sliding failure, the related strains can be found directly from the normality criterion (2.2.2).

$$(\epsilon_1, \epsilon_2, \epsilon_3) = (\lambda k, -\lambda, 0)$$

(2.3.12)

It is easy to find that

$$\epsilon_1 - \epsilon_2 = \lambda(k+1)$$
 (2.3.13)

$$\frac{\epsilon_1}{\epsilon_2} = -k = -\frac{1 + \sin\varphi}{1 - \sin\varphi}$$
 (2.3.14)

with (2.3.5), we have Comparing equation (2.3.13) with (2.3.4) and equation (2.3.14)

$$\lambda = \frac{V}{\Lambda(k+1)} = \frac{V}{\Lambda} \frac{1 - \sin \varphi}{2}$$
 (2.3.15)

$$\sin\alpha = \sin\varphi \tag{2.3.16}$$

Equation (2.3.16) yields

$$\alpha = \left\{ \begin{array}{l} \varphi \\ \pi - \varphi \end{array} \right. \tag{2.3.17}$$

Inserting equations (2.3.12) and (2.3.15) into (2.3.7), we get

$$W = \lambda (\sigma_1 k - \sigma_2) = \lambda f_C = \frac{V}{\Lambda} \frac{1 - \sin \varphi}{2} f_C$$
 (2.3.18)

For the separation failure, the related strains will be

$$(\epsilon_1, \epsilon_2, \epsilon_3) = (\lambda_1, 0, 0)$$

(2.3.19)

comparing equation (2.3.19) with (2.3.2) and (2.3.3), we find that

R H

NI4

$$\lambda_1 = \frac{V}{h} \tag{2.3.21}$$

Then the internal work is easily found by inserting equations (2.3.19) and (2.3.21) into (2.3.7)

$$W = \lambda_1 f_t = \frac{V}{\Delta} f_t \tag{2.3}$$

strains then can be found to be of the strain vectors belonging to the two surfaces. The related the strain vector will be an arbitrary positive linear combination equations (2.3.10) and (2.3.11) being satisfied simultaneously, At the intersection of the two yield surfaces, corresponding to

$$(\epsilon_1, \epsilon_2, \epsilon_3) = (\lambda k + \lambda_1, -\lambda, 0)$$
 (2.3.23)

Comparing equation (2.3.23) with (2.3.3) and (2.3.2), we find that

$$\lambda = \frac{V}{2\Delta} (1 - \sin \alpha) \tag{2.3.24}$$

$$\lambda_1 = \frac{V}{2\Lambda} (1 + \sin\alpha) - \lambda k = \frac{V}{\Lambda} \frac{\sin\alpha - \sin\phi}{1 - \sin\phi}$$

In this case, the ranges of α are

(2.3.24) and (2.3.25) into (2.3.7) The internal work will be found by inserting equations (2.3.23),

$$W = \frac{vf_c}{2\Lambda} \left[(1 - \sin\alpha) + 2 \frac{f_t}{f_c} \frac{\sin\alpha - \sin\varphi}{1 - \sin\varphi} \right]$$
 (2.3.27)

It should be noticed that when $\alpha=\frac{\pi}{2}$, equation (2.3.27) is identical to equation (2.3.22), and when $\alpha=\phi$ or $\pi-\phi$, equation (2.3.27) changes into equation (2.3.18).

The formula (2.3.27) thus gives the internal work per unit volume of concrete in plane strain field.

If σ_1 , σ_2 and σ_3 , and ϵ_1 , ϵ_2 and ϵ_3 can mean any of the principal stresses and strains, analogous derivations can be made and the solutions are the same.

The dissipation per unit length \mathbf{w}_1 is

$$W_{1} = Wb\Delta \tag{2.3.28}$$

where b is the dimension of the body in direction perpendicular to the plane. Inserting (2.3.27) into (2.3.28), we get the complete formula for the dissipation per unit length of a discontinuity line in concrete in plane strain:

$$W_1 = \frac{1}{2}b \ f_C v[1 - \sin\alpha + 2\frac{f_C}{f_C} \frac{\sin\alpha - \sin\phi}{1 - \sin\phi}] \text{ valid for } \phi \le \alpha \le \pi - \phi$$
(2.3.29)

Introducing the parameters

$$1 = 1 - 2 \frac{f_t}{f_t} \frac{\sin \varphi}{1 - \sin \varphi}$$
 (2.3.30)

$$m = 1 - 2 \frac{f_t}{f_c} \frac{1}{1 - \sin \varphi}$$

(2.3.31)

The brief version

$$W_1 = \frac{1}{2}b f_C V(1 - m \sin \alpha) \qquad \text{for } \phi \le \alpha \le \pi - \phi$$

(2.3.32)

The yield condition and the associated flow rule for concrete in

plane strain are shown in Fig. 2.13.

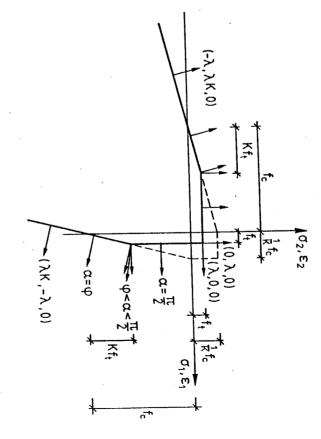


Figure 2.13. Yield condition and associated flow rule for concrete in plane strain.

2.3.2 The dissipation of concrete in plane stress fields

A plane stress field is defined as a stress field in which the stresses in sections parallel to a plane are zero; this plane is a principal section, and the normal to the plane is a principal direction with the corresponding principal stress equal to zero.

Denoting the principal stress in sections perpendicular to the plane as $\sigma_{
m I}$ and $\sigma_{
m II}\prime$ the yield condition can be expressed as follows:

For $\sigma_{\rm I} > \sigma_{\rm II} >$ 0, we get $\sigma_{\rm 1} = \sigma_{\rm I}$ and $\sigma_{\rm 3} =$ 0, which gives

as the condition for separation failure

for sliding failure is For $\sigma_{\rm I}$ > 0 > $\sigma_{\rm II}$, we have $\sigma_{\rm I}$ = $\sigma_{\rm I}$ and $\sigma_{\rm 3}$ = $\sigma_{\rm II}$, and the condition

$$\mathbf{r} = \sigma_{\mathsf{TT}} = \mathbf{f}_{\mathsf{C}} \tag{2.3.34}$$

(2.3.33). while the condition for separation failure is the same as

For 0 > $\sigma_{\rm I}$ > $\sigma_{\rm II}$, we have $\sigma_{\rm 1}$ = 0 and $\sigma_{\rm 3}$ = $\sigma_{\rm II}$, and only sliding failure is possible, the condition being

$$\sigma_{\text{II}} = f_{\text{C}} \tag{2.3.35}$$

conditions and the corresponding strains in plane stress. With the same procedure as in plane strain, we can find the yield

been drawn in a $(\sigma_1, \ \sigma_2)$ -coordinate system. In Fig. 2.14 the yield conditions and associated flow rule have

lines 3-4 in Fig. 2.14. The areas not covered by the formulae for plane strain are the

Along 3-4, the yield condition is

$$\sigma_2 = -f_c \tag{2.3.36}$$

The related strains are found to be

$$(\epsilon_1, \ \epsilon_2, \ \epsilon_3) = (0, -\lambda, \lambda k)$$

The strains in the deforming zone in Fig. 2.11 are again given by equations (2.3.2) and (2.3.3).

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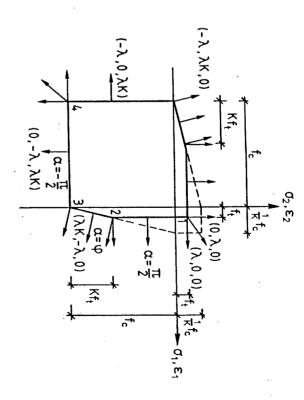


Figure 2.14. Yield condition and associated flow rule for concrete in plane stress.

Comparing (2.3.37) with (2.3.2) and (2.3.3), we get

$$\epsilon_1 = \frac{V}{2\Lambda}(\sin\alpha + 1) = 0$$

$$\epsilon_2 = \frac{V}{2\Lambda}(\sin\alpha - 1) = -\lambda$$

(2.3.39)

(2.3.38)

(2.3.40)

(2.3.41)

It yields

and

(2.3.37)

The internal work then is

At point 3, the strain vector can be an arbitrary positive linear combination of the strain vectors belonging to the two sides.

The related strains are

$$(\epsilon_1, \ \epsilon_2, \ \epsilon_3) = (\lambda k, -(\lambda + \lambda_1), \ \lambda_1 k)$$
 (2.3.43)

Comparing equation (2.3.43) with (2.3.3), we find that

$$\lambda + \lambda_1 = \frac{3}{2h}(1 - \sin\alpha) \tag{2.3.44}$$

The internal work at point 3 then is

$$W = \sigma_2 \epsilon_2 = \frac{vf_c}{2h}(1 - \sin\alpha) \tag{2.3.45}$$

In this case α will be

$$-\frac{\pi}{2} \le \alpha \le \varphi \quad \text{or} \quad \frac{3}{2}\pi \ge \alpha \ge \pi - \varphi \tag{2.3.46}$$

When $\alpha = -\frac{\pi}{2}$ or $\alpha = \frac{3}{2}\pi$, equation (2.3.45) is identical with equation (2.3.42).

The dissipation per unit length along line 3-4 and at point 3 can thus be expressed as

$$W_1 = WbA = \frac{bf_C}{2} v(1 - \sin \alpha)$$
 for $\alpha \le \varphi$ or $\alpha \ge \pi - \varphi$ (2.3.47)

It should be noted that in a plane stress field corresponding to line 3-4 and point 3 in Fig. 2.14, there must be inhomogeneous transitional zones between the homogeneous deformation field and the rigid part I and II in Fig. 2.11, because of the strain 63. The contribution from these transitional zones to the internal work is neglected in formula (2.3.47).

The complete dissipation per unit length of a discontinuity line for concrete in plane stress can be obtained by combining equations (2.3.32) and (2.3.47)

$$\eta_1 = \begin{cases} \frac{bf_C}{2} v(1 - \sin\alpha) & \text{for } \alpha \le \phi \text{ or } \alpha \ge \pi - \phi \\ \frac{bf_C}{2} v(1 - m \sin\alpha) & \text{for } \phi \le \alpha \le \pi - \phi \end{cases}$$
 (2.)

The expression (2.3.48) will be strongly simplified, if $f_{\rm t}$ = 0, as then in the whole $\alpha\text{-interval}$, we have

$$W_{1} = \frac{bf_{c}}{2} v(1 - \sin\alpha)$$
 (2.3.49)

In Fig. 2.15, the so-called square yield locus is depicted for concrete with a zero tension cut-off.

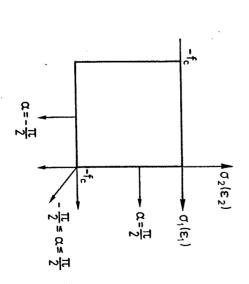


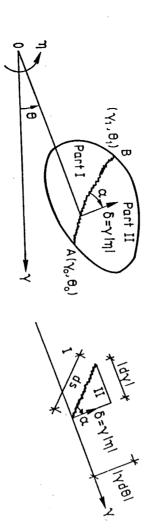
Figure. 2.15. Square yield locus for concrete in plane stress.

The dissipation formulae for yield lines in Coulomb materials and modified Coulomb materials were developed by B.C. Jensen [77.2] and Nielsen et al [78.3].

2.4 Dissipation of curved yield line corresponding to rotation around an arbitrary point in the plane.

In this section, all derivations are based on a zero tension cut-off in plane stress.

Fig. 2.16 shows a concrete disk with breadth b in a polar-coordinate system, r, θ .



a: Concrete disk with curved yield line.

b: Element of yield line.

Figure 2.16. Curved yield line corresponding to rotation around an arbitrary point.

A curved yield line $r=r(\theta)$ separates the disk into two rigid parts I and II. The two end points are $A(r_0, \theta_0)$ and $B(r_1, \theta_1)$ separately. The part II is rotating in the plane relatively to the part I around the pole 0 with rotation rate η .

By using (2.3.49), the dissipation of yield line element with length ds has been found to be

$$W_{I}ds = \frac{\eta}{2} bf_{c}r(\sqrt{r^{2} + r^{2}} - r^{2})d\theta$$
 (2.4.)

Here, $r' = \frac{dr}{d\theta}$ has been introduced.

The total dissipation of the curved yield line AB can be found by integrating along the curve

$$W_{I} = \frac{\eta}{2} bf_{C} \int_{\theta_{O}}^{\theta_{I}} F(x, x') d\theta$$
 (2.4.2)

where

$$F(r, r') = r(\sqrt{r'^2 + r^2} - r')$$
 (2.4.

The Euler equation, i.e. the necessary condition for getting an extreme value of W_{1} corresponding to the function $r(\theta)$, is

$$\frac{\partial \mathbf{F}}{\partial \mathbf{r}} - \frac{\mathbf{d}}{\mathbf{d}\theta} \left(\frac{\partial \mathbf{F}}{\partial \mathbf{r}} \right) = 0 \tag{2.4.4}$$

Inserting equation (2.4.3) into (2.4.4) yields

$$\mathbf{r} \cdot \mathbf{r}^{\parallel} - 3(\mathbf{r}^{\perp})^2 - 2\mathbf{r}^2 = 0$$
 (2.4.5)

The solution of this differential equation (2.4.5) is

$$r = \pm \sqrt{\frac{c_1}{\sin[2(\theta - c_2)]}}$$
, valid for $\frac{c_1}{\sin[2(\theta - c_2)]} > 0$ (2.4.6)

where ${\rm C_1}$ and ${\rm C_2}$ are arbitrary constants which must satisfy the edge conditions. Inserting equation (2.4.6) into (2.4.2), we get

$$W_{I} = \frac{\eta}{4} bf_{c} \left\{ \left[-c_{1} \frac{\cot^{2}(\theta - c_{2})^{-1}}{2 \cot(\theta - c_{2})} \right]_{\theta}^{\theta_{1}} - \left[r^{2}\right]_{r_{0}}^{r_{1}} \right\}$$
(2.4.7)

Transforming the polar coordinate system into Cartesian coordinates by

$$x = r \cdot \cos(\theta - C_2)$$

 $y = r \cdot \sin(\theta - C_2)$ (2.4.8)

We find a simple form of the curve determined by (2.4.6) (see Fig. 2.17)

$$x \cdot y = \frac{1}{2} C_1 \tag{2.4.9}$$

In the Cartesian coordinates, equation (2.4.7) can be simplified into

$$W_{I} = \frac{1}{2} |\eta| bf_{c} r_{m} d(1 + \frac{|\eta|}{\eta} sin(\beta - \gamma))$$
 (2.4.10)

Here, r_m , d, β and τ are 4 parameters in Cartesian coordinates (see Fig. 2.18).

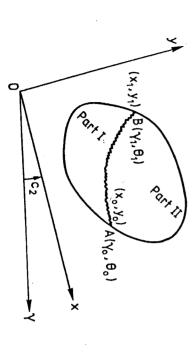


Figure 2.17. Transforming of coordinate systems.

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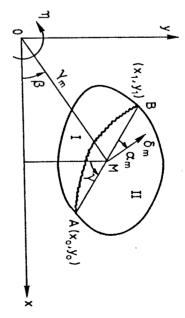


Figure 2.18. Introduction of 4 parameters.

d being the length of straight line AB, r_m the length of line OM,, where M is the centre of line AB, β the angle between line OM and x-axis, and τ is the angle between the line AB and the y-direction.

Introducing $\delta_{\rm m} = r_{\rm m} |\eta|$ and

$$\alpha_{\rm m} = \gamma - \beta$$
 for $\eta > 0$
 $\alpha_{\rm m} = -(\gamma - \beta)$ for $\eta < 0$

(2.4.11)

we get the internal work of the whole curved yield line

$$W_{I} = \frac{1}{2} \delta_{m} bdf_{c} (1 - \sin \alpha_{m})$$
 (2.4.12)

Here δ_m is the relative displacement of point M and α_m is an angle between line AB and the direction of δ_m .

It should be noted that equation (2.4.12) is only valid when

$$r_{\rm m} > \frac{1}{2} d$$
 (2.4.13)

The dissipation of curved yield lines corresponding to rotation

around an arbitrary point in the plane was developed by J.F. Jensen [81.3].

2.5 The plastic strengths of concrete

as a rigid, perfectly plastic material. But in fact, concrete is reinforced concrete structures, it is necessary to define concrete stress-strain relations for concrete are characterized by strain not a perfectly plastic material due to the fact that the As mentioned before, in order to apply the theory of plasticity to reached, the stresses decrease with increasing strains. softning; that is, as soon as the maximum value of the stresses is

has the form shown in Fig. 2.19 [87.2]. A typical stress-strain curve in uniaxial tension and compression

q (MPa)

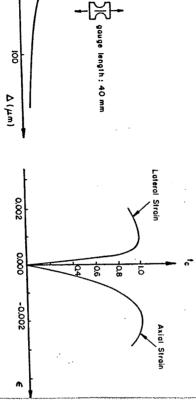


Figure 2.19. Uniaxial tensile curve (Peterson, 1981). stress-elongation

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stress-strain curve Uniaxial compressive (Kupfer, 1969).

In reality, the behaviour of concrete in tension is almost brittle and the ductility of concrete in compression is very limited.

Because of the limited deformability of the concrete and the

points of take place at the expense of losing some strength. redistribution of stresses, as assumed in plasticity, can only concrete stress to equal the maximum compressive strength at all unstable nature of the concrete failure we cannot expect the the yield lines at failure. It means

smaller than the strength measured by standard tests inserting in the theoretical formulae a concrete strength which is proved useful is to modify the exact plastic solutions by prediction of the observed shear capacity, a method that has been In order that the theory may give a reasonable quantitative

all other shortcomings of the theory as well. the compressive strength, and $\nu_{\mbox{\scriptsize t}}$, for the tensile strength, as two empirical measures of concrete ductility and as two absorbents of This suggests the introduction of the effectiveness factor $\boldsymbol{\mathfrak{v}}$, for

corresponding tensile strength of concrete f_{t}^{*} are corresponding plastic compressive strength of concrete $\mathbf{f}_{\mathbf{C}}^{\star}$ and the strength of concrete measured by some standard procedure, the a standard compression test and if f_{t} is the uniaxial tensile If $f_{\mathbf{C}}$ is the uniaxial compressive strength of concrete measured by

$$f_{c}^{*} = v f_{c}$$
 (2.5.2)
 $f_{t}^{*} = v_{t} f_{t}^{*}$

In the remaining sections, the notations f_{C}^{\star} and f_{L}^{\star} are always used in formulae. Since it is difficult to determine the tensile compressive strength $\mathbf{f}_{\mathbf{C}}^{m{x}}$ by introducing the factor $\mathbf{
ho}^{m{x}}$ defined by strength of concrete by an ordinary tensile test, the plastic tensile strength of concrete f_{t}^* will be compared to the plastic

values of v and ho^* normally must be determined by experiments. The quantities u and ho^* are called effectiveness factors.

2.6 Basic assumptions

Reinforced concrete beams (including deep beams and corbels) are very important structural members in concrete structures.

In the following chapters the plastic theory of reinforced concrete beams is treated based on the assumptions:

- The beam is in a state of plane stress.
- 2) Anchorage failure and bearing or supporting failure are prevented by special anchorage provisions, such as steel plates welded to the bar ends, and adequate design of the bearings and load platens so that the proper transfer of reinforcement tension to concrete compression will be possible.
- 3) The reinforcement is rigid, perfectly plastic with a stress-strain relation for tension and compression as shown in Fig. 2.20. The yield strength of reinforcement is denoted $\mathbf{f_y}$. For steels without a definite yield point, the $\mathbf{f_y}$ is defined in a suitable manner (e.g., as the 0.2% offset strength).

Furthermore, we assume that the reinforcing bars are only capable of carrying longitudinal tensile and compressive stresses. According to this assumption, the reinforcement bars are unable to resist any lateral forces, i.e., the dowel effects are neglected.

The assumption that the materials are rigid and perfect plastic means that any elastic deformations and work hardening effects are neglected and unlimited ductility is assumed.

4) The concrete is rigid, perfectly plastic with the modified Coulomb's failure criterion as yield condition and the associated flow rule (normality condition) as constitutive equations (see

The tensile and compressive strength of concrete are f_t^* and $f_{c'}^*$, respectively, where f_t^* and f_c^* are defined by (2.5.3) and (2.5.1).

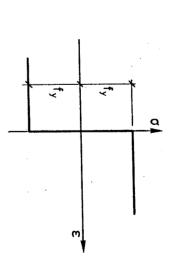


Figure 2.20. Uniaxial stress-strain relation for reinforcement.

5) The weight of the beam is disregarded.

CHAPTER III. BENDING CARRYING CAPACITY OF BEAMS, DEEP BEAMS

In this chapter, the plastic solution of pure bending is described only briefly, since the determination of the bending capacity of reinforced concrete beams is an old problem, which may be treated by many methods described in the literature and in design codes. by many methods described in the plastic theory in bending is to The purpose of reviewing the plastic theory in bending is to provide available tools to exclude the possible flexural failure when the shear problems are dealt with. The following solutions of bending were presented by M.P. Nielsen et al. [84.1], [84.6] and [78.3], except for solution of deep beams with horizontal web reinforcement, which is derived by the author at section 3.3.

3,1 Members with only tensile reinforcement

3.1.1 Beams and deep beams with rectangular section

The beams treated in this section are assumed to have a rectangular cross section b·h and to be loaded in pure bending. The tensile reinforcement with the yield strength \mathbf{f}_y and area \mathbf{A}_s is assumed to be concentrated in a stringer at a distance d from the top of the cross section as shown in Fig. 3.1.

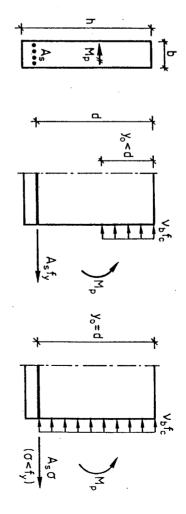
The concrete is assumed to be perfectly plastic with the compressive strength $f_c^*=\nu_b f_c$. Here ν_b is an effectiveness factor for bending and f_c is the compressive cylinder strength of

The complete solutions for the bending failure moment are given by

$$m_{p}^{*} = \begin{cases} (1 - \frac{1}{2} \phi_{b}^{*}) \phi_{b}^{*} & \text{valid for } \phi_{b}^{*} \le 1 \\ \frac{1}{2} & \text{valid for } \phi_{b}^{*} \ge 1 \end{cases}$$
 (3.1.1)

Here, the effective degree of reinforcement, defined by $\phi_b^*=A_sf_y/bdf_c^*$, and the dimensionless effective yield moment $m_{p'}$,

defined by $m_p^* = M_p/bd^2 f_c^*$, have been introduced. For details concerning the derivation of (3.1.1) see, for example, [84.1].



a) Normally reinforced section

b) Over reinforced section

Figure 3.1. Normal stress distribution at the yield moment.

3.1.2 Corbels with rectangular section

Fig. 3.2 shows a typical corbel with the rectangular cross section $\mathfrak{b} \cdot \mathfrak{h}$.

The complete solutions of bending capacity for such corbels are given by

$$\frac{P}{bdf_{c}^{*}} = \begin{cases} \sqrt{(\frac{a}{d})^{2} + \phi_{b}^{*}(2 - \phi_{b}^{*})} - \frac{a}{d} & \text{valid for } \phi_{b}^{*} \le 1 \\ \sqrt{(\frac{a}{d})^{2} + 1} - \frac{a}{d} & \text{valid for } \phi_{b}^{*} \ge 1 \end{cases}$$
(3.1.2)

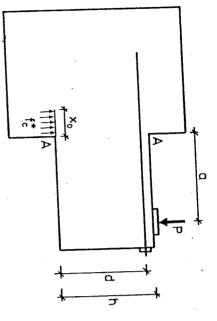


Figure 3.2. Bending capacity of rectangular corbel.

Equation (3.1.2) can easily be found by inserting equations

$$M_{p} = P \cdot (a + \frac{x_{0}}{2^{0}})$$
 (3.1.3)

$$x_{o} = \frac{P}{bf_{c}^{*}}$$
 (3.1.4)

into (3.1.1).

Solution (3.1.2) may be identified with the exact solution found by B.C. Jensen [79.7].

3.2 Beams with compression reinforcement

The compression reinforcement with the same tensile and compressive yield strength f_{yc} and area f_{sc} is assumed concentrated in a stringer at a distance f_{c} from the top of the cross section. All other data and assumptions are in agreement with those mentioned in section 3.1.1.

The dimensionless effective yield moment is found to be

$$\mathbf{m}_{\mathbf{p}}^{*} = \begin{cases} \phi_{\mathbf{b}}^{*} + \alpha \phi_{\mathbf{c}}^{*} - \frac{1}{2} (\phi_{\mathbf{b}}^{*} + \phi_{\mathbf{c}}^{*})^{2} & \text{valid for } (\phi_{\mathbf{b}}^{*} + \phi_{\mathbf{c}}^{*}) < \alpha \\ \frac{1}{2} \alpha^{2} + (1 - \alpha) \phi_{\mathbf{b}}^{*} & \text{valid for } (\phi_{\mathbf{b}}^{*} + \phi_{\mathbf{c}}^{*}) \geq \alpha \geq (\phi_{\mathbf{b}}^{*} - \phi_{\mathbf{c}}^{*}) \\ \phi_{\mathbf{b}}^{*} - \alpha \phi_{\mathbf{c}}^{*} - \frac{1}{2} (\phi_{\mathbf{b}}^{*} - \phi_{\mathbf{c}}^{*})^{2} & \text{valid for } (\phi_{\mathbf{b}}^{*} - \phi_{\mathbf{c}}^{*}) > \alpha \geq \alpha (\phi_{\mathbf{b}}^{*} - \phi_{\mathbf{c}}^{*}) \\ \frac{1}{2} + (1 - \alpha) \phi_{\mathbf{c}}^{*} & \text{valid for } (\phi_{\mathbf{b}}^{*} - \phi_{\mathbf{c}}^{*}) > 1 \end{cases}$$

$$(3.2.1)$$

Here, the following parameters have been introduced

$$\alpha = \frac{a_{\rm C}}{d} \tag{3.2..2}$$

$$\phi_{\mathbf{C}}^{*} = \frac{^{*}\mathbf{S}\mathbf{C}}{^{*}\mathbf{C}}$$
 (3.2.3)

For a more detailed presentation, see [84.1].

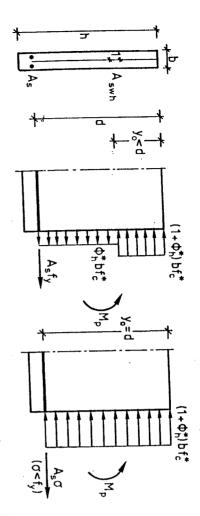
3.3 Deep beams with horizontal web reinforcement

Because the vertical web reinforcement has no contribution to the bending capacity of the beam, we will only discuss the influence of horizontal web reinforcement.

The horizontal web reinforcement with the same tensile and compressive yield strength f_{ywh} is assumed to be uniformly distributed in the cross section. The area of the horizontal web reinforcement per unit of height is denoted as h_{swh} . All other data and assumptions are the same as those mentioned in section f(x).

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When a flexural failure occurs, the stress distribution in the cross section will be as shown in Fig. 3.3.



a) Normally reinforced section. b) Ove

n. b) Over reinforced section

Figure 3.3. Normal stress distribution at the failure moment.

We easily obtain the dimensionless effective yield moment $m_{\widetilde{p}}^{\star}$ as

$$m_{p}^{*} = \begin{cases} \phi_{b}^{*} + \frac{1}{2}\phi_{hw}^{*} - \frac{1}{2}(\phi_{b}^{*} + \phi_{hw}^{*})^{2}/(1 + 2\phi_{hw}^{*}) & \text{for } \phi_{b}^{*} - \phi_{hw}^{*} \leq 1 \\ \\ \frac{1}{2}(1 + \phi_{hw}^{*}) & \text{for } \phi_{b}^{*} - \phi_{hw}^{*} \geq 1 \end{cases}$$

$$(3.3.1)$$

Here the effective degree of horizontal web reinforcement ϕ_{hw}^* defined as $\phi_{hw}^* = A_{swh} \cdot f_{ywh}/bf_c^*$ has been introduced.

The influence of the horizontal web reinforcement on the ultimate flexural load is shown graphically in Fig. 3.4, where $\alpha_1=\phi_{\text{hw}}^*/\phi_{\text{b}}$ has been introduced. The three curves in Fig. 3.4 are drawn for $\alpha_1=0$, which corresponds to the case without horizontal web reinforcement, formula (3.1.1), $\alpha_1=0.5$ and $\alpha_1=1.0$, respectively.

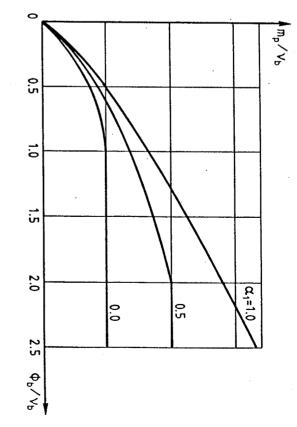


Figure. 3.4. Bending capacity of beams with horizontal web reinforcement.

3.4 The effectiveness factor for bending

For pure bending of a rectangular section with only tensile reinforcement, the effectiveness factor $v_{\rm b}$ has been analytically determined by Exner [79.5] using stress-strain curves measured by p.T. Wang et al. It turns out that $v_{\rm b}$ is a function of the uniaxial compressive strength $f_{\rm c}$, the yield stress of the reinforcement $f_{\rm y}$ and the reinforcement ratio $\phi_{\rm b}$, defined by $\phi_{\rm b} = \lambda_{\rm c}/{\rm bd}$.

For practical purpose, $\nu_{\mathbf{b}}$ can be calculated approximately by a simple empirical formula

$$v_{\rm b} = 0.97 - \frac{f_{\rm Y}}{5000} - \frac{f_{\rm C}}{300}$$
 for $\begin{cases} f_{\rm Y} < 900 \text{ MPa} \\ f_{\rm C} < 60 \text{ MPa} \end{cases}$ (3.4.1)

For most practical cases f_y will be less than 600 MPa, and conservatively we get

$$\mathbf{for} \left\{ \begin{array}{l} \mathbf{f_Y} < 600 \ \mathrm{MPa} \\ \mathbf{f_C} < 60 \ \mathrm{MPa} \end{array} \right. \tag{3.4.2}$$

 $v_b = 0.85 - \frac{f_c}{300}$

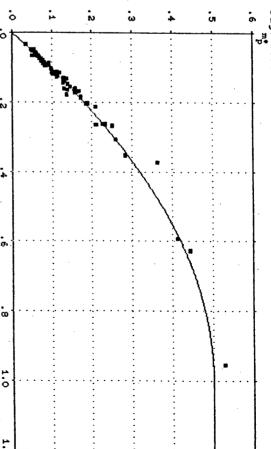
3.5 Experimental verification

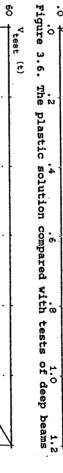
Comparison between almost 300 conventional beam tests from many different test series and formula (3.1.1), using the ν_b -value of (3.4.1), can be found in [84.1] and [84.6]. Here, we will only treat the deep beams with and without horizontal web reinforcement failing in flexure. The dimensions and material properties and measured ultimate loads of the tested deep beams are listed in Appendix A. The statistical values for the ratios of test to Appendix A. The statistical values for the ratios of test to appendix A. The statistical value of (3.4.2) are shown in Fig.

				- Tob reinforcement
0.072	29 0.998 0.072	0.998	29	deep beams with horizontal
			(-)	web reinforcement
0.069	0.069	0.999	74	deep beams without horizontal
				cases
C	a	×ι	3	Statistical varues
variation	deviation			/
coeff. of	standard	mean	number	1+ pm

Figure 3.5. The statistical values of the ratios of test to theory for deep beams with and without horizontal web reinforcement failed in flexure.

The good agreement can also be seen from Fig. 3.6, where the curve drawn is the solution (3.1.1) with the $v_{\rm b}$ -value from (3.4.2), and Fig. 3.7.





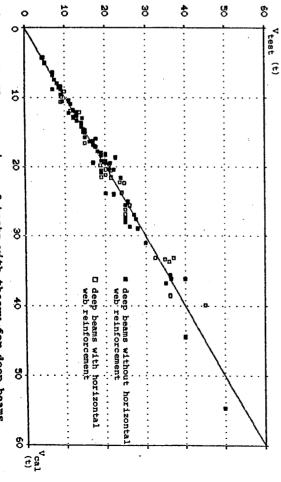


Figure 3.7. The comparison of tests with theory for deep beams.

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CHAPTER IV. SHEAR CARRYING CAPACITY OF BEAMS, DEEP BEAMS AND CORBELS WITHOUT WEB REINFORCEMENT

In this chapter we will only treat beams subjected to concentrated loading. Consider a horizontal, simply supported rectangular concrete beam loaded by two symmetrical loads P, (Fig. 4.1).

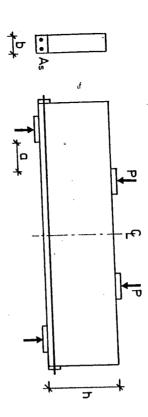


Figure 4.1. Rectangular reinforced concrete beam subjected to concentrated loads.

The beam is reinforced with the area $A_{\rm g}$ along the bottom face of the beam. The breadth and depth of the beam are termed b and h, respectively. The quantity a denotes the clear shear span, which is defined as the distance between the inside edge of support platen and the outside edge of load platen.

4.1 Shear capacity considering the tensile strength of concrete

In this section we will derive coinciding upper and lower bound solutions of non-webreinforced beams taking account of the tensile strength of concrete.

The coincident upper and lower bounds in the over reinforced case was found by J.F. Jensen in [81.3]. The complete solutions with their boundaries were derived by the author. These solutions can

be used not only for conventional beams but also for deep beams, corbels and joints. The solutions are compared with a great number of test series on non-webreinforced slender beams, deep beams and corbels subjected to concentrated loads on the top compression face.

A really good agreement has been found.

4.1.1 Lower bound solution

Let us consider the beam element A-B between the loading and supporting platen in a stringer beam.

The stress distribution is as shown in Fig. 4.2.

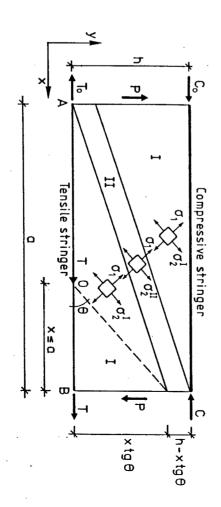


Figure 4.2. Stress distribution along the clear shear span.

We assume that if the tensile stringer is not strong enough, it will yield at a point 0. The distance between point 0 and the section B is denoted x. The tensile force at point 0 is T.

We must have

(4.1.0)

When point 0 reaches point A, i.e. x=a, the stress distribution in the shear span may change into the case shown in Fig. 4.3.

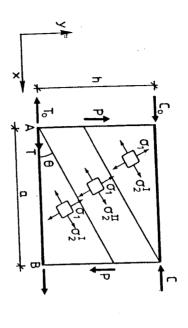


Figure 4.3. Stress distribution along the clear shear span of a stringer beam with small shear span.

In the areas marked I, the stresses are

$$\sigma_{\mathbf{X}}^{\mathbf{I}} = \sigma_{2}^{\mathbf{I}} \cos^{2}\theta + \sigma_{1} \sin^{2}\theta \tag{4.1.1}$$

$$\sigma_Y^{\rm I} = \sigma_1 \cos^2 \theta + \sigma_2^{\rm I} \sin^2 \theta \tag{4.1.2}$$

$$|\tau_{XY}^{I}| = |\sigma_1 - \sigma_2^{I}| \sin\theta \cos\theta$$
 (4.1.3)

As mentioned in our basic assumptions, the stringers can not resist any lateral forces.

This leads to

$$a^{\mathrm{T}} = 0 \tag{4.1.4}$$

Inserting (4.1.4) into equation (4.1.2) yields

$$\sigma_2^{\rm I} = -\sigma_1 \cot^2 \theta \tag{4.1}$$

From equations (4.1.1), (4.1.3) and (4.1.5) we easily get

$$\sigma_{\mathbf{X}}^{\mathbf{I}} = \sigma_{\mathbf{1}}(1 - \cot^2 \theta)$$
 (4.1.6)

$$|\tau_{xy}^{I}| = \sigma_{1} \cot \theta$$
 (4.1.7)

In area marked II, the stresses are

$$\sigma_2^{\text{II}} = k\sigma_1 - f_c^*$$
 (4.1.8)

$$\sigma_{\rm X}^{\rm II} = \sigma_{\rm 2}^{\rm II} \cos^2 \theta + \sigma_{\rm 1} \sin^2 \theta = \sigma_{\rm 1} + (k\sigma_{\rm 1} - f_{\rm c}^* - \sigma_{\rm 1}) \cos^2 \theta$$
 (4.1.9)

$$|\tau_{xy}^{II}| = |\sigma_1 - \sigma_2^{II}| \sin\theta \cos\theta = (\sigma_1 - k\sigma_1 + f_C^*) \sin\theta \cos\theta$$
 (4.1.10)

The first principal stress σ_1 in Fig. 4.2 and Fig. 4.3 is depending on the angle θ .

As mentioned in section 2.3, the first principal axis bisects the angle between the relative displacement rate vector and the yield line normal. It is demonstrated in Fig. 4.4.

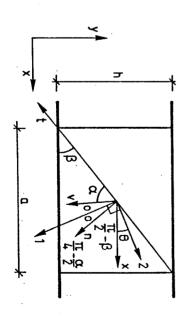


Figure 4.4. The relations between β , α and θ .

The relations between eta, lpha and heta can be easily found to be

$$\theta = \beta + \frac{\alpha}{2} - \frac{\pi}{4} \tag{4.1.11}$$

$$\cot \beta = \frac{a}{h} \tag{4.1.12}$$

The magnitude of the first principal stress σ_1 is

$$\sigma_1 = f_t^* = \rho^* f_c^*$$
 when $\alpha > \varphi$ (4.1.13)
 $0 \le \sigma_1 \le f_t^*$ when $\alpha = \varphi$ (4.1.14)

when
$$\alpha < \varphi$$
 (4.1.15)

The relation between σ_1 and θ may then be found to be

$$\sigma_1 = \rho^* f_C^*$$
 when $\theta > \beta + \frac{g}{2} - \frac{\pi}{4}$ (4.1.16)
 $0 \le \sigma_1 \le f_C^*$ when $\theta = \beta + \frac{g}{2} - \frac{\pi}{4}$ (4.1.17)

when
$$\theta < \beta + \frac{\mu}{2} - \frac{\pi}{4}$$
 (4.1.18)

At first, let us consider the case of x = a.

Using equilibrium conditions to section B in Fig. 4.3, we find

$$P = |\tau_{XY}^{I}| \text{ba tg}\theta + |\tau_{XY}^{II}| \text{b(h - a tg}\theta) \text{ valid for } x \ge \text{a (4.1.19)}$$

$$T + \text{ba tg}\theta \sigma_{X}^{I} + \text{b(h - a tg}\theta)\sigma_{X}^{II} + \text{ba}|\tau_{XY}^{I}| = 0 \qquad (4.1.20)$$

Inserting equations (4.1.6), (4.1.7), (4.1.9) and (4.1.10) into equation (4.1.19) and (4.1.20) yields

$$cot\theta = \frac{\left(\frac{a}{h}\right)^2 + 1 - \frac{D_0^2}{h} + \frac{a}{h}}{1 + D_0}$$
 (4.1.21)

$$\frac{\tau_{c}}{f_{c}^{*}} = \frac{\lambda_{o} \left(\frac{\dot{a}}{\dot{n}}\right)^{2} + 1 - D_{o}^{2} - \mu_{o} \frac{\dot{a}}{\dot{n}}}{2}$$
(4.1.22)

Here, $au = \frac{P}{bh}$ has been introduced as the average shear stress in the section in question.

The parameters $\lambda_{\rm O},~\mu_{\rm O}$ and $\rm D_{\rm O}$ in formulae (4.1.21) and (4.1.22) are as follows

$$\lambda_{O} = 1 - k \frac{\sigma_{1}}{f_{C}^{*}} + \frac{\sigma_{1}^{1}}{f_{C}^{*}}$$
 (4.1.23)

$$\mu_{0} = 1 - k \frac{\sigma_{1}}{f_{c}^{*}} - \frac{\sigma_{1}^{2}}{f_{c}^{*}}$$
 (4.1.24)

$$D_{0} = \frac{\mu_{0} - 2 \frac{T}{bhf_{c}^{*}}}{hf_{c}^{*}}$$
 (4.1.25)

Inserting equation (4.1.21) into equations (4.1.16), (4.1.17) and (4.1.18) yields

$$\sigma_1 = \rho^* f_C^*$$
 valid for $tg\phi - \frac{D}{\cos \phi} < \frac{a}{h}$ (4.1.26)

$$0 \le \sigma_1 \le f_{\mathsf{t}}^*$$
 valid for $\mathsf{tg}_{\theta} - \frac{\mathsf{c}}{\mathsf{cos}_{\theta}} \le \frac{\mathsf{a}}{\mathsf{h}} \le \mathsf{tg}_{\theta} - \frac{\mathsf{D}}{\mathsf{cos}_{\theta}} (4.1.27)$

$$_{1}$$
= 0 valid for $_{h}^{a} < tg \varphi - \frac{C}{\cos \varphi}$ (4.1.28)

Here, the parameters

$$C = 1 - 2 \frac{T}{bhf_c^*}$$
 (4.1.29)

and
$$\mu - 2 \frac{T}{bhf_c}$$

$$D = \frac{1}{bhf_c}$$

(4.1.30)

have been introduced.

The parameters λ and μ are as follows

$$\lambda = 1 - k\rho^* + \rho^* \tag{4.1.31}$$

(4.1.32)

In the case of
$$\frac{a}{h} > tg\phi - \frac{D}{\cos \phi}$$
, we have

μ = 1 - kρ* - ρ*

$$\cot \theta = \frac{\left(\frac{a}{h}\right)^2 + 1 - D^2 + \frac{a}{h}}{1 + D}$$
 (4.1.33)

$$\frac{T}{f_{\rm C}^*} = \frac{\sqrt{\left(\frac{a}{h}\right)^2 + 1 - D^2 - \mu \frac{a}{h}}}{2} \tag{4.1.34}$$

Now we determine the highest lower bound by maximizing equation (4.1.34) with respect to the tensile force ${f T}\cdot$

For T $\leq \frac{\mu}{2}$ bhf, the highest lower bound is obtained with the maximum tensile reinforcement force, i.e.

$$T = \phi^* bhf_c^* \tag{4.1.35}$$

Here, the effective reinforcement degree $\phi^*=\frac{A_S f_V}{b h f_C^*}$ has been introduced.

For $T > \frac{\mu}{2} bhf_C^*$, the highest lower bound is obtained with $T = \frac{\mu}{2} bhf_C^*$, i.e. D = 0.

Thus, the highest lower bound is

$$\frac{T}{f_{\rm c}^*} = \frac{\sqrt{\left(\frac{a}{h}\right)^2 + 1 - D^2 - \mu \frac{a}{h}}}{2}$$
 (4.1.36)

with

$$D = \begin{cases} \frac{\mu - 2\phi^2}{\lambda} & \text{valid for } \phi^* \le \frac{\mu}{2} \\ 0 & \text{valid for } \phi^* \ge \frac{\mu}{2} \end{cases}$$
 (4.1.37)

In the case of $\frac{a}{h} < tg\phi - \frac{c}{\cos \phi}$, we easily find that

$$\cot \theta = \frac{\left(\frac{a}{h}\right)^2 + 1 - c^2 + \frac{a}{h}}{1 + c}$$
 (4.1.38)

$$\frac{\tau_{\star}}{f_{\rm C}} = \frac{1}{2} \left[\sqrt{\left(\frac{a}{h}\right)^2 + 1 - c^2} - \frac{a}{h} \right] \tag{4.1.39}$$

ith

$$0 = \begin{cases} 1 - 2\phi^* & \text{valid for } \phi^* \le \frac{1}{2} \\ 0 & \text{valid for } \phi^* \ge \frac{1}{2} \end{cases}$$
 (4.1.40)

In the case of tgv - $\frac{C}{\cos\varphi}$ $\le \frac{a}{h} \le tgv$ - $\frac{D}{\cos\varphi}$, the highest lower bound may be found by maximizing equation (4.1.22) with respect to

o¦*

Through lengthy and tedious derivations, we get

$$\frac{\sigma_1}{f_C^*} = \frac{(1 - \sin\varphi)[(1 - 2\phi^*) - (\sin\varphi - \frac{a}{h}\cos\varphi)]}{2\cos\varphi[\cos\varphi - \frac{a}{h}\sin\varphi]}$$
(4.1.41)

Inserting equation (4.1.41) into equation (4.1.22) yields

$$\frac{\tau_{\star}}{f_{\rm C}^{\star}} = \frac{\left[\left(\frac{a}{h} \right)^2 + 1 \right] \left(1 - \sin \varphi \right) - 2\phi^{\star} \left(\frac{a}{h} \cos \varphi - \sin \varphi \right)}{2\left[\frac{a}{h} \sin \varphi + \cos \varphi \right]}$$
(4.1.42)

Solutions (4.1.36), (4.1.39) and (4.1.42) are only valid when the condition

$$\sigma_2^{\rm I} > {\rm kf}_{\rm t}^* - {\rm f}_{\rm c}^*$$
 (4.1.43)

is satisfied.

Inserting equations (4.1.5) and (4.1.33) into (4.1.43), we get the condition

$$\frac{a}{h} \le \mu \sqrt{\frac{1 - D^2}{\lambda^2 - \mu^2}} \tag{4.1.44}$$

When the equation (4.1.44) is not satisfied, the location of the yielding point of the tensile reinforcement 0 may depart from A and vary between A and B.

Inserting x=a into equations (4.1.33) and (4.1.34), we find solution in this case

$$\cot \theta = \frac{\left(\frac{X}{h}\right)^2 + 1 - D^2 + \frac{X}{h}}{1 + D} \quad \text{valid for } x < a \quad (4.1.45)$$

$$\frac{\tau_{x}}{\pi} = \frac{\lambda \left(\frac{X}{h}\right)^2 + 1 - D^2 - \mu \frac{X}{h}}{1 + D} \quad \text{valid for } x < a \quad (4.1.46)$$

The optimum value of $\frac{X}{h}$ can be found by maximizing equation (4.1.44) with respect to $\frac{X}{h}$. It's easy to find that

$$\frac{x}{h} = \mu \sqrt{\frac{1 - D^2}{\lambda^2 - \mu^2}}$$
 valid for x < a (4.1.47)

This value coincides with the equation (4.1.44). It means that when the shear span ratio reaches and surpasses some special value, the inclination of the principal direction and the lower bound will be independent of the shear span ratio.

Inserting equation (4.1.47) into (4.1.46) yields

$$\frac{1}{x} = \frac{1}{2} (\lambda^2 - \mu^2) (1 - D^2)$$
 (4.1.48)

It's worth noticing that when ϕ^* $\geq \frac{\mu}{2}$ (D = 0), i.e. in the case of an over reinforced beam, the stress distribution in the shear span may become homogeneous, as shown in Fig. 4.5.

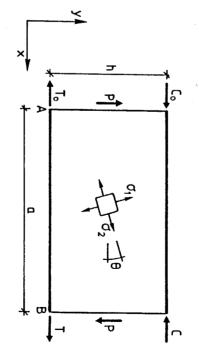


Figure 4.5. Stress distribution in the shear span of an over reinforced beam with large shear span ratio.

Solution (4.1.48) is valid only when the condition

$$\mu \sqrt{\frac{1 - D^2}{\lambda^2 - \mu^2}} > tg\phi - \frac{D}{\cos\phi}$$
 (4.1.49)

is satisfied. It yields

$$\lambda + \mu + D > 1$$
 (4.1.50)

Inserting equations (4.1.31), (4.1.32) and (4.1.37) into (4.1.50), we get

$$\rho^* \le \frac{k - \sqrt{k(1 + k\phi^* - \phi^*)}}{k(k - 1)}$$
 valid for $\phi^* \le \frac{k - 1}{4k}$ (4.1.51)

and

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(4.1.52)

(4.1.48) with respect to ρ *. The highest lower bound may be found by differentiating equation

 λ + μ + D < 1, the highest lower bound will be obtained when given by formulae (4.1.36), (4.1.39), (4.1.42), and (4.1.48). For It is found that for $\lambda + \mu + D \geq 1$, the highest lower bound is

$$k = \frac{k - \sqrt{(1 + k\phi^* - \phi^*)}}{k(k - 1)}$$

for
$$\phi^* \le \frac{k-1}{4k}$$
 (4.1.53)

for
$$\phi^* \ge \frac{k-1}{4k}$$
 (4.1.54)

It's coincident with the equations (4.1.51) and (4.1.52).

In this case the highest lower bound is

$$\frac{r}{f_{C}^{*}} = 2 \frac{\left[k - \sqrt{k[1 + \phi^{*}(k - 1)]} \right] \left[\sqrt{k(1 + \phi^{*}(k - 1))} - 1\right]}{\sqrt{k}(k - 1)^{2}}$$
valid for $\phi^{*} \le \frac{k - 1}{4k}$ (4.1.55)

valid for
$$\phi^* > \frac{k-1}{4k}$$
 (4.1.56)

when equation

$$\frac{a}{h} \geq \frac{\sqrt{1+2\frac{\sin\varphi}{1-\sin\varphi}} \phi^*}{\sin\varphi} - \cot\varphi \qquad (4.1.57)$$

is satisfied.

In the case of
$$\frac{a}{h} < \frac{1+2\frac{\sin\varphi}{1-\sin\varphi}}{\sin\varphi} - \cot\varphi$$
, solutions (4.1.39) and (4.1.42) are valid.

4.1.2 Upper bound solution

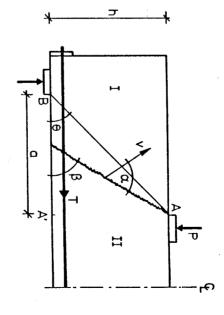


Figure 4.6. Failure mechanism in shear span for beams without web reinforcement.

concentrated load P. Consider a failure mechanism consisting of a Fig. 4.6 shows the shear span of a concrete beam subjected to the admissible, since it does not violate any support condition. angle α to yield line. The relative displacement rate is v, inclined at the and II in Fig. 4.6. Failure is assumed to take place along this This yield line subdivides the beam into two rigid parts marked I platen, to point 0, which may vary along the clear shear span. axis and running from A, which is the outside edge of the load single straight yield line inclined at the angle eta to the beam the yield line. This mechanism is kinematically

We assume that the reinforcement is yielding in tension, i.e.

$$\frac{\pi}{2} \le (\alpha + \beta) \le \pi \tag{4.1.58}$$

The range of the angle eta is

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) \ B \ 2

The lower limit of eta is imposed by the geometry of the beam, where

$$\cot \theta = \frac{a}{h} \tag{4.1.60}$$

The upper limit of eta ensures that the reinforcement is yielding in tension.

The rate of internal work dissipated by this failure mechanism is

$$W_{\rm I} = \frac{V}{2} (\lambda - \mu \sin \alpha) \frac{\text{bhf}_{\star}^{\star}}{\sin \beta} - \text{T V } \cos(\alpha + \beta) \qquad \text{valid for } \phi \le \alpha \le \pi - \phi$$
(4.1.6)

$$W_{\rm I} = \frac{V}{2}(1-\sin\alpha)\frac{{\rm bhf}_{\rm K}^*}{{\rm sin}\beta}$$
 - T $V\cos(\alpha+\beta)$ valid for $\alpha \le \phi$ or $\alpha \ge \pi - \phi$ (4.1.62)

where T is the tensile force of the reinforcement at the intersection with the yield line.

Here, the parameters λ and μ defined as

$$\lambda = 1 - 2 \frac{f_{c}^{\dagger}}{f_{c}^{*}} \frac{\sin \varphi}{1 - \sin \varphi} = 1 - \rho^{*}(K - 1)$$
 (4.1.63)

and

have been introduced.

 $\mu = 1 - 2 \frac{f_{\star}^{*}}{f_{c}^{*}} \frac{1}{1 - \sin \varphi} = 1 - \rho^{*}(k + 1)$

(4.1.64)

It is obvious that when $(\alpha + \beta) > \frac{\pi}{2}$, the reinforcement must be yielding. In this case we have

$$T = A_s f_y \tag{4.1.65}$$

When $\alpha + \beta = \frac{\pi}{2}$, we are in the over reinforced case.

The rate of external work done by the reaction is

$$W_{E} = P \cdot v \cdot \sin(\alpha + \beta) \qquad (4.1.66)$$

The work equation $W_{
m I}$ = $W_{
m E}$ yields the upper bound

$$\frac{\tau}{f_{\rm C}^*} = \frac{\lambda - \mu \sin(\alpha + \beta)\cos\beta + (\mu - 2\phi^*)\sin\beta \cos(\alpha + \beta)}{2\sin\beta \sin(\alpha + \beta)}$$

valid for $(\varphi \le \alpha \le \pi - \varphi)$

(4.1.67)

$$\frac{\tau}{f_c^*} = \frac{1 - \sin(\alpha + \beta)\cos\beta + (1 - 2\phi^*)\sin\beta\cos(\alpha + \beta)}{2\sin\beta\sin(\alpha + \beta)}$$

valid for
$$(\alpha \le \varphi \text{ or } \alpha \ge \pi - \varphi)$$
 (4.1.68)

Here, the effective reinforcement degree ϕ^* has been introduced

$$* = \begin{cases} \frac{\text{Agf}}{\text{bhf}_{c}^{*}} & \text{for } (\alpha + \beta) > \frac{\pi}{2} \\ \frac{\text{T}}{\text{bhf}_{c}^{*}} & \text{for } (\alpha + \beta) = \frac{\pi}{2} \end{cases}$$

$$(4.1.69)$$

It appears that
$$\frac{\partial (\tau/f_c^*)}{\partial \alpha} = 0$$
, for

$$\lambda \cos(\alpha + \beta) = -(\mu - 2\phi^*) \sin\beta \qquad (\phi \le \alpha \le \pi - \phi) \qquad (4.1.70)$$

$$\cos(\alpha + \beta) = -(1 - 2\phi^*)\sin\beta \quad (\alpha \le \varphi \text{ or } \alpha \ge \pi - \varphi) \quad (4.1.71)$$

$$\frac{\partial \left(\tau/\mathbf{f}_{\mathbf{C}}^{*}\right)}{\partial \beta} = 0, \text{ for }$$

$$(\lambda - \mu \sin \alpha) \sin (\alpha + 2\beta) = 2\phi^* \sin^2 \beta$$
 $(\phi \le \alpha \le \pi - \phi)$

(4.1.72)

$$(1 - \sin\alpha)\sin(\alpha + 2\beta) = 2\phi^*\sin^2\beta \quad (\alpha \le \phi \text{ or } \alpha \ge \pi - \phi) \quad (4.1.73)$$

Solving equations (4.1.70) and (4.1.72), we find that

$$\sin(\alpha + \beta) = \eta \cos\beta$$
 valid for $\phi^* \le \frac{\mu}{2}$ (4.1.74)

$$\cot \beta = \sqrt{\frac{1 - D^2}{\eta^2 - 1}} \qquad \beta \ge \theta \qquad (4.1.75)$$

$$\alpha = \frac{1+\eta}{\eta+D} \qquad \text{and} \quad \phi \leq \alpha \leq \pi-\phi \qquad (4.1.76)$$

Here, the parameters η and D are defined by

$$\eta = \frac{\lambda}{\mu} \tag{4.1.77}$$

$$\begin{cases} \mu - 2\phi^* & \text{for } \phi^* < \frac{\mu}{2} \\ 0 & \text{for } \phi^* \ge \frac{\mu}{2} \end{cases}$$
 (4.1.78)

Inserting equations (4.1.70), (4.1.74) and (4.1.75) into equation (4.1.67) we get the lowest upper bound

$$\frac{\tau}{f_{\rm c}^*} = \frac{1}{2} (\lambda^2 - \mu^2) (1 - D^2)$$
 (4.1.79)

It should be pointed out that equation (4.1.79) is only valid when

$$\beta > \theta \tag{4.1.80}$$

and

Inserting equations (4.1.75) and (4.1.60) into equation (4.1.80) and inserting equation (4.1.76) into equation (4.1.81) yield the following condition for equation (4.1.79) to be valid

$$\frac{1}{1} > \frac{1}{1} = \frac{D^2}{1}$$

(4.1.82)

and

70 %

$$\lambda + \mu + D \geq 1$$

(4.1.83)

If equation (4.1.82) is not satisfied, it means that $\beta \le \theta$. Thus β must be restricted to the value (4.1.60), i.e.

$$\cot \beta = \frac{a}{h} \tag{4.1.84}$$

Solving equations (4.1.70) and (4.1.84) yields

$$\sin \alpha = \frac{\frac{a}{h} \left(\frac{a}{h}\right)^2 + 1 - D^2 + D}{\left(\frac{a}{h}\right)^2 + 1}$$
 valid for $\varphi \le \alpha \le \pi - \varphi$ (4.1.85)

The lowest upper bound is obtained by inserting equations (4.1.70) and (4.1.84) into equation (4.1.67).

$$\frac{\tau_{\star}}{f_{c}} = \frac{1}{2} \left[\lambda \left(\frac{a}{h} \right)^{2} + 1 - D^{2} - \mu \frac{a}{h} \right]$$
 (4.1.86)

The condition for equation (4.1.86) to be valid can easily be found by inserting equation (4.1.85) into equation (4.1.81). This leads to

$$\frac{a}{h} \geq tg\varphi - \frac{D}{\cos\varphi} \tag{4.1.87}$$

When equation (4.1.87) is not satisfied, it means that α must be equal to or less than φ and equation (4.1.70) is not valid anymore.

Solving equations (4.1.71) and (4.1.84) yields the solutions.

$$\sin \alpha = \sin \varphi$$
 valid for tg $\varphi = \frac{C}{\cos \varphi} \le \frac{A}{h} \le \text{tg}\varphi = \frac{D}{\cos \varphi}$ (4.1.88)

$$\sin \alpha = \frac{\frac{1}{h}(\frac{1}{h})^2 + 1 - C^2 + C}{(\frac{1}{h})^2 + 1}$$
 valid for $\frac{1}{h} \le tg\phi - \frac{C}{cos\phi}$ (4.1.89)

Here, the parameter C has been introduced

$$C = \begin{cases} 1 - 2\phi^* & \text{when } \phi^* < \frac{1}{2} \\ 0 & \text{when } \phi^* \ge \frac{1}{2} \end{cases}$$
 (4.1.90)

As in the case corresponding to equation (4.1.88), the lowest upper bound can be easily found by inserting equations (4.1.84) and (4.1.88) into equation (4.1.68).

$$\frac{\tau}{f_{\rm c}^*} = \frac{\left[\left(\frac{{\rm a}}{{\rm h}}\right)^2 + 1 \right] (1 - \sin\varphi) - 2\phi^* \left(\frac{{\rm a}}{{\rm h}} \cos\varphi - \sin\varphi\right)}{2\left[\frac{{\rm a}}{{\rm h}} \sin\varphi + \cos\varphi\right]} \tag{4.1.91}$$

When $\frac{a}{h} \le tg_{\phi} - \frac{C}{\cos \phi}$, we find the lowest upper bound by inserting equations (4.1.84) and (4.1.89) into equation (4.1.68)

$$\frac{\tau}{f_c^*} = \frac{1}{2} \left[\left(\frac{a}{h} \right)^2 + 1 - c^2 - \frac{a}{h} \right]$$
 (4.1.92)

In the over reinforced case, i.e. when $\phi^* \geq \frac{\mu}{2}$ corresponding to $\varphi \leq \alpha \leq \pi - \varphi$, or when $\phi^* \geq \frac{1}{2}$ corresponding to $\alpha \leq \varphi$ or $\alpha \geq \pi - \varphi$, we may find the lowest upper bound in two ways.

One way is to use the same procedure as above which is done first.

When the beam is over reinforced, it appears that

$$\frac{\partial \left(\tau/f_{\mathbf{c}}^{*}\right)}{\partial \alpha} > 0 \tag{4.1.93}$$

It means that we find the lowest upper bound when α equals the smallest possible value. From the condition (4.1.58), we have

$$\alpha + \beta = \frac{\pi}{2} \tag{4.1.94}$$

Inserting equation (4.1.94) into equation (4.1.67) yields

$$\frac{\tau}{f_c^*} = \frac{\lambda - \mu \cos \beta}{2\sin \beta}$$

valid for $\varphi \le \alpha \le \pi - \varphi$

(4.1.95)

It appears that

$$\left[\frac{\partial \left(\tau/f_{\rm C}\right)}{\partial \beta}\right] = 0$$

for

$$\cot \beta = \frac{1}{1 - 2} = t g \alpha$$

(4.1.96)

Inserting equation (4.1.96) into equation (4.1.95) yields

$$\frac{\tau}{f_{\rm c}^*} = \frac{1}{2} \sqrt{\lambda^2 - \mu^2}$$
 valid for $\phi^* \ge \frac{\mu}{2}$ (4.1.97)

This solution is only valid if equations (4.1.80) and (4.1.81) are satisfied. From equations (4.1.80), (4.1.81) and (4.1.96), we get the condition for equation (4.1.97) to be valid:

$$\frac{a}{h} \ge \frac{1}{\eta^2 - 1} \tag{4.1.98}$$

and

When equation (4.1.98) is not fulfilled, we find the lowest upper bound by inserting equation (4.1.84) into equation (4.1.95)

$$\frac{r_{\star}}{f_{\rm C}^*} = \frac{1}{2} \left[\lambda \left[1 + \left(\frac{a}{h} \right)^2 - \mu \frac{a}{h} \right] \right] \quad \text{valif for } \phi^* \ge \frac{\mu}{2} \qquad (4.1.100)$$

The condition for equation (4.1.81) to be valid in this case becomes

$$\frac{a}{h} \ge tg\varphi$$
 (4.1.101)

In the case of $\frac{a}{h}$ \leq tg ϕ , we can easily find the solution by inserting equations (4.1.94) and (4.1.84) into equation (4.1.68).

$$\frac{\tau_*}{f_C} = \frac{1}{2} \left[\left(1 + \left(\frac{a}{h} \right)^2 - \frac{a}{h} \right] \quad \text{valid for } \phi^* \ge \frac{1}{2} \quad (4.1.102)$$

Another way to find the lowest upper bound in the case of an overreinforced beam is minimizing the equations (4.1.79), (4.1.86) and (4.1.92) with respect to ϕ^* . The same solutions as equations (4.1.97), (4.1.100) and (4.1.102) have been found in this way.

As mentioned above solutions (4.1.79) and (4.1.97) are valid only when equations (4.1.83) and (4.1.99) are fulfilled. When they are not satisfied, the lowest upper bound may be obtained by differentiating equations (4.1.79) and (4.1.97) with respect to *

$$\lambda + \mu + D = 1$$
 (4.1.10)

Inserting equations (4.1.63), (4.1.64) and (4.1.78) into equation (4.1.103) yields

$$\rho^* = \begin{cases} \frac{\sqrt{k} - \sqrt{1 + \phi^*(k - 1)}}{\sqrt{k}(k - 1)} & \text{valid for } \phi^* < \frac{k-1}{4k} \end{cases}$$

$$(4.1.1)$$

$$\frac{1}{2k}$$
valid for $\phi^* \ge \frac{k-1}{4k}$

In this case curves corresponding to equations (4.1.82) and (4.1.98) are overlapping the curves corresponding to equations (4.1.87) and (4.1.101). Inserting equation (4.1.104) into solutions (4.1.79) and (4.1.97), we find that

$$\frac{\tau_{\kappa}}{f_{C}^{*}} = 2 \frac{\left[\sqrt{k} - \sqrt{1 + \phi^{*}(k - 1)}\right] \left[\sqrt{k(1 + \phi^{*}(k - 1))} - 1\right]}{(k - 1)^{2}} \phi^{*} \leq \frac{k - 1}{4k}$$
(4.1.105)

$$= \frac{1}{2\sqrt{k}} \qquad \qquad \phi^* \ge \frac{k-1}{4k} \qquad (4.1.106)$$

Now equation (4.1.87) becomes

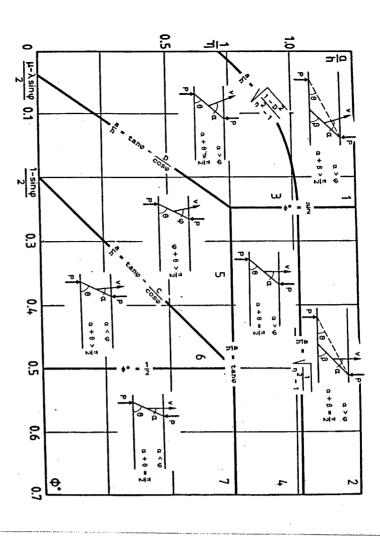
$$\frac{a}{h} \ge \frac{1 + \phi^*(k-1)}{\sinh 2} - \cot \varphi$$
 for $\phi^* \le \frac{k-1}{4k}$ (4.1.107)

$$\frac{a}{h} \geq tg\varphi \qquad \qquad for \ \phi^* \geq \frac{k-1}{4k} \qquad (4.1.108)$$

When equations (4.1.107) and (4.1.108) are not satisfied, solutions (4.1.91), (4.1.92) and (4.1.102) will still be valid.

Fig. 4.7 shows the domains of the different failure mechanisms in our upper bound analysis with the ρ^* -value satisfying equation (4.1.83). The boundaries are plotted for $\varphi=37^{\circ}$, i.e. k=4 and $\rho^*=0.1$.

In the case of large ρ^* -values, i.e. when condition (4.1.83) is violated, the failure mechanism marked 3 and 4 in Fig. 4.7 will disappear, because their upper and lower boundaries overlap, while the failure mechanism marked 1, 2, 5, 6 and 7 remain applicable.



 $tan\phi - \frac{c}{cos\phi} \le \frac{a}{h} \le tan\phi - \frac{D}{cos\phi}$

 $\frac{a}{h} \leq \tan \theta - \frac{c}{\cos \theta}$

 $\frac{1}{2} \left[\sqrt{\frac{a}{h}} \right]^2 + 1 - C^2 - \frac{a}{h}$

 $\frac{1}{2} \left[\sqrt{\frac{2}{h}} \right]^2 + 1 - \frac{2}{h}$

 $\tan \varphi - \frac{D}{\cos \varphi} \le \frac{a}{h} \le \sqrt{\frac{1 - D^2}{\eta^2 - 1}}$

2(1/(A)2+1-D2-4 A]

½ [λ /(ਜ)2+1-μ ਜ]

 $\left[\left(\frac{a}{h} \right)^2 + 1 \right] (1 - \sin \varphi) - 2\phi^* \left(\frac{a}{h} \cos \varphi - \sin \varphi \right)$

 $2\left(\frac{a}{h} \sin \phi + \cos \phi\right)$

 $\sqrt{\frac{1-D^2}{\eta^2-1}} \leq \frac{a}{h}$

 $\frac{1}{2}\sqrt{(\lambda^2-\mu^2)(1-D^2)}$

 $\frac{1}{2}\sqrt{\lambda^2-\mu^2}$

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2 / 0* / 2

ф* IV 21-1 λ + μ + D 2 1

Figure 4.7. Domains of the different failure mechanisms for small $\rho^{\star}\text{-values.}$

4.1:3 The complete plastic solutions

We have derived both the highest lower bound solutions and the lowest upper bound solutions. It is easy to note that they are identical. These solutions thus are exact according to our assumptions in section 2.6.

The complete plastic solutions are summarized in Fig. 4.8 and Fig.

Figure 4.8. The complete plastic solutions for small ho^* -values.

In Fig. 4.8 and 4.9, the parameters are defined as follows.

$$\rho^* = \frac{f_L}{f_C}, \qquad k = \frac{1 + \sin \varphi}{1 - \sin \varphi}$$

$$\lambda = 1 - \rho^*(k - 1) , \qquad \eta = \frac{\lambda}{\mu}$$

$$\mu = 1 - \rho^*(k + 1) , \qquad \phi^* = \frac{\lambda_S f_V}{bh f_C}$$

$$\begin{cases} 1 - 2\phi^* & \text{for } \phi^* \le \frac{1}{2} \\ 0 & \text{for } \phi^* > \frac{1}{2} \end{cases}$$

f* " bhf.*

$$D = \begin{cases} \frac{\mu - 2\phi^2}{\lambda} \\ 0 \end{cases}$$

n < tanφ = cosφ	$\tan \theta = \frac{c}{\cos \varphi} \le \frac{A}{h} \le \frac{\sqrt{1 + (k-1)} \phi^2}{\sin \varphi} = \cot \varphi$	$\frac{\sqrt{1+(k-1)\phi^2}}{\sin\phi} \cdot \cot\phi \le \frac{a}{h}$	TIP OF THE STATE O
1/21 /(1/2) 2+1-c2-4]	$\frac{\left(\frac{2}{h}\right)^{2}+11\left(1-\sin{\psi}\right)-2\phi^{*}\left(\frac{2}{h}\cos{\psi}-\sin{\psi}\right)}{2\left(\frac{2}{h}\sin{\psi}+\cos{\psi}\right)}$	$2\frac{[\sqrt{k}-\sqrt{1+(k-1)}\phi^*][\sqrt{k(1+(k-1)}\phi^*]-1]}{(k-1)^2}$	λ+μ+D<1 Φ* Δ 4k
· 기 / 류) ²+ 1 − 류]		2/2	

Figure 4.9. The complete plastic solutions for large ho * -values.

These solutions are also valid when $ho^*=0$, which correspond to the so-called square yield locus (Fig. 2.15). It is easy to find simplified and we get that in this case the complete solutions will be very much

$$\frac{\tau}{f_{c}^{*}} = \begin{cases} \frac{1}{2} \left[\left(\frac{a}{h} \right)^{2} + 1 - c^{2} - \frac{a}{h} \right] & \text{for } \phi^{*} \le \frac{1}{2} \\ \frac{1}{2} \left[\left(\frac{a}{h} \right)^{2} + 1 - \frac{a}{h} \right] & \text{for } \phi^{*} > \frac{1}{2} \end{cases}$$

(4.1.102)

(4.1.92)

considering the tensile strength of concrete. simple solutions to distinguish them from the complete solutions In what follows we will name equations (4.1.92) and (4.1.102) the

plotted for $\varphi = 37^{\circ}$, i.e. k = 4 and $\rho^* = 0.1$. Fig. 4.10 shows the domains of the complete plastic solutions for small ho^* -values in the $(\phi^*, \frac{1}{h})$ parameter plane. The boundaries are

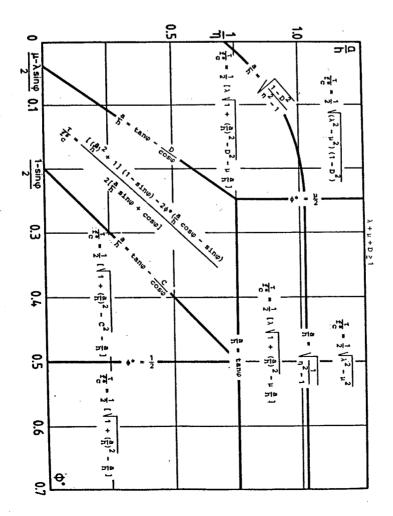


Figure. 4.10. Domains of the complete plastic solutions for beams reinforcement. subjected to point loading and without web

are plotted for $\varphi = 37^{\circ}$, i.e. k = 4 and $\rho^* = 0.2$. large ho^* -values in the $(\phi^*, \frac{a}{h})$ coordinate plane. The boundaries Fig. 4.11 shows the domains of the complete plastic solutions for

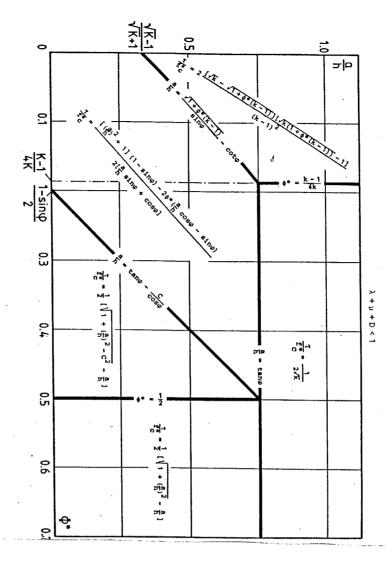
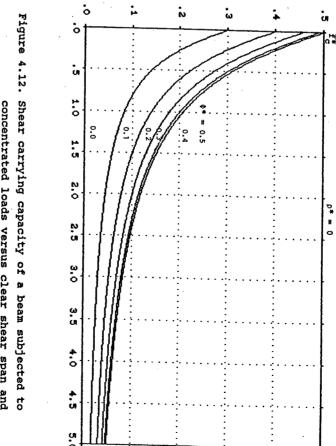
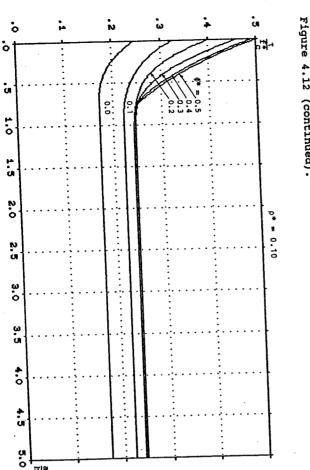


Figure 4.11. Domains of the complete plastic solutions for large ρ^* - values.

a/h for different values of ϕ^* and ρ^* In Fig. 4.12 the shear carrying capacity $au/f_{\mathbf{C}}^{*}$ is plotted against



concentrated loads versus clear shear span and longitudinal reinforcement degree.



and ρ^* - values.

Fig. 4.13 shows the relation between the shear capacity and the longitudinal reinforcement degree for different shear span ratios

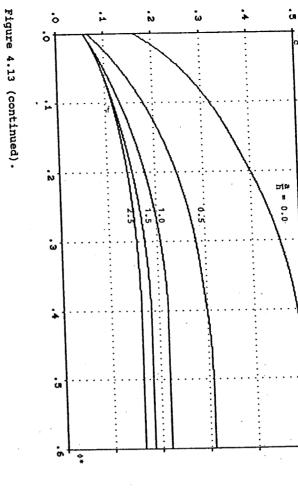
P* 0.0

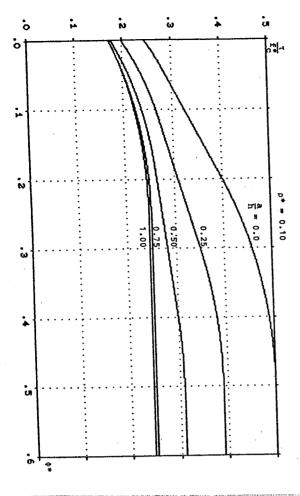
D 0 0 .

0.5

Figure 4.13. Shear capacity of a beam loaded by concentrated force versus longitudinal reinforcement degree.







has been shown in different values of $\frac{a}{h}$ and ϕ^* . capacity, In Fig. 4.14 the influence of the ho^*- value to the shear capacity ue. This special value is very much influenced by the ho^*- value. of the clear shear span ratio before it reaches some special vallongitudinal reinforcement degree and decreases with an increase in general, increases with an increase

From Fig. 4.12 and 4.13, it is very clear that the shear carrying

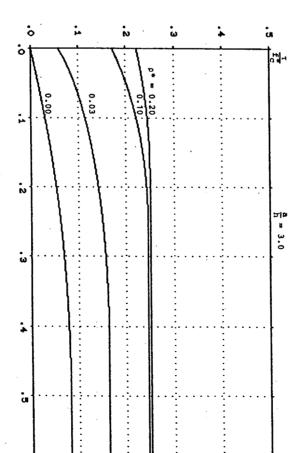
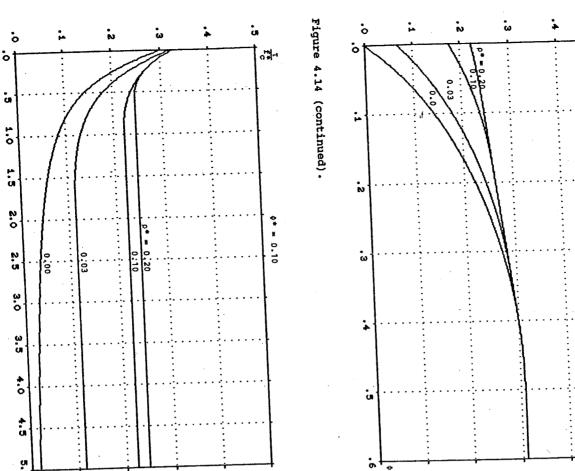


Figure 4.14. The influence of ho^* to shear capacity for different values of $\frac{a}{h}$ and ϕ^* .

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The difference between the cases of considering the tensile strength of concrete and neglecting it can be seen very clearly in Fig. 4.14, especially for beams with large shear span ratio and for corbels with low reinforcement degrees.

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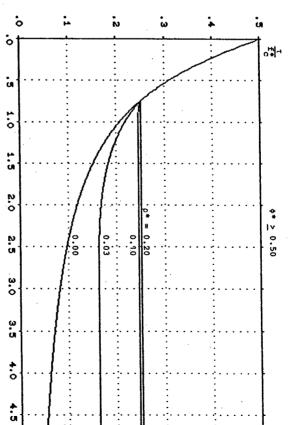
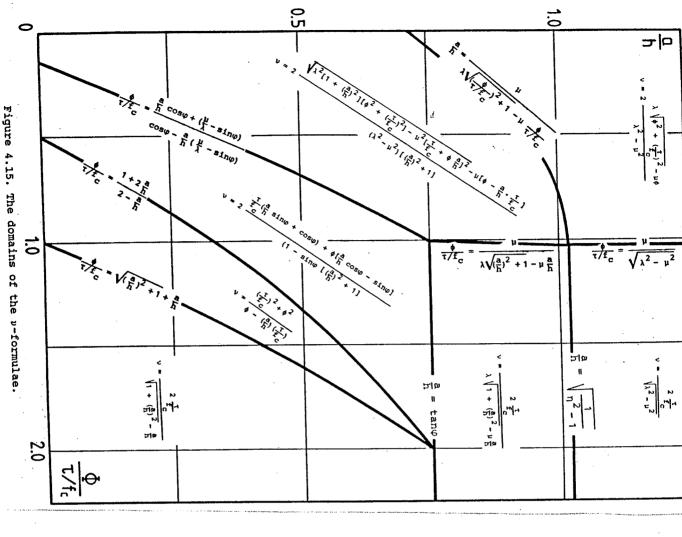


Figure 4.14 (continued).

4.1.4 The ρ^* and ν values in the plastic solutions

As mentioned before, the effectiveness factors ν and ρ^* in our complete solutions must be determined by experiments.

Inserting $f_C^* = \nu f_C$ into the complete plastic solutions shown in Fig. 4.8 and solving it with respect to ν , we find the corresponding domains of the ν -formulae as shown in Fig. 4.15. The boundaries are plotted for $\varphi = 37^{\circ}$, i.e. k = 4 and $\rho^* = 0.1$.



It is worth noticing that when we determine the v-value for shear from test results, we should exclude flexural failures and local failures and preferably use the test data from test specimen with high reinforcement ratios. Otherwise, it may produce very unreliable results especially in the case of tests with very low reinforcement ratios.

A great number of test data for non-shear reinforced beams, deep beams and corbels subjected to concentrated loads have been found in literature. The detailed test data, including the dimensions, material properties and measured ultimate loads of beams, deep beams and corbels, and the sources are shown in Appendix B.

It is obvious that the ν -value is very much dependent on the ρ^* -value that we have chosen. We may decide the ρ^* -value from the beam failure model. According to our plastic solutions, for the overreinforced case, when equation (4.1.98) is satisfied, the τ/f_C^* is independent from a/h. This conclusion has been verified by the tests. A great number of tests have shown that τ/f_C is almost a constant for beams with relatively big shear span ratios failed in shear. When the shear span ratio is less than a definite value, the τ/f_C is decreasing with increasing shear span ratio. The critical shear span ratio is about 2.5.

Inserting a/h = 2.5 into equation

$$\frac{a}{h} = \sqrt{\frac{1}{\eta^2 - 1}} \tag{4.1.109}$$

we get the ho^*- value about 0.03. This value satisfies the requirement of equations (4.1.51) and (4.1.52).

Parametric analysis indicates that the effectiveness factor ν mainly is a function of the uniaxial compressive strength f_C , the ratio of reinforcement φ , defined as $\varphi=A_S/bh$, the clear shear span ratio a/h and the absolute depth of the specimen.

Fig. 4.16 and Fig. 4.17 show some examples of the dependence of $\boldsymbol{\nu}$

on $\mathbf{f}_{\mathbf{C}}.$ It is very clear that ν decreases with the increase of the uniaxial compressive strength of concrete.

For simplicity, the relation between v and f_{C} may be described empirically by assuming that v is inversely proportional to $\sqrt{f_{C}}$.

$$v = \frac{k_1}{\sqrt{f_C}}$$
 (4.1.110)

Here, k_1 is a constant and f_C is in MPa.

For comparison, such kind of curves is drawn in Fig. 4.16 and Fig. 4.17 together with the tests. A good agreement is found.

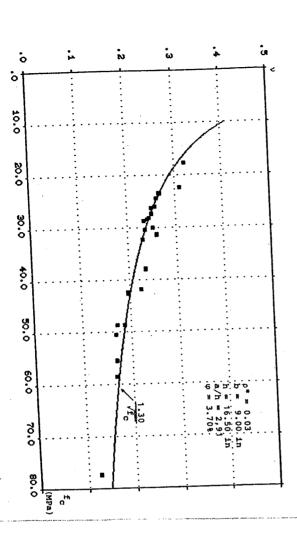


Figure 4.16. The v-dependence on f_C , [62.2].

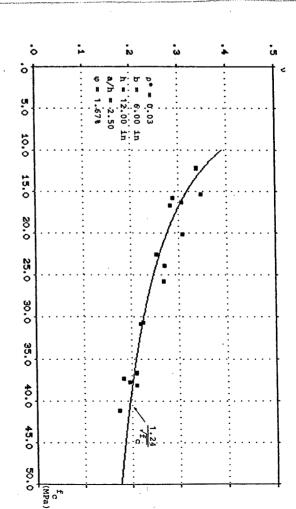


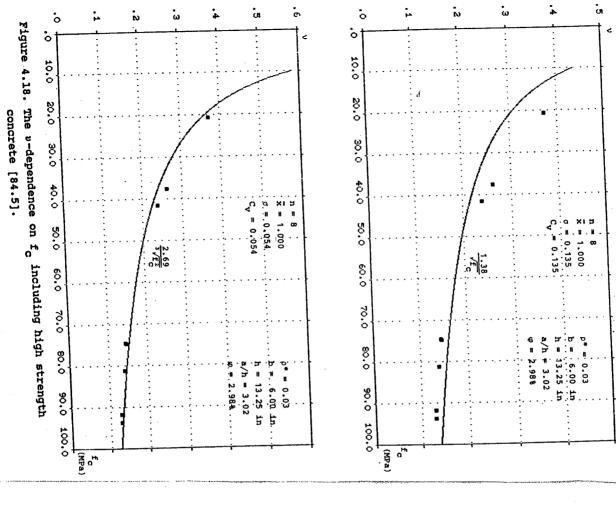
Figure 4.17. The ν -dependence on f_C , [54.2].

For high strengths of concrete, the $\nu\text{-dependence}$ on $\mathbf{f}_{\mathbf{C}}$ may be better expressed as

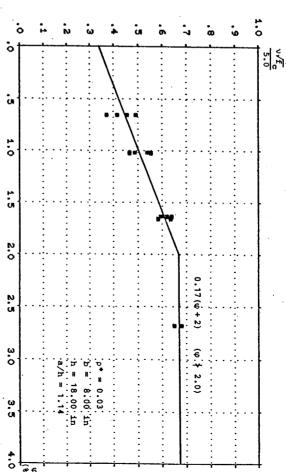
$$v = \frac{k}{3f_c^2}$$
 (4.1.111)

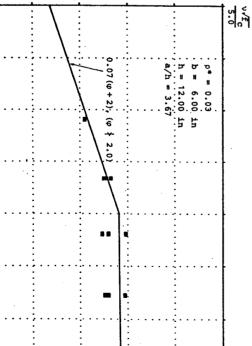
The comparison of equations (4.1.110) and (4.1.111) with some tests with high strength of concrete can be seen in Fig. 4.18. Unfortunately, we have only found a few shear tests with high strength of concrete. In the following sections we will only use (4.1.110) to describe approximately the ν -dependence on $\mathbf{f}_{\mathbf{C}}$.

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Figure 4.19. The ν -dependence on φ [63.2].

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linearly with ϕ for lower reinforcement ratios, for example ϕ < As clearly shown in Fig.4.19, the v-value is increasing almost 2.0, and is independent of φ for higher ratios, i.e. φ \geqslant 2.0.

The v-dependence on a/h may be expressed as

$$v = k_2(3.5 - \frac{a}{h})$$
, $(\frac{a}{h} \neq 2.5)$ (4.1.112)

Here k_2 is a constant

Some experimental evidence is shown in Fig. 4.20.

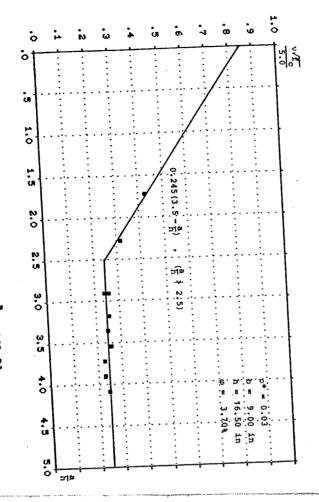


Figure 4.20. The v-dependence on $\frac{a}{h}$. [62.2].

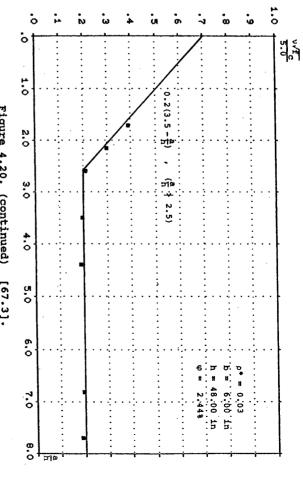
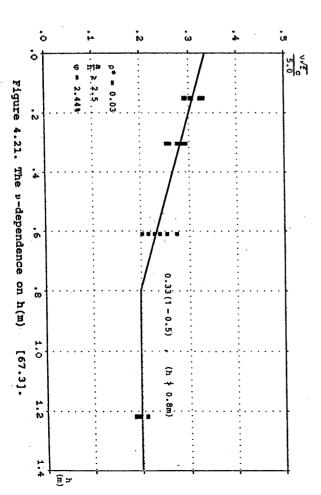


Figure 4.20. (continued) [67.3].

A parametric analysis has shown that the v-value decreases with increasing specimen depth. Some evidence is provided in Fig. 4.21.



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From the parametric analysis of 657 observed shear failures of beams, deep beams and corbels, the following empirical ν formula

with $\rho^* = 0.03$ is suggested

$$0.35(3.5 - \frac{a}{h})(\phi + 2)(1 - 0.2h) \le 1.0 \qquad \begin{cases} \frac{a}{h} > 2.5 \\ \phi > 2.0 \\ h > 1.25 \end{cases}$$

$$\sqrt{f_C}$$
(4.1.113)

Here f_{C} is in MPa, ϕ in percent and h in meter.

Formula (4.1.113) is verified by tests having $f_{\rm C} \le 100$ MPa, a/h ≤ 10.0 and h ≤ 1.20 m. Out of these ranges, more tests are needed. Results of 657 beams, deep beams and corbels from 30 available references were compared to the complete solutions in Fig. 4.8 using (4.1.113).

The statistical values of the ratios of test to theory are presented in Fig. 4.22. A very good agreement has been found.

		bserve	Observed shear failure	ure		Cheoreti	Theoretical shear failure	ilure
, Lean					aumhor.	mean	standard	coeff. of
State of the state	number	mean	standard deviation	coeff. of variation	TAGUIDET		deviation	variation
Sall Sall Sall Sall Sall Sall Sall Sall		×ι	Q	c C	p p	×ı	Q	°C√
Cases	,					245	0 148	0.146
	412	1.004	0.152	0.151	353	1.015		
beams				0 143		1.022	0.136	0.133
deep beams	127	1.018	0.146			2	0 142	0.141
1	118	1.002	0.140	0.140	110			0 1 43
COLUMN		1.006	0.149	0.148	564	1.015	0.145	

Figure 4.22. The statistical values of the ratios of test to theory.

The comparison between theory and test results is also shown in

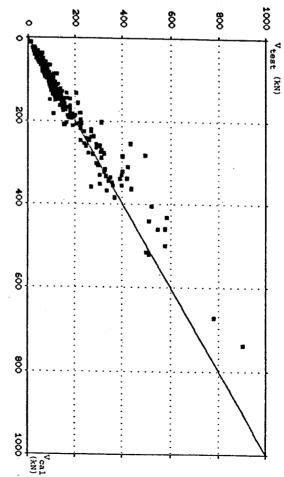


Figure 4.23. The comparison between theory and tests of beams.

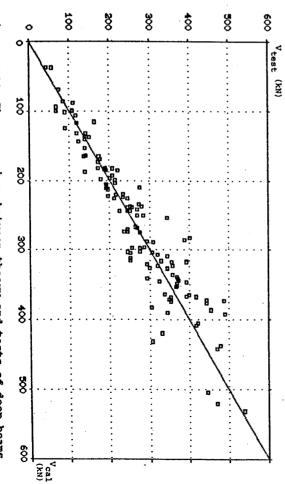


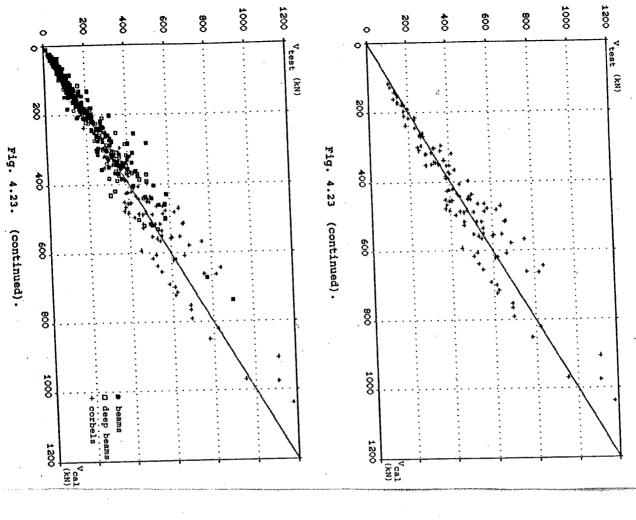
Figure 4.23. The comparison between theory and tests of deep beams (continued).



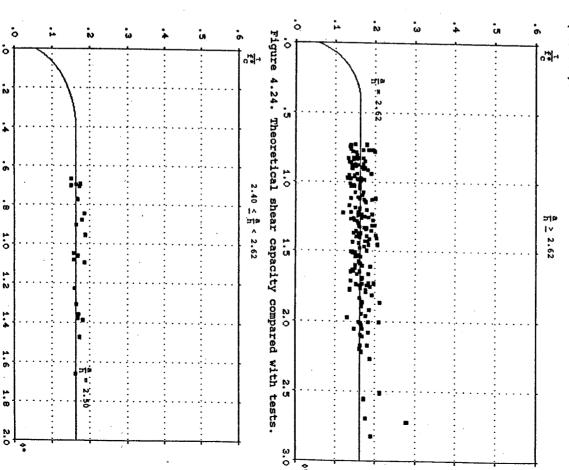
gure 4.24.

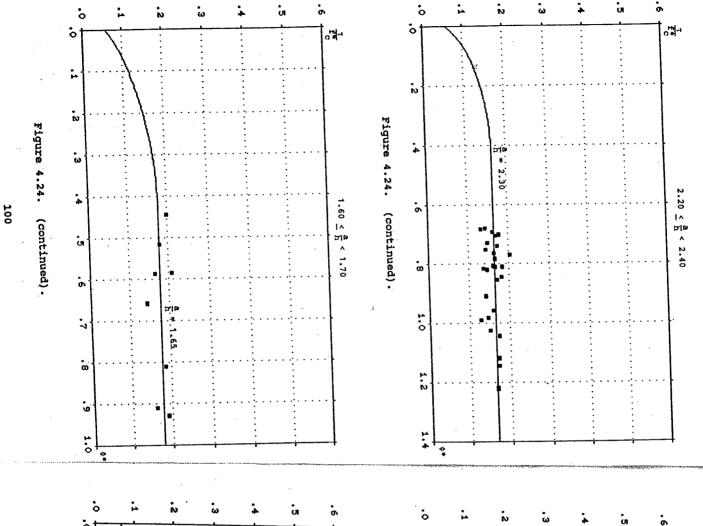
(continued).

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To demonstrate the general applicability of the complete plastic solutions, Fig. 4.24 shows the comparison of theory with the results of the shear tests of beams, deep beams and corbels for different shear span ratios. The ν -value has been calculated from (4.1.113).





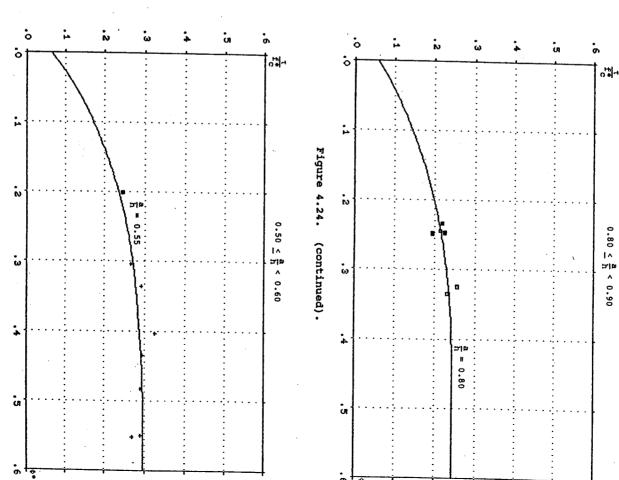
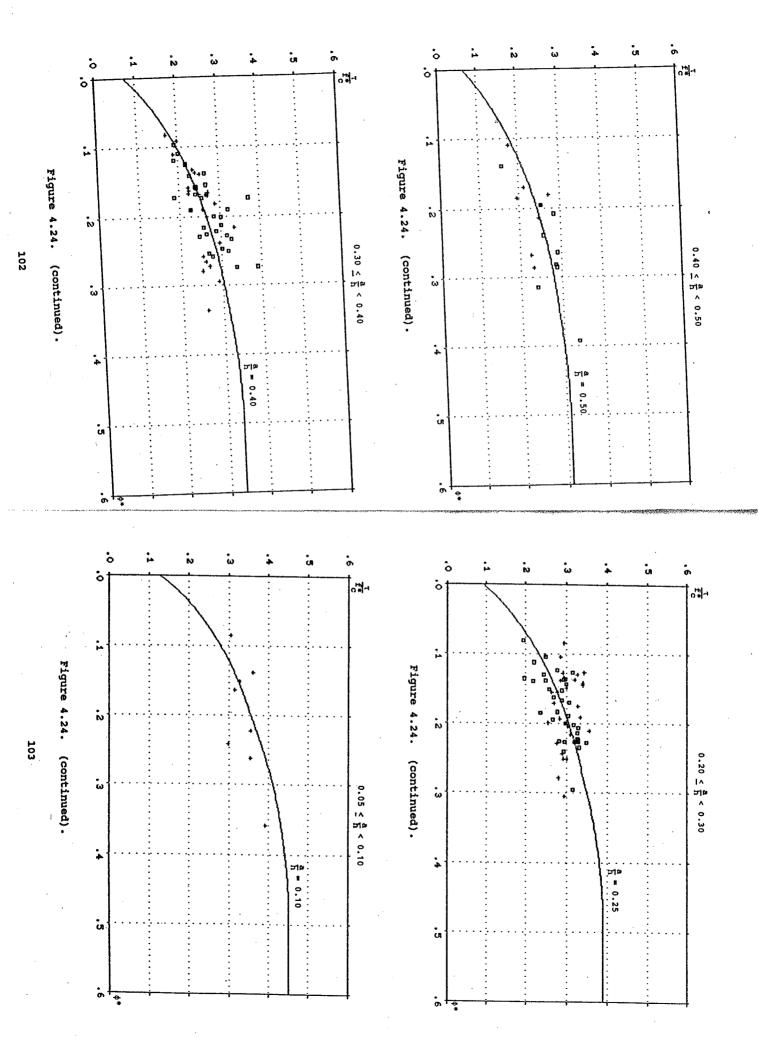


Figure 4.24. (continued).



4.2 Shear capacity neglecting the tensile strength of concrete

In section 4.1 we have derived the complete exact beam shear solutions for structural members without web reinforcement and subjected to concentrated loads. For practical use, it may often be considered to be too complicated to use. Further as a matter of fact, the behaviour of concrete in tension is almost brittle. Therefore, in most cases, it is reasonable to neglect the tensile strength of concrete. The exact simple solutions of beam shear strength of concrete. The exact simple solutions of beam shear based on the yield condition reducing to the so-called square based on the Yield condition reducing to the so-called square [78.2] and [79.3]. The corresponding empirical v formula were proposed by Roikjær [79.4].

In this section, an alternative derivation gaving identical solutions is developed.

4.2.1_Lower_bound_solution

a state of biaxial compression with horizontal stress $\sigma_{
m b}$ and hydrostatic compression with a stress termed $\sigma_{\mathbf{b}}.$ The part I is in subdivided into firmly connected with the lateral steel plate. The beam is beam with concentrated loading. The longitudinal reinforcement is Fig. 4.25 shows a stress distribution in a rectangular concrete uniaxial compression with stress $\sigma_{\mathbf{b}}$, but along line AB there is compression with stress $\sigma_{\mathbf{b}}.$ The part II is also in a state of vertical stress termed $\sigma_{\mathbf{V}}$. The part III is in a state of uniaxial conditions. some bond stress with the reinforcement. In the bending section, equations and the statical boundary conditions on the upper and force of concrete is termed C. This stress state is statically the reinforcement force is termed T and the resultant compression admissible ín The shaded regions are in a state of biaxial, the sense a number of areas with that j.t satisfies homogeneous stress the equilibrium

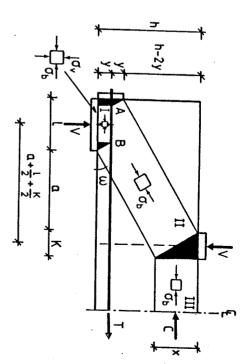


Fig. 4.25. Stress distribution in a rectangular beam subjected to symmetrical two-point loads.

The length of the load platen k is determined by the geometrical relation

$$\frac{k}{x} = \tan \omega = \frac{h-x}{a+x}$$
 (4.2.1)

with the solution

$$k = \frac{1}{2} \left[\left(a^2 + 4x(h-x) - a \right) \right]$$
 (4.2.2)

The depth ${\bf x}$ of the compressive concrete zone is determined by horizontal equilibrium

$$= \frac{T}{b\sigma_b} \tag{4.2.3}$$

The shear load V corresponding to this stress distribution is found by equilibrium of vertical forces

Inserting equation (4.2.3) into (4.2.4) yields

$$V = \frac{1}{2} \left[\sqrt{a^2 b^2 \sigma_b^2 + 4T(hb\sigma_b^{-T})} - ab\sigma_b \right]$$
 (4.2,5)

We now determine the highest lower bound by maximizing (4.2.5) with respect to the statical variables $\sigma_{\rm b}$ and ${
m T.}$ It appears that:

$$\frac{dv}{d\sigma_b} > 0$$
 always (4.2.6) $\frac{dv}{d\sigma_b} > 0$ for $T \le \frac{1}{2} hb\sigma_b$

Thus the highest lower bound is obtained with the maximum concrete stress, i.e. $\sigma_b = f_C^*$. For $T \le \frac{1}{2} \, hbf_C^*$, the highest lower bound is obtained when the reinforcement force is maximum, i.e. $T = \phi^*bhf_C^*$. Inserting into equation (4.2.5), we get

$$\frac{T}{f_c} = \frac{1}{2} \left[\left(\frac{a}{h} \right)^2 + 4\phi^* (1 - \phi^*) - \frac{a}{h} \right] \qquad \text{for } \phi^* \le \frac{1}{2}$$
 (4.2.7)

For T> $\frac{1}{2}$ bhf $_{\rm C}^*$, the highest lower bound is obtained with T= $\frac{1}{2}$ bhf $_{\rm C}^*$ and we find

$$\frac{\tau}{f_c^*} = \frac{1}{2} \left[\left(\frac{a}{h} \right)^2 + 1 - \frac{a}{h} \right] \qquad \text{for } \phi^* > \frac{1}{2} \qquad (4.2.8)$$

This solution coincides with the solution derived by Nielsen et alusing different stress distributions [78.2], [78.3].

4.2.2_Upper_bound_solution

Consider a failure mechanism consisting of a curved yield line in the shear span of the rectangular concrete beam subjected to the concentrated load V, (Fig. 4.26). The yield line is running from A, which is the outside edge of the load platen, to B, which is the inside edge of the support platen. The part I rotates relatively round an arbitrary point 0 with angular velocity of η . The inclination of the straight line AB is determined by the angle θ , where $\cot\theta = \frac{a}{h}$. M is the centre of line AB and the relative displacement of point M is δ_m , inclined at the angle α_m to the line AB. The length of line OM is termed γ_m , inclined at the angle θ to the horizontal direction. The means of the parameters t, 1 and y can be clearly found in Fig. 4.26.

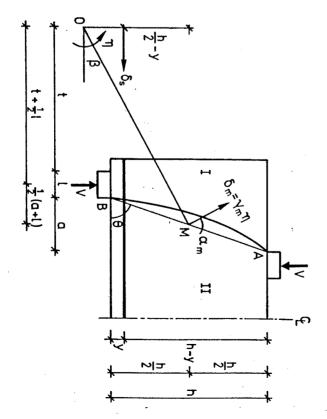


Fig. 4.26. Failure mechanism for shear span subjected to concentrated load.

We assume that the reinforcement is yielding in tension, i.e.

The rate of internal work dissipated by this mechanism is

$$W_{I} = \frac{1}{4} bf_{c}^{*} \eta \{ (2t+21+a) \left[\sqrt{a^{2}+h^{2}} \cdot \sqrt{1+tan^{2}\beta-htan\beta-a} \right]$$

$$+ 2 \frac{\frac{A_s f}{s_{\star}} Y_{[(2t+21+a) \tan \beta - (h-2y)]}}{b f_{c}}$$
 (4.2.10)

where we have used equation (2.4.12) to compute the contribution from the web concrete.

The external work is

$$W_{\rm E} = \frac{1}{2}(2t+1)\eta V \tag{4.2.11}$$

The work equation $\mathbf{W}_{\mathbf{I}} = \mathbf{W}_{\mathbf{E}}$ then yields an upper bound for the ultimate shear stress:

$$\frac{\tau}{f_{c}^{*}} = \frac{(2t+21+a)(\sqrt{a^{2}+h^{2}} \cdot \sqrt{1+\tan^{2}\beta-h\tan\beta-a})+2\phi^{*}h[(2t+21+a)\tan\beta-(h-2y)]}{2h(2t+1)}$$
(4.2.12)

The lowest upper bound can then be found by minimizing equation (4.2.12) with respect to the variables t and $an\!eta$.

It appears that

$$\frac{\partial (\tau/f_C^*)}{\partial (\tan \beta)} = 0 \quad \text{for } \sin \beta = \sin \theta \cdot (1-2\phi^*) \quad (\text{valid for } \phi^* \le \frac{1}{2}) \quad (4.2.13)$$

$$\frac{\partial (\tau/f_{C}^{*})}{\partial t} = 0 \quad \text{for } t \rightarrow \infty$$

(4.2, 14)

Inserting equations (4.2.13) and (4.2.14) and $\cot\theta = \frac{d}{h}$ into (4.2.12), we get the lowest upper bound

$$\frac{r_{\star}}{f_{\rm c}^*} = \frac{1}{2} \left[\sqrt{\left(\frac{a}{h}\right)^2 + 4\phi^* (1 - \phi^*)} - \frac{a}{h} \right] \qquad \text{for } \phi^* \le \frac{1}{2}$$
 (4.2.15)

For $\phi^*>\frac{1}{2}$, minimum is obtained with $\beta=0$ and $t\to\infty$ since the reinforcement can not be in compression.

In this case the lowest upper bound is easily found to be

$$\frac{\tau}{f_{\rm C}^*} = \frac{1}{2} \left[\left(\frac{a}{h} \right)^2 + 1 - \frac{a}{h} \right] \qquad \text{for } \phi^* > \frac{1}{2}$$
 (4.2.16)

This solution is valid when the longitudinal reinforcement is sufficiently strong, it no longer contributes to the shear capacity, which is determined by the web concrete alone.

Solutions (4.2.15) and (4.2.16) coincide with the solutions derived by Nielsen et al assuming a different failure mechanism, which consists of a straight yield line running from the load to the support and with a relative displacement inclined at an angle to the yield line, [78.2]. We note that the lowest upper bound is identical with the highest lower bound. This means that we have determined the correct value of ultimate shear strength corresponding to our assumptions.

It should be noticed, as only the shear problem is concerned in our analysis here, that the ultimate shear capacity is independent from the location of the horizontal reinforcement. But, the distance y in Fig. 4.25 and Fig. 4.26, which is measured from the bottom of the cross section to the centroid of horizontal reinforcement, does strongly influence the bending carrying capacity of the beam, because the increasing of y leads to decrease of the effective depth d of the beams and then to decrease of the ultimate bending capacity.

Compare equations (4.2.7) and (4.2.8), or (4.2.15) and (4.2.16), with solution (3.1.1), which is the moment capacity formula, we may find the critical position of the horizontal reinforcement Yo. may find the critical position of the horizontal reinforcement Yo. may find the critical position of the carrying capcity of the beam, while for y > yo, the beam will fail in flexure. Same beam, while for y > yo, the length of supporting platen & in influence can be found for the length of supporting platen & in flexure. Same Fig. 4.25 and Fig. 4.26. In [81.2] and [81.3], some solutions had Fig. 4.26. In [81.2] and [81.3], some solutions had been derived by J.F. Jensen and B.C. Jensen for the combinations of bending, shear and bearing carrying capacity of the beam. Those of bending, shear and bearing carrying capacity of the beam depth h, shear span a, the length of supporting platen &, the beam depth h, the distance y and the reinforcement degree \(\phi^* \), and are somewhat

complicateu. It has been found that Jensen's solution are fully identical with It has been found that Jensen's solution are fully identical with the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (3.1.1) and (4.2.7), (4.2.8). But the way the combining equations (4.2.1) and (4.2.7), (4.2.8). But the way the combining equations (4.2.1) and (4.2.7), (4.2.8). But the way the combining equations (4.2.1) and (4.2.7), (4.2.8). But the way the combining equations (4.2.1) and (4.2.7), (4.2.8). But the way the combining equations (4.2.1) and (4.2.1) and

4.2.3 Modified simple formula

The comparison of the simple solutions with the complete solutions can be seen in Fig. 4.14. It appears that the simple solutions have some obvious disadvantages. In the first place, it declares have some obvious disadvantages. In the first place, it declares that the shear capacity will continuously decrease along with the increase of the shear span to depth ratio. According to the simple formula in the case of an overreinforced beam, the shear capacity of beams with shear span ratio equal to 5 will only be half that of beams with shear span ratio equal to 2.5. It does not agree with the tests. In order to fill up the shortage of the simple with the tests. In order to fill up the shortage of the simple formula, the corresponding empirical v formula [79.4] consists of a parabolic contribution with the purpose of increasing the value in the case of big shear span ratios. The v value is very often bigger than unit in such cases.

Secondly, it gives zero shear capacity for beams without any reinforcement and underestimate the shear carrying capacity of beams with low reinforcement degrees, especially when the shear span ratio is rather small. In conventional beams it perhaps makes no problem at all since the requirement of reinforcement is governed by the bending design. But for specimens with very small shear span ratio, such as deep beams, corbels and joints, it may result in too conservative results.

Furthermore, for the reason given above, it is not possible to extend the simple solutions into fibre concrete beams and to shear in corbels with tensile normal force, in which an example can be found in section 4.3.

By utilizing the properties of the complete solutions, we may modify the simple solutions to eliminate the first mentioned shortage of the simple solutions. The simplest way is giving a limit of 2.5 to the shear span ratio when using the simple solutions, i.e. when the shear span ratio is larger than 2.5 it is put equal to 2.5. It means that the modified simple solutions now are

$$\frac{T}{f_c^*} = \begin{cases} \frac{1}{2} \left[\left(\frac{a}{h} \right)^2 + 4\phi^* (1 - \phi^*) - \frac{a}{h} \right] & \text{for } \phi^* \le \frac{1}{2} \\ \frac{1}{2} \left[\left(\frac{a}{h} \right)^2 + 1 - \frac{a}{h} \right] & \text{for } \phi^* > \frac{1}{2} \end{cases}$$

The v formula corresponding to the modified simple solutions (4.2.17) has been found by the parametric analysis of 412 conventional beams from 21 references in the literature (see Appendix B). The result is

$$v = \frac{0.60(2-0.4\frac{a}{h})(\varphi+2)(1-0.25h)}{\sqrt{f_c}} \le 1.0 \qquad \begin{bmatrix} \frac{a}{h} \ne 2.5\\ \varphi \ne 2.0\\ h \ne 1.0 \end{bmatrix}$$
 (4.2.18)

As in equation (4.1.123), f_c is in MPa, ϕ in percent and h in

meter.

The statistical data of the ratios of tests to modified simple theory are presented in Fig. 4.27. If we compare them to the values shown in Fig. 4.22, the equations (4.2.17) and (4.2.18) are seen to be satisfactory for beams. For deep beams and corbels, the complete solutions are somewhat better. Comparisons of the modified simple formulae with test results can also be seen in Fig. 4.28.

total	corbels	deep beams	beams	Cases	1,00%	A Solution of the second of th	State /	Item	
657		127	412		3		number		
1.000	1.086	1.103	1.028		×ı		mean	1000	horva
	0.100				a		mean standard deviation		observed shear failure
	0.163	0.171	0.154	2 451	۷,		variation	es of	lure
	614	116	118	380		3		number	
	1.058	1.088	1.094	1.038		χı		mean	Theoreti
	0.163	0.187	0.151	0.155		a a	DEATURE.	standard	Theoretical shear raiture
na na na	0. 104	0.172	0.138	0.149	Hara	€.	Moreov (Sec	coeff.	1144

Fig. 4.27. The statistics of the ratios of tests to the modified simple theory.

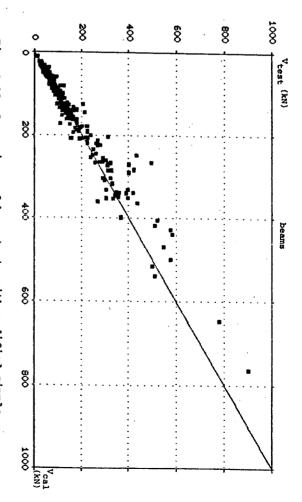


Fig. 4.28. Comparison of beam tests with modified simple theory.

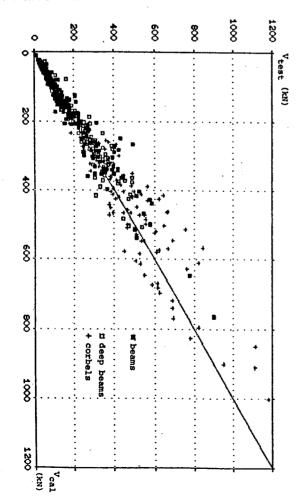


Fig. 4.28. The comparison of beam, deep beam and corbel tests with modified simple theory (continued).

4.3 The influence of normal force on the shear capacity

In this section the problem how normal forces influence the shear capacity of reinforced concrete beams is treated. The best upper bound solution considering the tensile strength of concrete for rectangular beams is presented in what follows. The theory then is compared with the tests of corbels with tensile normal force and beams with compressive normal force. A rather good agreement has been found.

4.3.1 Theoretical solutions

Fig. 4.29 shows a shear span of a rectangular beam, subjected to the concentrated load P and the normal force N.

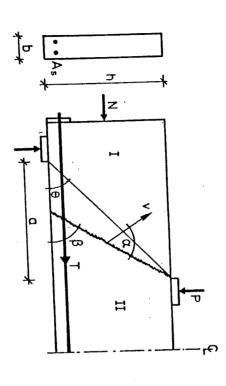


Fig. 4.29. Failure mechanism for a shear span with normal force.

The yield line pattern is the same as in 4.1.2.

The range of the variables α and β in this case is

and

(4.3.2)

(4.3.1)

The rate of external work is

$$W_{E} = Pvsin(\alpha+\beta) + Nvcos(\alpha+\beta)$$
 (4.3.3)

Here, N is the normal force, positive as a compressive force.

The rate of internal work dissipated by this failure mechanism is

$$W_{\rm I} = \frac{V}{2}(\lambda - \mu \sin \alpha) \frac{\text{bhf}_{\star}^{\star}}{\sin \beta} + \text{Tv} |\cos(\alpha + \beta)| \quad \text{for } \phi \le \alpha \le \pi - \phi \qquad (4.3.4)$$

$$W_{I} = \frac{V}{2}(1-\sin\alpha) \frac{bhf_{c}^{*}}{\sin\beta} + Tv|\cos(\alpha+\beta)| \quad \text{for } \alpha \le \varphi \text{ or } \alpha \ge \pi-\varphi \quad (4.3.5)$$

The work equation $W_{I} = W_{E}$ yields the upper bound

$$\frac{\tau}{f_{\rm C}^*} = \frac{\lambda - \mu \sin(\alpha + \beta) \cos\beta + (\mu - 2\chi) \sin\beta \cos(\beta + \alpha)}{2 \sin\beta \sin(\alpha + \beta)} \quad \text{for } \phi \le \alpha \le \pi - \phi \quad (4.3.6)$$

$$\frac{\tau}{f_c^*} = \frac{1-\sin(\alpha+\beta)\cos\beta+(1-2\chi)\sin\beta\cos(\beta+\alpha)}{2\sin\beta\sin(\alpha+\beta)} \quad \text{for } \alpha \le \phi \text{ or } \alpha \ge \pi-\phi$$

$$(4.3.7)$$

Here the parameter χ and the normal force degree $n^* = \frac{N}{bhf_C^*}$ has been introduced. The parameter χ is defined by

$$\chi = \begin{cases} \frac{T}{bhf_{C}^{*}} + \frac{N}{bhf_{C}^{*}} = \phi^{*} + n^{*} & \text{for } (\alpha + \beta) \ge \frac{\pi}{2} \\ \frac{N}{bhf_{C}^{*}} - \frac{T}{bh^{*}} = n^{*} - \phi^{*} & \text{for } (\alpha + \beta) \le \frac{\pi}{2} \end{cases}$$
(4.3.8)

In the case of $\beta>\theta$ and $\alpha>\varphi$, we may find the lowest upper bound by minimizing equation (4.3.6) with respect to α and β . Following the same procedure as in 4.1.2, we find the lowest upper bound

$$\frac{\tau_{*}}{f_{c}} = \frac{1}{2} \left(\lambda^{2} - \mu^{2} \right) (1 - D_{1}^{2})$$
 (4.3.10)

Here the parameter $\mathbf{D}_{\mathbf{I}}$ is defined as

$$D_{1} = \begin{cases} \frac{\mu - 2(n^{*} + \phi^{*})}{\lambda} & \text{for } n^{*} + \phi^{*} < \frac{\mu}{2} \\ 0 & \text{for } n^{*} + \phi^{*} \ge \frac{\mu}{2} \ge n^{*} - \phi^{*} \end{cases}$$
 (4.3.1)

The range of validity of equation (4.3.10) has been found to be

$$\frac{1-D_1^2}{h} > \frac{1-D_2^2}{\frac{2}{2-1}}$$
 (4.3.12)

$$n^{*}+\phi^{*}\geq -\rho^{*}$$
 (it means $D_{1}\leq 1.0$) (4.3.13)

and

$$\lambda + \mu + D_1 > 1 \cdot 0 \tag{4}$$

For the case of $\beta=\theta$ and $\alpha>\phi$, the lowest upper bound may be obtained by differentiation of equation (4.3.6) with respect to α and making β equal to θ . It yields

$$\frac{\tau_{\star}}{f_{c}^{\star}} = \frac{1}{2} \left[\lambda \left(\frac{a}{h} \right)^{2} + 1 - D_{1}^{2} - \mu_{h}^{a} \right]$$

(4.3.15)

which is valid for

$$\left\{\begin{array}{c} 1-D_1^2 \\ \eta^2-1 \end{array}\right.$$

(4.3.16)

and

$$\frac{a}{h} + \tan \theta - \frac{D_1}{\cos \theta}$$
 (4.3.18)

Using the condition $\beta=\theta$ and $\alpha=\varphi$, the lowest upper bound is easily found by inserting $\beta=\theta$ and $\alpha=\varphi$ into equation (4.3.7). It appears that

$$\frac{\tau}{f_{\rm C}^*} = \frac{\left[\left(\frac{a}{h}\right)^2 + 1\right] (1 - \sin\varphi) - 2\chi \left(\frac{a}{h}\cos\varphi - \sin\varphi\right)}{2\left(\frac{a}{h}\sin\varphi + \cos\varphi\right)}$$
(4.3.19)

In this case χ becomes

$$\chi = \begin{cases} n^{+} + \phi^{+} & \text{for } \frac{a}{h} \le \tan \phi \\ n^{+} + \phi^{+} & \text{for } \frac{a}{h} \ge \tan \phi \end{cases}$$
 (4.3.20)

The necessary condition for equation (4.3.19) to be applicable is

$$tan\varphi - \frac{c_1}{\cos\varphi} < \frac{a}{h} \le tan\varphi - \frac{D_1}{\cos\varphi}$$
 (4.3.22)

and the conditions (4.3.13) and (4.3.14). Here the parameter c_1 has been introduced as

$$C_{1} = \begin{cases} 1-2(n^{*}+\phi^{*}); & n^{*}+\phi^{*} < \frac{1}{2} \\ 0 & ; & n^{*}+\phi^{*} > \frac{1}{2} > n^{*}-\phi^{*} \end{cases}$$

$$\begin{cases} 1-2(n^{*}-\phi^{*}); & \frac{1}{2} < n^{*}-\phi^{*} \end{cases}$$

$$(4.3.23)$$

be found by minimizing equation (4.3.7) with respect to α and When we are in the case of $\beta=\theta$ and $\alpha<\varphi$, the lowest upper bound can keeping $\beta = \theta$. Following the same step as in 4.1.2, we get

$$\frac{\tau_{\star}}{f_{c}^{\star}} = \frac{1}{2} \left[\left[\left(\frac{a}{h} \right)^{2} + 1 - c_{1}^{2} - \frac{a}{h} \right] \right]$$
 (4.3.24)

which is identical to the expressions derived by Roikjær [78.3] using the so-called square yield locus as yield condition.

The condition for solution (4.3.24) to be valid is:

$$\frac{a}{h} \le \tan \varphi - \frac{C_1}{\cos \varphi} \tag{4.3.25}$$

n*-+* < 1.0

$$n^{+} + + \times - \rho^{+}$$
 (4.3.26)

In the situation when $\lambda + \mu + D_1 \le 1.0$, the lowest upper bound can only $\alpha = \varphi$. It has been found that differentiation of equation (4.3.7) with respect to eta and keeping be obtained with $lpha \ \leqq \ arphi$. We may find the lowest upper bound by

$$\frac{\tau}{f_c^*} = 2 \frac{[\sqrt{k} - \sqrt{1 + (k-1)} \xi] [\sqrt{k} (1 + (k-1) \xi] - 1]}{(k-1)^2}$$

(4.3.28)

Here the parameter ; is defined as

$$\xi = \begin{cases} n^{*} + \phi^{*}; & n^{*} + \phi^{*} < \frac{K-1}{4K} \\ \frac{K-1}{4K} & n^{*} + \phi^{*} \ge \frac{K-1}{4K} \ge n^{*} - \phi^{*} \end{cases}$$

$$(4.3.29)$$

$$n^{*} - \phi^{*} \frac{K-1}{4K} < n^{*} - \phi^{*}$$

The condition for solution (4.3.28) to be valid is

$$\frac{1}{h} > \frac{1}{\sin \varphi} [\sqrt{1 + (k-1)\chi} - \cos \varphi]$$
 (4.3.30)

$$n^{*} + \phi^{*} \geq -\rho^{*}$$
 (4.3.31)

and

$$n^* - \phi^* \le 1.0$$
 (4.3.32)

condition is changed into this situation, solution (4.3.19) is valid, but the necessary geometrical condition and it is necessary that it equals heta. In If equation (4.3.30) is not satisfied, eta is limited by the

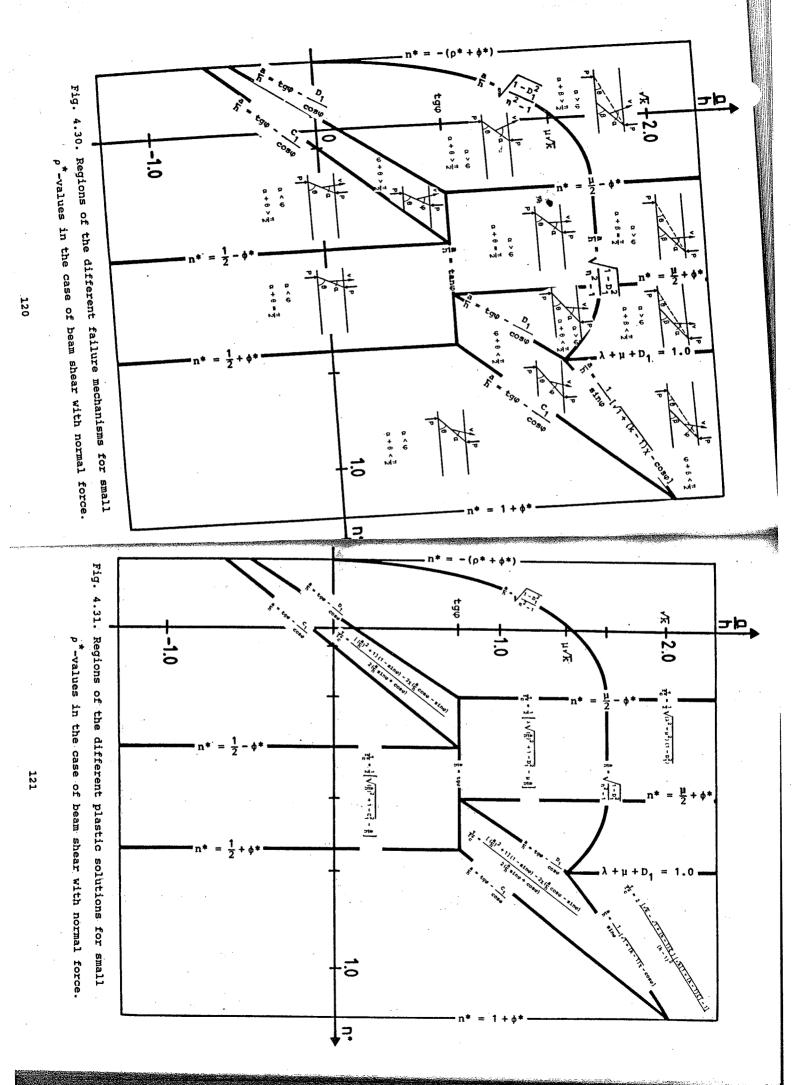
$$tan \varphi = \frac{c_1}{\cos \varphi} < \frac{a}{h} < \frac{1}{\sin \varphi} [\sqrt{1 + (k-1)\chi} - \cos \varphi]$$
 (4.3.33)

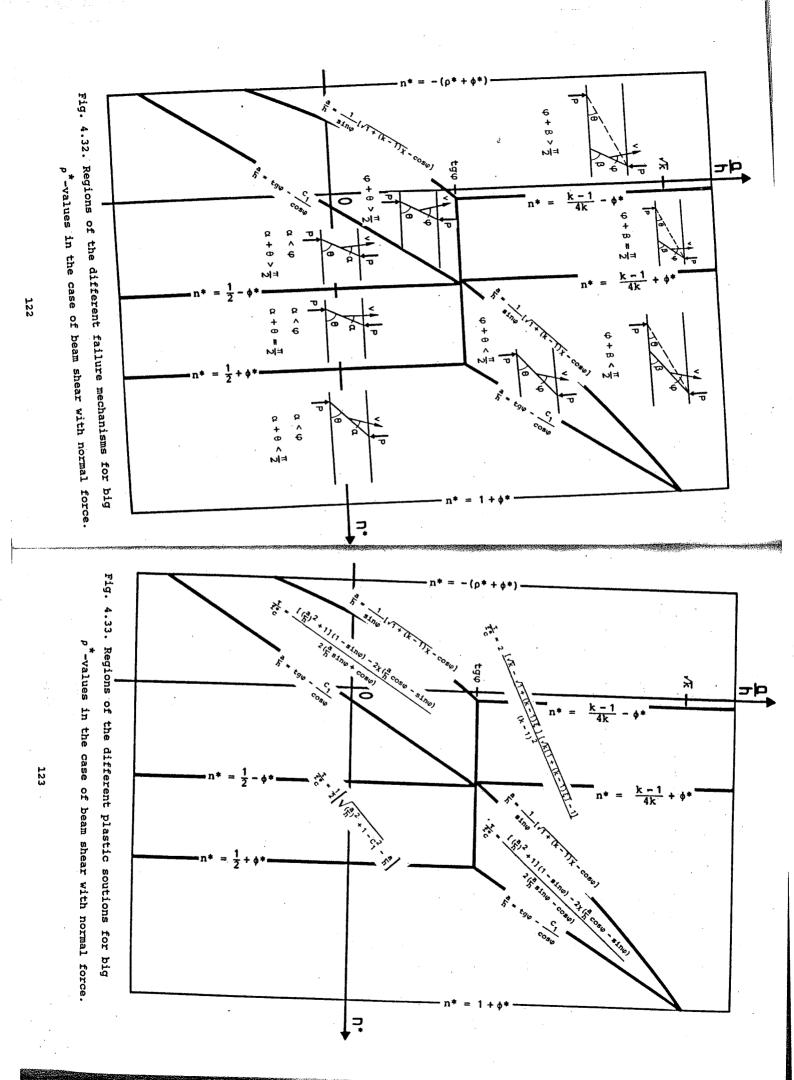
and equation (4.3.13).

validity of equations (4.3.25), (4.3.26) and (4.3.27) are still In the case of $\beta=\theta$ and lpha<arphi , the solution (4.3.24) and the range of

 $0.06 \text{ and } \phi^{2} = 0.15.$ $(n^*, \frac{a}{h})$ coordinate system. The boundaries are plotted for ho^* = plastic solutions are depicted in Fig. 4.30 and Fig. 4.31 in a ho * -values in the upper bound analysis and the corresponding The regions for the different failure mechanisms for small

are plotted for $\rho^* = 0.25$ and $\phi^* = 0.15$. solutions are depicted in Fig. 4.32 and Fig. 4.33. The boundaries The conditions for big ho^* -values and the corresponding plastic





The solutions are shown in Fig. 4.34, Fig. 4.35 and Fig. 4.36 for some values of the parameters. For low values of the normal force, the shear capacity increases with increasing normal force, while for high values of the normal force the shear capacity decreases with increasing normal force.

In Fig. 4.36 the influence of the ρ^* -values to the shear capacity has been shown for different $\frac{a}{h}$ and ϕ^* .

The difference between the case of considering the tensile strength of concrete or neglecting it can be seen very clearly in Fig. 4.37.

In the case of beam shear with normal tensile forces, the simple solutions ($\rho^*=0$) may underestimate the shear carrying capacity throughout all ranges of shear span ratio. In the case of big normal compressive forces, the simple solutions may overestimate the shear capacity in beams for large shear span ratios.

The flexural capacity of course might govern the load carrying capacity in some regions of the parameters. For the calculation of bending with normal force, see [84.1].

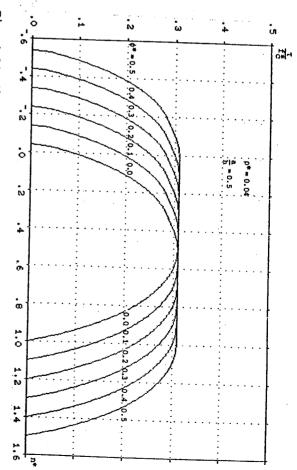


Fig. 4.34. Shear carrying capacity of beams subjected to concentrated loads and normal force versus the normal force degree for different longitudinal reinforcement degrees.

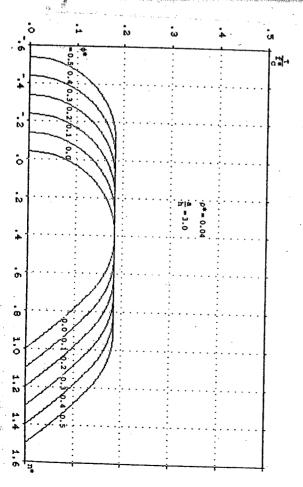
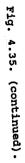


Fig. 4.34. (continued).



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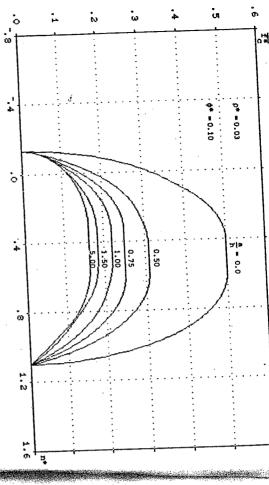
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0.75

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♦* ■ 0.2 h = 3.0 꺴

Fig. 4.35. Shear carrying capacity versus the normal force degree for different shear span ratios.

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p* = 0.03

1 0.0

 $\phi^* = 0.50$

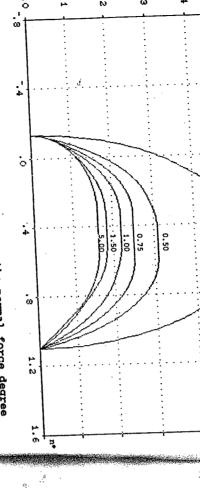


Fig. 4.36. The influence of ho^* -values to the shear capacity of beams subjected to combined shear forces and normal forces.

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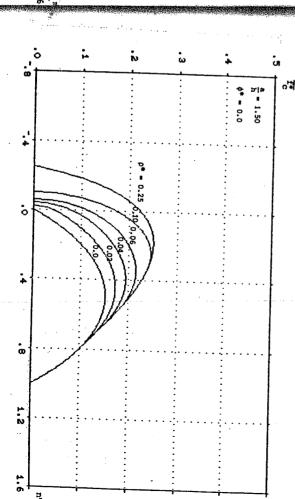
1.6

0.00

0.02

. .0.06 . . 0.04

0.10



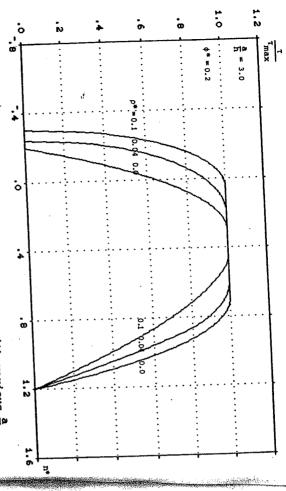


Fig. 4.37. The influence of ρ^* on shear capacity with various $\frac{a}{h}$ and ϕ^* .

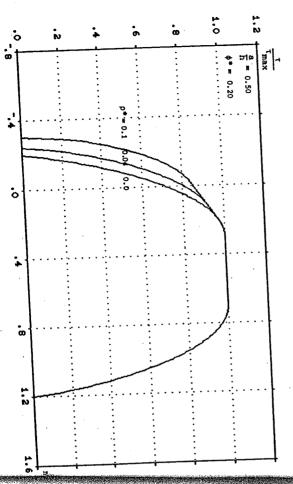


Fig. 4.37. (continued).

4.3.2_Experimental_verification

To test the validity of the complete set of equations in Fig. 4.31 for shear specimens with normal force, an investigation has been made for corbels with tensile normal force and beams with compressive normal force reported in the literature, [65.1], [60.1] and [57.1].

The detailed test data are shown in Appendix C.

A parametric analysis of 24 observed shear failures of corbels with normal force has given the following best empirical ν formula with $\rho^*=0.1$ as

$$v = 0.055(2.0-0.8 \frac{d}{h})(\varphi+4)(1-0.25 h) \le 1.0$$
 (4.3.34)

Here φ should be in percent and h in meter.

The mean value \bar{x} of the ratios of test to theory is 1.000, while the standard deviation σ is 0.076 and the coefficient of variation c_V is 0.076.

The comparison of test results with theory is shown on Fig. 4.38.

The very good agreement between tests and theory can also be seen from Fig. 4.39.

The complete solutions shown in Fig. 4.31 may be a little complicated for practical use. For the shear specimens with tensile normal force, we are forced to consider the tensile strength of concrete to get the agreement with tests, especially when the tensile normal force is bigger than the force that the reinforcement can take.

However, in some cases we may simplify the calculation according to the actual situation. For example we may only use equation (4.3.15) in stead of the complete solutions to calculate the shear capacity of corbels subjected to the combined shear and normal tensile force.

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becomes more easy and simple. The results is still quite satisfactory and the calculation deviation σ is 0.092 and the coefficient of variation C_V is 0.090. failures of corbels, the mean value $\bar{\mathbf{x}}$ is 1.014, while the standard Using (4.3.15) with $\rho^*=0.1$ and (4.3.34) for 24 observed shear

equation (4.3.24) for $\frac{a}{h} \neq 2.5$, just as we did in section 4.2. normal force, we may simplify the complete solutions by modifying For the specimens subjected to combined shear and compressive

frames from 2 reports in the literature, [60.1], [57.1]. parametric analysis of 47 stub beams with axial loads and knee In this case the corresponding ν formula has been found by the

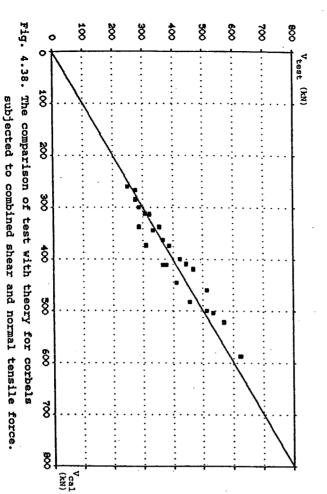
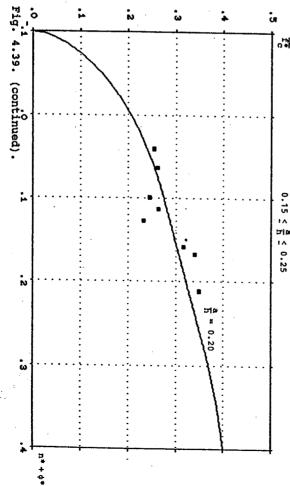


Fig. 4.39. The plastic solutions for beam shear with normal force ò ŗ 'n **. ;** Ġ chi $0.045 \le \frac{8}{h} \le 0.07$ 'n 90.0 n* + o*

compared with corbel test results.



$$= \frac{0.58(\varphi+2)(2-0.4\frac{a}{h})(1-0.25h)(1+0.75\frac{N}{A_{S}f_{Y}})}{\sqrt{f_{C}}}$$
 (4.3.35)

 $\left(\frac{a}{h} \neq 2.5, \varphi \neq 2.0 \text{ and } h \neq 1.0\right]$

As before f_{C} is in MPa, φ in percent and h in meter.

coefficient of variation c_V is 0.095. theory is 0.997, while the standard deviation σ is 0.094 and the The mean value \bar{x} of the ratios of 47 observed shear tests to

The comparison of test results with theory is shown on Fig. 4.40.

been calculated from (4.3.35). and knee frames for different shear span ratios. The v value has shows the results of shear tests of stub beams with axial loads To demonstrate the general applicability of the theory, Fig.4.41

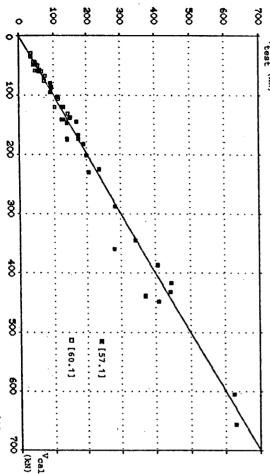


Fig. 4.40. The comparison of tests with theory of stub bemas with axial loads and knee frames.

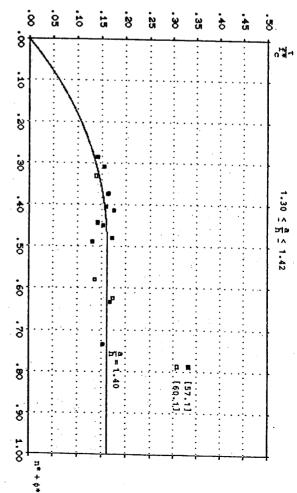


Fig. 4.41. Theoretical shear capacity compared with test results of stub beams with axial loads and knee frames.

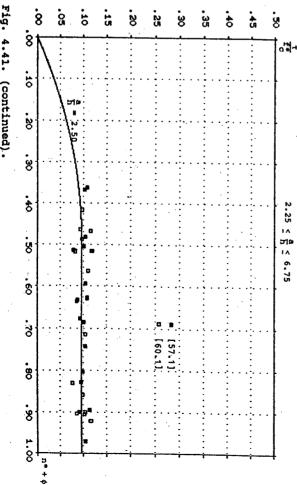


Fig. 4.41. (continued).

4.4. The influence of prestress to shear capacity of beams

The problem how prestress influences the shear capacity of concrete beams without web reinforcement has only been very preliminarily studied. According to plastic theory, prestress does not influence the theoretical shear solution in any way. The reason is that the prestress introduces an internal, self-equilibrated stress system, which does not affect either the failure mechanism, nor the stress distribution at collapse.

However, the existence of an initial, non-zero stress state does influence the amount of stress redistribution which must take place before failure, and this has an effect on the ultimate load.

In the following, we will use the complete plastic solutions in Fig. 4.10 and we will try to modify the ν expression for prestressed beams.

The effect of prestress seems to be that it increases the effectiveness factor with increasing level of effective prestress.

An investigation how the prestress influences the effectiveness factor ν has been made for 286 prestressed beams without web reinforcement reported in the literature [58.2], [59.2], [60.3], [60.9], [63.3], [65.5] and [69.4].

The detailed test data are shown in Appendix D.

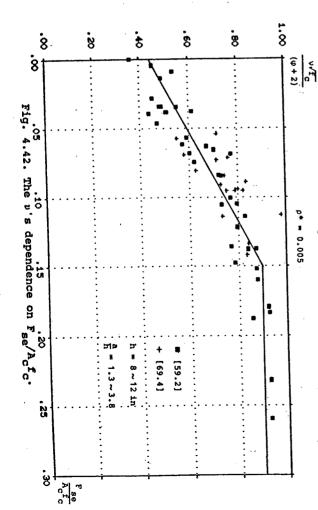
A treatment similar to that reported in 4.1.4 has shown that the effective prestress in the section increases the ν value and that prestressed beams do not exhibit an a/h dependence to a degree that has been found for non-prestressed beams. In fact for a carefully chosen ρ^* -value, we may consider the influence of the shear span to the ν value for prestressed beams to be non-existent. A proper ρ^* value is found to be 0.003 _ 0.005 for prestressed rectangular beams reported in [59.2].

Fig. 4.42 shows an example of the dependence of ν on $F_{\rm ge}/A_{\rm c}f_{\rm c}$, which is the average effective prestress ratio in the section of the beams. Here $A_{\rm c}$ is the area of the section and $F_{\rm ge}$ is the effective prestress acting on the section. Similar to the ν - ψ relation shown in Fig. 4.42, the ν value is increasing almost linearly with increasing the average effective prestress ratio $F_{\rm ge}/A_{\rm c}f_{\rm c}$. Beyond a certain limit of $F_{\rm ge}/A_{\rm c}f_{\rm c}$ the ν value seems to be independent of $F_{\rm ge}/A_{\rm c}f_{\rm c}$. This limit may be chosen as 0.15 for simplicity.

A parametric analysis of 54 observed shear failures of prestressed beams with rectangular section [59.2] and [69.4] has given the following empirical ν formula assuming $\rho^*=0.003$

$$0.575(\varphi+1.5)(1-0.2h)\begin{bmatrix} 1+8 & \frac{F_{Se}}{A_{C}f_{C}} \\ h + 1.0 \\ \frac{F_{Se}}{A_{C}f_{C}} + 0.15 \end{bmatrix}$$

Here φ should be in percent, f_C in MPa and h in meter.



 $c_{
m V}$ is 0.088. The comparison of tests with theory is shown in Fig. the standard deviation σ is 0.089 and the coefficient of variation The mean value \bar{x} of the ratios of test to theory is 1.007, while

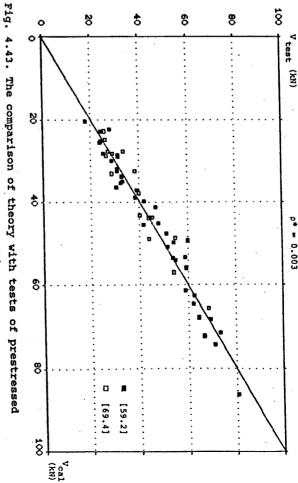
The good agreement between tests and theory can also be seen from

formulae (4.2.15) and (4.2.16). The corresponding ν formula then tensile strength of concrete, i.e. let $\rho^*=0$ and use the simple For practical use, we may simplify our theory by neglecting the

$$\frac{0.52(\varphi+2)(1-0.25h)\left[1+6\frac{F_{Se}}{A_{C}f_{C}}\right]}{\sqrt{f_{C}}} \qquad \begin{pmatrix} \varphi \ * \ 2.0 \\ h \ * \ 1.0 \\ \frac{F_{Se}}{A_{C}f_{C}} \ * \ 0.15 \end{pmatrix}$$
(4.4.2)

Units are as before.

also be seen from Fig. 4.46. shown in Fig. 4.45. The agreement between tests and theory can of variation $c_{
m V}$ is 0.132. The comparison of tests to theory is 1.007, while the standard deviation σ is 0.133 and the coefficient In this case the mean value $\bar{\mathbf{x}}$ of the ratios of test to theory is



beams with rectangular section.

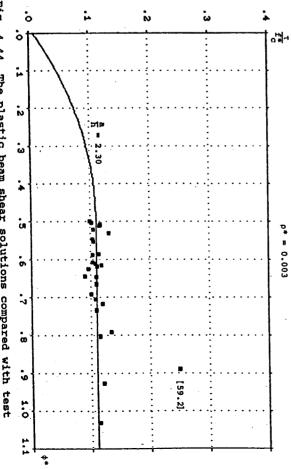


Fig. 4.44. The plastic beam shear solutions compared with test results of prestressed beam with rectangular section.

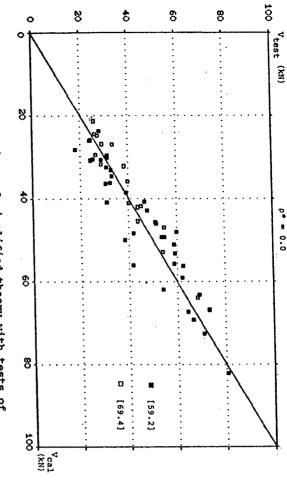


Fig. 4.45. The comparison of simplified theory with tests of prestressed beams.

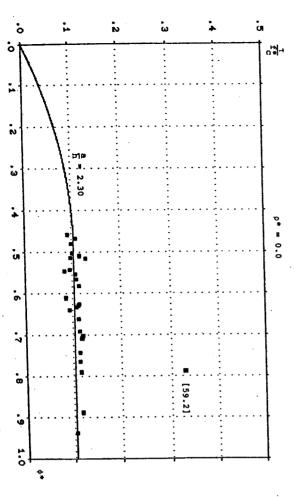


Fig. 4.46. The simplified plastic solutions compared with tests.

In practice, a lot of prestressed beams are with T cross section or I cross section. The considerable influence of tension or compression flanges on the maximum shear capacity of beams without shear reinforcement has been studied in [78.3]. The upper bound technique for such beams is somewhat complicated. An alternative method has been developed in this section. We may employ the complete plastic shear solutions shown in Fig. 4.10 for the web and take account of the influence of the flanges in the ν formula.

A parametric analysis of 286 observed shear failure of prestressed beams without web reinforcement has shown a relatively simple and quite good empirical v formula assuming $\rho^{\frac{1}{n}}=0.007$

$$= \frac{0.445(\varphi+2)\left[1+8\frac{F_{Se}}{A_{C}f_{C}}\right]\left[\frac{A_{C}}{bh}\right]}{\left[\frac{F_{Se}}{A_{C}f_{C}} + 0.10\right]} \frac{F_{Se}}{A_{C}f_{C}} + 0.10$$

$$\left[\frac{A_{C}}{bh} + 1.8\right]$$
(4.4.3)

Units are as before.

The mean value \bar{x} of the 286 ratios of test to theory is 1.000, while the standard deviation σ and the coefficient of variation C_V are 0.166. The comparison of tests with theory is shown in Fig. 4.47.

In practice, it's more convenient to use the modified simple solutions (4.2.17) to calculate the shear capacity of prestressed beams. The corresponding ν formula can be taken as

$$v = \frac{0.505(\varphi+2)\left[1+6\frac{F_{Se}}{A_{C}f_{C}}\right]\left[\frac{A_{C}}{bh}\right]}{\sqrt[4]{f_{C}}} \left[\frac{A_{C}}{bh} + 1.8\right]} \left[\frac{A_{C}}{bh} + 1.8\right] (4.4.4)$$

Using formulae (4.2.17) and (4.4.4), the mean value \bar{x} of the 286 ratios of the test to theory is 1.004, while the standard deviation is 0.188 and the coefficient of variation C_V is 0.187. The comparison of tests with theory is depicted in Fig. 4.48.

The simplified plastic solution (4.2.15) and (4.2.16) may also be used for prestressed beams. In this case we may use the corresponding ν formula as

$$0.655(\varphi \pm 1.5) \left[1 + 6 \frac{F_{\text{Se}}}{A_{\text{C}} f_{\text{C}}} \right] \left[\frac{A_{\text{C}}}{bh} \right] \left[\frac{A_{\text{C}}}{bh} + 1.8 \right]$$

$$\left[\frac{F_{\text{Se}}}{A_{\text{C}} f_{\text{C}}} + 0.12 \right]$$

$$\left[\frac{A_{\text{C}}}{A_{\text{C}} f_{\text{C}}} + 0.12 \right]$$

The calculation using formulae (4.2.15), (4.2.16) and (4.4.5) has been carried out for 286 prestressed beams. The mean value \bar{x} is 1.005, while the standard deviation σ is 0.200 and the coefficient of variation C_V is 0.199. Fig. 4.49 shows the comparison of theory with tests.

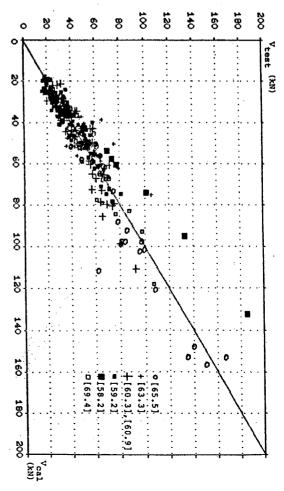


Fig. 4.47. The comparison of theory with tests of 286 prestressed beams without web reinforcement.

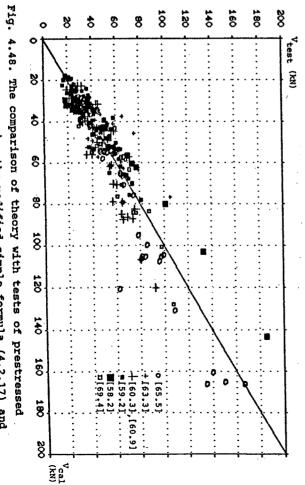


Fig. 4.48. The comparison of theory with tests of prestressed beams using the modified simple formula (4.2.17) and corresponding ν formula (4.4.4).

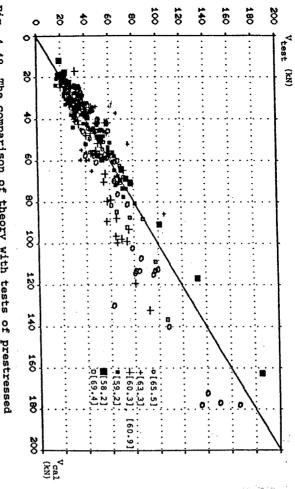


Fig. 4.49. The comparison of theory with tests of prestressed beams using the simplified formulae (4.2.15), (4.2.16) and the corresponding v formula (4.4.5).

4.5. Shear capacity of joints

The complete plastic shear solutions listed in Fig. 4.8 can be easily extended to solve the shear problem of concrete joints, which are shown in Fig. 4.50.

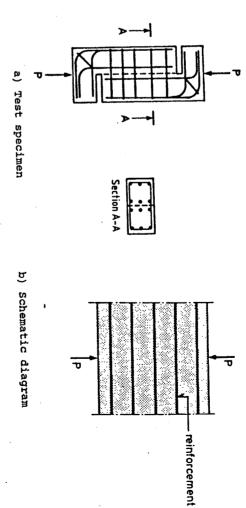


Fig. 4.50. Joints of monolithic concrete.

Inserting $\frac{a}{h}=0$ into the relevant formulae along the line $\frac{a}{h}=0$ in Fig. 4.10, we get the shear carrying capacity of reinforced monolithic concrete as follows:

$$\frac{\frac{1}{2} \cdot \lambda \sqrt{1-D^2}}{\frac{1-\sin\varphi}{2 \cos\varphi} + \varphi^* \operatorname{tg}\varphi} \qquad \frac{\psi^* \leq \frac{\mu-\lambda \sin\varphi}{2}}{\frac{1-\sin\varphi}{2}} \qquad (4.5.1)$$

$$\frac{\frac{1}{f_*^*}}{\frac{1}{f_*^*}} = \begin{cases} \frac{1-\sin\varphi}{2} + \varphi^* \operatorname{tg}\varphi & \frac{\mu-\lambda \sin\varphi}{2} < \varphi^* \leq \frac{1-\sin\varphi}{2} & (4.5.2) \\ \frac{1}{2}\sqrt{1-C^2} & \frac{1-\sin\varphi}{2} < \varphi^* \leq \frac{1}{2} & (4.5.3) \\ \frac{1}{2}\sqrt{1-C^2} & \frac{1}{2} < \varphi^* & (4.5.4) \end{cases}$$

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Inserting equations (4.1.31), (4.1.32), (4.1.37) and (4.1.40) into equations (4.5.1) - (4.5.4), we may rewrite the shear capacity into

$$\frac{\tau}{f_{C}^{*}} = \begin{cases} \sqrt{(\rho^{*} + \phi^{*})[(1 - 3\rho^{*}) - (\rho^{*} + \phi^{*})]} & \phi^{*} \leq 0.2 - 1.6\rho^{*} & (4.5.5) \\ 0.25 + 0.75\phi^{*} & 0.2 - 1.6\rho^{*} < \phi^{*} \leq 0.2 & (4.5.6) \\ \sqrt{\phi^{*}(1 - \phi^{*})} & 0.2 < \phi^{*} \leq 0.5 & (4.5.7) \\ 0.5 & 0.5 < \phi^{*} & (4.5.8) \end{cases}$$

Solutions (4.5.5) - (4.5.8) are identical with the solutions given firstly by B.C. Jensen [77.2].

The shear carrying capacity of joints versus effective reinforcement degree for some ρ^* -values is depicted in Fig. 4.51.

It is obvious that the tensile strength of concrete plays an important part in shear carrying capacity when the reinforcement degree is rather low.

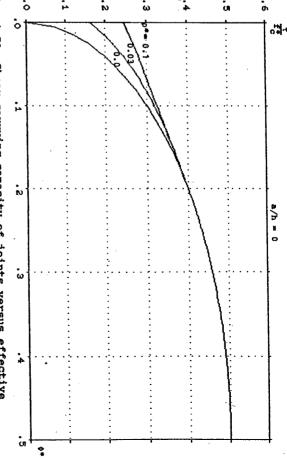


Fig. 4.51. Shear carrying capacity of joints versus effective reinforcement degree for some ρ^* -values.

The theory with $\rho^*=0.035$ and $\nu=0.665$ has been compared with joint tests of monolithic concrete performed by Hofbeck et al. [69.1]. The agreement between the formulae for the carrying capacity and the test results is exceptionally good.

The comparison between theory and test results is depicted in Fig. 4.5.2. To demonstrate the applicability of the complete plastic shear solutions to the monolithic concrete joints, Fig. 4.53 shows how better the test results fit the theoretical curves.

In this chapter we will only treat beams with vertical stirrups and subjected to point loading. Besides the assumptions that we have made in section 2.6, we further assume:

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V_{test} (kips)

[69.1]

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ν = 0.665

 $\rho^* = 0.035$

The web reinforcements is rigid, perfectly plastic. The stirrup spacing is sufficiently small to permit a continuous distribution of the stirrup forces. The cross-sectional steel area per unit length in both directions is $^{\rm A}_{\rm SWh}$ and $^{\rm A}_{\rm SWV}$, respectively. The yield stress of the web reinforcement is $^{\rm f}_{\rm yw}$.

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Fig. 5.1 shows a shear span of a concrete beam subjected to the shear force V. The breadth and depth of the beam are termed b and h^{\star} , respectively. The failure mechanism is assumed to be consisting of a single yield line inclined at angle β to the beam axis.

The relative displacement rate is ν inclined at the angle α to the beam normal.

Fig. 4.52. The comparison between theoretical calculations and

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. n = 15.

 $C_{\mathbf{V}} = 0.037$

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For the beams with web reinforcement, it is reasonable to neglect the tensile strength of concrete.

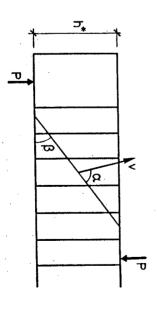
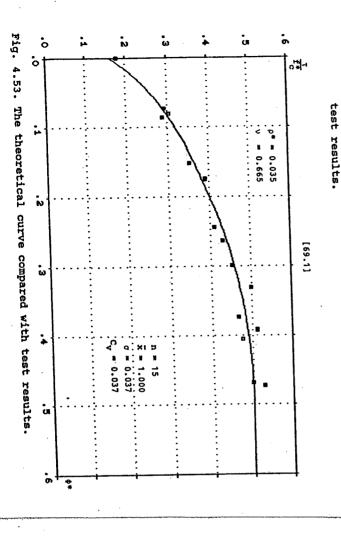


Fig. 5.1. Failure mechanism of beam with vertical stirrups subjected to point loading.

On the above assumptions, the following upper bound solutions were



are still missing the lower bound solutions [79.1]. are backed by coincident lower bound solutions, but some of them derived by Nielsen and Bræstrup [76.7]. Most of these upper bounds

5.1 Theoretical solutions

The complete lowest upper bounds can be summarized as following:

\$	$\phi_{VO}^* \le \phi_{V}^* \le \frac{1}{2}$	ф* Ч^ф* V0	φ _V γ/f _c φ _h *
$\sqrt{\phi_h^*(1-\phi_h^*)}$	$2\sqrt{\phi_{\rm h}^*(1-\phi_{\rm h}^*)\phi_{\rm v}^*(1-\phi_{\rm v}^*)}$	$\frac{1}{2} \left[\left(\left(\frac{a}{h^*} \right)^2 + 4\phi_h^* (1 - \phi_h^*) - \frac{a}{h^*} \right) + \phi_V^* \cdot \frac{a}{h^*} \right]$	**
ЫH	$\sqrt{\frac{\phi_{\mathrm{V}}^{*}(1-\phi_{\mathrm{V}}^{*})}{}}$	$\frac{1}{2} \left[\left(\left(\frac{a}{h^*} \right)^2 + 1 - \frac{a}{h^*} \right) + \phi_V^* \cdot \frac{a}{h^*} \right]$	ው ከ * ። 2] ተ

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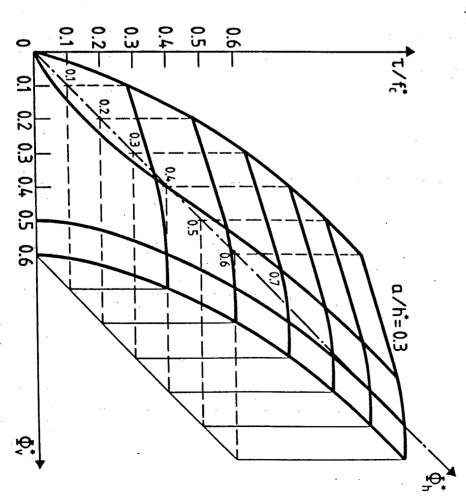
$$\phi_{\text{VO}}^{*} = \begin{cases} \frac{\left(\frac{a}{h^{*}}\right)^{2} + 4\phi_{\text{h}}^{*}(1 - \phi_{\text{h}}^{*}) - \frac{a}{h^{*}}}{2\left(\frac{a}{h^{*}}\right)^{2} + 4\phi_{\text{h}}^{*}(1 - \phi_{\text{h}}^{*})} & \text{valid for } \phi_{\text{h}}^{*} \leq \frac{1}{2} \\ \frac{\left(\frac{a}{h^{*}}\right)^{2} + 1 - \frac{a}{h^{*}}}{\left(\frac{a}{h^{*}}\right)^{2} + 1} & \text{for } \phi_{\text{h}}^{*} > \frac{1}{2} \end{cases}$$

$$(5.2)$$

(5.3)

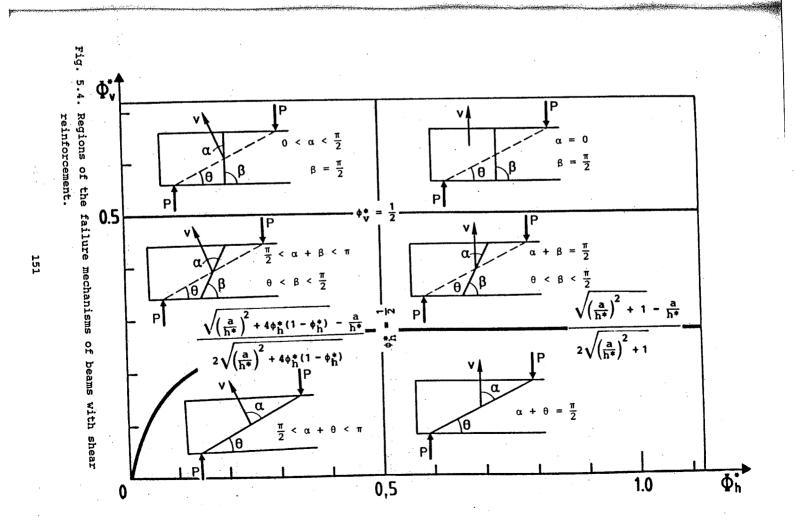
$$\phi_{V}^{*} = \frac{A_{SWV}f_{VW}}{bf_{C}^{*}}$$

 $a/h^* = 0.3.$ solutions (5.1) is also visualized in Fig. 5.2 for deep beams with



(5.1)

Fig. 5.2. Shear capacity of beams with vertical stirrups versus the reinforcement degree in both directions.



 $\sqrt{\left(\frac{a}{h^*}\right)^2 + 1} - \frac{a}{h^*}$

 Φ_h

 $2\sqrt{\left(\frac{a}{h^*}\right)^2+1}$

The regions of complete upper bound solutions and the corresponding failure mechanisms are illustrated in Fig. 5.3 and Fig. 5.4, respectively. The boundaries are plotted for $a/h^*=0.5$.

5.2 The effective shear depth h^{\star} and the effectiveness factor ν

In order to be able to use the theoretical solutions (5.1) in practice we need to assess the values of the effectiveness factor ν and the effective shear depth h*. These two parameters are interrelated in the sense that an assessment of the effectiveness factor ν from experimental results of shear tests is influenced by the choice of h* and vice versa.

In our plastic analysis, the effective shear depth h is defined as the distance between the tension and compression stringers. It's rational to consider the tensile reinforcement as the tension stringer. Concerning the compression stringer, things become more complicated and vague. For T-beams, we may take the compression stringer as being identical to the compression flange. If for conventional rectangular beams, the compression zone is choiced, iteration must be used in the analysis of existing tests.

er possibility of defining the effective shear depth h*. The usual practice and standards are to relate the nominal shear stress to the effective depth of the beam. The Danish Code of Practice requires the use of the internal moment lever arm Z calculated at the section of maximum moment. It has been found that the internal moment lever arm may not be suitable for the effective shear depth h* in the analysis of existing shear tests of rectangular beams with both strong tensile reinforcement and shear reinforcement. A more rational way may be to define the effective shear depth h* as the stirrup depth d_S, as illustrated in Fig. 5.5. Unfortunately, up to now we cannot propose a rational and convenient way to de-

cide h^{\star} for rectangular beams, since the experimental documentation is lacking.

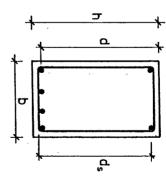


Fig. 5.5. Illustration of effective shear depth h .

with the view of combining the physics of our theory with practice in the analysis of test results described below, the effective shear depth h^{\star} is defined as the distance measured from the centroid of the tensile reinforcement to the centre of the compressive flange for T-section beams and as the 0.85 d for the rectangular beams.

Having decided upon h*, we can proceed to determine v. Towards this end we need test results covering a wide range of various parameters such as concrete strengths from weak to strong, both tensile reinforcement degrees and shear reinforcement degrees from lower to higher, shear span ratio from small to large and so on.

Unfortunately, not very many available web failure tests have been found in the literature.

In order to assess the effectiveness factor v from the existing experimental results of shear tests, inserting $f_{\rm C}^* = v f_{\rm C}$ into the complete plastic equations (5.1) and solving them with respect to v, we get the v formulae and the corresponding regions as shown in Fig. 5.6. The boundaries are plotted for $a/h^* = 0.5$.

5.3 Experimental verification of the theory

5.3.1 Conventional thin-webbed beams

In order to verify the plastic solutions for the web failure mechanism and to investigate the dependence of the web effectiveness factor ν on the concrete strength, the width of the web and the reinforcement details, (i.e. the number and dimensions of the longitudinal bars, their distances from the edges of the beam, and whether they are supported by stirrups), several series of T-beams with shear reinforcement had been tested at the Structural Research Laboratory of the Technical University of Denmark since 1973, [76.6], [80.2].

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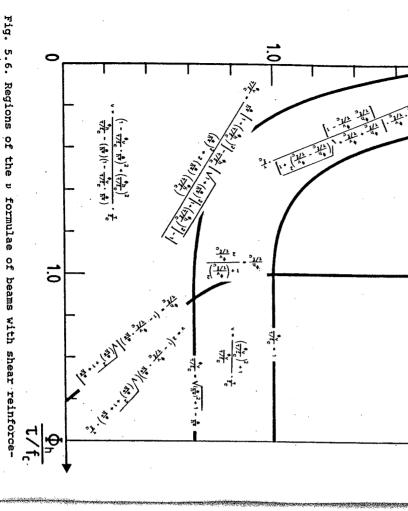
The test data are summarized in Appendix E.

The results are generally in good agreement with the web crushing criterion (Eq. $(5\cdot1)$). Because all these tested beams are with the same shear span ratio, it's impossible to investigate whether the web effectiveness factor ν is dependent or independent upon the shear span ratio from these tests.

The concrete strengths of most beams in these tests are between 5 $_{\rm 25~MPa}$, only very few beams are with concrete strength around 35 MPa, wherefore even using the constant ν value the coefficient of variation $C_{\rm V}$ is nearly the same as the case considering the ν dependence on $f_{\rm C}$.

For 106 observed shear failures, the best correspondance between test results and the theory are obtained for v=0.725, when the shear effective depth is defined as in 5.2. The mean value \bar{x} is 1.001, while the standard deviation σ is 0.120 and the coefficient of variation C_V is 0.120.

The tendency of ν to decrease with increasing concrete strength has been found even though it's not very clear since the tests



didn't cover a wide range of concrete strength. A linear relation-ship

$$v = 0.79 - \frac{f_{\rm c}({\rm MPa})}{200}$$
 (5.4)

may be used when h is defined as in 5.2.

If the experimental results are compared with this model, the correspondance is reasonably good, (see Fig. 5.7).

The mean value is 1.003, while both the standard deviation σ and the coefficient of variation $c_{_{f V}}$ are 0.113.

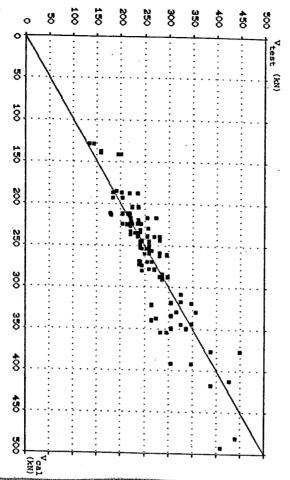
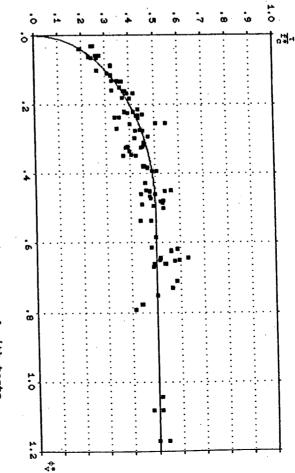


Fig. 5.7. The comparison of theory with tests of T-beams with shear reinforcement and subjected to point loads.

To demonstrate the general applicability of the web crushing criterion (Eq. (5.1)), Fig. 5.8 shows how the test results fit the theoretical curve.



rig. 5.8. Theoretical shear capacity compared with tests.

crete strength $f_c > 20$ MPa when h shear in beams without shear reinforcement have shown that the dependence But, if the whole beam depth h is used instead of h^{π} , just like we above that such model may underestimate the v-value for the con-However, it has been found from the shear test series mentioned variation of v may be described by an inverse square root depen-Results from previous applications of the plasticity theory to between test results and the theory is obtained for tests. For all these series of tests, do for beams without shear reinforcement, this inverse square root dence $v = k_1/\sqrt{f_C}$ where k_1 is a constant (see Fig. 4.16 to 4.18). model may also give a satisfactory agreement with is defined as in section 5.2. the best correspondance

$$v = 1.86/\sqrt{f_G}$$
 (f_G in MPa) (5.5)

With ν according to (5.5), and comparing tests with theory (5.1), we find the mean value is 1.001, while both the standard deviation σ and the coefficient of variation C_V are 0.107. It's a little better than using the ν model (5.4).

To establish a final conclusion, more test data which cover a wide range of concrete strength are needed.

5.3.2 Deep_beams_with_web_reinforcement

Here only 35 deep beam tests reported by Prof. F.K. Kong et al [70.2] are analyzed. These tests consist of seven different types of web reinforcement. The detailed test data are shown in Appendix F. Twenty beams subdivided into 4 types, have only the horizontal web reinforcement. For such deep beams, which have no vertical web reinforcement, the failure mechanism shown in Fig. 4.7 and the corresponding complete theoretical solutions presented in Fig. 4.8 together with the empirical v equation (4.1.123) are available. The comparisons of test to theoretical calculation prove our confidence. A rather good agreement has been found in the statistics of ratios of test to theory.

If we exclude the horizontal web reinforcement in calculation of the mechanical reinforcement degree ϕ defined by

$$\phi = \frac{A_{\rm S} f_{\rm Y}}{b h f_{\rm C}} \tag{5.6}$$

for 20 beams, the mean value is 1.051, while the standard deviation and the coefficient of variation $C_{\mathbf{V}}$ are 0.113 and 0.108, respectively. The comparison of tests with calculated results is shown in Fig. 5.9.

However, if we include the horizontal web reinforcement in the reinforcement degree ϕ defined by

$$\Rightarrow = \frac{\text{Asfy}}{\text{bhfc}} + \frac{\text{Aswh} \cdot \text{fywh}}{\text{bfc}}$$
 (5.7)

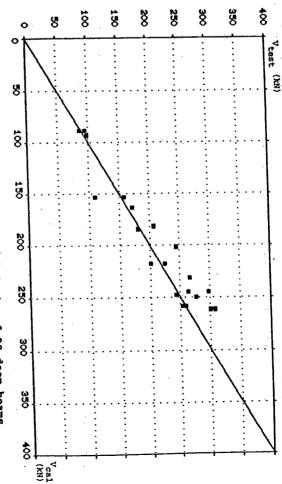


Fig. 5.9. The comparison of theory with tests of 20 deep beams with only horizontal web reinforcement.

where \mathbf{A}_{SWh} and \mathbf{f}_{YWh} are the area per unit length and the yield respectively. We note no considerable difference between the two and the coefficient of variation $C_{
m V}$ are 0.107 and 0.114, case, the mean value $\bar{\mathbf{x}}$ is 0.936, while the standard deviation σ calculated results are a little higher than the tests. In this strength of horizontal web reinforcement, respectively, then the to use the former way to calculate the shear carrying capacity of decreasing the shear crack width in the web). Until more dicisive conservative and the latter is too unsafe. It may be that the statistics, except the mean value. The former calculation is too experimental evidence is available, for the safe reason, we prefer as 'the main reinforcement, horizontal web reinforcement is not so effective on shear capacity in the calculation of the reinforcement degree. such deep beams, i.e. excluding the horizontal web reinforcement (of course it may be useful in

For the other 15 deep beams with vertical web reinforcement, which consist of 3 types of vertical web reinforcement arrangements, a

statistical analysis has been shown that the best agreement between theory and test may be reached by using

$$= \frac{2.35}{\sqrt{f_C}} \qquad (f_C \text{ in MPa}) \qquad (5.8)$$

when the whole depth of the deep beam is used as the shear depth. The mean value is 0.999, while both the standard deviation σ and the coefficient of variation C_V are 0.100.

The comparison of theoretical calculations with tests is shown in Fig. 5.1. In these tests the concrete strength only varied in a small range, whereby the ν dependence on f_C is also not quite clear.

If we put the effective shear depth of rectangular beams equal to 0.85 d, the statistics for different v expressions can be seen in Table 5.3.1.

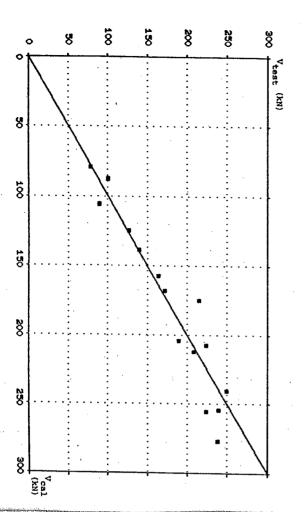


Fig. 5.10. The plastic solution for beam shear with web reinforcement compared with deep beam test results.

—г		r	 1
cv	a	ΧI	e .
0.145	0.146	1.006	0.72
Ó.142	0.146	1.030	0.8 - [£] C
0.139	0.139	1.002	रही इ.उ.

Table 5.3.1. Statistics of the ratios of tests to theoretical calculations for different v expressions.

It seems that the difference can hardly be said to be striking. However, in order to be unique and identical with the T-beams described in 5.3.1, equation (5.4) appears preferable to the inverse square root model $v=k_1/\sqrt{f_C}$, where k_1 is a constant, since k_1 , in these two cases deviate too much from each other.

5.3.3 Prestressed_beams_with_web_reinforcement

As we have mentioned in Section 4.4, according to the theory of plasticity, prestressing should have no effect on the shear carry-ing capacity of a perfectly plastic structure. However, for concrete beams, the existence of prestressing does influence the stress redistribution before failure and results in the change of the ultimate load.

How the prestress influences the effectiveness factor ν of concrete beams with web reinforcement has only been very preliminarily investigated for 93 observed shear failure tests of prestressed beams reported in the literature [64.2], [71.2], [73.3], [73,4], and [76.3].

The detailed test data are shown in Appendix G.

The experimental investigation of the shear capacity of prestressed concrete beams with shear reinforcement has shown that the effective prestress in the section does increase the ν value but the dependence of ν on the average effective prestress in the section is smaller than for beams without shear reinforcement. For the available test results, defining the effective shear depth h as in Section 5.2, the best correspondence between tests and theoretical calculations is obtained for

$$v = \left[0.8 - \frac{f_{\rm C}}{200}\right] \left[1 + 2.2 \frac{f_{\rm Se}}{A_{\rm C} f_{\rm C}}\right] \qquad \left(\frac{f_{\rm Se}}{A_{\rm C} f_{\rm C}} \neq 0.5\right]$$
 (5.9)

The mean value \bar{x} is 1.002, while both the standard deviation σ and the coefficient of variation C_V are 0.151.

The comparison of tests with calculations is shown in Fig. 5.11.

The applicability of the plastic solutions for beams with shear reinforcement can also be seen in Fig. 5.12, although the scatter is considereable. However, it is believed that as long as the effectiveness factor ν is only considered to be depending on the concrete compressive strength and the average effective prestress on the section, it will be difficult to obtain better results from any theory.

If the inverse square root model of concrete strength is considered in the ν expression, the best correspondence between tests and theory is found to be

$$v = \frac{3 \cdot 8}{\sqrt{f_C}} \left[1 + 2 \frac{F_{Se}}{A_C f_C} \right] \qquad \left[\frac{F_{Se}}{A_C f_C} \right\} \text{ 0.5}$$
 (5.10)

Now the statistical results are: mean value \bar{x} is 1.006, and the standard deviation σ and the coefficient of variation C_V are 0.159 and 0.158, separately.

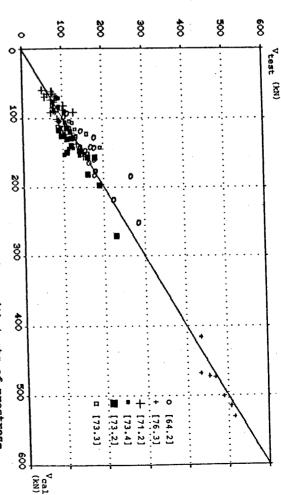


Fig. 5.11. The comparison of calculations with tests of prestressed beams with web reinforcement.

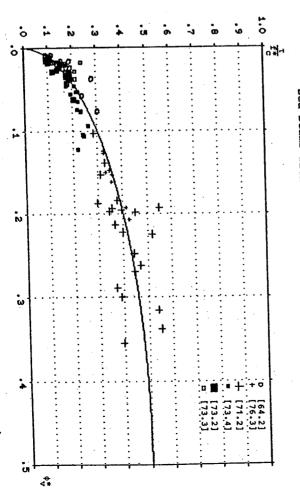


Fig. 5.12. Theoretical shear capacity for beams with shear reinforcement compared with tests of prestressed beams.

5.4 Some discussion on v

In Chapter IV it is shown that the main parameters governing the ν -value for ordinary beams with rectangular section and without shear reinforcement are the concrete strength, the amount of longitudinal reinforcement and the shear span/depth-ratio. Besides the above major factors, the scale effect, represented by the absolute value of the depth, also gives a smaller effect on the

The dependence on the concrete strength is easily understood since increasing concrete strength leads to decreasing ductility.

The dependence on the reinforcement ratio is probably due to the dowel action of the reinforcement, which is neglected in the plastic theory model according to our assumptions in Section 2.6, and to the fact, that increasing reinforcement ratios leads to decreasing crack widths at rupture.

The decreasing ν -value for increasing a/h-values up to about 2.5 is probably due to the dowel action and to the fact that the required stress redistribution is strongly increasing for increasing a/h-values in this interval.

The slight decreasing v-value for increasing depth of the section may be explained as a result of an unstable crack propagation at a lower load if the crack is longer, therefore the shear compression capacity is not proportional to the depth of the section but less.

For prestressed beams with rectangular section and without shear reinforcement, the main parameters governing the ν -value are the same as for ordinary beams except that the shear span/depth-ratio has only a small influence and that the average effective prestress on the section has an influence.

The experimental investigation has shown the ν -value is increasing almost proportional to the average effective prestressing of the section up to about 0.15. This phenomenon is probably due to the

fact that the prestressing effectively delays the formation and prevents the propagation of the inclined shear cracks. Therefore prestressing increases the depth of shear compressive zone of concrete and raises the shear carrying capacity of the beams.

Indeed, pre-tensioned prestressed beams do not exhibit such a strong a/h-ratio dependence like that found for ordinary beams.

The explanation may be that small span-to-depth ratios enable the dowel action to develop in ordinary beams, provided that there is a well anchored reinforcement. However, in pretensioned prestressed beams the anchoring capacity of the strands will limit the dowel effects. But some tests still show more or less the dependence of v on the a/h-ratio. To draw a more convincing and definite conclusion, more available test data are needed. Before these are available, and in order to be on the safe side and finally to be convenient in practice, the slight increasing of v for small a/h-ratio should not be taken into account.

The precise plastic shear solutions for beams with T-section or I-section and without web reinforcement are somewhat complicated. To find a more rational and precise v dependence is even more difficult. For the sake of simplicity and in order to utilize the exsisting conclusions for beams with rectangular section, we may employ the plastic solutions to the web area of such beams but include the favourable effects of the compressive flange into the v expression. We must point out that the physical meaning of v here is somewhat different from it's original difinition, i.e. it's not only an empirical measure of concrete ductility and the absorbent of other shortcomings of the theory (for example, the neglecting of dowel action), but also a conversion factor of an alternative simple method. This is the reason why in such cases the v-value may become larger than unit.

A parametric analysis of prestressed beams without web reinforcement has shown the v-value in such cases increase approximatively proportional to the (whole area)/(web area) ratio until about 1.8.

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reasons for the difference found, we would like to mention the beams with shear reinforcement in Chapter $extsf{V}$. Among the possible reinforcement are rather different from the v expressions for that the resulting empirical formulae of $\boldsymbol{\nu}$ for beams without shear It's an unfortunate consequence of the above mentioned conclusions following:

- ۳ The ductility is probably increased and the dowel action may be decreased by supplying shear reinforcement.
- 2) depth of the beam as the shear depth, while for beams with i.e. for beams without shear reinforcement we use the whole Different measures of the shear depth are used for beams with shear reinforcement we use the effective depth $oldsymbol{h}^{oldsymbol{\star}}$ as the shear shear reinforcement and beams without shear reinforcement,
- ω strengths and shear span/depth ratios more available data are In order to determine the v dependence for high concrete

ıs: In order to unify the ν expression in various cases, more research needed in this field.

SUMMARY

joints are mainly treated. crete theory. Shear resistance of beams, deep beams, corbels and This paper deals with some elementary problems in reinforced con-

cut-off as yield condition and the associated flow rule (normality of the modified Coulomb failure criterion with non-zero tension The classical plastic theory has been used under the assumptions condition) as constitutive equations for the concrete.

For plane structural members without shear reinforcement an exact found for beams, deep beams, corbels and joints. solution (i.e. with corresponding upper and lower bounds) has been

identical to the exact shear solution without normal force. ted to combined shear and normal force, which is proved to Moreover an upper bound solution is presented for members subjec-

strength of concrete f_t by the reduced values $v \cdot f_c$ and $\rho \overset{*}{\cdot} \cdot f_c$ re-Due to the well-known fact that concrete is not the perfect plasreplacing the compressive strength of concrete $f_{\mathbf{C}}$ and the tensile tic material the exact plastic solutions are then modified by spectively. The v and ho^{\star} are called effectiveness factors.

A lot of work has been done to develop a more general and rational method for the assessment of approximate $\upsilon ext{-value}$ in various shear cases, especially for members without shear reinforcement.

The theoretical predictions are lastly compared with a great number of experimental results from the literature.

good agreement has been found.

RESUME

Rapporten omhandler en beskrivelse af nogle elementære teoretiske problemer i armeret beton.

Der behandles forskydningsproblemer i slanke bjælker, høje bjælker, konsoller og fuger.

Den klassiske plasticitetsteori anvendes under forudsætning af, at den associerede flydelov gælder. (Normalitetsbetingelsen).

Betonen antages at være et modificeret Coulomb materiale med en trækstyrke, der er forskellig fra nul.

For plane elementer uden forskydningsarmering udledes en eksakt løsning (dvs. med sammenfaldende øvre- og nedreværdi), der finder anvendelse både for slanke bjælker, konsoller samt fuger.

Ligeledes præsenteres en øvreværdiløsning for konstruktioner belastet med kombineret forskydning og normalkraft, der ses at være identisk med den eksakte forskydningsløsning uden normalkraft.

På grund af forskellen mellem den rigtige spændings- tøjningssammenhæng for beton og det sammenhæng der eksisterer for et stift plastisk materiale, skal de eksakte plastiske løsninger reduceres.

Dette gøres ved at indsætte de reducerede værdier $v\cdot f_c$ og $\rho^*\cdot f_c$ for henholdsvis den plastiske trykstyrke f_c og trækstyrken $f_t\cdot$

v og p^* kaldes effektivitetsfaktorer.

Der er blevet lagt et stort arbejde i at udvikle en mere generel og rationel metode til at vurdere tilnærmede v-værdier for forskydningstilfældet, især for forskydningspåvirkede legemer uden forskydningsarmering.

De teoretiske beregninger er herefter sammenlignet med et stort antal forsøgsresultater fra litteraturen omhandlende slanke bjælker, høje bjælker og konsoller.

Der er fundet udmærket overensstemmelse.

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APPENDIX A (CONTINUED)

TEST DATA OF DEEP BEAMS FAILED IN FLEXURE

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APPENDIX A

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APPENDIX B

TEST DATA OF CONVENTIONAL BEAMS FAILED IN SHEAR

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	B12 B13 B14 B15	B7 B8 B9 B10 B11	8 8 8 8 8 8 1 1 1 1 1 1 1 1 1 1 1 1 1 1	AC1 AC2 AC3 AC4 BC4	AA4 AB1 AB3 AB4	>>> H H H		III24A III24B III25A III25B III25B III26A	TE BEAM MARK
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	12.00 12.00 12.00 12.00	12.00 12.00 12.00 12.00	12.00 12.00 12.00 12.00	12.00 12.00 12.00 12.00	12.00 12.00 12.00 12.00	24.00 24.00 12.00 12.00 12.00	24.00 24.00 24.00 24.00 24.00	24.00 24.00 24.00 24.00 24.00	IN O
	10.56	10.56	10.56	10.55 10.70 10.75 10.80 10.56	10.63 10.50 10.55 10.63	21.00 21.00 10.30 10.50 10.55	21.00 21.00 21.00 21.00 21.00	21.00 21.00 21.00 21.00 21.00	IN d CONVENT
	36.00 36.00 36.00	36.00 36.00 36.00	00000	31.50 31.50 31.50 31.50 36.00	31.50 31.50 31.50 31.50	32.00 32.00 31.50 31.50 31.50	32.00 32.00 32.00 32.00	154.2 32.00 32.00 32.00 32.00 32.00	acc IN
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	5.48 3.27 2.42 3.42	4.48 1.77 5.97 3.47 5.53	4644	0.92 0.88 0.98 5.32	4.57 3.07 3.13 2.79 2.43	3.62 4.40 4.50	3.10 3.32 3.25	22.00 20.00	fc
	1.200 1.200 1.200 1.200	1.200		0.600 0.600 0.600 1.200	1.760 1.190 1.200 1.190 1.240	6.250 6.250 1.560 1.580 1.640	4.000 4.000 5.090 5.090	4.000 5.000 5.090 6.250	As IN2
	45.00 45.00 45.00	455.00	n 00000	45.00 45.00 45.00 45.00	45.00 45.00 45.00	00000	43.80 45.70 45.70 45.40	45.70 45.70 45.40 45.40	FY
	10.60 12.50 9.70 11.50	31.0	7.71.70	4.50 5.50 5.70 13.00	16.00 12.65 13.50 12.50	75387	89.00 78.00 80.00 68.00 76.50	66.50 65.00 65.00	Vt KIPS

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G1.5---5 G2----1 G2----2 G2.5---1

72 96 96 120

12.0 12.0 12.2 12.2

48.0 48.0 48.3

43.2

24.0 32.0 40.0

100

246 246 250 246

C---4 E1.5---1 G1.5---1 G1.5---2 G1.5---2

68 72 72

15.0 12.0 12.0

45.0 48.6 48.7

40.5 43.7 43.8 43.2

22.5 24.0 24.0

100

349 345 340 305

C1.5---1 C1.5---2 C1.5---3 C2----1 C2----3

75 75 100

16.5 15.6 15.8 16.3

50.5 50.3 50.6

45.3 45.3 45.6

25.0 25.0 25.0 33.3

10010

317 309 415 415 317

4.52 4.52 4.52

3490 3890 3500 3770 3730

11.9 11.5 12.2 12.0

32.50 24.00 36.00 25.00 25.00

Z----506 Z----507 Z----501 Z----505

135 135 135 135

12.0 12.0 12.0 12.0

90.0

86.4 84.0 81.0

22.5 33.8 54.8 57.5

200

289 328 328 230 230

2.26 2.26 4.52 6.78 8.54

3500 3500 3500 3500 3500

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34.80 23.50 45.00 37.00 36.25

M-0.4--2 M-0.4--3 M-0.4--4 M-0.6-7A M-0.6-8A

120 120 120 80,

0 11.0 0 11.5 0 11.0

48.0 48.0 48.0 5.5

46.7 45.0 46.0 46.5

40.0 40.0 25.0 25.0

10 10 10 10 10

209 207 196 277 279

2858 2960 2960

11.50 12.00 13.50

10.4 7.3 10.8 9.9 7.3

2957 2957

18.75 19.50

III-4--2 III-6--H III-6--I M-0.33-2 M-0.33-3

12.3 12.1 12.1 11.0 10.8

60.5 60.2 40.2

37.5 38.5 37.5

55.0 40.0

100

198 198 198 282 269

8.04 8.04 8.04 3.14 3.14

2800 2800 2800 2938 2953

14.1 10.7 5.3 9.9 7.5

21.40 9.25 9.75

BEAM MARK

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TEST DATA OF DEEP BEAMS WITH HORIZONTAL WEB REINFORCEMENT AND FAILED IN FLEXURE [82.4]

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As IN2	L BEAMS FAILED IN SHEA
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	67.50				68.00	•	3	63.20	7.	7.	3.6	2.2	3.9	47.80	7	8.3	ט מ ט מ	7 O	62.20		48.20			4	4.9	6.2	9.4	61.30	α		63.40	GI (n	61.20	٠	7	7	KSI	ı
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				2.000		2.000			٠,		2.000	2.000	2.000	2.000		2,000				594	0.594	067	067	067	18	18	N	0.951			0.954			1.47	1.47	2.361	2.35	IN2	Ď
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2.12 1.84 3.27 3.82 4.19

16.00 16.00 16.00 16.00 16.00 16.00 16.00 16.00 14.75 14.25 14.63 14.50 14.50 14.50 14.38 14.50 14.56 14.00 24.50 31.50 31.50 31.50 31.50 24.50 24.50 24.50 24.50 24.50 5.50 5.50 5.50 6.60 2.13 1.99 3.99 4.69 3.51 4.63 4.56 6.57

14.50 13.88 14.50 14.50 14.50 31.50 31.50 31.50 43.50 59.50 5.50 5.50 4.80 6.84 6.36 5.04 2.13

B---28E4 B---28A6 B---28B6

12.00 12.13 12.00 12.00 12.00

16.00 16.00 16.00 16.00

B---2186 B---28B2 B---28E2 B---28A4 B---28B4

12.00 12.00 12.13 12.00 12.00

B---21E4 B---21E4 B---21E5 B---21F4 B---21A6

12.00 12.00 12.00 12.00

B---14A6 B---14B6 B---21B2 B---21E2 B---21A4

12.00 12.00 12.00 12.00

16.00 16.00 16.00 16.00

14.00 14.50 14.44 14.75 14.50

17.50 17.50 24.50 24.50 24.50

5.50 5.50 5.00

6.59 6.78 2.01 1.64 4.32

B---14B2 B---14A4 B---14B4 B---14B4 B---14E4

12.00 12.00 12.00 12.00 12.00

16.00 16.13 16.00 16.00

[57.1] 17.50 17.50 17.50 17.50 17.50

5.50 5.50 5.50

BEAM MARK

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TEST DATA OF CONVENTIONAL BEAMS

APPENDIX B (CONTINUED)

B---56B2

59.50 59.50 59.50 59.50 5.50 5.50 5.50

2.13 3.62 3.95 4.12 5.78

B---56B6 B---70B2 B---70A4 B---70A6 B---84B4

12.00 12.00 12.00 12.00

16.00 16.00 16.00 16.00

14.63 14.38 14.50 14.00 14.31

59.50 73.50 73.50 73.50 87.50

5.50 5.50 5.50

B--113B4

12.00

16.00

14.38

116.50

G

B---56A4 B---56B4 B---56E4 B---56A6

12.00 12.00 12.00 12.00 12.13

16.00 16.00 16.00 16.00

14.50 14.75 14.50 14.50 14.00

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6.63 2.37 3.95 6.52 3.95

50

3.68 3.33 3.17 3.83 3.73

[63.2] 24.00 24.00 24.00 24.00

3.50 3.50 3.50 3.50

15.86 15.86 15.86 15.86

8.00

18.00 18.00 18.00 18.00

I-----2 II----3 II----4 III----5

71.00

64.85

193

195

B----74 B----75 B----76

6.00

24.00 24.00 24.00 12.00

B----79

B----66 B----69 B----71 B----72

6.15 6.19 6.11 6.10

24.00 24.00 24.00 24.00 24.00

B----61 B----63 B----64 B----64

6.10 6.08 6.15 5.89

6.00 24.00 24.00 24.00 24.00

B----57 B----57 B----58 B----59

5.92 6.03 6.00

6.000

5.95 9.95 9.50 9.50 9.50

6.00

B-----43 B-----45

5.97 5.96 5.95 5.95

6.00

B-0--3 C-0--1 C-0--2 C-0--3 D-0---1

8.00 8.00 8.00

18.00 18.00 18.00

A-0---1 A-0---2 A-0---3 B-0----1 B-0----2

8.00

18.00 18.00 18.00 18.00

BEAM MARK

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D-0---2 D-0---3

8.00

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	9.8	ຜູ	9.4	53.00	,	55.70	, ,			,	4.2	1.0	1.0	50.60	٥ ۵	0	60.40	4.4	8.4	6.8		6.8	6.8	56.80	20 0	מ ע	0	٠ . د د	56.80) (i	6.2	,		53.71	3.7	٠.	ى د 1 -	ι ω 1 - 1	3.7	53.71		3.7	53.71	ن د د د) i.) 1	KSI	Fy	AR
٠	٦. ٦	ω ω		24.200	•	44.250	0.00) i	20.40	,	5.25	7.75	. 95	36.700	83		1 /1	:	:.	•			4.90	6.495	9 1		0.06	л U 3 К	. 6. 545	1 0	7.19)		50.19	8) . 0 U	י פינ	9.1	28.78		21.1	27.21	* 4			KIPS	۷t	
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	- 1		'n	K) I	TV19A2	IV18E2	-17	<u>.</u>	Ļ,	<u>.</u> J	* 1 4 A	. L	ָבָּיבָיבָ בַּיבָיבָיבָ	180			TI-1	7	ייייייייייייייייייייייייייייייייייייי	7 A			304	B3046		304	B3044	304	304	304		7.7	7.7	/ 2	ł		9	9	9	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1		9		9	8	1 1 1 1 8	MARK	BEAM	,
	0.0	6.0	8.0	6.00	•				6.00						6.00		0	0	5 6	» «	,		÷	6.10		÷	6.00	•	0	•		 H	4	4 <i>«</i>	• 0		•			6.03	5	0	000	.08	95	14	Ė	ָ ֓֞֓֞֞֓֞֓֞֓֓֓֓֞֓֞֓֓֓֓֞֞֓֓֓֓֓֞֡	
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19))	• •	36.00	6.0	•	6.0	6.0	36.00	6.0		6	6.0	6.0	36.00	۱ د	6.0	6.0	6.0	36.00	6.00	λ γ γ		331.00	3	15.0	171.80	29.0	07.7	6.4		2.1	2	ω. σ	64.20		6.7	6.7	2.1	42.72	6.7	}	4. L	4.0	2.7	[67.3] 32.04		acc N	
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TEST DATA OF CONVENTIONAL BEAMS FAILED IN SHEAR APPENDIX B (CONTINUED)

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		5CC 6.00 12.00 9.86 60.00 2.43 2	6.00 12.00 9.94 60.00	CC 6.00 12.00 10.00 60.00 2.48 3	RCC 6.00 12.00 10.06 60.00 1.77 1	0.00 T7.00 3:00 #0:00	.00 12.00 9.94 40.00 1.80	6.00 12.00 10.00 #6.00 2.23 2:	3AC 6.00 12.00 10.06 40.00 1.87	SAAC 6.00 12.00 9.86 36.00 1.99 1.	1.94 2.	6.00 12.00 9.94 36.00	6.00 12.00 10.00 36.00 2.23 2.	00 10.06 36.00	6.00 12.00 9.86 /2.00	SEC 6.00 12.00 9.94 /2.00 4.90 2.	5,43 2,	-6CC 6.00 12.00 9.00 80.00	00 12.00 9.94 00.00 5.57	-4CC 6.00 12.00 10.00 60.00 5.3 2.00	6.00 12.00 9.86 48.00 4	-5AC 6.00 12.00 9.94 48.00 4./6 2.00		6.00 12.00 10.00 48.00 4.42 1.	6.00 12.00 10.06 48.00 4.62	2	6.00 12.00 9.94 36.00 4.76 2.	4.24 1.58	3C 6.00 12.00 10.00 30.00	-6C 6.00 12.00 9.94 36.00 5.01 1.20	3C 6.00 12:00 9:86 84:00 3:10 2	C 6.00 12.00 9.94 84.00 3.16 2.	GC 6.00 12.00 10.00 84.00 3.05		6EC 6.00 12.00 9.86 72.00 2.77 2.		6.00 12.00 10.00 72.00 3.08 1.	600 6 00 12 00 9.86 60.00 2.98 2.53	2.95 2.	4CC 6.00 12.00 10.00 60.00 2.98 1	3CC 6.00 12.00 10.06 60.00 2.97 1.2	6AC 6.00 12.00 9.86 48.00 3.31 2.53	5AC 6.00 12.00 9.94 48.00 2.66 2.00	.00 12.00 10.00 48.00 2.39	C 99.1	BEAM b h d acc 1 fc As Fy MARK IN IN IN IN KSI IN2 KSI	TEST DATA OF CONVENTIONAL BEAMS FAILED IN SHEAR	1	ADDENDTY B (CONTINUED)
		19.000	8.900	7 700	7.000		9.200	9.800	9.000	8.300	14.000		11.300	9.600	9.100	.00	12.000		8	9	11.800	3 6	Š	12.100	12.000	13.500	12.800	13.000		12.500	ពោ ។	9.100	9.400	2000 2000	9.500	8.900	9.400	10.000	10.000		000	- C-	4.000	8.500	rda la colo	y Vt I KIPS	Le∵ con		
A-410	A-48	A-47	A-46	4	A-44	-4	A-42	A-41		1	E4	1	1		. 1	D20	D19	D18	D17	D16	,	7 - 1 - 1 - 1 - 1 - 1		D13	1		D10		1	1	J6			D	D2	D1			S-II-OCB	S-TT-CV	S-11-008	S-LCA	XIPCB	XIPCA		BEAM MARK	F		
.00 16.50 14.	00 16.50 14.	9.00 16.50 14.	.00 16.50 14.	. 00 to 00	9.00 16.50 14.	.00 16.50 14.	.00 16.50 14.	.00 16.50 14.		.00 16.50 14.	9.00 16.50 14.	.00 16.50 14.	.00 16.50 14.	.00 16.50 14.		.00 16.50 14.	.00 16.50 14.	.00 16.50 14.	9.00 16.50 14.2	.00 16.50 14.		00 16.50 14.	00 16.50 14.	9.00 16.50 14.1	00 16 50 14	000000000000000000000000000000000000000	.00 16.50 14.	9.00 16.50 14.1	.00 16.50 14.	.00 16.50 14.	.00 16.50 14.		.00 16.50 14.	9.00 16.50 14.1	00 16 50 14.	.00 L6.50 L4.		-	10.00 20.00 17.9		00 20 00 17.	00 12 00 10.	00 10 00 10.	20 12		IN IN IN	I DATA OF	DE CE	
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A0--7--1 A0-11--1 A0-15-1A A0-15-1B

A0--7--2 A0-11--2 A0-15-2A A0-15-2B A0--3--1

A0-11-3B A0-15-3A A0-15-3B A0-15-3C A0--3--2

A0--3-3B A0--3-3C A0--7-3A A0--7-3B A0--11-3A

A-5----4 A-5----5 A-5----6 A-5----7

A-4---11 A-4---12 A-5----1 A-5----2 A-5----3

BEAM MARK

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27 258.00 27 258.00 118 258.00 60 258.00	80 95.00 95.00 95.00 76.00	370 370 370	70 70 70 70	70 70 70	70 40 70 70	0000	.500 .500 .660	FY KSI	IN SHEAR
7.550 5.500 6.150 5.750	9.000 8.400 6.990	70.000 97.200 62.000 111.200	00.000 00.000 000.000 0001.000	88888	14.500 15.000 18.500 18.600 20.200	21.500 23.000 27.000 34.500	55.000 40.000 41.500 22.500	Vt KIPS	docina.
R3 R3 24	0-A1 0-A2 0-A3	3 G 3	3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3	2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	NNNN N	2	18 30 42 48 248	BEAM MARK	
4.00 8.00 6.8 4.00 8.00 7.2 4.00 8.00 6.7 4.00 8.00 6.7	12.20 21.90 18.15 12.00 22.10 18.35 12.10 21.90 18.17	7.50 10.00 8.75 7.50 10.00 8.75 7.50 10.00 8.75 7.50 10.00 8.75	7.50 10.00 8.75 7.50 10.00 8.75 7.50 10.00 8.75 7.50 10.00 8.75 7.50 10.00 8.75	7.50 10.00 8.75 7.50 10.00 8.75 7.50 10.30 9.05 7.50 10.30 9.05 7.50 10.00 8.75	.50 10.00 8.7 .50 10.00 8.7 .50 10.00 8.7 .50 10.00 8.7 .50 10.00 8.7	7.50 10.00 8.75 7.50 10.00 8.75 7.50 10.00 8.75 7.50 10.00 8.75	6.00 12.00 10.50 6.00 12.00 10.50 6.00 12.00 10.50 6.00 12.00 10.50 7.50 10.00 8.75	TEST DATA OF CONVEN b h d IN IN IN	
0000	[63.1] 72.00 90.00 126.00	36.00	36.00 36.00	36.00 36.00 36.00	00000	36.00 36.00	[60.8] 18.00 30.00 42.00 48.00 36.00	CONVENTIONAL BEARING IN IN IN IN	TOWN BEA
2.00 4 2.00 3 2.00 4	უ უ უ უ უ უ უ დ დ უ დ დ	തത്ത്	ພພ4.៧∢	44400	44440	4444	<u> </u>		MS FAILED
. 75 . 88 . 40	. 27 . 44 5	.18 .34 .36	. 18 1 . 83 1 . 76 1	.47 1 .45 1 .41 1 .91 1	00448	16 0 10 0 72 0 56 0	35 8 2 1 1		
1.200	.00	1.178 1.227 1.178 1.227	1.227 1.190 1.178 1.227 1.190	. 553 173333	.982 .982 .178 .178	.785 .785 .785 .982	. 5500 8000 9000	As IN2	IN SHEAR
) Pun			÷				Fy	
5.75 5.75 5.75		11.13 12.46 11.68 12.13	9.81 10.25 10.25 11.02		9.150 9.700 10.030 10.030 10.250	9.260 11.020 8.930 10.140 10.140	$\circ \circ \circ \circ \circ$	Vt KIPS	

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\$103/1 \$103/2 \$103/3 \$103/4	\$1C1/4 \$1C2/1 \$1C2/2 \$1C2/2 \$1C2/3 \$1C2/4	C	C 1 C 1 C 2 C 2 C 2	A 1 A 2 B 2 B 2 B 3	1 4 7	IIB1 IIB2 IIB3 IA1A	381	R1 R2 R3 R3 R7	BEAM MARK
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7.87 7.87 7.87 7.87	7.87 7.87 7.87 7.87 7.87	5, 5, 5, 5, 5, 5, 5, 5, 5, 5, 5, 5, 5, 5	9 9 9 9 9 8 8 8 8 8 9 9 9 9	39.40 39.40 19.70 19.70 19.70	13.00 13.00 13.00	14.00 14.00 14.00 14.00	12.00	12.00 12.00 12.00 12.00	DATA OF h N IN
6 6 6 6 6 6 6 6 7 7 7 7		55.55 55.55	9.20 9.20 9.20 9.20	36.80 36.80 18.40 18.40 18.40	10.70 10.70 10.70 10.70	12.50 12.25 12.12	10.03	10.70 10.70 10.70 10.70	CONVEN d IN
5.99 6.65 7.32 7.98	[74.2] 2.66 3.33 3.99 4.66 5.32	27.58, 16.55 16.55 16.55	27.58 27.58 27.58 27.58 27.58	[72.5 110.30 110.30 55.15 55.15 55.15	[71.3 16.00 16.00 21.40 21.40	[60. 42.00 48.00 54.00 30.00	[81. 30.00	171. 36.00 36.00 36.00 36.00	CONVENTIONAL d acc
	. •	1.50 0.90 0.90 0.90	1.50 1.50 1.50 1.50	6.00 3.00 3.00	1.50 1.50 1.50 1.50	4] 4.00 4.00 4.00 4.00	1] 3.00	7.1	BEAMS 1 IN
3.67 3.67 3.67 3.67	3.36 3.89 3.89 89	00000 00000	4.64 4.64 4.98 3.77 4.06	5.28 4.62 4.93 5.80	4.56 5.19 4.29 4.36	3.58 3.03 3.03	5.40	3.80 3.80 3.60 4.07	FAILED fc KSI
0.380 0.380 0.380 0.380	0.380 0.380 0.380 0.380	0.484 0.174 0.174 0.174 0.174	0.484 0.484 0.484 0.484 0.484	7.744 7.744 1.936 1.936 1.936	1.570 0.614 1.570 0.614	1.148 1.148 1.148 1.148	0.114	0.630 0.940 0.940 0.940	IN SHEAR As IN2
39.85 855	39.85 39.85 85	62.50 62.50 62.50 62.50	62.50 62.50 62.50 62.50	62.50 62.50 62.50 62.50	56.00 56.00	46.80 46.40 46.60 47.00	169.00	e e	AR Fy KSI
13.47 12.34 14.12 10.75	22.48 17.92 19.58 18.39 15.22	2.56 2.38	5.06 5.06 5.06	80.70 73.90 23.40 19.60	42.37 21.37 22.24 15.28	10.20 9.15 8.30 13.30	6.50	10.10 10.60 10.10 12.20 11.80	Vt KIPS

D-18---2 E-18---1 E-18---2

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18.00

15.90 15.90

24.00 24.00 24.00

3.50 3.50

1.476 0.954 0.954

96.90 99.50 99.50

60.00 49.65 50.00 B-18---1 B-18---2 C-18---1 C-18---2

8.000

18.00

15.90 15.90 15.90 15.90

[58.3] 24.00 24.00 24.00 24.00 24.00

3.68 3.33 3.71 3.83 3.72

3.880 3.880 2.353 2.391 1.488

38.70 38.70 71.00 67.50

70.00 69.50 65.00 70.00

D-18---1

T-B---1 T-B---5 T-B---5 T-B---6 T-B---6

9.84 9.84 9.84 9.84

[84.3] 19.69 19.69 19.69 19.69 29.53

2.55 3.00 3.40 3.40

0.616 0.616 0.889 0.889 0.889

55.47 55.47 55.47 55.47 55.47

20.72 18.74 24.80 27.56 13.01

T-B---10

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TEST DATA OF CONVENTIONAL BEAMS FAILED IN SHEAR

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APPENDIX B (CONTINUED)

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	88888	8.000	88.000	88888	8.00	8.00 8.00 8.00	8.00 8.00 16.00 8.00	8.000 8.000	IN p	
	18.00 26.00 26.00 26.00 18.00	26.00 26.00 18.00 18.00	26.00 26.00 18.00 18.00 22.00	18.00 22.00 26.00 26.00 22.00	26.00 26.00 26.00 18.00	18.00 18.00 18.00 22.00 22.00	18.00 22.00 22.00 14.00 26.00	22.0 26.0 18.0 22.0	NI	TEST
	16.10 24.10 24.10 24.10 16.10	24.10 24.10 16.10 16.10 16.10	24.00 24.00 16.10 16.10 20.10	16.10 20.10 24.10 24.10 20.00	24.20 24.20 24.20 16.10 16.10	16.20 16.20 20.20 20.20	0 16.10 0 20.10 0 20.10 0 20.10 0 12.10 0 24.20	0 20.20 0 24.20 0 16.10 0 20.10 0 24.10	N G	DATA OF
	9.50	9.50 9.50	6.00	6.000 0000	6.00 6.00 6.00	6.00 6.00	2.75 2.75 2.75 4.00	[65. 2.75 2.75 2.75 2.75 2.75	acc	APPENDIX F CORBELS
	22.50 22.50 25.50 25.50	1.50 2.50 2.50 2.50	1.50 1.50 1.50 1.50	1.50 1.50 1.50 1.50	1.50 1.50 1.50	1.50 1.50 1.50	1.50 2.50 1.50	1.50 1.50 1.50 1.50	IN 1	*
	4.20 4.85 4.14 3.84	3.96 3.77 4.70 4.49 4.34	4.04 4.39 3.83 4.07 3.82	4.28 4.32 4.63 3.73 4.26	3.92 3.74 3.95 4.25 6.41	4.50 3.99 4.21 3.79 3.79	3.26 6.50 3.90	3.05.07 3.05.07 3.05.07 97.44	fc KSI	B (CONTINUED)
	1.198 1.195 1.195 1.195 2.396	2.391 2.391 1.198 1.198 1.198	1.574 1.574 2.396 2.396 2.396	1.198 1.206 1.195 1.195 1.584	0.620 0.620 0.620 1.198 1.198	0.622 0.622 0.622 0.614	2.396 2.396 2.396 2.401 0.407	0.614 0.620 1.198 1.206 1.195	As IN2	(NUED)
	44.40 52.50 45.70 45.40 50.50	47.30 54.30 53.00 54.50 44.30	46.60 45.60 47.30 53.30 47.30	53.30 47.30 47.30 47.50 45.60	43.20 95.80 45.00 47.30 46.60	48.10 95.80 47.30 43.20 95.80	43.30 45.80 45.00 47.70 51.00	45.30 45.30 43.60 43.30 47.00	Fy KSI	
	78.05 150.00 119.92 112.02	155.01 155.98 85.52 86.81 86.94	150.14 139.97 113.99 123.52 132.18	110.64 114.97 124.93 123.39 135.04	95.0 104.9 88.44 89.0 130.0	72.60 85.54 81.00 92.44 86.13	140.3 175.2 265.3 184.6 84.0	100.0 109.0 100.0 136.5 137.4	Vt Kips	
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	CO98 CO99 CO102 CO103	CO91 CO93 CO94 CO95 CO96	CO	CO 80 CO 81 CO 82 CO 83	CO67 CO69 CO75 CO76	CO61 CO62 CO64 CO65 CO66	CO55 CO56 CO59 CO59	CO46 CO47 CO48 CO49 CO54	BEAM MARK	
	CO-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1	CO91 8.0 CO93 8.0 CO94 8.0 CO95 8.0	CO85 8.0 CO87 8.0 CO89 8.0	CO - 1 - 1 - 1 - 8 - 0 8 - 0 CO - 1 - 1 - 1 - 1 - 8 3 8 - 0 CO - 1 - 1 - 1 - 1 - 8 3 8 - 0 CO - 1 - 1 - 1 - 1 - 8 3 8 - 0 CO - 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1 -	CO75 8.0 CO75 8.0 CO76 8.0 CO77 8.0	CO - 1 - 1 - 1 - 1 - 6 - 6 - 6 - 6 - 6 - 6		46 8.0 47 8.0 48 8.0 54 8.0	BEAM b MARK IN	
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APPENDIX B (CONTINUED)

APPENDIX B (CONTINUED)

TEST	
DATA	
O.F	
CORBELS	
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N	
SHEAR	

CO----7E CO----9E CO----10E

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8.00 8.00 8.00

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24.10 24.10 24.10 16.10 16.20

[65.1] (CONTINUED)
9.50 2.50 4.44 1
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44.90 43.00 44.90 44.50 48.10 BEAM MARK

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acc IN

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KSI KSI

As IN2

FY KSI

					Vt XIPS 0 144.79 0 124.93 0 145.56 0 114.50 0 63.89 0 109.48 0 139.97 0 148.84 0 127.06	
OAO44 OAO48 OBO49 OCO50 ODO47	G33S11 G33S31	K11		A A A A A A A A A B B B B B B B B B B B	BEAM MARK R101 R102 R103 R104 R111 R115 R115 R111 R1110	
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31.37 30.60 33.50 26.00	19.20 24.05	12.55	20.25 31.60 26.65 31.15	24.40 15.25 20.25 27.80	Vt KIPS 100.00 100.00 82.50 90.00 87.50 109.50 109.50 66.00	

APPENDIX	
8	
(CONTINUED)	

TEST DATA OF DEEP BEAMS FAILED IN SHEAR

		· ·	9.05 3589 40.38 10.05 4864 54.04 10.06 4864 53.53 10.05 4864 45.89 10.09 4864 38.24	.03 3518 27.3 .04 4864 28.5 .05 4864 28.5 .26 3426 33.1 .39 3426 40.7 .52 3589 33.1	3.97 3848 15. 1.51 3212 15. 1.02 4864 17. 1.03 4864 23. 1.04 4864 23. 1.04 4864 25.	3.97 3848 12.90 11.27 3325 13.50 2.53 3325 20.00 2.53 3325 15.00 2.53 3325 11.50 3.97 3848 26.00 3.97 3848 18.10	3325 1 3325 1 3325 1 3326 1 3848 1	As Fy Vt
to de la								
-	206 207 208 700A 700B	2002 2003 2004	M1.08 M1.09 M1.010 M1.011 M1.012			· · · · · · · · · · · · · · · · · · ·	MARK III0 III0 MV1 MV2	BEAM
• •	12.0 12.0 12.0 10.0	12.0		12.4 10.7 10.3 10.7	1100 1100 1100 1100 1100 1100 1100 110	84444 0 0	12.0 12.1 12.2 12.2	ੇ ਹ
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	45.0 60.0 90.0 37.5	45.0 67.5 90.0	o nnnnn	32.0 32.0 32.0	30.0 30.0 30.0 30.0 30.0 30.0	0.07775	100 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	OM C
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	294.0 294.0 253.0 341.0 305.0	214.0 214.0 214.0 294.0	50.	359.0 337.0 337.0 353.0 353.0	239. 258. 258. 344. 344. 3599. 0	55. 55.	206.0 198.0 257.0 257.0	fc KG/CM2
٠.	12.06 12.06 12.56 7.60 7.60	12.06 12.06 12.06 12.06		9.10 3.08 4.02 8.16	8.502 8.502 8.502 8.502	4.64 6.16 10.18 6.04 6.16 7.82 10.10	12.06 8.04 3.65 4.64	As CM2
-	22850 2850 2850	2222	N N W N N N	Nooon	2944 2944 2859 2780 2732 2632 2632	2712 2850 3820 3040 2920 3660 3680	2790 2800 3392 2712 2712	Fy KG/CM2
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fc KG/CM2 TEST DATA OF CORBELS FAILED IN SHEAR

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B12 B13 B1.51	II23 II24 A1.57 A1.57 B11	II9 II10 II11 II21 II22	II2 II3 II4 II5	I11 I12 I13 I14 I15	I	I	801A 804A 805A 806A	801 802 803 803A 802A	BEAM MARK	
16.2 5 16.0 5 16.0 5	10.0 10.0 13.0 12.0	10.0	10.0	10.0	10.0 10.0 10.0	10.0	10.7 10.7 10.4 10.5	10.3 10.5 10.3 10.0	S A	
0.50	75.0 75.0 48.0 48.0	75.0 75.0 75.0 75.0 75.0	75.0 75.0 75.0 75.0 75.0	60.0 60.0 60.0	60.0 60.0 60.0	60.0	75.0 75.0 75.0	75.0 75.0 75.0 75.0	S P	TEST
45.0 45.0 45.3	69.0 69.0 43.2 43.2	69.0 69.0 69.0	69.0 69.0 69.0	57.0 54.0 57.0 57.0 57.0	54.0 57.0 57.0 57.0	57.0 57.0 57.0 57.0	67.5 67.5 67.5	67.5 67.5 5.5 5.5	čk g	DATA
25.0 25.0 25.0	50.0 24.0 24.0	50.0 50.0 50.0	25.0 75.0 90.0 102.5 25.0	50.0 50.0 110.0 80.0 60.0	30.0 75.0 50.0 37.5	30.0	40.0 70.0	400. 400. 400.	- CM C	APPEN OF DEEP
10.0	15.0 12.0 10.0 10.0	12.0 12.0 12.0 15.0 21.0	12.0 12.0 12.0 12.0	12.0 12.0 12.0 12.0	12.0 12.0 12.0 12.0	12.0 12.0 12.0 12.0	20.0 20.0 15.0	200	CM CM	APPENDIX B
309.0 415.0 317.0	389.0 285.0 252.0 252.0 317.0	362.0 277.0 345.0 225.0 317.0	317.0 308.0 301.0 326.0 288.0	297.0 318.0 286.0 348.0 310.0	298.0 277.0 276.0 331.0 348.0	236.0 284.0 283.0 330.0	324. 267. 273. 252.	CONTINUED) 0 269.0 0 319.0 0 278.0 0 263.0 0 317.0	fc KG/CM2	S FAILED IN
4.27 4.27 4.27	12.57 6.16 4.71 4.71 4.27	10.17 12.57 15.20 12.57 12.57	10.17 15.20 15.20 15.20 15.20	9.78 8.85 7.72 7.72 7.72	8.85 7.72 7.72 7.72 6.34	0 7.72 0 7.72 0 7.72 0 6.34 0 9.78	0 7.77 0 7.77 0 7.77	7777	As 12 CM2	ENUED)
3888 3918 3692	2670 2770 3090 3090 3984	2920 2670 2680 2670 2670	2920 2620 2620 2620 2900	2440 2630 2620 2620 2620 2620	2630 2620 2620 2620 2580	2620 2620 2520 2580 2440	9 3067 9 3067 9 3067 9 3067	3067 9 3067 9 3067 9 3067 9 3067	FY 12 KG/CM2	SHEAR
25.00 26.00 27.00	36.00 28.00 22.00 20.00 28.00	25.50 25.50 31.00 25.50	31.25 27.50 21.60 19.50 34.40	23.75 23.70 6.90 14.20	25.00 14.00 21.00 27.50 17.40	18.10 24.05 40.10 35.00 29.25	7 37.50 7 35.00 7 21.80 7 21.80	7 36.40 7 38.00 7 35.00 7 33.50 7 36.50	W2 T	
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CO-110 CO-111 CO-112 CO-116

CO-104 CO-105 CO-106 CO-107 CO-108

BEAM MARK

APPENDIX C (CONTINUED)

APPENDIX D

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36.00 18.20 7.500 P13 2.00 9.92 6.97 41.90 14.60 4.480 0.314 36.00 18.20 18.20 7.500 P13 2.00 9.92 6.97 41.90 30.00 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.00 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.00 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.00 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.30 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.30 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.30 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.30 4.480 0.314 36.00 18.20 18.20 18.20 9.92 6.97 41.90 30.30 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.30 4.480 0.314 36.00 18.20 18.20 9.92 6.97 41.90 30.30 4.480 0.314 36.00 18.20 1 | 86.00 45.00 4.040 0.442 21.90 9.650 P==10 1.50 9.92 7.05 36.70 29.60 4.400 0.314 38.00 16.53 18.00 45.00 4.930 0.661 22.10 10.750 P==11 2.0 9.80 6.93 43.40 0.314 36.45 20.80 12.70 45.00 4.350 0.661 20.70 7.400 P==112 2.45 9.80 6.93 43.40 29.10 4.720 0.314 36.45 21.80 12.40 45.00 37.00 4.350 0.681 20.70 7.400 P==138 200 9.22 6.97 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14.80 4.00 0.314 36.00 7.25 86.00 28.00 4.91 0.602 18.30 9.10 9.00 30.00 4.91 14.90 9.10 9.10 4.00 30.30 4.00 0.314</td> <td>23.40 55.00</td> <td>\$\begin{align*}{35.00} \begin{align*}{55.00} \begin{align*}{4.150} \oldsymbol{0.755} \begin{align*}{25.00} \begin{align*}{25.00} \begin{align*}{4.120} \oldsymbol{0.7755} \begin{align*}{25.00} \begin{align*}{25.00} \begin{align*}{25.00} \begin{align*}{4.240} \oldsymbol{0.7755} \begin{align*}{25.00} al</td> <td>\$5.00 \$5.00 \$4.150 0.601 \$29.20 8.200 \$\begin{array}{cccccccccccccccccccccccccccccccccccc</td> <td> 18.90 36.00 1.900 0.785 18.90 16.950 29.20 36.00 25.00 1.900 0.34 36.00 2.34 36.00 2.35 36.</td> <td>## 15:00 42:50 5:00 0.601 18:00 13:000 P==== 1:50 10:04 7:05 40:00 20:00 4.200 0.314 36:00 7:05 36:00 36:00 4.300 0.785 28:00 11:000 P==== 1:50 10:04 7:05 40:00 20:00 4.200 0.314 36:00 7:05 36:00 30:00 4.400 0.314 36:00 7:05 36:00 35:00 4.430 0.785 28:00 11:000 P==== 1:50 10:04 7:05 40:00 20:00 4.200 0.314 36:00 7:05 36:00 35:00 4.430 0.785 27:00 7:000 P==== 1:50 10:04 7:05 40:00 30:00 4.200 0.314 36:00 7:05 36:00 35:00 4.430 0.785 27:00 7:000 P==== 1:50 10:00 7:05 36:00 35:00 4.200 0.314 36:00 12:80 35:00 4.730 0.442 21:30 8.450 P==== 1:50 10:00 7:09 34:50 11:30 4.160 0.314 36:00 12:80 35:00 4.730 0.742 21:30 8.450 P==== 1:50 10:00 7:09 34:50 11:30 4.160 0.314 36:00 12:80 35:00 4.200 0.785 21:80 11:400 P==== 1:50 10:00 7:09 34:50 11:30 4.200 0.314 36:00 7:03 36:00 15:00 4.200 0.314 36:00 12:80 35:00 4.200 0.785 18:20 7:500 P==== 1:50 10:00 7:09 34:50 11:30 4.200 0.314 36:00 7:03 36:00 15:00 4.200 0.314 36:00 7:05 36:00 12:80 37:00 7:00 36:00 12:80 37:00 12:80 37:00 12:80 37:00 12:</td> <td>\$8.00 42.50 5.000 0.994 12.20 14.750</td> <td> 13.00 14.55 5.060 0.944 11.20 14.056 14.250 15.00 15</td> <td> 13.10 23.756 13.10 23.756 13.10 23.756 13.10 23.756 13.10 23.756 13.10 23.756 13.10 23.10 13.10 23.10 13.10 23.10 13.10 23.10 13.10 23.10 13.10 13.10 23.10</td> <td> 13-30 13-10 23-750 13-10 23-750 13-10 23-750
23-750 23-75</td> <td> 1.00 </td> <td>42.00 30.00 7.300 0.222 32.80 15.000 32.40 8.00 5.740 0.441 22.70 9.05 1.401 180 9.00 31.00 32.40 8.00 5.740 0.441 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 9.05 1.401 0.20 10.00 31.00 28.00 5.200 0.443 22.70 10.00 31.00 28.00 5.200 0.443 22.70 10.00 31.00 28.00 5.200 0.443 22.70 10.00 31.00 28.00 5.200 0.443 22.70 10.00 31.00 28.00 5.200 0.443 22.70 10.00 31.00 28.00 5.200 0.443 22.70 10.00 31.00 28.00 5.200 0.443 22.70 10.00 31.00 28.00 5.200 0.443 20.00 10.22 22.70 10.00 5.200 0.442</td> <td>44.00 42.00 44.00 30.00 7.200 0.222 32.00 12.000 44.00 30.00 7.200 0.222 32.00 12.000 32.000 32.00 12.000 32.0000 32.000 32.000 32.000 32.</td> <td>44.00 42.00</td> <td>44.00 36.00</td> <td>44.00 10.00 6.880 0.222 33.00 13.900 846 31.00 4.790 0.119 11.00 3.95 44.00 31.00 6.260 0.222 33.00 13.900 846 31.00 4.790 0.119 11.00 3.95 44.00 30.00 6.260 0.222 33.00 13.900 846 31.00 4.790 0.119 11.00 3.95 44.00 30.00 7.200 0.222 33.00 13.00 846 31.00 10.00 33.40 28.00 5.260 0.443 24.00 12.00 10.00 31.00 28.00 5.260 0.443 24.00 12.00 12.00 10.00 31.00 28.00 5.260 0.443 22.70 9.70 12.00 12.00 10.00 31.00 28.00 5.260 0.443 22.70 9.70 12.00 12</td> <td>40.00 18.00 6.880 0.222 33.00 18.900 844 32.00 6.00 4.790 0.118 11.00 6.886 14.00 44.00 36.00 5.760 0.222 33.00 11.300 844 32.00 6.00 4.790 0.118 11.00 6.886 14.00 44.00 36.00 5.760 0.222 33.00 11.300 846 3.10 6.00 4.00 10.00 33.40 28.00 6.260 0.443 24.00 10.00 12.200 0.222 33.00 11.300 846 11.300 12.200 10.00 33.40 28.00 5.760 0.443 22.20 33.00 11.300 846 11.300 12.200 10.00 33.40 28.00 5.760 0.443 22.20 33.00 11.300 846 11.300 12.200 10.00 33.40 28.00 5.760 0.443 22.20 33.00 11.200 10.200 11.80 28.00 5.760 0.443 22.20 33.00 12.200 10.200 11.80 28.00 6.200 0.443 22.20 33.00 12.200 10.200 11.80 28.00 6.200 0.443 22.20 33.00 12.200 10.200 11.80 28.00 6.200 0.443 22.20 33.00 12.200 10.200 11.80 28.00 12.200 10.200 11.80 28.00 12.200 10.200 11.80 28.00 12.200 10.200 11.80 28.00 12.200 10.200 11.80 28.00 12.200 10.200 11.80 28.00 12.200
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13-16 3.00 12.00 10.38 53.40 28.0 13-26 2.94 12.00 10.03 53.03 28.0 13-41 2.90 12.00 10.04 52.78 28.0 21-26 2.96 12.00 10.21 53.15 54.0 22-09 2.96 12.00 11.07 53.15 36.0	12-29 3.00 12.00 9.18 53.0 12-34 3.08 12.00 9.99 55.6 12-35 3.08 12.00 9.99 55.6 12-50 2.96 12.00 10.20 53.1 12-61 3.00 12.00 9.90 53.4	B12-10 3.06 12.00 11.11 53.77 36.0 B12-12 3.00 12.00 11.13 53.40 36.0 B12-14 3.00 12.00 11.14 53.40 36.0 B12-19 2.98 12.00 11.09 53.28 36.0 B12-26 3.03 12.00 10.06 54.40 36.0	A32-37 6.00 12.00 8.20 72.0 A32-49 6.00 12.00 8.20 72.0 B11-20 2.95 12.00 10.21 52.6 B11-29 2.95 12.00 10.00 52.8 B11-40 2.95 12.00 10.00 52.8	A22-40 6.00 12.00 8.20 72.00 36.0 A22-49 6.00 12.00 8.20 72.00 36.0 A32-19 6.10 12.00 9.03 73.20 36.0 A32-22 6.00 12.00 9.38 72.00 36.0 A32-27 6.00 12.00 9.16 72.00 36.0	A22-28 6.10 12.00 8.75 73.20 36.0 A22-31 6.00 12.00 8.06 72.00 36.0 A22-34 6.00 12.00 8.35 72.00 36.0 A22-36 6.00 12.00 8.35 72.00 36.0 A22-36 6.00 12.00 8.35 72.00 36.0	A21-39 6.00 12.00 8.95 72.00 54.0 A21-51 6.00 12.00 8.12 72.00 54.0 A22-20 6.00 12.00 8.45 72.00 36.0 A22-24 6.00 12.00 8.85 72.00 36.0	114-39 6.00 12.00 8.35 72.00 24.00 A14-44 6.00 12.00 8.50 72.00 24.00 A14-55 6.00 12.00 8.53 72.00 24.00 A14-58 6.00 12.00 8.42 72.00 24.00 A14-59 6.00 12.00 8.45 72.00 54.00	IN IN IN d	ST DATA OF PRESTRESSED
13-16 3.00 12.00 10.38 53.40 28.00 8.0 13-26 2.94 12.00 10.03 53.03 28.00 8.0 13-41 2.90 12.00 10.04 52.78 28.00 8.0 21-26 2.96 12.00 10.21 53.15 54.00 8.0 22-09 2.96 12.00 11.07 53.15 36.00 8.0	12-29 3.00 12.00 9.70 55.00 36.00 8.0 12-34 3.08 12.00 9.99 55.64 36.00 8.0 12-50 2.96 12.00 10.20 53.15 36.00 8.0 12-61 3.00 12.00 9.90 53.40 36.00 8.0	B12-10 3.06 12.00 11.11 53.77 36.00 8.0 B12-12 3.00 12.00 11.13 53.40 36.00 8.0 B12-14 3.00 12.00 11.14 53.40 36.00 8.0 B12-19 2.98 12.00 11.09 53.28 36.00 8.0 B12-26 3.03 12.00 10.06 54.40 36.00 8.0	A32-37 6.00 12.00 8.20 72.00 36.00 8.0 A32-49 6.00 12.00 8.20 72.00 36.00 8.0 B11-20 2.95 12.00 10.21 52.63 54.00 8.0 B11-29 2.95 12.00 10.00 52.80 54.00 8.0 B11-40 2.95 12.00 10.00 52.80 54.00 8.0	A22-40 6.00 12.00 8.20 72.00 36.00 8.0 A22-49 6.00 12.00 8.20 72.00 36.00 8.0 A32-19 6.10 12.00 9.03 73.20 36.00 8.0 A32-22 6.00 12.00 9.38 72.00 36.00 8.0 A32-27 6.00 12.00 9.16 72.00 36.00 8.0	A22-28 6.10 12.00 8.75 73.20 36.00 8.0 A22-31 6.00 12.00 8.06 72.00 36.00 8.0 A22-34 6.00 12.00 8.35 72.00 36.00 8.0 A22-36 6.00 12.00 8.35 72.00 36.00 8.0 A22-36 6.00 12.00 8.35 72.00 36.00 8.0	A21-39 6.00 12.00 8.95 72.00 54.00 8.0 A21-51 6.00 12.00 8.12 72.00 54.00 8.0 A22-20 6.00 12.00 8.15 72.00 36.00 8.0 A22-24 6.00 12.00 8.80 72.00 36.00 8.0	159.2] (0 A14-39 6.00 12.00 8.35 72.00 24.00 8.0 A14-44 6.00 12.00 8.50 72.00 24.00 8.0 A14-55 6.00 12.00 8.53 72.00 24.00 8.0 A14-68 6.00 12.00 8.42 72.00 24.00 8.0 A14-68 6.00 12.00 8.45 72.00 54.00 8.0	b h d Ac acc l IN IN IN IN2 IN IN	ST DATA OF PRESTRESSED
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APPENDIX D (CONTINUED)

4A 8A 7B 7A	C32-22 C32-37 C32-42 C32-50 C32-80	C22-40 C22-46 C22-62 C22-62 C22-73 C32-11	C12-57 C22-29 C22-31 C22-36 C22-36	NNNNN	B32-41 B32-54 C12-09 C12-18 C12-19	B31-15 B32-11 B32-19 B32-31 B32-34	B22-23 B22-30 B22-41 B22-65 B22-68	BEAM MARK		
1.500	1.882	1.75 1.79 1.89 1.75	1.84 1.84 1.77	1.86 1.88 1.75	2.96 2.78 1.75 1.75	3.20 3.12 3.12 3.05 3.05 3.05 3.05 3.05 3.05 3.05 3.05	3.00 1	N Q	TEST	
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I15 3.00 12.00 8.50 55.50 21.25 5.640 0.240 116 3.00 12.00 9.00 55.50 17.00 4.900 0.300 118 3.00 12.00 9.00 55.50 17.00 4.900 0.300 119 3.00 12.00 9.00 55.50 29.75 4.900 0.300 120 3.00 12.00 8.50 55.50 24.00 5.100 0.300 120 3.00 12.00 8.50 55.50 24.00	B9 4.00 8.00 6.00 32.00 30.00 5.140 0.240 B10 4.00 8.00 6.00 32.00 30.00 5.140 0.240 B10 3.00 12.00 9.00 55.50 29.75 5.600 0.300 I11 3.00 12.00 9.00 55.50 29.75 5.000 0.300 I14 3.00 12.00 9.00 55.50 29.75	B3 4.00 8.00 6.00 32.00 21.00 4.230 0.240 B4 4.00 8.00 6.00 32.00 27.00 4.065 0.240 B5 4.00 8.00 6.00 32.00 27.00 4.065 0.240 B6 4.00 8.00 6.00 32.00 28.00 4.380 0.300 B7 4.00 8.00 6.00 32.00 21.00 4.810 0.240	A7 5.00 10.00 7.00 50.00 27.00 4.380 0.300 A8 5.00 10.00 7.00 50.00 29.00 4.950 0.300 A9 5.00 10.00 7.00 50.00 27.00 4.900 0.300 A10 5.00 10.00 7.00 50.00 35.00 4.610 0.300 A12 5.00 10.00 7.00 50.00 28.00 5.060 0.240	A1 5.00 10.00 7.00 50.00 35.00 5.210 0.125 A2 5.00 10.00 7.00 50.00 35.00 50.50 0.126 A4 5.00 10.00 7.00 50.00 24.50 4.180 0.300 A5 5.00 10.00 7.00 50.00 21.00 5.000 0.300 A6 5.00 10.00 7.00 50.00 28.00 4.065 0.240	1.50 10.00 7.09 39.75 47.50 5.040 0.314 2 15A 1.50 10.00 7.09 39.75 15.00 5.040 0.314 2 15B 1.50 10.00 7.09 39.75 30.00 5.040 0.314 2 16 1.50 10.00 7.09 39.75 47.50 5.040 0.314 2	111A 1.50 10.00 7.09 39.75 15.00 5.520 0.314 2 1.1B 1.50 10.00 7.09 39.75 47.50 5.520 0.314 2 1.50 10.00 7.09 39.75 15.00 3.760 0.314 2 1.3A 1.50 10.00 7.09 39.75 15.00 3.760 0.314 2 1.50 10.00 7.09 39.75 30.00 3.760 0.314 2 1.50 10.00 7.09 39.75 30.00 3.760 0.314 2 1.50 10.00 7.09 39.75 30.00	2 1.50 10.00 7.09 39.75 20.00 4.400 0.314 2 1.50 10.00 7.09 39.75 25.00 4.400 0.314 2 1.50 10.00 7.09 39.75 30.00 4.400 0.314 2 1.50 10.00 7.09 39.75 43.00 4.400 0.314 2 98 1.50 10.00 7.09 39.75 85.00 4.400 0.314 2	b h d AC IN IN KSI IN2 IN IN IN2 IN IN KSI IN2 [58.2] (CONTINUED)	TEST DATA OF PRESTRESSED CONCRETE BEAMS	(0)
I15 3.00 12.00 8.50 55.50 21.25 5.640 0.2 I16 3.00 12.00 9.00 55.50 21.25 4.900 0.3 I18 3.00 12.00 8.50 55.50 17.00 4.900 0.3 I19 3.00 12.00 9.00 55.50 29.75 4.900 0.3 I20 3.00 12.00 8.50 55.50 24.00 5.100 0.3	B9 4.00 8.00 6.00 32.00 30.00 5.140 0.240 214.0 B10 4.00 8.00 6.00 32.00 30.00 5.140 0.240 214.0 I10 3.00 12.00 9.00 55.50 29.75 5.000 0.300 214.0 I11 3.00 12.00 9.00 55.50 29.75 5.000 0.300 214.0 I14 3.00 12.00 9.00 55.50 29.75 5.000 0.300 214.0	B3 4.00 8.00 6.00 32.00 21.00 4.230 0.240 214.0 B4 4.00 8.00 6.00 32.00 27.00 4.640 0.240 214.0 B5 4.00 8.00 6.00 32.00 27.00 4.065 0.240 214.0 B6 4.00 8.00 6.00 32.00 28.00 4.380 0.300 214.0 B7 4.00 8.00 6.00 32.00 21.00 4.810 0.240 214.0	A7 5.00 10.00 7.00 50.00 27.00 4.380 0.300 214.0 A8 5.00 10.00 7.00 50.00 29.00 4.900 0.300 214.0 A9 5.00 10.00 7.00 50.00 27.00 4.900 0.300 214.0 A10 5.00 10.00 7.00 50.00 35.00 4.610 0.300 214.0 A12 5.00 10.00 7.00 50.00 28.00 5.060 0.240 214.0	A1 5.00 10.00 7.00 50.00 35.00 5.210 0.126 201.01 A2 5.00 10.00 7.00 50.00 35.00 4.180 0.300 214.0 1 A5 5.00 10.00 7.00 50.00 21.00 5.000 0.300 214.0 1 A6 5.00 10.00 7.00 50.00 28.00 4.065 0.240 214.0 1 A6 5.00 10.00 7.00 50.00 28.00 4.065 0.240 214.0 1 A6 5.00 10.00 7.00 50.00 28.00	1.50 10.00 7.09 39.75 47.50 5.040 0.314 211.6 2 15A 1.50 10.00 7.09 39.75 15.00 5.040 0.314 211.6 2 15B 1.50 10.00 7.09 39.75 30.00 5.040 0.314 211.6 2 16 1.50 10.00 7.09 39.75 47.50 5.040 0.314 211.6 2	111 1.50 10.00 7.09 39.75 15.00 5.520 0.314 211.6 2 111 1.50 10.00 7.09 39.75 47.50 5.520 0.314 211.6 2 12 1.50 10.00 7.09 39.75 15.00 3.760 0.314 211.6 2 138 1.50 10.00 7.09 39.75 30.00 3.760 0.314 211.6 2	2 1.50 10.00 7.09 39.75 20.00 4.400 0.314 211.6 2 1.50 10.00 7.09 39.75 25.00 4.400 0.314 211.6 2 1.50 10.00 7.09 39.75 30.00 4.400 0.314 211.6 2 5 1.50 10.00 7.09 39.75 43.00 4.400 0.314 211.6 2 9B 1.50 10.00 7.09 39.75 85.00 4.400 0.314 211.6 2	h d AC IN IN KSI IN2 KSI IN IN IN IN IN [58.2] (CONTINUED)	TEST DATA OF PRESTRESSED CONCRETE BEAMS FAILED IN	
I15 3.00 12.00 8.50 55.50 21.25 5.640 0.240 214. I16 3.00 12.00 9.00 55.50 17.00 4.900 0.300 214. I18 3.00 12.00 9.00 55.50 17.00 4.900 0.300 214. I19 3.00 12.00 9.00 55.50 29.75 4.900 0.300 214. I20 3.00 12.00 8.50 55.50 24.00 5.100 0.300 214.	B9 4.00 8.00 6.00 32.00 30.00 5.140 0.240 214.0 18.5 B10 4.00 8.00 6.00 32.00 30.00 5.125 0.300 214.0 23.4 B10 3.00 12.00 9.00 55.50 29.75 5.600 0.300 214.0 28.3 I11 3.00 12.00 9.00 55.50 29.75 5.000 0.300 214.0 32.7 I14 3.00 12.00 9.00 55.50 29.75 5.000 0.300 214.0 18.7	B3 4.00 8.00 6.00 32.00 21.00 4.230 0.240 214.0 10.8 B4 4.00 8.00 6.00 32.00 24.00 4.640 0.240 214.0 9.2 B5 4.00 8.00 6.00 32.00 27.00 4.065 0.240 214.0 13.4 B6 4.00 8.00 6.00 32.00 28.00 4.380 0.300 214.0 14.7 B7 4.00 8.00 6.00 32.00 21.00 4.810 0.240 214.0 14.7	A7 5.00 10.00 7.00 50.00 27.00 4.380 0.300 214.0 28.3 A8 5.00 10.00 7.00 50.00 29.00 4.900 0.300 214.0 32.7 A9 5.00 10.00 7.00 50.00 27.00 4.900 0.300 214.0 32.7 A10 5.00 10.00 7.00 50.00 35.00 4.610 0.300 214.0 32.7 A12 5.00 10.00 7.00 50.00 28.00 5.060 0.240 214.0 23.1 A12 5.00 10.00 7.00 50.00 28.00	A1 5.00 10.00 7.00 50.00 35.00 50.50 0.126 201.0 14.55 A2 5.00 10.00 7.00 50.00 35.00 50.50 0.126 201.0 14.55 A4 5.00 10.00 7.00 50.00 24.50 4.180 0.300 214.0 23.5 A5 5.00 10.00 7.00 50.00 21.00 50.00 0.300 214.0 11.0 A6 5.00 10.00 7.00 50.00 28.00 4.065 0.240 214.0 11.0 A6 5.00 10.00 7.00 50.00 28.00	1.50 10.00 7.09 39.75 47.50 5.040 0.314 211.6 25.1 15A 1.50 10.00 7.09 39.75 15.00 5.040 0.314 211.6 25.1 15B 1.50 10.00 7.09 39.75 30.00 5.040 0.314 211.6 25.1 16 1.50 10.00 7.09 39.75 47.50 5.040 0.314 211.6 25.1	11A 1.50 10.00 7.09 39.75 15.00 5.520 0.314 211.6 25.18 1.50 10.00 7.09 39.75 47.50 5.520 0.314 211.6 25.18 1.50 10.00 7.09 39.75 47.50 5.520 0.314 211.6 25.18 1.50 10.00 7.09 39.75 15.00 3.760 0.314 211.6 25.18 1.50 10.00 7.09 39.75 30.00 3.760 0.314 211.6 25.18 1.50 10.00 7.09 39.75 30.00 3.760 0.314 211.6 25.18	2 1.50 10.00 7.09 39.75 20.00 4.400 0.314 211.6 25.18 1.50 10.00 7.09 39.75 25.00 4.400 0.314 211.6 25.18 1.50 10.00 7.09 39.75 30.00 4.400 0.314 211.6 25.18 1.50 10.00 7.09 39.75 43.00 4.400 0.314 211.6 25.18 1.50 10.00 7.09 39.75 55.00 4.400 0.314 211.6 25.18 1.50 10.00 7.09 39.75 81.50 4.400 0.314 211.6 25.18 1.50 10.00 7.00 4.400 0.314 211.6 25.18 1.50 10.00 7.00 39.75 81.50 10.00 7.00 39.75 81.50 10.00 7.00 39.75 81.50 10.00 7.00 7.00 39.75 81.50 10.00 7.00 7.00 39.75 81.50 10.00 7.00 7.00 7.00 39.75 81.50 10.00 7.00 7.00 7.00 7.00 7.00 7.00 7.	h h d AC III IN KSI IN2 KSI KIPS IN IN IN IN2 IN IN KSI IN2 KSI KIPS [58.2] (CONTINUED)	TEST DATA OF PRESTRESSED CONCRETE BEAMS FAILED IN SH	

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APPENDIX E

TEST DATA OF R.C.T-BEAMS WITH WEB REINFORCEMENT AND FAILED IN SHEAR

APPENDIX D (CONTINUED)
TEST DATA OF PRESTRESSED CONCRETE BEAMS FAILED IN SHEAR

BEAM MARK

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Ac IN2

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U6010CE U6013CW U6013CE U6017CW

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U6007HE U6010HE

APPENDIX E (CONTINUED)

BEAM MARK	TEST
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V6004-E U6002-W

V6002-W V6002-E V6004-W

22222

800

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T6018MW T6018ME T9029MW T9029ME

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U6010-W

U6017-E U6017-W

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34.33 34.33 34.33

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R.C.T-BEAMS WITH WEB REINFORCEMENT AND FAILED IN SHEAR 1 fc As Fy Asc Fyc dc Sy Nt CM KG/CM2 CM2 KG/CM2 CM2 KG/CM2CM KG/CM2CM

APPENDIX E (CONTINUED)

OM OM

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	.0 3.0 3.640 0.884 41.60 0.00 0.00 44.0 56.80 .0 3.0 3.782 0.884 41.60 0.17 0.00 44.0 67.40 .0 3.0 3.640 0.884 41.60 0.34 0.00 44.0 58.30 .0 3.0 3.085 0.884 41.60 0.68 0.00 44.0 59.30 .0 3.0 3.085 0.884 41.60 0.85 0.00 44.0 66.80	1.0 3.0 3.782 0.884 41.60 0.51 0.00 44.0 69.20 1.0 3.0 3.640 0.884 41.60 0.61 0.00 44.0 59.80 1.0 3.0 3.782 0.884 41.60 0.77 0.00 44.0 55.00 1.0 3.0 3.782 0.884 41.60 1.02 0.00 44.0 38.80 1.0 3.0 3.640 0.884 41.60 1.53 0.00 44.0 22.10	0.0 3.0 2.690 0.884 41.60 0.61 0.61 40.6 53.80 0.0 3.0 2.790 0.884 41.60 0.61 0.61 40.6 46.80 0.0 3.0 2.920 0.884 41.60 0.61 0.61 40.6 38.80 0.0 3.0 3.180 0.884 41.60 0.61 0.61 40.6 28.60 0.0 3.0 3.180 0.884 41.60 0.61 0.61 40.6 28.60 0.0 3.0 3.270 0.884 41.60 0.61 0.61 40.6 17.50	0.0 3.0 3.040 0.884 41.60 0.86 0.00 44.0 45.20 0.0 3.0 2.920 0.884 41.60 0.86 0.00 44.0 40.60 0.0 3.0 2.920 0.884 41.60 0.86 0.00 44.0 24.60 0.0 3.0 3.190 0.884 41.60 0.86 0.00 44.0 24.60 0.0 3.0 3.280 0.884 41.60 0.86 0.00 44.0 21.50	10 3.0 3.040 0.884 41.60 2.45 0.00 40.6 50.7 1.0 3.0 3.040 0.884 41.60 2.45 0.00 40.6 46.7 1.0 3.0 3.180 0.884 41.60 2.45 0.00 40.6 35.8 1.0 3.0 3.180 0.884 41.60 2.45 0.00 40.6 35.8 1.0 3.0 3.280 0.884 41.60 2.45 0.00 40.6 19.4 1.0 3.0 3.0 3.280 0.884 41.60 2.45 0.00 40.6 19.4 1.0 3.0 3.0 3.280 0.884 41.60 2.45 0.00 40.6 19.4 1.0 3.0 3.0 3.0 3.0 3.0 3.0 3.0 3.0 3.0 3	.0 3.0 2.785 0.884 41.60 0.00 0.86 44.0 56.0 .0 3.0 2.700 0.884 41.60 0.00 0.86 44.0 50.4 .0 3.0 2.880 0.884 41.60 0.00 0.86 44.0 48.4 .0 3.0 3.300 0.884 41.60 0.00 0.86 44.0 31.4 .0 3.0 2.920 0.884 41.60 0.00 0.86 44.0 22.4 .0 3.0 2.920 0.884 41.60 0.00 0.86 44.0 22.4	.0 3.0 3.120 0.884 41.60 0.00 2.45 40.6 53.7 .0 3.0 3.560 0.884 41.60 0.00 2.45 40.6 50.4 .0 3.0 3.080 0.884 41.60 0.00 2.45 40.6 42.6 .0 3.0 3.080 0.884 41.60 0.00 2.45 40.6 36.9 .0 3.0 3.140 0.884 41.60 0.00 2.45 40.6 20.1	l fc As Fy wh wv fwy IN KSI IN2 KSI % % KS	DEEP BEAMS WITH WEB REINFORCEMENT [70.2]
2-B5 1.00 8.88 32.5 17.75 5.65 0.359 1.259 14.0 17.02 2-B6 1.00 8.88 32.5 17.75 6.46 0.359 1.259 0.0 16.25 3-C2 1.00 8.88 32.5 26.63 4.82 0.359 1.259 66.4 17.02 3-C3 1.00 8.88 32.5 26.63 4.88 0.359 1.259 46.6 17.70 3-C4 1.00 8.88 32.5 26.63 4.42 0.359 1.259 34.0 12.67	[71.2] [72-A3 1.00 8.88 32.5 17.75 4.90 0.359 1.259 80.7 18.20 [71-A4 1.00 8.88 32.5 8.88 5.03 0.359 1.259 79.8 28.35 [72-B2 1.00 8.88 32.5 17.75 6.23 0.359 1.259 68.0 22.53 [72-B3 1.00 8.88 32.5 17.75 6.37 0.359 1.259 48.4 22.93 [72-B4 1.00 8.88 32.5 17.75 5.52 0.359 1.259 34.9 17.78	4A1 4.72 21.26 247.4 59.06 5.9 4.57 3.253 119.9 0.491 40.0 100.7 4B1 4.72 21.26 247.4 59.06 5.9 4.41 3.253 120.4 0.506 47.0 102.1 5A0 4.72 21.26 247.4 59.06 5.9 3.73 3.898 105.2 0.518 0.0 97.7 5B0 4.72 21.26 247.4 59.06 5.9 3.86 3.898 106.3 0.497 0.0 97.7	[76.3] 2A3 4.72 21.26 247.4 59.06 5.9 4.73 2.006 185.9 0.474 141.9 113. 2B3 4.72 21.26 247.4 59.06 5.9 4.92 2.006 185.8 0.494 141.4 115. 2B2 4.72 21.26 247.4 59.06 5.9 4.51 2.610 142.6 0.510 94.6 109. 3B2 4.72 21.26 247.4 59.06 5.9 3.99 2.610 142.8 0.480 94.6 97.6	明明	3B 3.00 15.80 102.0 40.00 3.0 6.84 0.902 260.0 0.287 87.7 504B 3.00 15.80 102.0 50.00 3.0 6.34 0.902 260.0 0.150 94.6 395B 3.00 15.80 102.0 50.00 3.0 6.41 0.902 260.0 0.188 87.0 407B 3.00 15.80 102.0 60.00 3.0 6.62 0.902 260.0 0.125 93.7 34.	6 3.00 15.80 102.0 100.00 3.0 6.23 0.902 2 7A 3.00 15.80 102.0 60.00 3.0 6.62 0.902 2 10A 3.00 15.80 102.0 70.00 3.0 7.05 0.902 2 B 3.00 15.80 102.0 48.00 3.0 6.65 0.902 2 1B 3.00 15.80 102.0 30.00 3.0 6.82 0.902 2	[64.2] **PXIA 3.00 15.80 102.0 48.00 3.0 6.65 0.902 260.0 0.117 91.7 32.00 **F-21A 3.00 15.80 102.0 30.00 3.0 6.82 0.902 260.0 0.188 92.3 60.00 **F-22A 3.00 15.80 102.0 40.00 3.0 6.55 0.902 260.0 0.117 86.3 40.00 **F-23A 3.00 15.80 102.0 40.00 3.0 6.84 0.902 260.0 0.084 87.7 40.00 **F-25A 3.00 15.80 102.0 50.00 3.0 6.41 0.902 260.0 0.081 87.0 32.20	l fc As IN KSI IN2	Z F G
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TEST DATA OF DEEP BEAMS WITH WEB REINFORCEMENT [70.2]

APPENDIX F

APPENDIX G

	. '	14 5.90 11.1 15 5.90 10.7 16 5.90 10.7 17 5.90 10.6	12 5.90 11.1 13 5.90 11.1	5.90 10.6 5.90 11.1	8 5.90 10.7	CM-680 2.00 12.0 CH-680 2.00 12.0	000	680 2.00 12.0 6-240 2.00 12.0	000	680 2.00 12.0 0-160 2.00 12.0	6-160 2.00 12 6-80 2.00 12	-240 2.00 12. -240 2.00 12.	3-E3 1.50 8.8 2-F1 1.00 16.6 2-F2 1.00 16.6 2-F3 1.00 16.6 2-F4 1.00 16.6	3-C5 1.00 8.8 3-D1 1.00 8.8 3-D2 1.00 8.8 3-D3 1.00 8.8 3-D4 1.00 8.8	BEAM b d MARK IN IN	TEST DATA
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					2007年10月1日	STORY CONTRACTOR			200 BANKS							
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APPENDIX G (CONTINUED)

AFDELINGEN FOR BERENDE KONSTRUKTIONER DANMARKS TEKNISKE HØJSKOLE

Department of Structural Engineering Technical University of Denmark, DK-2800 Lyngby

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