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Integrated circuits for multimedia applications

Luca Vandi

This thesis is submitted in partial fulfillment of the requirements for obtaining the Ph.D. degree at:

Kongens Lyngby, 27 September 2006

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To Daniela

She made it possible, I just had to get it done.

Abstract

This work presents several key aspects in the design of RF integrated circuits for portable multimedia devices.

One chapter is dedicated to the application of negative-feedback topologies to receiver frontends. A novel feedback technique suitable for common multiplier-based mixers is described, and it is applied to a broad-band dual-loop receiver architecture in order to boost the linearity performances of the stage. A simplified noise- and linearity analysis of the circuit is derived, and a comparison is provided with a more traditional dual-loop topology (a broad-band stage based on shunt-series feedback), showing a difference in compression point in the order of 10dBm for the same power consumption. The same principle is also applied to a more conventional narrow-band stage in which a single loop is employed in order to enhance noise performances. Noise analysis shows sensible improvements in the noise figure (up to $\sim 1 \, \mathrm{dB}$) in low-performance technologies, when stringent specifications are considered in terms of power consumption.

A recently-reported current-reuse technique, applied to a complete RX frontend, is examined in the following chapter in order to sketch a simplified numerical analysis for the performances of the stage. Semi-ideal models are used in simulations to validate the derived calculations, and the fundamental limits of the basic structure are discussed. The design of a current-mode base-band output stage implemented in a $0.13\mu m$ technology is presented: the amplifier draws $\sim 500\mu A$ from a 1.2V supply,

providing 35dB gain and 135MHz GBWP.

The integration of high-performance passive components is studied in the last chapter which presents the first reported toroidal inductor fabricated in a standard CMOS process. Field-confinement properties of the structure are exploited in order to reduce the impact of substrate-induced currents. Basic models are derived in the design phase, and the technological limits of the device are considered. Measurement results show that a very compact coil can provide $\sim 1 \, \mathrm{nH}$ inductance up to $20 \, \mathrm{GHz}$ (physical limit for the measurement equipment), with a peak quality factor around $10 \, \mathrm{at} \, 15 \, \mathrm{GHz}$.

Resumé

I afhandlingen præsenteres forskellige nøgle-aspekter om design af RF integrerede kredsløb for bærbare multimedia-apparater.

Et kapitel omhandler negativ-feedback topologier til modtager-frontends. En ny feedback-teknik, der passer til almindelige multiplikator-baserede mixere, beskrives og anvendes til en bredbåndsdobbeltsløjfe-modtager-arkitektur for at forbedre trinets linearitet. En forenklet støj- og linearitetsanalyse af kredsløbet udledes, og en sammenligning gives med en mere traditionel dobbeltsløjfe-konfiguration (et bredbåndstrin baseret på shuntseries feedback), der viser en forskel i kompressionspunktet i nærheden af 10dBm med samme strømforbrug. Det samme princip anvendes til et mere konventionelt smalbåndstrin: en enkeltsløjfe benyttes for at forbedre støjegenskaber. En støjanalyse giver mærkbare forbedringer i støjtal (op til $\sim 1 \, \mathrm{dB}$) i lavtydende teknologier, når strenge specifikationer med hensyn til strømforbrug tages i betragtning.

En nyligt rapporteret strømgenbrugsteknik (anvendt til en komplet RX-frontend) undersøges i det næste kapitel for at skitsere en simplificeret numerisk analyse af trinets ydelser. Semi-ideelle modeller bruges i simuleringer for at underbygge udledte beregninger, og strukturens grundlæggende grænser klarlægges. Designet af et current-mode base-band udgangstrin realiseret i 0.13μ m-teknologi fremvises: forstærkeren trækker $\sim 500\mu$ A fra en 1.2V-forsyning og giver 35dB forstærkning sammen med 135MHz GB-WP.

Integrationen af højtydende passive komponenter undersøges i det sidste kapitel, der introducerer den første rapporterede toroidal-induktionsspole fabrikeret i standard CMOS teknologi. Strukturens feltsbegrænsningsegenskaber udnyttes for at nedsætte virkningen på substrat-induceret strøm. Grundmodeller udledes i designfasen, og komponentens teknologiske grænser behandles. Måleresultater viser, at en meget kompakt spole kan levere $\sim 1 \, \mathrm{nH}$ induktans op til 20GHz (fysisk grænse for måleapparatet) med en maksimum kvalitetsfaktor på omkring 10 ved 15GHz.

Acknowledgements

"Beggar that I am, I am even poor in thanks; but I thank you. [...]" William Shakespeare, *Hamlet*

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Despite the fact that there is only one name on the front page, this is not a one-man job. I take full responsibility over everything I wrote, but nothing could have been accomplished without help. I'd never be able to mention all the people who have contributed to this work in some way or another, but each of them deserves my complete gratitude.

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- the Microelectronics Lab at University of Pavia (Italy), led by prof. Rinaldo Castello, which welcomed me as a visiting Ph.D. student during my external sojourn,

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List of Symbols and Acronyms

2.5D	Quasi three-dimensional	
3D	Three-dimensional	
AC	Alternate current	
ASIC	Application-specific integrated circuit	
$BW_{\text{-3dB}}$	Bandwidth at -3dB	
C_{gd}	Gate-drain capacitance of a MOSFET in saturation region	[F]
C_{gs}	Gate-source capacitance of a MOSFET in saturation region	[F]
CMFB	Common-mode feedback	
CMOS	Complementary MOS	
C'_{ox}	Field-oxide specific capacitance of a MOS transistor	$[F \cdot m^{-2}]$
$\mathit{CP}_{\mathrm{1dB,in}}$	Input-referred 1dB-compression point	[dBm]
$CP_{1\mathrm{dB,out}}$	Output-referred 1dB-compression point	$[\mathrm{dBm}]$
DC	Direct current	

Units are indicated according to the International System (SI), apart from electrical power which is indicated in dBm.

DSB	Double sideband	
F	Noise factor	
f_{\max}	Maximum oscillation frequency of a MOS device	[Hz]
$f_{ m self}$	Self-resonance frequency of an integrated inductor	[Hz]
f_T	Unity-gain frequency of a MOSFET	[Hz]
GBWP	Gain-bandwidth product	[Hz]
g_{ds}	Drain-source output conductance of a MOSFET	[A/V]
G_{loop}	Closed-loop gain for a feedback system	[dB]
g_m	Transconductance of a MOSFET	[A/V]
GSG	Ground-Signal-Ground	
IC	Integrated circuit	
I_D	Drain current of a MOSFET	[A]
IF	Intermediate frequency	
IIP_3	Input-referred $3^{\rm rd}$ -order intercept point	[dBm]
$\langle i_{\rm n,in}^2\rangle$	Input-referred current-noise density	$[{\rm A^2/Hz}]$
J	Imaginary unit	
k_B	Boltzmann constant	$[\mathrm{J/K}]$
L	Gate length of a MOS transistor	[m]
LC	Inductor-capacitor	
$L_{ m min}$	Minimum gate length of a MOS transistor for a given technology	[m]
LMV	LNA-Mixer-VCO	
LNA	Low-noise amplifier	
LO	Local oscillator	
MGTR	Multiple-gated transistors	
MOS	Metal-Oxide-Semiconductor	
OSFET	MOS field-oxide transistor	

N_a	Acceptors' concentration in pn -junctions	$[m^{-3}]$		
N_d	N_d Donors' concentration in pn -junctions			
NF	Noise figure	[dB]		
nMOS	n-channel MOS transistor			
$Op ext{-}Amp$	Operational amplifier			
PAC	Periodic alternate current			
pMOS	p-channel MOS transistor			
PN	Phase noise	[dBc/Hz		
PSS	Periodic steady-state			
Q	Quality factor of a resonator			
q	Charge of the electron	[C]		
RF	Radio frequency			
RX	Receiver			
s	Generic complex variable in the Laplace domain	[Hz]		
SC	Switched capacitors			
SOI	Silicon on insulator			
SOM	Self-oscillating mixer			
SSB	Single sideband			
T	Temperature	[K]		
t_{ox}	Gate-oxide thickness	[m]		
VCO	Voltage-controlled oscillator			
V_{DD}	Supply voltage of a circuit	[V]		
V_{GS}	Gate-source voltage of a MOSFET \dots	[V]		
VLSI	Very large scale of integration			
$\langle v_{\rm n,in}^2 \rangle$	Input-referred voltage-noise density \dots	$[{ m V}^2/{ m Hz}]$		
V_T	Threshold voltage of a MOS transistor	[V]		
W	Gate width of a MOS transistor	[m]		
WLP	Wafer-level packaging			

x_j	pn-junction depth of MOS systems	[m]
z	Generic complex variable in the ${\mathcal Z}$ -domain	
γ	Drain-current thermal noise parameter for MOSFETs	
ϵ_{ox}	Dielectric constant in oxide	[F/m]
θ	Vertical-field mobility-reduction parameter of a MOSFET	$[V^{-1}]$
λ	Scaling factor for CMOS technology	
μ_0	Magnetic permeability in vacuum	[H/m]
μ_{n0}	Nominal carrier mobility for nMOS transistors	$[\mathrm{m}^2 \cdot (\mathrm{V \cdot s})^{-1}]$
ϕ_m	Phase-margin of a feedback loop	[deg]
ϕ_{th}	Thermal potential	[V]
ω_T	Unity-gain angular frequency of a MOSFET	[rad/s]

CHAPTER 1

Introduction

"Gallia est omnis divisa in partes tres [...]"
"All Gaul is divided into three parts [...]"
Caius Iulius Caesar, Commentarii De Bello Gallico

Many times since its formulation, *Moore's Law* [2] has been under discussion, but it proved itself quite impervious to attacks. The main reason behind the exponentially-growing number of components¹ which can be crammed onto the same IC is the fact that semiconductor technologies offer devices with smaller and smaller minimum dimensions. Major crisis points have been faced and overcome in multiple occasions, and the improvement rate has been constant [3]; in recent years, the progress of photolithography (among other things) has been so fast that technology enhancements have repeatedly beaten expectations, and forecasts for the near-term future are more optimistic than ever [4], as shown in Figure 1.1. The reduction of minimum dimensions is not enough to explain the astounding progress in VLSI technology: many other parameters have to scale in order to guarantee a reliable development path. It is interesting to notice, though, that

• Moore's Law had been conceived almost a decade before a proper scaling theory was devised [5], and

¹Transistors, in many practical cases.

2 Introduction

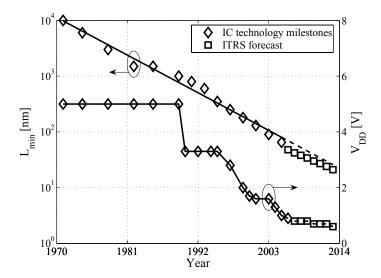


Figure 1.1: Historical data and projections for IC technology development in terms of minimum gate length (L_{min}) and supply voltage (V_{DD}) . For L_{min} a linear fit is indicated for reference.

• the scaling theory itself has not been followed thoroughly, since for example supply voltage has been kept constant for a very long time for no reason but compatibility, as shown in Figure 1.1.

The key-item in scaling theory is the concept of constant field, which ensures faster switching of logic gates, lower energy per transition, cheaper devices. In Table 1.1 the main parameters are associated with their respective scaling factor, function of the common factor $\lambda > 1$. RF CMOS designers have to face the benefits and the drawbacks of a scaling technique which is targeting purely digital applications.² Key aspects like operating frequency, noise and linearity (among others) are heavily influenced by the scaling strategy.

1. Faster logic gates means for certainty higher cut-off frequency f_T and maximum oscillation frequency f_{max} , and this has a major positive

² "[...] the process is optimized and characterized primarily for one tradeoff: that between speed and power dissipation." [6]

Parameter	Scaling factor
t_{ox}	$\frac{1}{\lambda}$
L	$\frac{1}{\lambda}$
W	$\frac{1}{\lambda}$
x_{j}	$\frac{1}{\lambda}$
N_a	λ
N_d	λ
V_T	$\frac{1}{\lambda}$
V_{DD}	$\frac{1}{\lambda}$

Table 1.1: Scaling behaviour according to the constant-field theory.

impact on RF performances [7]. There is also a very peculiar side-effect, though: in inductively-degenerated LNAs the inductance on the source must be inversely proportional to f_T in order to provide correct input impedance matching, and today designers are finding increasing difficulties in tuning this passive component without sacrificing noise performances.

- 2. Noise performances can be also expected to be positively influenced by scaling, in terms of minimum noise figure as well as sensitivity to deviation from optimal noise matching [8]. Therefore the design of LNAs should be easier, because lower noise levels can be achieved over a wider range of design parameters (bias, transistors' dimensions and so on). It should be noted however that VCO designers will experience serious troubles coping with $\frac{1}{f}$ noise. Gate-oxide thickness must be scaled together with L_{\min} ; being a few percent of the channel length, t_{ox} is approaching the sub-nanometre scale, but a thickness below 1nm is not enough to keep direct tunnelling currents at reasonable levels [9]. The introduction of materials with high dielectric constant (so-called "high- κ materials") can alleviate this problem, but integration of such structures is still difficult [10], and the poor quality of these insulators leads to bad noise performances [11].
- 3. Linearity is indeed affected by scaling, but in a negative way. It can be shown that optimal biasing for maximum linearity moves toward higher and higher currents [12], and this is in contrast with

4 Introduction

the natural tendency to reduce power consumption. The reduction of supply voltage has also a dramatic negative impact on linearity performances.

4. Scaling of threshold voltage has not been very aggressive, and it is probably close to its limits, due to the fact that the value of V_T is intimately connected to thermal potential $\phi_{th} = k_B T/q$ [13], therefore apart from reducing operating temperature³ there are very few chances to overcome this limitation on the long run. Subthreshold operation is widely popular in low-frequency analogue circuits, but it still must be fully exploited in RF, and this is probably one of the main challenges for the incoming years.

The amount of bandwidth in excess and the improvement of high-frequency noise performances suggest the possibility to employ negative feedback in order to relax linearity design constraints, particularly in the design of RX front-ends. This goes in the direction of implementing a broadband stage that is capable of handling several communication standards at a time, thereby providing a solution that can serve different applications and is easy to migrate from a technology node to the next, since it relies on non-tuned passives (i.e. resistors) and on transistors' performances that are continuously improving. In Chapter 2 we describe a novel approach to negative feedback applied to Gilbert-cell mixers, that leads to a dual-loop RX front-end with broad-band operation and high linearity. The same principle is also applied to a more traditional single-loop narrow-band stage.

Reduction of supply voltage is simply not enough in order to guarantee a continuous path of lower and lower power consumption, especially considering battery-operated implementions: the number of features increases at a pace hardly sustainable by digital CMOS roadmap. Several design techniques try to address this issue, and current-reusing is definitely one of the most interesting. Recently the first implementation of a stacked structure comprising LNA, mixer and VCO biased by the same current has been proposed [1]. The difficulties of providing the correct trade-offs between all different aspects of RF design just sum up when three different blocks are sharing the same devices (active and passive); the robustness of the design must be thoroughly evaluated, too. In Chapter 3 a simplified analysis is proposed in order to grasp the basic behaviour of the stage and

³This is also a viable option [14], but not for mainstream consumer products.

verify the functionality of the circuit at first order.

The great breakthrough of RF CMOS technology is the possibility of integrating a large amount of digital functions on the same chip together with analogue circuitry, but integration of passive components remains a key necessity in modern radios. Even though we have not examined the subject in detail, we can assume that newer technologies will provide better performances in terms of inductors, varactors and the like [8]. Sharing the same silicon substrate is always a matter of isolating critical parts of the circuit from each other, but this turns out particularly difficult when several coils are integrated on the same die. In the quest for the integration of the perfect component (compact, scalable, not prone to inducing currents in the substrate), an implementation of a toroidal inductor is proposed in Chapter 4 in order to demonstrate the superior field-confinement properties of this structure, which has never been implemented in standard CMOS technologies before. Design considerations are presented together with basic modelling and measurement results.

Some conclusions are drawn in the final chapter.

Chapter 2

Receiver front-ends based on negative feedback

Traditional active mixers based on the Gilbert cell [15], like the one depicted in Figure 2.1a, are ubiquitous in modern RF architectures. The design of such cells is carried out under the general assumption that the linearity of the block is intimately related to the linearity of the transconductor (M_1 in the figure). To improve linearity performances, several solutions have been devised, most of which are based on the following methods [16]: feedback, feedforward, predistortion and piecewise approximation. In practice many applications exploit a combination of these techniques. A non-exhaustive overview of the most popular approaches includes:

- source degeneration, which is the most basic form of negative feedback.
- 2. Stages based on multiple-gated transistors (MGTR) combine opposing nonlinear contributions in order to cancel non-linearities, thereby resembling the feedforward approach [17].
- 3. Several implementations rely on the MOS I-V characteristic to pre-

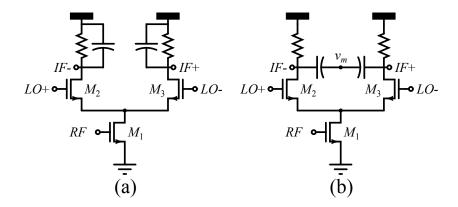


Figure 2.1: A down-conversion mixer based on the Gilbert cell: conventional arrangement (a) and modified topology (b).

distort the input signal and rebuild it at the output, by using for example a transconductor based on a cross-quad cell [16].

4. The so-called CMOS " g_m -cell", which is a variant of the bipolar-based "multi-tanh" technique [18] that exploits perfectly the piecewise approximation approach.

In Table 2.1 several solutions found in the literature are presented. Whenever possible, simulation results have been listed, rather than measurement results. Some implementations exploit specific linearization techniques, others are chosen among the present state-of-the-art, in order to provide useful comparisons. Feedback-based solutions have been traditionally relegated to low-frequency applications, because of the intrinsic bandwidth reduction associated with the feedback. As more advanced technologies become available, bandwidth limitations become less and less critical. To fully exploit the capabilities of negative feedback, it would be beneficial to embrace the complete cell (transconductor and switching pair) into the feedback loop, rather than closing the loop locally by means of source degeneration.

In Figure 2.1b a slightly-modified multiplier cell is depicted: as explained in the following, the node v_m offers an unmodulated amplified replica of the input RF signal. This node will be fed back into the loop. The actual implementation of the load might differ, since many different alternatives are feasible, but the basic concept is well illustrated in the figure.

Table 2.1: Reported implementations of down-convertion mixers based on the Gilbert cell.

	Method		Gain	NF	$I_{ m BIAS}@V_{ m DD}$	${ m IIP_3@freq.}$		
	1	2	3	4	[dB]	[dB]	[mA] [V]	[dBm] [GHz]
[19]	•				5	18	3.8@2	13@0.9
[20]	•				9.48	17.69	4.7@2	3@2.4
[21]	•				10	24.2	0.6@1	6@2
[22]	•		•		-2.8	N/A	1.4@2.5	-20@N/A
[23]	•				16	13.1	4@1.8	9@2.1
[24]				•	26.6	11.9	1.2@1.5	-0.1@2.4
[25]				•	16.6	16.5	5@1.8	12.12@2.1
[17]		•			16.5	17.2	3@1.8	9@2.4
[26]		•			1.7	14.5	4.4@2	4.9@2.4
[27]		•			-6.5	N/A	3.2@1.25	15@0.9
[28]			•		3.3	14.87	3.7@1.5	5.46@2.4
[29]			•		-9.1	N/A	2.5@1.2	9.5@0.9
[30]		•			6.5	11	6.9@1	3.5@8
[31]			•		2	13.5	5.2@0.9	3.5@0.9
[32]		•			4.7	N/A	1.8@1	10.3@5.5
[33]	•	•			6.6	21.4	N/A@1.8	1.5@1.69
[34]					6	18.5	7@3.3	11.5@1.9
[35]					2	N/A	N/A@2	20@24
[36]					8.3	24.5	5.5@0.9	0@5.25
[37]					12.76	12.4	1.8@1.8	-6.9@5.7
[38]					7	N/A	N/A@3.3	4.5@30
[39]					11.9	13.9	3.2@1	-3@2.4

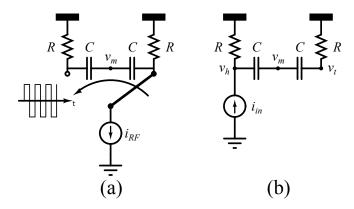


Figure 2.2: Full-switched model of a differential pair used as a mixer (a) and R-C network resembling the passive load (b).

In Figure 2.2a we have reproduced the cell in Figure 2.1b by modelling the differential pair $(M_1 - M_2)$ as a couple of switches, controlled by a time-varying periodic function at frequency f_{LO} ; we indicate the function as $m_{LO}(t)$. The computation of the output voltage can be developed with the aid of Figure 2.2b. We can easily show that the following relations hold:

$$v_h = i_{in} R \frac{1 + s \frac{CR}{2}}{1 + sCR}, (2.1)$$

$$v_t = i_{in} \frac{R}{2} \frac{sCR}{1 + sCR}, \tag{2.2}$$

$$v_m = i_{in} \frac{R}{2}. (2.3)$$

We can then define equivalent impedances, as follows:

$$Z_h \triangleq \frac{v_h}{i_{in}} = R \frac{1 + \frac{sCR}{2}}{1 + sCR},$$
 (2.4)

$$Z_t \triangleq \frac{v_t}{i_{in}} = \frac{R}{2} \frac{sCR}{1 + sCR}, \tag{2.5}$$

$$Z_m \triangleq \frac{v_m}{i_{in}} = \frac{R}{2}. (2.6)$$

The application to the structure in Figure 2.2a is straightforward, due to the symmetry properties:

$$v_{IF} = -i_{RF} \frac{1 + m_{LO}}{2} Z_h - i_{RF} \frac{1 - m_{LO}}{2} Z_t - \left(-i_{RF} \frac{1 + m_{LO}}{2} Z_t - i_{RF} \frac{1 - m_{LO}}{2} Z_h\right). \tag{2.7}$$

By simplification, we get:

$$v_{IF} = -i_{RF} \cdot m_{LO} \left(Z_h - Z_t \right)$$

$$= -i_{RF} \cdot m_{LO} \cdot \frac{R}{1 + sCR}, \qquad (2.8)$$

which is exactly the result that we expected: the structures in Figure 2.1a and 2.1b are equivalent with respect to the down-converted signal. The fundamental property of middle point v_m is revealed by the following equation:

$$v_m = -i_{RF} \frac{1 + m_{LO}}{2} Z_m - i_{RF} \frac{1 - m_{LO}}{2} Z_m$$
$$= -i_{RF} \frac{R}{2}. \tag{2.9}$$

The input RF current is amplified without any kind of filtering nor mixing. This voltage can be used in active feedback circuit topologies, as we will show in the rest of the chapter. More detailed calculations related to the properties of node v_m can be found in Appendix A.

2.1 Dual-loop implementation

Traditional single-stage broad-band LNAs for RF applications present a fundamental trade-off between input matching, noise figure and linearity. The best possible compromise might turn out to be simply not good enough for the particular application. To break the trade-off, several solutions have been proposed, among which feed-forward [40] and negative feedback should be mentioned. The latter is quite popular, but certain limitations occur when a single-loop feedback is implemented. As an example, we can examine the simplified schematic in Figure 2.3. By inspection we can easily get the following result:

$$Z_{in} = \frac{R_f}{1+A}. (2.10)$$

In order to provide correct input matching, the forward gain A must be carefully controlled. By imposing $Z_{in} = R_S$ we get:

$$G_{loop} = -\frac{A}{A+2}, \tag{2.11}$$

$$\frac{v_{out}}{v_S} = -\frac{A}{2}. (2.12)$$

The loop gain can not be arbitrarily large, in fact it is going to be smaller than unity. In these conditions, negative feedback does not provide benefits in terms of linearity. For all these reasons, a dual-loop solution is a more sensible choice, but without proper care for the loop properties a dual-loop configuration can not guarantee any benefit by itself. As a practical case, we can study the shunt-series stage depicted in Figure 2.4, which resembles the topology in Figure 2.3 and does exhibit a dual loop: the source degeneration (R_1 in the figure) provides a stable forward gain (in principle $A \approx \frac{R_L}{R_1}$) and the shunt feedback (given by R_f) defines the correct input matching. However, some limitations occur in practical implementations. We can easily determine the following results:

$$Z_{in} = \frac{(1 + g_{m_1} R_1)(R_f + R_L)}{1 + g_{m_1}(R_1 + R_L)}, \tag{2.13}$$

$$\frac{v_{out}}{v_S} = \frac{Z_{in}}{Z_{in} + R_S} R_L \frac{1 + g_{m_1}(R_1 - R_f)}{(1 + g_{m_1}R_1)(R_f + R_L)}, \tag{2.14}$$

$$G_{loop} = -\frac{g_{m_1} R_L R_S}{(1 + g_{m_1} R_1)(R_f + R_L + R_S)}. (2.15)$$

If we impose perfect matching $(Z_{in} = R_S)$, we can determine the value of R_f as a function of R_S , R_1 , R_L and g_{m_1} . By imposing $R_f > 0$, we can easily determine a practical limitation for the value of the degeneration resistance: $R_1 < R_S$. Since the loop-gain of the series feedback is equal

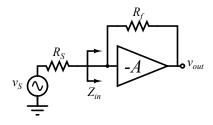


Figure 2.3: A simple negative-feedback stage.

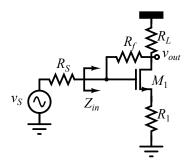


Figure 2.4: A shunt-series input stage.

to $g_{m_1}R_1$, given for example $R_S = 50\Omega > R_1$ it turns out quite difficult to boost such loop-gain way above unity. So the series feedback alone can not provide much benefit in terms of linearization. On the other hand, the same matching condition $Z_{in} = R_S$ affects G_{loop} as well, and the expression in (2.15) becomes:

$$G_{loop} = -\frac{g_{m_1}R_L}{g_{m_1}R_L + 2 + 2g_{m_1}R_1},$$
 (2.16)
 $|G_{loop}| < 1.$

We can conclude that any dual-loop topology must provide more degrees of freedom in order to accommodate strict input-matching requirements and highly-demanding linearity specifications.

The proposed topology is shown in Figure 2.5. Both input stage and mixer exhibit local feedback implemented as source degeneration (R_{s1}, R_2) , but the circuit exploits also a dual-loop feedback involving both stages, provided by the node v_m and the resistor R_3 . Node v_m can be directly fed back into the loop, since there is no DC coupling. Loads R and R_L can be easily implemented by means of active devices, but in the following we will treat them as purely resistive loads. We will show that the dual loop has a high and controllable loop-gain, leading to a robust configuration with good properties in terms of linearity and input matching.

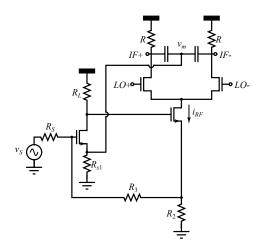


Figure 2.5: A dual-loop high-linearity stage.

2.1.1 Circuit analysis

To simplify the analysis of the stage, we consider a stripped-down version operating at a single frequency: as explained in Appendix A, if we neglect the impedance related to capacitance C we can define an equivalent resistance $R_1 \triangleq R_{s1} / \frac{R}{2}$ and disregard the mixing operation with respect to RF signals. The simplified schematic is shown in Figure 2.6: the output signal is actually i_2 , which is the RF current that is going to be down-mixed and converted into a differential voltage signal. It might be useful in this case to derive a semi-ideal model based on a single nullor (see Appendix B). The nullor-based schematic is shown in Figure 2.7. Due to the (ideal) properties of the nullor, we easily find:

$$i_2 = \frac{v_a}{R_1},$$
 (2.17)

$$i_S = \frac{v_a - R_3 \cdot i_S}{R_2} + i_2,$$
 (2.18)

from which we can obtain the basic results:

$$Z_{in} = \frac{v_a}{i_S} = \left(1 + \frac{R_3}{R_2}\right) \frac{R_1 R_2}{R_1 + R_2},$$
 (2.19)

$$Z_{in} = \frac{v_a}{i_S} = \left(1 + \frac{R_3}{R_2}\right) \frac{R_1 R_2}{R_1 + R_2}, \qquad (2.19)$$

$$i_2 = v_S \frac{1}{R_1 \left(1 + \frac{R_S}{Z_{in}}\right)}. \qquad (2.20)$$

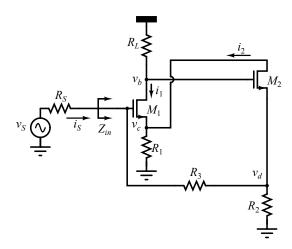


Figure 2.6: Simplified stage for first-order analysis.

Reverting to Figure 2.6, we can derive a proper set of equations:

$$v_a = v_S - R_S i_S, (2.21)$$

$$v_b = -R_L i_1, (2.22)$$

$$v_c = R_1(i_1 - i_2),$$
 (2.23)

$$v_d = R_2(i_S + i_2), (2.24)$$

$$i_S = \frac{v_a - v_d}{R_3},$$
 (2.25)

$$i_1 = g_{m_1}(v_a - v_c),$$
 (2.26)

$$i_2 = -g_{m_2}(v_b - v_d). (2.27)$$

It is important to notice that these equations, although they do not include but the most ideal parameters, can be easily modified to take parasitics into account, since the number of relevant signals do not change. By solving the equations, we find:

$$R_{\Delta} \triangleq (R_2 + R_3) \{ (1 + g_{m_1} R_1) [1 + g_{m_2} (R_2 /\!\!/ R_3)] + g_{m_1} g_{m_2} R_1 R_L \},$$
(2.28)

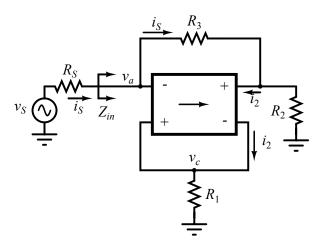


Figure 2.7: Dual-loop stage based on the nullor representation.

$$Z_{in} = \frac{R_{\Delta}}{(1 + g_{m_1}R_1)(1 + g_{m_2}R_2) + g_{m_1}g_{m_2}R_L(R_1 + R_2)},$$
 (2.29)

$$v_a = v_S \frac{Z_{in}}{Z_{in} + R_S}, (2.30)$$

$$v_b = \frac{-v_a g_{m_1} R_L \left\{ R_3 + R_2 \left[1 + g_{m_2} (R_3 - R_1) \right] \right\}}{R_{\Delta}}, \tag{2.31}$$

$$v_c = \frac{v_a R_1 \left\{ g_{m_2} R_2 + g_{m_1} (R_2 + R_3) \left[1 + g_{m_2} (R_L + R_2 / / R_3) \right] \right\}}{R_{\Delta}}, (2.32)$$

$$v_d = \frac{v_a R_2 \left\{ 1 - g_{m_1} \left[g_{m_2} R_L (R_3 - R_1) - R_1 \right] \right\}}{R_{\Delta}}, \tag{2.33}$$

$$i_1 = \frac{v_a g_{m_1} \left[R_2 + R_3 + g_{m_2} R_2 (R_3 - R_1) \right]}{R_{\Delta}}, \tag{2.34}$$

$$i_2 = \frac{g_{m_2} v_a \left\{ g_{m_1} R_3 R_L + R_2 \left[1 + g_{m_1} (R_1 + R_L) \right] \right\}}{R_{\Delta}}.$$
 (2.35)

Simplifications can be made assuming that R_L is adequately large, so that $g_{m_1}R_L, g_{m_2}R_L \gg 1$, and $R_1, R_2 \ll R_L$. Under these assumptions, we get:

$$Z_{in} \approx \left(1 + \frac{R_3}{R_2}\right) \frac{R_1 R_2}{R_1 + R_2},$$
 (2.36)

$$i_2 \approx v_S \frac{1}{R_1 \left(1 + \frac{R_S}{Z_{in}}\right)},$$
 (2.37)

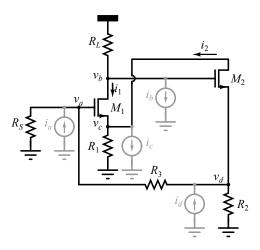


Figure 2.8: Simplified stage for first-order noise calculations.

which can be directly compared with (2.19) and (2.20). Deriving the loopgain is not obvious, since there are actually two feedback loops. It is anyway easy to notice that by removing the connection from node v_b to the gate of transistor M_2 we can cut the outer loop. The total loop-gain can be determined directly:

$$G_{loop} = -\frac{g_{m_1}g_{m_2}R_L[R_2R_S + R_1(R_2 + R_3 + R_S)]}{(1 + g_{m_1}R_1)\{R_3 + R_S + R_2[1 + g_{m_2}(R_3 + R_S)]\}}.$$
 (2.38)

2.1.2 Noise

Some rough noise calculations can be performed with the help of Figure 2.8. It is easy to demonstrate that all relevant noise sources can be seen as linear combination of the currents depicted in the figure, namely i_a , i_b , i_c and i_d . First of all we have to calculate the transfer with respect to these sources:

$$i_2 = G_a \cdot i_a + G_b \cdot i_b + G_c \cdot i_c + G_d \cdot i_d,$$
 (2.39)

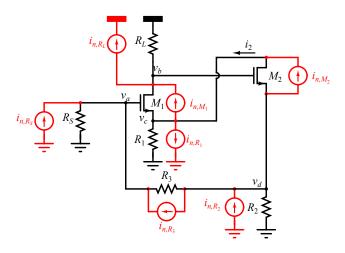


Figure 2.9: Simplified stage with noise sources.

$$G_a = \frac{\frac{R_S Z_{in}}{Z_{in} + R_S} \cdot g_{m_2} \left\{ g_{m_1} R_3 R_L + R_2 \left[1 + g_{m_1} (R_1 + R_L) \right] \right\}}{R_{\Delta}}, \quad (2.40)$$

$$G_b = \frac{Z_{in}}{Z_{in} + R_S} \cdot g_{m_2} R_L (R_2 + R_3 + R_S) (1 + g_{m_1} R_1)}{R_{\Delta}}, \tag{2.41}$$

$$G_{c} = \frac{\frac{Z_{in}}{Z_{in} + R_{S}} \cdot g_{m_{1}} g_{m_{2}} R_{1} R_{L} (R_{2} + R_{3} + R_{S})}{R_{\Delta}}, \qquad (2.42)$$

$$G_{d} = \frac{\frac{Z_{in}}{Z_{in} + R_{S}} \cdot g_{m_{2}} R_{2} \left[(R_{3} + R_{S})(1 + g_{m_{1}} R_{1}) + g_{m_{1}} R_{S} R_{L} \right]}{R_{\Delta}}, \qquad (2.43)$$

$$G_d = \frac{\frac{Z_{in}}{Z_{in} + R_S} \cdot g_{m_2} R_2 \left[(R_3 + R_S)(1 + g_{m_1} R_1) + g_{m_1} R_S R_L \right]}{R_{\Delta}}, \quad (2.43)$$

where of course R_{Δ} and Z_{in} are defined as in (2.28)-(2.29). In Figure 2.9 the main noise sources are depicted. It is then easy to evaluate the different contributions:

$$i_{2} = G_{a} \cdot i_{n,R_{S}} + G_{c} \cdot i_{n,R_{1}} + G_{d} \cdot i_{n,R_{2}} + (G_{a} - G_{d}) \cdot i_{n,R_{3}} + (G_{c} - G_{b}) \cdot i_{n,M_{1}} - (G_{c} + G_{d}) \cdot i_{n,M_{2}} + G_{b} \cdot i_{n,R_{L}}.$$
(2.44)

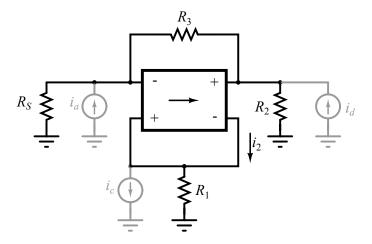


Figure 2.10: Noise analysis of the stage based on the nullor representation.

The noise factor can then be derived directly, considering that all these noise contributions will sum up with the signal:

$$F = 1 + \frac{G_c^2 \cdot \overline{i_{n,R_1}^2} + G_d^2 \cdot \overline{i_{n,R_2}^2} + (G_a - G_d)^2 \cdot \overline{i_{n,R_3}^2} + G_b^2 \cdot \overline{i_{n,R_L}^2}}{G_a^2 \cdot \overline{i_{n,R_S}^2}} + \frac{(G_c - G_b)^2 \cdot \overline{i_{n,M_1}^2} + (G_c + G_d)^2 \cdot \overline{i_{n,M_2}^2}}{G_a^2 \cdot \overline{i_{n,R_S}^2}}.$$

$$(2.45)$$

The above expression can be greatly simplified by means of nullor approximation. First of all, by inspection we find out that node v_b is actually *inside* the nullor itself. From (2.39) and the following, we derive:

$$\frac{G_b}{G_c} = \frac{1 + g_{m_1} R_1}{g_{m_1} R_1}. (2.46)$$

We can refer to Figure 2.10 to derive new expressions for current gains:

Source	Contribution
R_S	1
R_1	$\frac{R_1}{R_S} \left(1 + \frac{R_S}{R_2 + R_3} \right)^2$
R_2	$\frac{R_2}{R_S} \left(\frac{R_S}{R_2 + R_3} \right)^2$
R_3	$\frac{R_3}{R_S} \left(\frac{R_S}{R_2 + R_3} \right)^2$
R_L	$\frac{1}{g_{m_1}^2 R_L R_S} \left[\left(1 + g_{m_1} R_1 \right) \cdot \left(1 + \frac{R_S}{R_2 + R_3} \right) \right]^2$
M_1	$\frac{\gamma}{g_{m_1}R_S} \left(1 + \frac{R_S}{R_2 + R_3}\right)^2$
M_2	$\gamma \cdot \frac{g_{m_2}}{R_S} \left[\frac{R_1 \cdot (R_2 + R_3 + R_S) + R_S R_2}{R_2 + R_3} \right]^2$

Table 2.2: Different noise contributions in the dual-loop topology.

$$G_a \approx \frac{R_S \cdot (R_2 + R_3)}{R_S R_2 + R_1 \cdot (R_2 + R_3 + R_S)},$$
 (2.47)

$$G_b \approx \frac{1 + g_{m_1} R_1}{g_{m_1} R_1} \cdot \frac{R_1 \cdot (R_2 + R_3 + R_S)}{R_S R_2 + R_1 \cdot (R_2 + R_3 + R_S)}, \qquad (2.48)$$

$$G_c \approx \frac{R_1 \cdot (R_2 + R_3 + R_S)}{R_S R_2 + R_1 \cdot (R_2 + R_3 + R_S)},$$
 (2.49)

$$G_d \approx \frac{R_S R_2}{R_S R_2 + R_1 \cdot (R_2 + R_3 + R_S)}.$$
 (2.50)

The expression found in (2.45) is still valid. We can write the noise density for a MOS transistor as follows: $\overline{i_{n,M}^2} = 4k_BT\gamma g_m$ (like e.g. in [41], where the typical value of $\gamma = \frac{2}{3}$ valid for long-channel devices is used). Different contributions are summarized in Table 2.2 (the contribution from the source is unitary). Some useful indications come from this approximated analysis:

- R_1 and R_3 should be the dominant noise sources.
- To minimize the overall noise figure, R_2 should be somewhat larger than R_1 , and R_1 should be as small as possible. Reducing R_1 value will eventually lead to:

- a) higher and higher noise contribution from M_1 ,
- b) lower and lower feed-back loop-gain.
- The limit condition $R_3 \to 0$ leads to a theoretical NF> 3dB.

2.1.3 Linearity

The linearity analysis can not be performed by *brute force*: introducing arbitrarily non-linear terms in equations (2.21)-(2.27) lead to unmanageable results. A *softer* approach, based on a semi-iterative method, is then preferred. Since our simplified schematic is virtually operating at a single frequency, we want ultimately to derive a polynomial approximation like:

$$i_2 \approx g_S \cdot v_S + \alpha_2 \cdot v_S^2 + \alpha_3 \cdot v_S^3, \tag{2.51}$$

where g_s is the small-signal transconductance resulting from (2.21)-(2.27), and $\alpha_2 - \alpha_3$ are appropriate non-linear coefficients. In fact, we will derive only α_3 , since input signal v_S is a harmonic signal and even-order distortion doesn't affect the transfer properties at the frequency of interest:

$$v_{S} = v_{0} \cdot \cos \omega_{0} t, \qquad (2.52)$$

$$i_{2} \approx \frac{\alpha_{2} v_{0}^{2}}{2} + \left(g_{s} v_{0} + \frac{3\alpha_{3} v_{0}^{3}}{4}\right) \cos \omega_{0} t +$$

$$+ \frac{\alpha_{2} v_{0}^{2}}{2} \cos 2\omega_{0} t + \frac{\alpha_{3} v_{0}^{3}}{4} \cos 3\omega_{0} t. \qquad (2.53)$$

The following assumptions are considered:

1. the non-linearities are concentrated on the transconductors, i.e. M_1 and M_2 in Figure 2.6. Equations (2.26)-(2.27) are modified as follows:

$$i_1 = g_{m_1}(v_a - v_c) + \hat{i_1},$$
 (2.54)

$$i_2 = -g_{m_2}(v_b - v_d) - \hat{i_2}.$$
 (2.55)

The new system is solved with respect to non-linear terms $\hat{i_1}$ and $\hat{i_2}$:

$$i_2 = g_S \cdot v_S + \widehat{G}_1 \cdot \widehat{i}_1 + \widehat{G}_2 \cdot \widehat{i}_2. \tag{2.56}$$

2. Third-order non-linearities are implemented through empirical factors that depend only on the gate-source voltage:

$$\widehat{i_1} \triangleq \alpha_{31}(v_a - v_c)^3, \tag{2.57}$$

$$\widehat{i_2} \triangleq \alpha_{32}(v_b - v_d)^3. \tag{2.58}$$

First we can solve the (still linear) system (2.26)-(2.27) with respect to the non-linear terms defined in (2.54)-(2.55):

$$\widehat{G}_{1} = \frac{\frac{Z_{in}}{Z_{in} + R_{S}} \cdot g_{m_{2}} R_{L} (R_{2} + R_{3} + R_{S})}{R_{\Delta}}, \tag{2.59}$$

$$\widehat{G}_2 = -\frac{Z_{in}}{Z_{in} + R_S} \cdot (1 + g_{m_1} R_1) (R_2 + R_3 + R_S)}{R_{\Lambda}}.$$
 (2.60)

By substitution of (2.57)-(2.58) in (2.56) we find:

$$i_2 = g_S \cdot v_S + \widehat{G}_1 \cdot \alpha_{31} (v_a - v_c)^3 + \widehat{G}_2 \cdot \alpha_{32} (v_b - v_d)^3.$$
 (2.61)

The expressions of v_a , v_b , v_c and v_d here used are the solutions of the linear system, i.e. (2.30)-(2.33). We introduce those expressions in (2.61) and we radically simplify to get the following result:

$$\widehat{R}_{T} \triangleq (R_{2} + R_{3} + R_{S}) \cdot \left(\frac{Z_{in}}{Z_{in} + R_{S}}\right)^{4}, \qquad (2.62)$$

$$\alpha_{3} = \alpha_{31} \frac{g_{m_{2}} \widehat{R}_{T} R_{L} [R_{2} + R_{3} + g_{m_{2}} R_{2} (R_{3} - R_{1})]^{3}}{R_{\Delta}^{4}} + \alpha_{32} \frac{\widehat{R}_{T} (1 + g_{m_{1}} R_{1}) \{g_{m_{2}} R_{3} R_{L} + R_{2} [1 + g_{m_{1}} (R_{1} + R_{L})]\}^{3}}{R_{\Delta}^{4}}.$$

$$(2.63)$$

To quantify non-linear terms α_{31} and α_{32} we have to consider second-order effects in the MOS I-V characteristic: the simple quadratic expression $I_D \propto (V_{GS} - V_T)^2$ does not provide any high-order non-linearity. Many mechanisms have been studied and implemented in advanced MOS models [42], among which:

• threshold voltage modulation (due to different mechanisms like body effect),

- mobility reduction due to vertical field,
- velocity saturation,
- bulk charge effect.

For the sake of simplicity, only one of them will be examined, and since we have assumed that non-linear terms are functions of the gate-source voltage, mobility reduction due to vertical field is the preferred choice. We rewrite the I-V characteristic as follows:

$$I_D = \frac{1}{2} \frac{\mu_{n0}}{1 + \theta(V_{GS} - V_T)} C'_{ox} \frac{W}{L} (V_{GS} - V_T)^2$$
 (2.64)

$$= \frac{k_0}{1 + \theta(V_{GS} - V_T)} (V_{GS} - V_T)^2, \qquad (2.65)$$

where

 μ_{n0} is the effective mobility at $V_{GS} = V_T$,

 θ is the vertical-field reduction factor and

 k_0 is equal to $\frac{1}{2}\mu_{n0}C'_{ox}\frac{W}{L}$.

Given an input signal v_{gs} on top of the bias voltage V_{gs} , the output current can be determined as a sum of bias current I_d , a linear term i_d and a non-linear term $\hat{i_d}$:

$$I_{d} + i_{d} + \hat{i}_{d} = \frac{k_{0}(V_{gs} + v_{gs} - V_{T})^{2}}{1 + \theta(V_{gs} + v_{gs} - V_{T})}$$

$$= \frac{k_{0}}{1 + \theta(V_{gs} - V_{T})} \frac{(V_{gs} + v_{gs} - V_{T})^{2}}{1 + \frac{\theta v_{gs}}{1 + \theta(V_{gs} - V_{T})}}.$$
 (2.66)

We can approximate the following term:

$$\frac{1}{1 + \frac{\theta v_{gs}}{1 + \theta(V_{gs} - V_T)}} \approx 1 - \frac{\theta v_{gs}}{1 + \theta(V_{gs} - V_T)},$$
(2.67)

and (2.66) becomes:

$$I_{d} + i_{d} + \hat{i}_{d} \simeq \frac{k_{0}(V_{gs} + v_{gs} - V_{T})^{2}}{1 + \theta(V_{gs} - V_{T})} \left[1 - \frac{\theta v_{gs}}{1 + \theta(V_{gs} - V_{T})} \right]$$

$$\simeq \frac{k_{0}(V_{gs} - V_{T})^{2}}{1 + \theta(V_{gs} - V_{T})} + \frac{k_{0}(V_{gs} - V_{T})}{1 + \theta(V_{gs} - V_{T})} v_{gs} + \frac{1 - \theta(V_{gs} - V_{T})}{[1 + \theta(V_{gs} - V_{T})]^{2}} v_{gs}^{2} - \frac{k_{0} \cdot \theta}{[1 + \theta(V_{gs} - V_{T})]^{2}} v_{gs}^{3}, \qquad (2.68)$$

$$I_{d} = \frac{k_{0}(V_{gs} - V_{T})^{2}}{1 + \theta(V_{gs} - V_{T})}, \qquad (2.69)$$

$$i_{d} \simeq k_{0}(V_{gs} - V_{T}) \frac{2 + \theta(V_{gs} - V_{T})}{[1 + \theta(V_{gs} - V_{T})]^{2}} v_{gs}, \qquad (2.70)$$

$$\hat{i}_{d} \simeq k_{0} \frac{1 - \theta(V_{gs} - V_{T})}{[1 + \theta(V_{gs} - V_{T})]^{2}} v_{gs}^{2} - \frac{1 - \theta(V_{gs} - V_{T})}{[1 + \theta(V_{gs} - V_{T})]^{2}} v_{gs}^{2}, \qquad (2.70)$$

$$\widehat{i_d} \simeq k_0 \frac{1 - \theta(V_{gs} - V_T)}{[1 + \theta(V_{gs} - V_T)]^2} v_{gs}^2 - \frac{k_0 \cdot \theta}{[1 + \theta(V_{gs} - V_T)]^2} v_{gs}^3.$$
(2.71)

From (2.71) we can derive the third-order polynomial coefficient:

$$\alpha_{3i} \triangleq -\frac{k_{0_i} \cdot \theta_i}{\left[1 + \theta(V_{qs_i} - V_{T_i})\right]^2},$$
(2.72)

where the index i refers to the specific transistor M_i .

2.1.4 Numerical results

The simulation of the stage is quite problematic because the different parasitics disseminated in the circuit reduce the loop gain (in magnitude and phase) very rapidly. We can identify a few parasitics among the most critical ones:

1. the finite output resistance of the input transistor (g_{ds_1}) ,

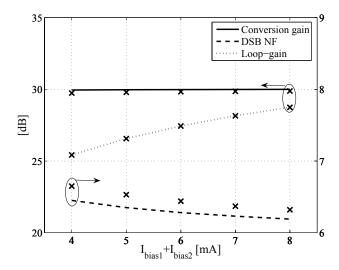


Figure 2.11: Main linear parameters: conversion gain, noise figure and loop-gain. Results from simulations are indicated with "×"-markers.

- 2. the gate-drain capacitance of the input transistor (C_{gd_1}) ,
- 3. the gate-drain capacitance of the switching pair in the mixer,
- 4. the total capacitance from the drain of transistor M_2 to ground.

Many solutions can be devised to reduce the effects of these parasitics (for example, an obvious improvement for 1. and 2. would be a cascode transistor), but in this context we want to validate the results derived in the previous sections. Therefore we make use of semi-ideal transistor models, where the main parasitics have been stripped down. These simplifications enable us to perform simulations of critical parameters based on linear analyses (e.g. conversion gain and noise figure), but unfortunately prevent any non-linear simulation from completing successfully. For this reason, non-linear analysis will be entirely based on numerical calculations. Semi-ideal models are based on a $0.35\mu m$ process. The bias currents of the input stage (I_{bias_1}) and the mixer (I_{bias_2}) are swept at the same time. Calculated values for conversion gain and loop-gain exhibit excellent matching with simulations, whereas NF shows a difference around 0.2dB, due to the fact that the contribution from M_2 at low frequencies has been disregarded.

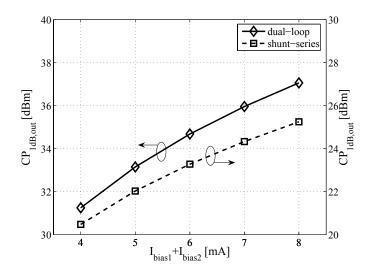


Figure 2.12: Output-referred 1-dB compression point for the proposed circuit and a front-end based on a shunt-series input stage.

Linearity analysis is performed on the basis of the calculations from the previous section. It is useful to have a term of comparison: the same approach that we described in Section 2.1.3 is applied to a front-end based on a shunt-series input stage (see Figure 2.4) and a conventional multiplier-based active mixer (like the one in Figure 2.1). The results are shown in Figure 2.12 in terms of output-referred 1-dB compression point. In Table 2.3 the main design parameters are summarized.

2.2 Application to a narrow-band receiver

Narrow-band LNAs based on tuned LC-tanks are extensively used in state-of-the-art receiver chains, where specifications are highly demanding; in Table 2.4 a selection of existing implementations is presented. The selective nature of LC resonators determines a high rejection of adjacent blockers, and the relatively low number or critical noise sources has a positive impact on total noise figure. A simplified configuration is shown in Figure 2.13a: the inductive degeneration provides the correct matching

R_S 50Ω 50Ω R_1 10Ω 10Ω R_2 10Ω 50Ω	
D 100 500	
$R_2 10\Omega 50\Omega$	
R_3 90 Ω N/A	
R_f N/A $1k\Omega$	
R_L 3.7k Ω 3.7k Ω	
R 1k Ω 1k Ω	
I_{bias_1} 2÷4mA 2÷4mA	
I_{bias_2} 2÷4mA 2÷4mA	
V_{DD} 2.5V 2.5V	

Table 2.3: Main parameters related to results in Figures 2.11 and 2.12.

[41], and a cascode stage reduces the impact of the Miller effect related to gate-drain capacitance at transistor M_1 . The mixer is built as a direct implementation of the Gilbert multiplier.

The open loop configuration can be easily modified to exploit negative feedback as shown in Figure 2.13b; to simplify the structure the following assumptions are made:

- 1. The contribution of the cascode transistor in the LNA is omitted.
- 2. Since we are interested in the operation around resonance, we represent the load LC tank with the equivalent parallel resistance R_d .
- 3. No parasitics are considered except for the gate-source capacitance of transistor M_1 . Switches and transconductors are ideal.
- 4. The quality factor of tuning inductor L_q is considered infinite.
- 5. At high frequency, the load resistors of the mixer combine in parallel with the degeneration inductor L_s . We assume that the equivalent parallel resistance of the inductor is marginally affected.

Table 2.4: Reported implementations of narrow-band receiver front-ends based on tuned LC-tanks.

	Gain	NF	Freq.	I _{BIAS} @V _{DD}	Tech.	
	[dB]	[dB]	[GHz]	[mA] $[V]$	[nm]	
[43]	51.5	2.9	0.894	29@2.6	250	
[44]	7	11.77	1.1	22@1	350	
[45]	28.4	1.4	5	11.2@1.8	180	
[46]	47	5.6	2.1	9@1.8	180	
[47]	92	4.8	1.6	17@1.6	180	1
[48]	110	4	1.6	15@1.6	180	2
[49]	44	7	1.6	5@1.8	180	
[50]	31	3.5	2.4	3.9@1.8	180	
[51]	17.8	1	2.4	2.3@1.2	180	
[52]	26	3.5	5.825	30@1.2	130	
[53]	30	3	0.915	2@1.8	180	
[54]	28	4.7	2.4	20.3@1.8	180	
ib.	30	5.1	5	25.4@1.8	180	
[55]	28	12.5	60	7.5@1.2	130	3
[56]	38	1.8	1.6	60@1.4	90	4
[57]	50	3.9	1.9	11@1.2	130	
[58]	20.4	19	2.4	0.5@1	180	5
ib.	30.5	10.1	2.4	0.5@1	180	6
[59]	33	1.7	0.85	18.6@1.5	65	
[60]	16	8.5	5	6.3@1.8	180	7
[61]	31.5	2.1	1.9	22@1.5	90	8

¹ performances refer to the complete radio

² gain and current consumption include the whole chain

³ based on common-gate topology

⁴ current consumption refers to the complete radio

⁵ LNA and Mixer are stacked

⁶ folded-cascode LNA+Mixer

⁷ current consumption includes VCO

⁸ current consumption includes LO drivers

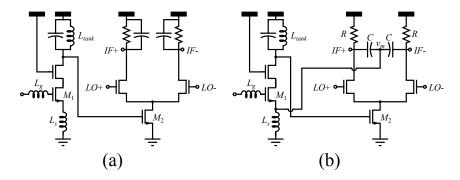


Figure 2.13: A typical LNA-Mixer chain based on LC-tank resonators (a), and the proposed topology (b). Biasing details are omitted.

Concerning the last assumption, we observe that the source-degeneration inductor can be conveniently represented at resonance frequency by an inductance and a series resistance (like in Figure 2.14a) or by an inductance and a parallel resistance (see Figure 2.14b) [16]. The impedance at node v_m (equal to $\frac{R}{2} + \frac{1}{s2C}$) is connected in parallel, and if we neglect the contribution from capacitance 2C we can define the total impedance as

$$Z_s = sL_p /\!\!/ R_{p1} /\!\!/ \frac{R}{2}. \tag{2.73}$$

According to assumption 5. we consider $R_{p1} \approx R_{p1} / \frac{R}{2}$. The simplified circuit results as in Figure 2.15. We will express the degeneration impedance Z_s in the Laplace domain in the following forms, using well-known seriesparallel transformations:

$$Z_s(s) \approx R_1 + sL_s \tag{2.74}$$

$$Z_s(s) \approx R_1 + sL_s$$
 (2.74)

$$\approx \frac{sL_pR_{p1}}{R_{p1} + sL_p}.$$
 (2.75)

2.2.1Circuit analysis

Once again, we study the behaviour of the simplified circuit at a single frequency, and the output signal is the current i_2 . From the figure we can

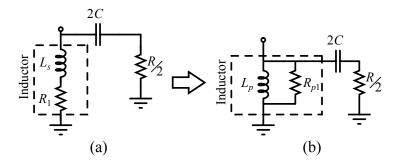


Figure 2.14: Series-parallel transformation of the source-degeneration inductor.

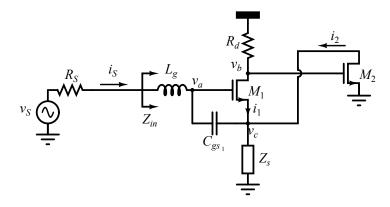


Figure 2.15: Simplified circuit derived from Figure 2.13b.

easily define a fundamental set of equations, like we did for the dual-loop stage:

$$v_a = v_S - i_S(R_S + sL_g),$$
 (2.76)

$$v_b = -R_d i_1, (2.77)$$

$$v_c = Z_s(s) \cdot (i_S + i_1 + i_2),$$
 (2.78)

$$i_S = sC_{gs_1}(v_a - v_c),$$
 (2.79)

$$i_1 = g_{m_1}(v_a - v_c),$$
 (2.80)

$$i_2 = -g_{m_1} v_b. (2.81)$$

First of all, we can easily determine the input impedance as follows:

$$Z_{in}(s) = \frac{1}{sC_{gs_1}} + sL_g + Z_s(s) + \frac{g_{m_1}Z_s(s)}{sC_{gs_1}}(1 + g_{m_2}R_d).$$
 (2.82)

The result does not differ from the usual one [41] except for the multiplying factor $(1 + g_{m_2}R_d)$. By using (2.74) we find:

$$Z_{in}(s) = R_1 + \frac{g_{m_1} L_s}{C_{gs_1}} (1 + g_{m_2} R_d) + s \cdot (L_g + L_s) + \frac{1 + g_{m_1} R_1 (1 + g_{m_2} R_d)}{s C_{gs_1}}.$$
(2.83)

We define the unity-gain frequency $\omega_T \triangleq \frac{g_{m_1}}{C_{g_{s_1}}} = 2\pi f_T$ for transistor M_1 , and we can determine the degeneration inductance L_s and the tuning inductance L_g by imposing input matching at the operating frequency ω_0 :

$$\Re[Z_{in}(\jmath\omega_0)] = R_S \quad \Rightarrow \quad L_s = \frac{R_S - R_1}{\omega_T(1 + g_{m_2}R_d)},$$
 (2.84)

$$\Im[Z_{in}(\jmath\omega_0)] = 0 \quad \Rightarrow \quad L_g = \frac{1 + g_{m_1}R_1(1 + g_{m_2}R_d)}{\omega_0^2 C_{gs_1}} - L_s. \quad (2.85)$$

The optimal value for L_s will be effectively smaller, compared to the traditional open-loop case, by a factor equal to $(1+g_{m_2}R_d)$; this can be a significant problem in high-performance technologies, since the low value of L_s can be difficult to control accurately. The tuning inductance L_g will not differ substantially, since $g_{m_1}R_1 \ll 1$. We can determine the transconductance of the stage as follows:

$$g_s \triangleq \frac{i_2}{v_S},\tag{2.86}$$

$$g_s(s) = \frac{g_{m_1}g_{m_2}R_d \frac{Z_{in}(s)}{Z_{in}(s) + R_S}}{1 + s \cdot C_{qs_1}[s \cdot L_q + Z_s(s)] + g_{m_1}Z_s(s)(1 + g_{m_2}R_d)}. \quad (2.87)$$

The expression in (2.87) can be greatly simplified if we assume perfect matching at the resonance frequency:

$$g_s(s) = \frac{1}{2} \frac{g_{m_1} g_{m_2} R_d}{s \cdot C_{as_1} R_S},\tag{2.88}$$

which is the same result that we would obtain with an open-loop stage. Therefore the total conversion gain is not affected by the feedback.

From equations (2.76)-(2.81) we can derive an expression for the loop-gain. By using the same technique described at the end of Section 2.1.1 we find:

$$G_{loop}(s) = -\frac{g_{m_1}g_{m_2}R_dZ_s(s)}{1 + g_{m_1}Z_s(s) + s \cdot C_{gs_1}[R_S + s \cdot L_g + Z_s(s)]}.$$
 (2.89)

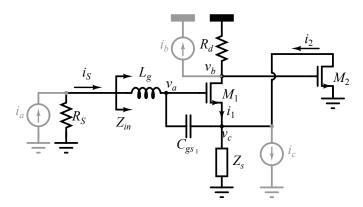


Figure 2.16: Simplified stage for first-order noise calculations.

To simplify the expression we might try to evaluate the limit for $R_1 \longrightarrow 0$:

$$\lim_{R_1 \to 0} G_{loop}(j\omega_0) = -\frac{g_{m_2} R_d}{2 + g_{m_2} R_d},\tag{2.90}$$

which shows that the loop-gain is smaller than unity in magnitude around the operating frequency.

2.2.2 Noise

Noise analysis is performed by means of fictitious current sources i_a , i_b and i_c shown in Figure 2.16, like in Section 2.1.2. The transfer functions are easily determined:

$$i_2 = G_a \cdot i_a + G_b \cdot i_b + G_c \cdot i_c,$$
 (2.91)

$$G_a(s) = \frac{g_{m_1}g_{m_2}R_dR_S}{1 + sC_{qs_1}[sL_q + Z_s(s)] + g_{m_1}Z_s(s)(1 + g_{m_2}R_d)},$$
(2.92)

$$G_b(s) = \frac{g_{m_2}R_d\{1 + g_{m_1}Z_s(s) + sC_{gs_1}[sL_g + R_S + Z_s(s)]\}}{1 + sC_{gs_1}[sL_g + R_S + Z_s(s)] + g_{m_1}Z_s(s)(1 + g_{m_2}R_d)}, (2.93)$$

$$G_c(s) = \frac{g_{m_1}g_{m_2}R_dZ_s(s)}{1 + sC_{gs_1}[sL_g + R_S + Z_s(s)] + g_{m_1}Z_s(s)(1 + g_{m_2}R_d)}. (2.94)$$

In this simplified analysis, induced gate noise is neglected: even though it plays an important role in inductively-degenerated LNAs [62], its contribution will be the same for both structures in Figure 2.13. In Figure 2.17

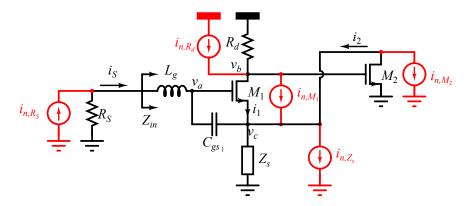


Figure 2.17: Simplified stage with noise sources.

the main noise sources are depicted. The different contributions around operating frequency are evaluated as follows:

$$F = 1 + \frac{\overline{i_{n,M_{1}}^{2}} \cdot |G_{c}(\jmath\omega_{0}) - G_{b}(\jmath\omega_{0})|^{2} + \overline{i_{n,M_{2}}^{2}} \cdot |G_{c}(\jmath\omega_{0})|^{2}}{\overline{i_{n,R_{S}}^{2}} \cdot |G_{a}(\jmath\omega_{0})|^{2}} + \frac{\overline{i_{n,R_{S}}^{2}} \cdot |G_{c}(\jmath\omega_{0})|^{2} + \overline{i_{n,R_{d}}^{2}} \cdot |G_{b}(\jmath\omega_{0})|^{2}}{\overline{i_{n,R_{S}}^{2}} \cdot |G_{a}(\jmath\omega_{0})|^{2}}$$

$$= 1 + \gamma g_{m_{1}} R_{S} \left| \frac{G_{c}(\jmath\omega_{0}) - G_{b}(\jmath\omega_{0})}{G_{a}(\jmath\omega_{0})} \right|^{2} + \gamma g_{m_{2}} R_{S} \left| \frac{G_{c}(\jmath\omega_{0})}{G_{a}(\jmath\omega_{0})} \right|^{2} + \frac{R_{S}}{R_{p_{1}}} \left| \frac{G_{c}(\jmath\omega_{0})}{G_{a}(\jmath\omega_{0})} \right|^{2} + \frac{R_{S}}{R_{d}} \left| \frac{G_{b}(\jmath\omega_{0})}{G_{a}(\jmath\omega_{0})} \right|^{2}.$$

$$(2.95)$$

To provide a useful comparison with the traditional open-loop stage, we summarize the different contributions to the total noise factor in Table 2.5. The results are calculated around operating frequency for $R_1 \to 0$. The factor related to M_2 includes two terms, because the transistor generates noise at high frequency (and it gets down-converted together with the signal) and at low frequency (where $\frac{1}{f}$ -noise will be dominant, even though we have chosen not to include it explicitly in the calculations); the simplified formula overestimates the latter term quite a bit. In the closed-loop solution several terms exhibit improvements in noise contributions, specifically M_2 , R_d and R_{p1} .

	Closed loop	Open loop
R_S	1	1
M_1	$\gamma g_{m_1} R_S \left(\frac{\omega_0}{\omega_T}\right)^2$	$\gamma g_{m_1} R_S \left(\frac{\omega_0}{\omega_T}\right)^2$
M_2	$\gamma g_{m_2} R_S \left(\frac{\omega_0}{\omega_T}\right)^2 \frac{1}{(1+g_{m_2}R_d)^2} +$	$\gamma g_{m_2} R_S \left(\frac{\omega_0}{\omega_T} \frac{1}{g_{m_2} R_d} \right)^2 \left(4 + \frac{\pi^2}{2} \right)$
	$+\gamma g_{m_2} R_S \left(\frac{\omega_0}{\omega_T}\right)^2 \frac{1}{2} \left(\frac{\pi}{g_{m_2} R_d}\right)^2$	
R_{p1}	$\frac{R_S}{R_{p1}} \left(\frac{\omega_0}{\omega_T} \frac{1}{1 + g_{m_2} R_d} \right)^2$	$\frac{R_S}{R_{p1}} \left(\frac{\omega_0}{\omega_T} \right)^2$
R_d	$\frac{R_S}{R_d} \left(\frac{\omega_0}{\omega_T} \frac{2 + g_{m_2} R_d}{1 + g_{m_2} R_d} \right)^2$	$\frac{R_S}{R_d} \left(2 \frac{\omega_0}{\omega_T} \right)^2$
R	$\frac{R_S}{R} \left(\frac{\omega_0}{\omega_T} \frac{\pi}{q_{m_2} R_d} \right)^2$	$rac{R_S}{R} \left(rac{\omega_0}{\omega_T} rac{\pi}{q_{mo}R_d} ight)^2$

Table 2.5: Noise contributions for the proposed stage compared to the traditional case.

2.2.3 Numerical analysis

Like we did in Section 2.1.4, we use simplified models to perform simulations. In this case, only linear parameters will be extracted, so we can compare simulations results with previously derived calculations. From a quick glance at Table 2.5 we can conclude that the most critical parameters in noise analysis are the unity-gain frequency f_T and the transconductance of the second stage g_{m_2} ; therefore sweeping the bias current of the second stage (I_{bias_2}) and the unity-gain frequency sounds like a sensible choice. Figure 2.18 shows that the noise contribution from transistor M_2 is actually overestimated, but the general trend of the calculations match the behaviour seen in simulations. The difference between open-loop and closed-loop solution can be as high as $\sim 0.76 \mathrm{dB}$. Simulations are based on a $0.35 \mu\mathrm{m}$ process, and the main parameters are summarized in Table 2.6.

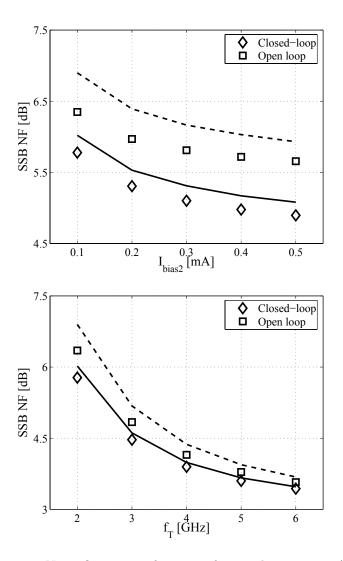


Figure 2.18: Noise figure as a function of mixer bias current (top) and unity-gain frequency (bottom) for the proposed stage and the traditional open-loop stage. Calculations are shown as lines (continuous line and dashed line respectively), whereas simulations are indicated with markers.

Table 2.6: Main parameters related to results in Figure 2.18.

	Closed loop	Open loop	
R_S	50Ω	50Ω	
I_{bias_1}	$2\mathrm{mA}$	2 mA	
I_{bias_2}	$0.1\text{-}0.5\mathrm{mA}$	$0.1\text{-}0.5\mathrm{mA}$	1
	0.1 mA	$0.1 \mathrm{mA}$	2
f_T	$2\mathrm{GHz}$	$2\mathrm{GHz}$	1
	$2\text{-}6\mathrm{GHz}$	$2\text{-}6\mathrm{GHz}$	2
V_{DD}	2.5V	2.5V	
f_{LO}	$900 \mathrm{MHz}$	$900 \mathrm{MHz}$	
$Q@f_{LO}$	10	10	
L_{tank}	2nH	2nH	

 $^{^1}$ conditions for results in Figure 2.18 (top) 2 conditions for results in Figure 2.18 (bottom)

Chapter 3

Simplified analysis of a current-reuse LNA-Mixer-VCO architecture

In the past years, battery-operated portable devices have been pushing the development of power-saving techniques to the limit. Approaches based on current reuse are extremely effective, for different reasons:

- the supply voltage of large systems is determined by complex architecture specifications, if not by the voltage provided by the battery itself. Therefore even if a single block *could* be operated at a much lower voltage, in practice this is not a viable option to save power.
- Given a certain supply voltage, current consumption of different blocks is determined by a number of non-trivial trade-offs, but it is generally possible to identify one or very few *critical* blocks that determine a lower bound for power usage.

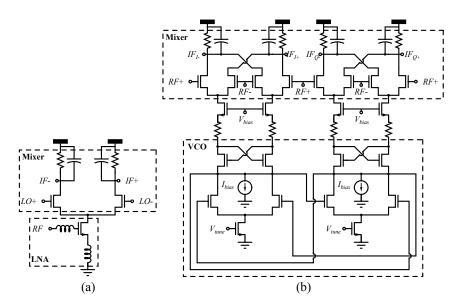


Figure 3.1: Examples of current-reuse topologies: a stacked LNA-Mixer (a) and a self-oscillating mixer (b).

To get as close as possible to the aforementioned lower limit, stacking multiple stages over the same bias source sounds like a very sensible choice. A couple of popular examples are given in Figure 3.1:

- (a) a narrow-band LC-based LNA can provide bias current for a conventional Gilbert-based mixer [63][64],
- (b) the bias current of a VCO can be used to supply a common-base double-balanced mixer [65]. These architectures are commonly referred to as self-oscillating mixers (SOM).

The LNA-Mixer-VCO (*LMV*) architecture presented in [1] is capable of stacking an LNA, a mixer and a VCO over the same bias current. LC-tank VCOs and active mixers share the same principle (i.e. a pair of switching MOS transistors), but the LC-tank can not offer high impedance at low frequencies (where the down-converted signal should be collected): an additional stage must be inserted, thereby providing isolation between the LO output and the down-converted signal. A schematic of the LMV

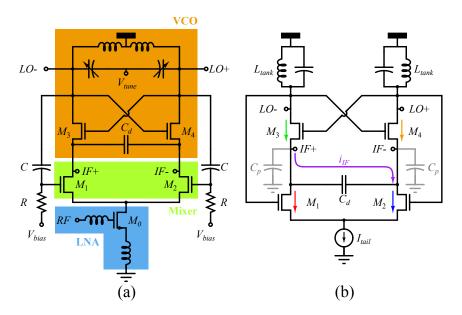


Figure 3.2: The LMV cell: complete circuit (a) and simplified schematic with parasitics at output nodes (b).

cell is reproduced in Figure 3.2a:

- transistor M_0 is a transconductor configured as a narrow-band LNA, which is also supplying bias current for the rest of the circuit.
- Transistors $M_3 M_4$ provide the necessary negative resistance to induce oscillation; capacitor C_d should behave as a short-circuit at radio frequencies.
- Transistors M_1-M_2 are operated as switches to synthesize the mixing function. The R-C network is just a decoupling high-pass filter that carries the LO signal back to the input of the mixer.

The adopted description is very simplified and somehow inaccurate (e.g. we can say that $M_3 - M_4$ are also involved in the mixing operation), but it can be useful to a large extent for a first-order analysis of the stage. Calculating the transfer characteristic of this cell (i.e. the conversion gain and the output bandwidth, among other things) is not obvious, since there

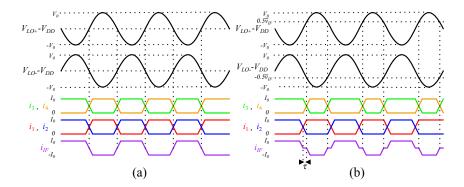


Figure 3.3: Waveforms for the simplified schematic in Figure 3.2b: steady-state operation for $v_{IF} = 0$ (a) and perturbed waveforms when a signal v_{IF} is applied (b).

is apparently no load to convert the IF current signal into a differential voltage (like the R-C low-pass filter in Figure 3.1a, for instance). In particular, the influence of parasitics at the output nodes must be taken in serious consideration. Based on the simplified schematic in Figure 3.2b, we will develop a first-order analysis of the stage, in order to quantify the impact of parasitic capacitance C_p . As a fundamental precondition, we assume that transistors M_1-M_4 behave as ideal switches.

3.1 Effective resistance at output nodes

First of all, we consider the limit case $C_p \to 0$. Referring to Figure 3.2b, we define:

$$I_{tail} = I_0 + i_{RF} \cdot \sin(\omega_{RF}t), \tag{3.1}$$

$$V_{LO+} = V_{DD} + V_0 \cdot \sin(\omega_0 t), \qquad (3.2)$$

$$V_{LO-} = V_{DD} - V_0 \cdot \sin(\omega_0 t). \tag{3.3}$$

For $i_{RF} = 0$, transistors $M_1 - M_4$ are switching on and off periodically, and we can easily conclude that the couple M_1, M_4 is on half of the time, whereas the couple M_2, M_3 is on the rest of the time. All relevant waveforms are summarized in Figure 3.3a. We will show that by applying a

probe signal $v_{IF} = v_{IF+} - v_{IF-}$ at low frequency and evaluating the current i_{IF} associated with it, we can calculate an equivalent resistance R_{IF} .

We assume that the test signal does not interfere with the oscillation, therefore waveforms V_{LO+} and V_{LO-} do not change. We can also safely assume that the behaviour of the couple M_1, M_2 is not affected at the first order, since a low-frequency voltage at the drains does not influence normal operation; finally, we observe that the couple M_3, M_4 will indeed be affected, since one transistor will be on for a longer time and the other will be on for a shorter time: the reason is related to the fact that the switching threshold is reduced for one transistor and increased for the other. The concept is illustrated by the waveforms in Figure 3.3b: in every half-period, there is a time-slot τ in which the tail current I_0 is flowing entirely either through M_1, M_3 or through M_2, M_4 . We can easily derive the following equations:

$$\frac{\overline{v_{IF}}}{2} = V_0 \sin \omega_0 \tau, \tag{3.4}$$

$$\overline{i_{IF}} = \frac{I_0 \frac{T_0}{2} - I_0 \left(\frac{T_0}{2} - 2\tau\right)}{T_0},$$
(3.5)

where $T_0 = \frac{2\pi}{\omega_0}$. From the previous equations, we can determine

$$\tau = \frac{1}{\omega_0} \arcsin \frac{\overline{v_{IF}}}{2V_0}$$

$$\approx \frac{\overline{v_{IF}}}{2\omega_0 V_0}, \qquad (3.6)$$

$$\overline{i_{IF}} = \frac{2I_0 \tau}{T_0}$$

$$\approx \frac{\overline{v_{IF}}}{2\omega_0 V_0} \frac{2I_0}{T_0} = \frac{I_0}{2\pi V_0} \overline{v_{IF}}, \qquad (3.7)$$

$$R_{IF} \triangleq \frac{\overline{v_{IF}}}{\overline{i_{IF}}}$$

$$\approx \frac{2\pi V_0}{I_0}. \qquad (3.8)$$

The last equation shows that there is an *effective* resistance acting at low frequencies, which is related to the amplitude of the oscillation and to the bias current. We can also express (3.8) in another form:

$$\frac{V_0}{I_0} \approx \frac{1}{\pi} \omega_0 L_{tank} Q \Longrightarrow R_{IF} \approx 2\omega_0 L_{tank} Q,$$
 (3.9)

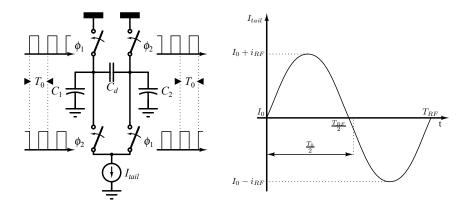


Figure 3.4: SC-theory applied to the LMV cell.

where Q is the quality factor of inductor L_{tank} . To complete the analysis, we notice that the capacitance C_d is in parallel to the effective resistance R_{IF} , therefore the complete transfer function can be expressed as follows:

$$\frac{v_{IF}}{i_{RF}} \approx \frac{2}{\pi} \frac{R_{IF}}{1 + sC_dR_{IF}}.$$
(3.10)

The last equation helps us to determine an upper limit for the GBWP of the stage.

3.2 Parasitics-related losses

We want to evaluate the impact of parasitic capacitance C_p . Since we treat transistors $M_1 - M_4$ as ideal switches, we can greatly simplify the circuit by applying the theory of switched capacitors (SC). As a further simplification, we consider the impedance of the tank negligible. The simplified structure is depicted in Figure 3.4, where obviously we have $C_1 = C_2 = C_p$. In this case we are interested in the RF signal injected by I_{tail} , therefore the DC component I_0 is disregarded. During a phase (for example ϕ_1) the

input current injects a certain amount of charge Q_{in} in the system:

$$Q_{in} = \int_0^{\frac{T_0}{2}} (I_{tail}(t) - I_0) dt$$
 (3.11)

$$= i_{RF} \int_0^{\frac{T_0}{2}} \sin(\omega_{RF} t) dt = i_{RF} \int_0^{\frac{T_0}{2}} \sin\frac{2\pi t}{T_{RF}} dt \qquad (3.12)$$

$$\approx i_{RF} \int_0^{\frac{T_0}{2}} \sin(\omega_0 t) dt = \frac{2i_{RF}}{\omega_0}.$$
 (3.13)

This charge is going to be shared by C_d and C_2 . During the (instantaneous) transition $\phi_1 \longrightarrow \phi_2$, capacitor C_2 is discharged, and the charge Q_d , already present on C_d , is redistributed partially on C_1 . At this point we can repeat for phase ϕ_2 the same steps (of course C_1 and C_2 have opposite roles): the input RF current has opposite sign now, but it is injected on the opposite side. A detailed sequence is depicted in Figure 3.5. Given the periodic nature of the system, the entire behavioural analysis boils down to a single discrete-time dynamic equation:

redistribution sharing of the of the charge input charge
$$Q_d(n+1) = \overbrace{(1-\alpha)Q_d(n)}^{\text{redistribution}} + \overbrace{(1-\alpha)Q_{in}(n)}^{\text{sharing of the}}, \qquad (3.14)$$

where $\alpha \triangleq \frac{C_p}{C_p + C_d}$. In the \mathcal{Z} -domain we can rewrite (3.14) as follows:

$$zQ_d(z) = (1-\alpha)[Q_d(z) + Q_{in}(z)],$$
 (3.15)

$$\mathcal{H}(z) \triangleq \frac{Q_d(z)}{Q_{in}(z)} = \frac{1-\alpha}{z-1+\alpha}.$$
 (3.16)

We can easily perform a proper mapping from the \mathcal{Z} -domain to the Laplacedomain; since we are interested in frequencies close to DC, we can approximate the mapping function as follows:

$$z = e^{\frac{sT_0}{2}} \approx 1 + \frac{sT_0}{2} \Longrightarrow \mathcal{H}(s) = \frac{1 - \alpha}{\alpha + \frac{sT_0}{2}}.$$
 (3.17)

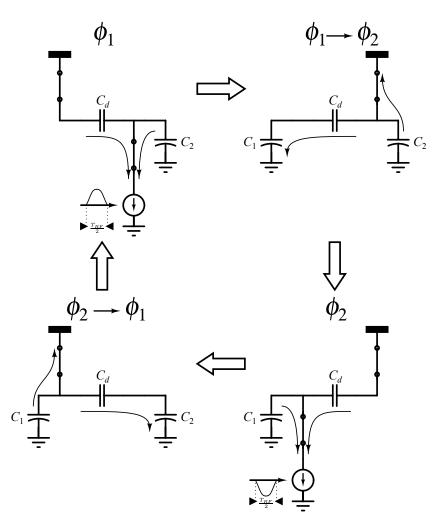


Figure 3.5: Complete cycle $\phi_1 - \phi_2$. Different charge-discharge processes are indicated.

From (3.13) and (3.17) we derive the total transfer function:

$$\begin{pmatrix}
Q_{in} & \approx & \frac{2i_{RF}}{\omega_0} \\
Q_d & = & C_d \cdot v_{IF}
\end{pmatrix} \frac{v_{IF}}{i_{RF}} \approx \frac{2}{\omega_0 C_d} \cdot \frac{1 - \alpha}{\alpha + \frac{sT_0}{2}} \tag{3.18}$$

$$\approx \frac{2}{\omega_0} \cdot \frac{1 - \alpha}{\alpha C_d} \cdot \frac{1}{1 + \frac{sT_0}{2\alpha}}$$
 (3.19)
$$\approx \frac{2}{\omega_0 C_p} \cdot \frac{1}{1 + \frac{sT_0}{2\alpha}}.$$
 (3.20)

$$\approx \frac{2}{\omega_0 C_p} \cdot \frac{1}{1 + \frac{sT_0}{2\alpha}}.$$
 (3.20)

We notice that the transfer function in (3.20) has a diverging behaviour for $C_p \to 0$: this stems from the fact that the lossy mechanism described in Section 3.1 is not taken into consideration here.

All in all, equations (3.10) and (3.20) provide reasonable asymptotic bounds for the GBWP of the stage, as it is shown in Section 3.3. In Appendix C we try to provide a unified formulation which embraces the SC-based approach leading to (3.20), and to give a more rigorous explanation for the effective resistance described in Section 3.1.

3.3 Numerical results

To prove the validity of limits (3.10) and (3.20) we perform simulations of the LMV structure in semi-ideal conditions:

- MOS transistors are stripped-down of their main capacitive parasitics, in order to resemble ideal switches as much as possible,
- current source I_{tail} is ideal,
- the LC tank is implemented as an ideal centre-tapped inductor, in which the mutual coupling between the two branches is equal to unity.

We write the transfer function as follows:

$$\frac{v_{IF}}{i_{RF}} = \frac{R_{\text{gain}}}{1 + \frac{s}{BW_{\text{adB}}}}.$$
(3.21)

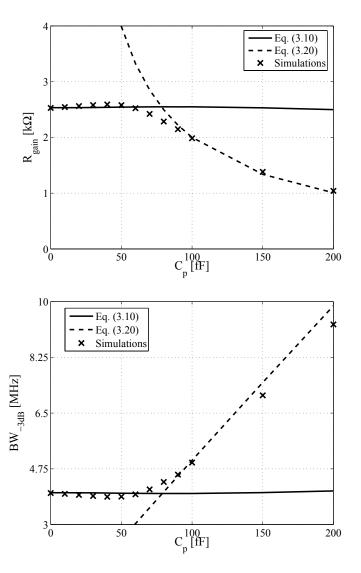


Figure 3.6: Calculation results from (3.10) and (3.20) (continuous and dashed lines respectively) together with simulation results (" \times " markers). The low-frequency gain (on top) and the -3dB bandwidth (bottom) are shown.

 f_0 1.6GHz I_0 1mA L_{tank} 5nH $Q@f_0$ 20 C_d 10pF

Table 3.1: Main parameters for results in Figure 3.6.

It is then easy to extract parameters $R_{\rm gain}$ and $BW_{-3{\rm dB}}$ from simulations using PSS+PAC analyses in Cadence. Results are plotted in Figure 3.6, whereas main parameters are summarized in Table 3.1: we can appreciate that asymptotic curves from (3.10) and (3.20) follow the simulations quite closely in the peripheral regions, and the transition between (3.10)- and (3.20)-regime (which happens for $C_p \sim 50 \div 100 fF$ in the figure) is very sharp. As predicted by (3.20), the roll-off of $R_{\rm gain}$ as a function of C_p is extremely fast, and this makes this structure rather prone to degradation in performance due to parasitics. To alleviate the problem substantially, a current-mode output is chosen in practical implementations: it can be proven that performance degradation due to parasitic capacitance C_p is radically reduced [66].

3.4 Current-mode output stage

A simple current-mode output stage is depicted in Figure 3.7a; the fully-differential operational amplifier should be able to drive a relatively low resistance, therefore a topology based on a super-cascode stage is preferred (see Figure 3.7b). In order to reduce the required voltage headroom, a self-biased active load is chosen: the final configuration is drawn in Figure 3.7c, where capacitor C_{out} provides additional low-pass filtering. The fully-differential op-amp schematic is reported in Figure 3.8 (biasing details are simplified). The basic circuit (a) consists of

• a differential pair $M_7 - M_8$,

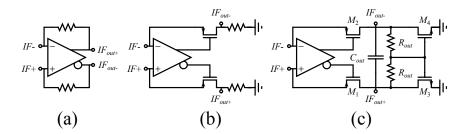


Figure 3.7: A current-mode output stage: simple solution (a), super-cascode with passive load (b) and active load (c).

- an active load $M_5 M_6$,
- a cascode stage $M_9 M_{10}$ followed by
- a self-biased active load $M_{11} M_{12} R_{load}$.

This folded-cascode topology makes the most out of the reduced supply voltage, but the fully differential circuitry requires a common-mode feedback (CMFB) loop, which is sketched in Figure 3.8b. Very large resistors $(R_a > 100k\Omega)$ are used in order to sense the DC voltage at the inputs and feed it into the loop. A basic amplifier OpAmp2 performs the comparison with the reference voltage V_{ref} in order to provide the required bias voltage V_{bias_2} to pair $M_5 - M_6$. This solution provides reasonable phase margin, even though the gate of $M_5 - M_6$ is a high-gain node; the filter $R_b - C_b$ offers one extra degree of freedom to adjust the open-loop frequency-response of the CMFB. Referring again to Figure 3.7c, we notice one additional problem: super-cascode transistors $M_1 - M_4$ must be biased with a current which is directly *spilled* from the bias current of VCO's core. Therefore, such bias current must be determined accurately, and for this reason reference voltage V_{ref} must be generated internally. The complete schematic is displayed in Figure 3.8c: devices $M_{1a} - M_{4a}$ are a replica of the super-cascode stage $M_1 - M_4$, they are biased with an adjustable reference current I_{bias_2} , and the source of $M_{1a} - M_{2a}$ provides the correct V_{ref} . OpAmp2 is a simple differential pair with an active load $(M_{13} - M_{16})$ in the figure). A summary of measured results for the LMV in [1] is presented in Table 3.2, together with simulated parameters for the amplifier in Figure 3.8; the micrograph of the chip is reproduced in Figure 3.9, and some measurement results are reported in Figure 3.10.

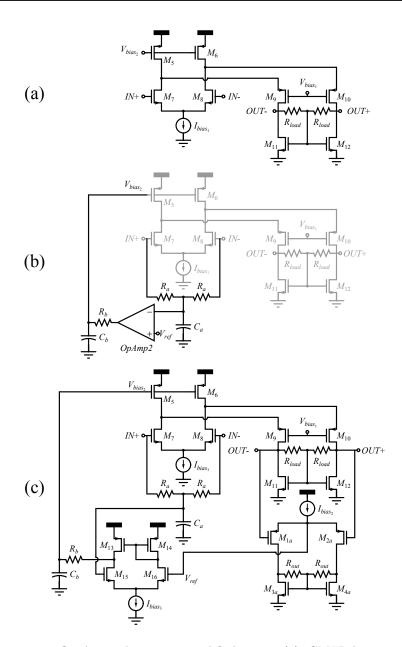


Figure 3.8: Op-Amp schematic: simplified circuit (a), CMFB loop in evidence (b) and complete schematic (c).

Table 3.2: Summary of measured parameters.

5.2. Sammary or m	casarca param
LMV cell in [1]	_
Gain	36dB
NF $(3 \div 5 \text{MHz})$	4.8 dB
IIP_3	-19 dBm
$CP_{1dB,out}$	-31 dBm
PN@1MHz	$-104 \mathrm{dBc/Hz}$
$I_{\rm bias}$	$4.5 \mathrm{mA}$
V_{DD}	1.2V
Amplifier in Fig	gure 3.8
Gain	35 dB
GBWP	$135 \mathrm{MHz}$
$\sqrt{\langle i_{ m n,in}^2 angle}$	$4.25 \mathrm{nA}/\sqrt{\mathrm{Hz}}$
$\dot{\phi}_m$ (CMFB loop)	68^o

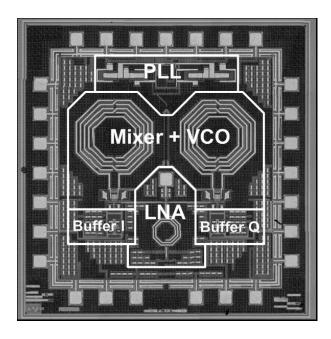


Figure 3.9: Micrograph of the chip presented in [1].

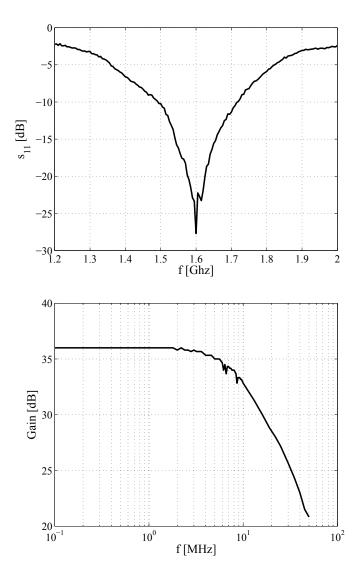


Figure 3.10: Measurement results from [1]: s_{11} plot (top) and conversion gain (bottom). Input reflection coefficient is lower than -10dB in the range $1.5 \div 1.7 \mathrm{GHz}$, and $\mathrm{BW}_{-3\mathrm{dB}} \approx 10 \mathrm{MHz}$.

Chapter 4

Toroidal inductors in standard CMOS processes

Planar spiral inductors are key passive components in any integrated RF design [16]. In Table 4.1 several implementations from the present stateof-the-art are summarized, and we can see that these components are extremely area-hungry; recently it has been proposed to mitigate this issue by integrating active components underneath [83][84], but such coils have anyway the strong tendency to induce RF currents into the substrate. In some applications, magnetic field confinement might be needed, especially when large systems with several different coils are integrated on the same die: mutual interference between coils is likely going to jeopardize overall performances. In this sense, solenoid-based inductors can be a handy choice, and toroidal geometry is definitely preferable because it is very compact and symmetrical: even the internal area can be actively used. Good properties of toroidal inductors have been exploited in advanced micromachining processes [85], but CMOS processes present much different issues. In the following pages a possible approach to toroids' integration is presented: the aim of this implementation is not competing in terms of performances, but rather demonstrating the feasibility of this kind of

Table 4.1: Performance of integrated inductors in CMOS-compatible technologies.

	Area	L	${ m Q@f_{ m peak}}$	Technology	
	$[\mu \mathbf{m}^2]$	[nH]			
[67]	180×180	0.127	$48@10\mathrm{GHz}$	$0.25 \mu\mathrm{m}$	
[68]	240×240	4	$22@5\mathrm{GHz}$	N/A	
[69]	$120{\times}120$	2	$11.2@5\mathrm{GHz}$	$0.25 \mu\mathrm{m}$	
[70]	N/A	2	$8@2\mathrm{GHz}$	$0.25 \mu\mathrm{m}$	1
[71]	300×300	5.6	$9@2.4{\rm GHz}$	$0.18 \mu \mathrm{m}$	2
[72]	160×200	1.3	$15@8\mathrm{GHz}$	$0.35 \mu\mathrm{m}$	
[73]	21×21	0.42	${>}15@50{\rm GHz}$	$0.13 \mu \mathrm{m}$ SiGe	
[74]	$390{\times}390$	2.3	$34@7.5\mathrm{GHz}$	$90\mathrm{nm}$	1
[75]	N/A	3	$10.8@3.5\mathrm{GHz}$	$0.18 \mu \mathrm{m}$	
[76]	N/A	3.4	$15.7@2.45{\rm GHz}$	$0.13 \mu \mathrm{m}$	3
[77]	N/A	0.93	$24.9@5\mathrm{GHz}$	$0.13 \mu \mathrm{m}$	
ib.	N/A	3.5	$16@2.4\mathrm{GHz}$	$0.13 \mu \mathrm{m}$	
[78]	N/A	N/A	$16@6\mathrm{GHz}$	$0.35 \mu\mathrm{m}$	
[79]	380×380	1.8	$41@5\mathrm{GHz}$	$90\mathrm{nm}$	1
[80]	825×825	1	$51@1{\rm GHz}$	N/A	4
[81]	$165{\times}165$	2	$12@10.5\mathrm{GHz}$	$90\mathrm{nm}$	
[82]	$232{\times}232$	1.1	${>}13@20\mathrm{GHz}$	$0.13 \mu \mathrm{m}$	

¹ The coil is not exactly integrated on the die, therefore area consumption is not an issue, but the WLP technology is nonetheless compatible with CMOS processing.

² The dimensions of the coil are estimated by inspecting the die micrograph.

³ SOI processing is applied.

⁴ A CMOS-compatible post-processing is applied.

solution, with respect to obvious technology limitations. The figures summarized in Table 4.1 reveal that planar coils are quiet effective in terms of quality factor, which is the ultimate bottle-neck in the vast majority of RF designs. Nevertheless, cross-talking can become a serious issue in the future, since nanometre-scale technologies offer the possibility to integrate larger and larger systems on the same die. Under this perspective, field-confinement properties of toroidal inductors become appealing, also because internal area might be effectively utilized.

4.1 Physical implementation

The structure resembles a rectangular-section toroid, and the windings are evenly spaced along the circle. The single winding is built by using the top metal layer running along the radius, a stack of vias going down to the bottom metal level, the bottom metal layer itself running along the radius in the opposite direction, and another stack of vias to get back at top metal level. A graphical representation can be seen in Figure 4.1. Planar inductors can be accurately characterized by means of 2.5D electro-magnetic simulations: the structure is basically extended over 2 dimensions, and this hugely simplifies the analysis. In the case of toroids, pure 3D analysis is necessary, since the structure itself is not planar. Nevertheless the dimensions involved are extreme, since the total area is in the order of fractions of mm^2 , whereas the distance between close windings can be as small as a few μm : simulations might be extremely slow and compromises must be accepted in terms of accuracy. Some qualitative analysis can be performed, though. Figure 4.2 shows a possible setup: the 3D structure is enclosed into a box (resembling the oxide), and the two terminals are driven out for probing on opposite sides of the box, on top and bottom. The negative terminal (together with the bottom plane of the box) is grounded, whereas the positive terminal is driven with a probing signal (basically a constant voltage). The toroid (on the left) is compared with a 2-turn square planar inductor (on the right), which occupies roughly the same area and should provide (according to design kit models) approximately the same inductance. This setup has been simulated with a 3D finite-elements simulator (COMSOL MultiphysicsTM) in order to provide a qualitative analysis of the structure with respect to the most significant property: field confinement. The normal component of the magnetic field is plotted in magnitude on the silicon plane (in the middle of the picture) and on the cross section

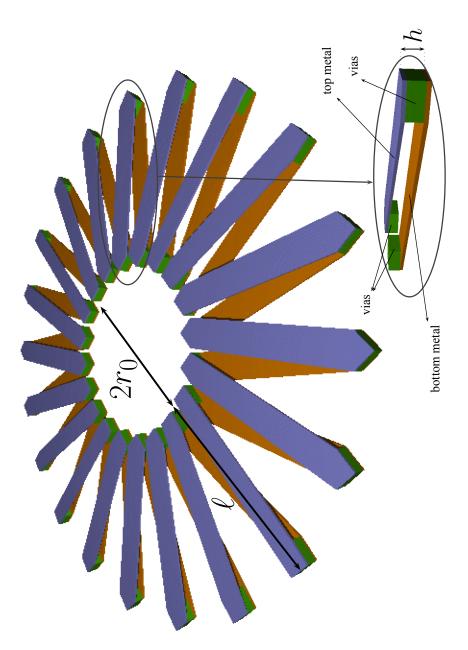


Figure 4.1: 3D representation of the complete structure.

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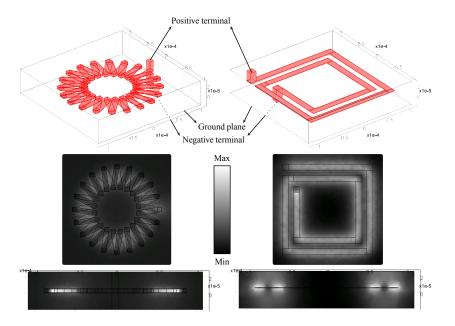


Figure 4.2: 3D representation of a finite-elements simulation setup (on top) together with the results, that show the distribution of the magnetic field magnitude along the silicon plane (middle) and the cross-section (bottom), for the toroidal structure (left) and a reference planar inductor (right).

(at the bottom): the simulation shows that the field is reasonably enclosed in the toroid, especially when compared with the planar spiral inductor. A close-up plot, shown in Figure 4.3, reveals that the field is actually running along the path inside the toroid. In the following, we try to exploit these properties and investigate the behaviour of the component through simple straightforward calculations.

4.2 Modelling

Two different approaches are examined in order to analyze the structure and provide reasonable design guidelines. The aim is not providing an accurate model, but rather clear indications about different trade-offs.

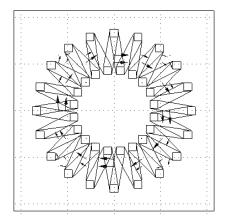


Figure 4.3: Magnetic flux density simulation along the silicon plane for the toroidal inductor.

4.2.1 Classical approach

An air-core rectangular-section toroidal inductor can be easily modelled by means of static Maxwell's equations. A straightforward calculation leads to a well known value for the inductance:

$$L_{coil} = \frac{\mu_0 h}{2\pi} n^2 \log\left(\frac{r_0 + \ell}{r_0}\right),\tag{4.1}$$

where

 μ_0 is the magnetic permeability in air (the core is actually built with multilayer oxide, but magnetic properties do not change significantly);

h is the height of the rectangular section of the toroid, as distance between top and bottom metal layers;

n is the number of windings;

 r_0 is the internal radius of the toroid;

 ℓ is the difference between external and internal radius.

The calculations are carried on under rather restrictive assumptions:

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- 1. wire dimensions are negligible;
- 2. flux dispersion outside the toroid is very low;
- 3. magnetic field is constant along toroid symmetry axis.

None of them is actually verified in this case, but we might get some useful indications anyway. The square dependency shown in (4.1) is very appealing: very high Q values could be achieved with many windings, since the resistance is growing linearly with n. However, in practical implementations the physical dimensions must be considered. A major shortcoming for this kind of structure would be the dominant impact of contact resistance associated with vias. In fact, the availability of copper vias is a necessary precondition for the implementation of the coil. In the rest of this section we will consider only the contribution associated with the metal lines, but we will bear in mind that in practice that contribution can be as important as the contact resistance (if not even less important). With this simplification we get a total DC resistance

$$R_{\rm DC} \approx 2\overline{R}_{\square} \frac{\ell}{w} n,$$
 (4.2)

where \overline{R}_{\square} is the average metal sheet resistance in Ω/\square , and w is the width of the metal stripe. Furthermore, windings cannot be infinitely close to each other; if s is the minimum spacing between close lines, we can pack the windings as much as possible, but ultimately we get:

$$2\pi r_0 \approx n(w+s) \Rightarrow n \approx \frac{2\pi r_0}{w+s}.$$
 (4.3)

Combining equations (4.1), (4.2) and (4.3) and assuming a simple expression for the impedance $(Z(j\omega) \approx R_{\rm DC} + j\omega L_{coil})$, we can find a compact and elegant expression for the quality factor normalized with respect to the frequency:

$$\frac{Q}{\omega} \approx \frac{L_{coil}}{R_{\rm DC}} \approx \frac{\mu_0 h}{2\overline{R}_{\square}} \cdot \frac{w}{w+s} \cdot \frac{\log\left(1 + \frac{\ell}{r_0}\right)}{\frac{\ell}{r_0}}.$$
 (4.4)

This final result shows that performances are basically limited by technology, because the upper limit for (4.4) is in any case $\mu_0 h/2\overline{R}_{\square}$. Some marginal optimization can be achieved using wide and short metal stripes. The design parameters in Table 4.2 lead to an inductance value of $\sim 0.3nH$.

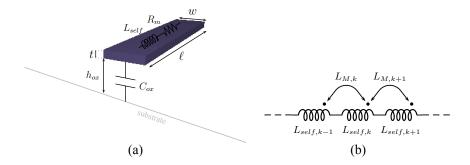


Figure 4.4: A single metal stripe with related parameters (a) and a series of mutually-coupled inductors (b).

4.2.2 Local approach

A self-inductance is associated with the field generated locally by each metal line [86]. This self-inductance can be approximated by means of a widely used expression:

$$L_{\text{self}} = \frac{\mu_0}{2\pi} \ell \left[\log \left(\frac{2\ell}{w+t} \right) + 0.50049 + \frac{w+t}{3\ell} \right], \tag{4.5}$$

where ℓ is the length of the line, w is the width and t is the thickness, as shown in Figure 4.4(a). This approach is extremely versatile, because mutual correlations between neighbours can be calculated as well; the final expression for the total inductance will be in the form:

$$L = \sum_{k=1}^{N} L_k + \sum_{k \neq j}^{\text{mutual}} L_{k,j}. \tag{4.6}$$

where N is the number of metal lines. For example, in Figure 4.4(b) we can see a special case where mutual coupling is considered only between adjacent lines, leading to a total inductance which is still linearly dependent with respect to the number of windings. The expressions in [86] are tailored around a planar structure, nevertheless it might be possible to consider special correction factors in order to accommodate a full 3D structure like the toroid. Such kind of approach has been followed [85] with a good matching between experimental results and numerical model. To get quick

4.3 Design 61

indications, we shall try to focus on fundamental results. The most important fact is that the total inductance is basically linear with respect to the number of metal stripes, and therefore to the number of windings n. We choose not to consider any kind of mutual effect, as a first-order approximation. Referring to a single metal line as in Figure 4.4(a), we define L_{self} as in (4.5), and

$$R_m \approx R_{\Box} \frac{\ell}{w},$$
 (4.7)

where R_{\square} has appropriate values for top and bottom metal layers. Based on these assumptions, we find out some indications for the final design: metal lines should be wide and long.

4.3 Design

As previously shown, the two modelling approaches provide different indications, but the first one is not applicable in this case, because it's based on assumptions that are not verified; we shall focus on the indications from the second approach. We can simplify the calculations considering a single turn, since no mutual coupling between turns is considered. Resistance from vias stacks should be added to the total resistance, and then some capacitance should be considered in order to estimate self-resonance. Stray capacitance between turns can be easily neglected, as the main contribution comes from bottom metal-to-substrate capacitance. The expression for the single stripe is straightforward:

$$C_{ox} = \epsilon_{ox} \frac{w\ell}{h_{ox}},\tag{4.8}$$

where ϵ_{ox} is the effective dielectric constant of the oxide, and h_{ox} is the distance between the bottom plate and the substrate. The expression is valid for bottom metal lines, since we consider $C_{ox} = 0$ for top metal lines. In order to dimension a test-coil, we decide to follow some general guidelines:

• the total inductance should be relatively large $(L \gtrsim 1nH)$, otherwise de-embedding the parasitics can produce misleading results,

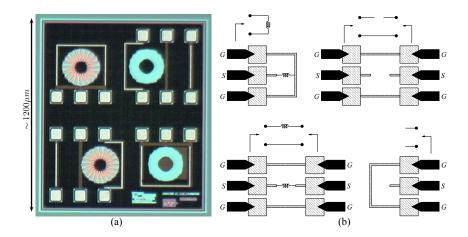


Figure 4.5: Micro photograph of the die (a) and measurement setups (b). From the upper-left quadrant, clockwise: 1-port test setup, 2-port deembedding setup, 1-port de-embedding setup, 2-port test setup.

- parasitic capacitance should not limit the self-resonance too much (a reasonable target is $f_{\rm self} \gtrsim 5 GHz$),
- DC resistance should of course be minimized.

In addition, regular design rules for CMOS processes must be followed. For example, arbitrarily wide metal lines cannot be fabricated. Design choices and estimated parameters are summarized in Table 4.2. The physical layout of the device is particularly complicated because of the geometry of the toroid itself. It is useful to build a parametric automated procedure to perform all the drawing steps; for the details, see Appendix D. The complete silicon die can be seen in Figure 4.5a. It has been fabricated in a standard CMOS $0.13\mu m$ process with 6 copper metal lines and copper vias.

4.4 Measurement approach and results

The measurements are carried directly on the die by means of standard Ground-Signal-Ground (GSG) microprobes. A 1-port and a 2-port mea-

surement setup is provided on each die (even though only 2-port measurements are used in practice for data analysis, 1-port measurements are useful for quick checks), together with de-embedding counterparts, in order to rule out contributions from pads and parasitics. De-embedding is accomplished by means of an open-circuit structure, as shown in Figure 4.5b. Data analysis is carried on through manipulation of admittance parameters, since the de-embedded values are given by a simple difference:

$$\mathbf{Y_{dut}}(\omega) = \mathbf{Y_{dir}}(\omega) - \mathbf{Y_{de}}(\omega),$$
 (4.9)

where

 $\mathbf{Y_{dut}}(\omega)$ is the de-embedded result,

 $\mathbf{Y}_{dir}(\omega)$ is the measured admittance matrix of the coil,

 $\mathbf{Y}_{\mathbf{de}}(\omega)$ is the measured admittance matrix of the open-circuit de-embedding structure.

This approach provides reasonable results, though some limitations exist:

- short-circuit de-embedding is not performed, and this might have a significant impact especially on Q factor measurements [87]. The quality factor is going to be underestimated, but previous implementations in older technologies demonstrated that the underestimation is not dramatic;
- the inductor itself can shield the pads [88], but for the sake of simplicity we chose not to perform any calibration aimed to alleviate this issue. Because of this effect, the Q factor will suffer further underestimation.

To get a closer insight, the idea is to build a very simple Π model like the one showed in Figure 4.6, based on measurement results. Such model is extremely rough, since it does not contain frequency-dependent elements nor mutual couplings, but it is useful in order to make comparisons with hand calculations. More sophisticated approaches can be found in the literature [89].

Table 4.2: Design parameters and estimated values.

Table 4.3:	Equivalent	П	model	pa-
rameters.				

Param.	Value
r_0	$\sim 40 \mu m$
ℓ	$\sim 60 \mu m$
w	$\sim 12 \mu m$
n	20
\overline{L}	$\sim 1.4nH$
R	$\sim 6.2\Omega$
C_{ox}	$\sim \! 0.45 pF$

Param.	Value
\overline{L}	~1.1nH
R	$\sim 6.13\Omega$
C_{ox}	$\sim 0.42 \mathrm{pF}$
C_{sub}	$\sim 5 \mathrm{fF}$
R_{sub}	${\sim}115\Omega$
C_{stray}	$\sim 2 \mathrm{fF}$

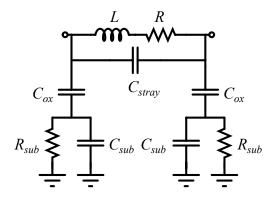


Figure 4.6: Simple Π model of the inductor.

The procedure is totally empirical: model and measurements data are fed into a simulation tool which applies a straightforward optimization routine to extract the model parameters. The results are summarized in Table 4.3.

It should be easily appreciated the fact that extracted parameters show a good matching with hand calculations already seen in Table 4.2. The inductance itself is easily the most critical parameter, but this is no surprise since we did not account for any mutual correlation between windings. The de-embedded data can be easily manipulated in order to extract relevant performance information. In Figure 4.7 a plot of the Q factor is shown, together with the effective inductance at different frequencies. We can see a clear trend: the quality factor is increasing as the frequency grows, as

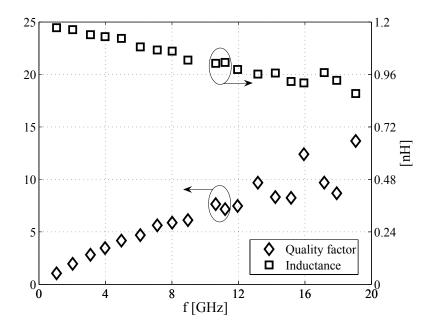


Figure 4.7: Measured Q factor and inductance of the coil.

anticipated. This means that eddy currents induced in the substrate do not interfere heavily. The inductance value is reasonably constant over the whole range. In Figure 4.8 we can see the real and imaginary part of the impedance, whereas in Figure 4.9 we show the S-parameters on a Smith chart, for 1-port and 2-port measurements. The model is of course not very accurate at high frequencies.

The practical manufacturability of a toroidal inductor with good areaconsumption characteristics has been demonstrated. Results show that this implementation cannot challenge a traditional spiral inductors on a performance level, but field confinement properties fulfil the expectations, and high-frequency behaviour suggests that certain applications can greatly benefit from this kind of approach. First-order modelling has been really helpful in the design phase, showing also decent matching with measurements, but nonetheless great improvements can be achieved in this area.

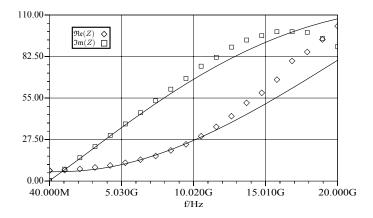


Figure 4.8: One-port measurement of the impedance, shown in terms of real and imaginary part, together with the Π -model approximation (solid lines).

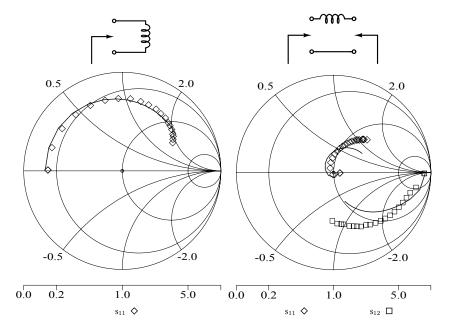


Figure 4.9: S-parameter plots for 1-port and 2-port measurements (on the left and on the right side respectively), together with the Π -model approximation (solid lines).

Chapter 5

Conclusions

"Σκόπει δή, [...] ໕ Κέβης, εἰ ἐκ πάντων τῶν εἰρημένων τάδε ἡμῖν συμβαίνει [...];"
"Then reflect, [...] Cebes: is not this the conclusion of the whole matter [...]?"
Plato, Phaedo

The aim of this work is to exploit the capabilities of CMOS processes for RF applications. The promise of an ever-growing level of integration and interaction with the digital domain is very appealing, and the goal of a full-fledged single-chip radio looks close [90][91]. Portable devices will especially benefit from this (r)evolutionary path: a very compact and low-power structure will embrace several wireless standards at a time, thereby providing multiple functions (e.g. GPS, UMTS, Bluetooth and so on...) with a respectable battery life. In order to offer a wide perspective of the subject, different aspects of RF IC design have been taken into consideration.

The application of negative feedback to the design of RX front-ends has been approached in Chapter 2 on a simulation level. The possibility to embrace a traditional Gilbert-based mixer in the loop without the need for radical circuit modifications has been presented. From this starting point, the concept of dual-loop feedback has been proposed as a viable solution for the very demanding trade-off between input matching, noise and

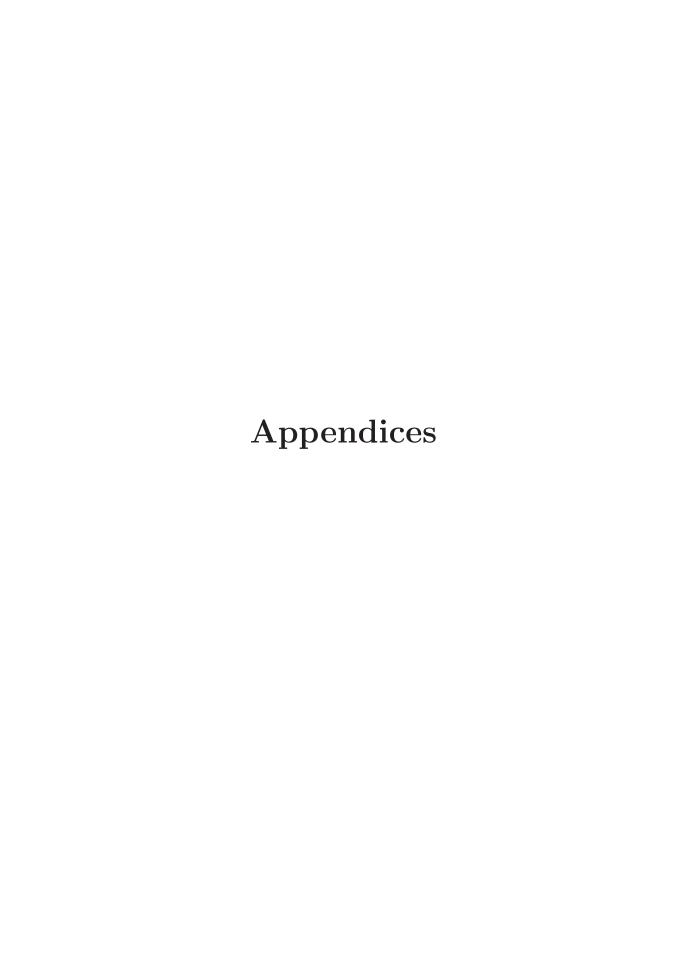
68 Conclusions

linearity. A first-order analysis has been performed with the aid of simplified BSIM3 models and detailed calculations: the analysis demonstrates radical linearity improvements, stable performance and reliable noise behaviour. A single-loop narrow-band solution has been introduced as a straight-forward application of the feedback approach: the possibility to improve noise performances (especially for low-performances cheap technologies) has been shown, and a reliable (though very simplified) noise analysis has been sketched and compared with simulations based on the aforementioned models.

Current-reuse techniques can be successfully applied to ultra low-power systems in order to shrink power consumption to the lowest possible limit. A very effective approach has been presented in Chapter 3: the LMV cell efficiently performs the difficult task of integrating an LNA, a mixer and a VCO over the same bias current. A simplified analysis of the structure has been derived, in order to fit the complex trade-offs related to the different blocks sharing the same components. A detailed description of the current-mode output stage has been offered, and measurements results have been summarized, in order to show the remarkable properties of the LMV stage.

In top-performance RF systems the use of integrated inductors is not only a necessity, but an ultimate bottle-neck.¹ As CMOS technology moves on, we can expect improvements in performance and characterization of integrated planar inductors, but the devices remain bulky and difficult to scale, and they present the unpleasant *feature* of inducing strong RF currents in the substrate. An alternative solution has been portrayed in Chapter 4 in the form of a toroidal inductor. The first implementation in CMOS process has been presented, together with a first-order modelling approach and measurements. Overall performances show difficulties in competing with more traditional planar inductors, but the ultimate target of field-enclosure has been achieved without penalties in terms of area consumption.

 $^{^1\,\}rm ``From$ the point of view of RF circuits, the lack of a good inductor is by far the most conspicuous shortcoming of standard IC processes." [16]



Appendix A

Multiplier-based mixers with passive load

As we have shown in Chapter 2, node v_m in Figure 2.1 (the picture is reproduced in Figure A.1 together with the simplified schematic based on the full-switched model) presents an amplified unmodulated replica of the RF

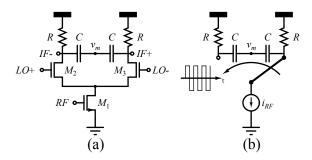


Figure A.1: A down-conversion mixer based on the Gilbert cell: schematic (a) and full-switched model (b).

current signal provided by transconductor M_1 . The fundamental results concerning the down-converted signal and the node v_m are summarized in equations (2.8)-(2.9), which are reported here for convenience:

$$v_{IF} = -i_{RF} \cdot m_{LO} \cdot \frac{R}{1 + sCR}, \tag{A.1}$$

$$v_m = -i_{RF} \frac{R}{2}. \tag{A.2}$$

The node v_m is actively used into a feedback topology: if we connect it to another node v_b in the loop, we must consider the influence of the impedance Z_b at that node. Referring to Figure A.2, the calculations get a little bit involved, but they can be summarized as follows:

$$Z_{h} \triangleq \frac{v_{h}}{i_{in}} = \frac{R\left(\frac{1}{sC} + \frac{Z_{b}\left(R + \frac{1}{sC}\right)}{Z_{b} + R + \frac{1}{sC}}\right)}{R + \frac{1}{sC} + \frac{Z_{b}\left(R + \frac{1}{sC}\right)}{Z_{b} + R + \frac{1}{sC}}},$$
(A.3)

$$Z_t \triangleq \frac{v_t}{i_{in}} = \frac{R\left(\frac{R \cdot Z_b}{Z_b + R + \frac{1}{sC}}\right)}{R + \frac{1}{sC} + \frac{Z_b\left(R + \frac{1}{sC}\right)}{Z_b + R + \frac{1}{sC}}},\tag{A.4}$$

$$Z_{m} \triangleq \frac{v_{m}}{i_{in}} = \frac{R \frac{Z_{b}(R + \frac{1}{sC})}{Z_{b} + R + \frac{1}{sC}}}{R + \frac{1}{sC} + \frac{Z_{b}(R + \frac{1}{sC})}{Z_{b} + R + \frac{1}{sC}}}.$$
(A.5)

After some manipulation we find out that (A.1) still holds: the output mixing product is unaffected. We can then calculate the voltage at node v_m :

$$v_m = -i_{RF} \cdot Z_m = -i_{RF} \frac{sCRZ_b}{1 + sC(R + 2Z_b)}.$$
 (A.6)

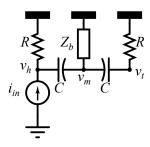


Figure A.2: R-C low-pass network with buffer input impedance.

Obviously the limit of the previous expression is the same that we found in (A.2):

$$\lim_{Z_h \to \infty} Z_m = \frac{R}{2}.\tag{A.7}$$

We can evaluate two different situations:

$$Z_b = R_b \quad \Rightarrow \quad Z_m = \frac{sCRR_b}{1 + sC(R + 2R_b)},$$
 (A.8)

$$Z_b = \frac{1}{sC_b} \Rightarrow Z_m = \frac{\frac{C}{C_b}}{1 + \frac{2C}{C_b}} \cdot \frac{R}{1 + \frac{sCR}{1 + \frac{2C}{C_b}}}.$$
 (A.9)

In the first case the transfer function resembles a first-order high-pass filter, whose cut-off frequency is smaller than $\omega_0 \triangleq \frac{1}{RC}$. Since $\omega_0 \ll \omega_{RF}$, we can conclude that the transfer at node v_m is basically flat at the frequencies of interest. This is the most relevant case for the applications presented in Chapter 2. In the second case $(Z_b = \frac{1}{sC_b})$ we have a first-order low-pass filter; the transfer is acceptable as long as the cut-off frequency is reasonably larger than the RF carrier, which means:

$$\omega_1 \triangleq \frac{1 + \frac{2C}{C_b}}{CR} \gg \omega_{RF} \Rightarrow 1 + \frac{2C}{C_b} \gg \omega_{RF} CR.$$
 (A.10)

Since $\omega_{RF}CR \gg 1$, we conclude that C_b must be extremely small compared to C.

In practical implementations, the value of resistance R will be such that the voltage drop across the resistor will be too high. In order to provide a more compact biasing solution, we might think about draining DC current from the tail generator, as in Figure A.3. Solutions (a) and (b) are perfectly equivalent in principle to the original solution shown in Figure A.1a, and therefore all results presented in equations (A.1)-(A.2) hold. There are some differences that can not be seen at a first glance, and solution (b) should perform better in terms of

switching: the differential pair must switch a smaller amount of current;

conversion gain: the output resistance of the DC current source should be compared with the switching pair input resistance rather than the load resistance R.

 $^{^{1}\}omega_{RF}$ is the RF carrier frequency.

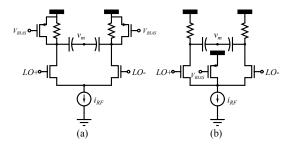


Figure A.3: Ideal mixers with active DC current drain.

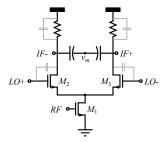


Figure A.4: A resistive-load mixer with parasitic components.

There is another point, that is more subtle. In Figure A.4 we have indicated some critical parasitics that can affect the functionality of the stage. From the point of view of the differential signal $i_{M_2} - i_{M_3}$ (which is going to be down-converted) parasitic capacitors at nodes IF+ and IF- provide simply an additional low-pass filtering (at first order); however, the same capacitors attenuate the common-mode signal $i_{M_2} + i_{M_3}$ and affect the transfer at node v_m directly. Given a certain capacitance C_{par} at both nodes, equation (A.2) becomes

$$v_m = -\frac{1}{2} \frac{R}{1 + sRC_{par}}. (A.11)$$

All these considerations lead us to prefer the solution in Figure A.3b, since the influence of parasitics at output nodes should be minimized.

Appendix B

Nullor-based representation of amplifiers

Transmission parameters provide a very convenient way to represent the behaviour of 2-port linear circuits. We can refer to Figure B.1 to define the relationships between relevant signals as follows:

$$\begin{pmatrix} v_{in} \\ i_{in} \end{pmatrix} = \underbrace{\begin{pmatrix} A & B \\ C & D \end{pmatrix}}_{\text{chain matrix}} \cdot \begin{pmatrix} v_{out} \\ i_{out} \end{pmatrix}. \tag{B.1}$$

Transmission parameters A, B, C and D can be reciprocated to obtain transfer parameters, somewhat more familiar in electronic design:

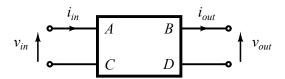


Figure B.1: A 2-port circuit represented by its transmission parameters.

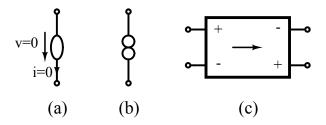


Figure B.2: Symbols of a nullator (a), a norator (b) and a nullor (c).

An ideal amplifier can be represented by means of two components, a *nullator* and a *norator* [92]:

- a nullator (see Figure B.2a) is an ideal 2-terminal element which enforces both a zero-voltage and a zero-current between its terminals,
- a norator (see Figure B.2b) is an ideal 2-terminal element which imposes no constraints on its branch current and voltage.

The *nullor* is a combination of a nullator and a norator. Such element is associated with an all-zeros chain matrix:

$$\begin{pmatrix} v_{in} \\ i_{in} \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \cdot \begin{pmatrix} v_{out} \\ i_{out} \end{pmatrix}. \tag{B.2}$$

In Figure B.2c the common symbol for the nullor is given. An example of a single-amplifier topology represented with the use of a nullor is given in Figure B.3 [93]. The nullator imposes the same voltage v_{in} at both its terminals, but no current flows into them. Therefore the current flowing into R_1 is determined entirely by v_{in} , and the norator determines the

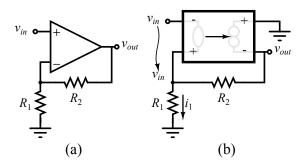


Figure B.3: A conventional operational amplifier used in a non-inverting feedback topology (a) and its representation based on the nullor (b).

output voltage accordingly. The following equations describe the behaviour of the circuit and determines familiar relationships:

$$\begin{split} i_1 &= \frac{v_{in}}{R_1}, \\ v_{out} &= v_{in} + i_1 \cdot R_2 = v_{in} \left(1 + \frac{R_2}{R_1}\right). \end{split}$$

Appendix C

A unified approach to the analysis of the LMV cell

In Chapter 3 we have performed a simplified analysis of the LMV cell based on two approaches, leading to equations (3.10) and (3.20); such equations describe the behaviour of the stage by means of two different asymptotic scenarios. In the following, we will try to provide a *single* model that accounts for both aspects of the problem. We will modify the approach followed in Section 3.2 in order to accommodate additional lossy mechanisms that limit the GBWP of the stage for low values of parasitic capacitance. The starting point is again Figure 3.4 and Figure 3.5. First of all, we have to examine carefully what happens during transition $\phi_1 \longrightarrow \phi_2$ (and similarly $\phi_2 \longrightarrow \phi_1$). Referring to Figure C.1, we consider only the core of the oscillator and we apply a test current i_{probe} . During the transition, transistors $M_3 - M_4$ work as current generators, and we can

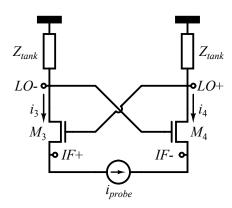


Figure C.1: Evaluation of resistance at the sources of M_3 and M_4 .

sum up all relevant equations as follows:

$$i_3 = g_{m_3}(v_{LO+} - v_{IF+})$$
 (C.1)

$$= g_{m_3}(-Z_{tank} \cdot i_4 - v_{IF+}) \tag{C.2}$$

$$= g_{m_3}(Z_{tank} \cdot i_{probe} - v_{IF+}), \qquad (C.3)$$

$$i_4 = g_{m_4}(v_{LO-} - v_{IF-})$$
 (C.4)

$$= g_{m_4}(-Z_{tank} \cdot i_3 - v_{IF-}) \tag{C.5}$$

$$= g_{m_4}(-Z_{tank} \cdot i_{probe} - v_{IF-}), \qquad (C.6)$$

$$G_{d} \triangleq \frac{i_{probe}}{v_{IF-} - v_{IF+}} = \frac{g_{m_{4}}(-Z_{tank} \cdot i_{probe} - v_{IF-})}{g_{m_{3}} + g_{m_{4}} - 2g_{m_{3}}g_{m_{4}}Z_{tank}}.$$
 (C.7)

Capacitor C_d is discharged almost instantaneously, therefore we can approximate (C.7) as follows:

$$G_d \approx \frac{g_{m_3}g_{m_4}}{g_{m_3} + g_{m_4}}.$$
 (C.8)

Assuming that the discharge is taking place in a very short time interval $[-\Delta t, \Delta t]$ around the transition instant, the charge loss can be easily

determined:

$$-C_d \frac{\mathrm{d}v_{IF}}{\mathrm{d}t} = G_d v_{IF} \tag{C.9}$$

$$\downarrow \qquad \qquad \downarrow$$

$$v_{IF}(\Delta t) = v_{IF}(-\Delta t)e^{-\frac{1}{C_d} - \frac{\Delta t}{-\Delta t}G_d(t)dt};$$
 (C.10)

$$\Delta Q_d = C_d v_{IF}(-\Delta t) e^{-\frac{1}{C_d} - \frac{\Delta t}{C_d} G_d(t) dt} - C_d v_{IF}(-\Delta t) (C.11)$$

$$\approx -\frac{Q_d(-\Delta t)}{C_d} \int_{-\Delta t}^{\Delta t} G_d(t) dt, \qquad (C.12)$$

We can calculate G_d explicitly by means of equations (53)-(54) in [94]: they have been derived for a common-source switching pair, but they can fit this analysis at first order.

$$g_{m_3} = k_0 V_0 \left(-\sin \varphi + \sqrt{2\sin^2 \Phi - \sin^2 \varphi} \right);$$
 (C.13)

$$g_{m_4} = k_0 V_0 \left(\sin \varphi + \sqrt{2 \sin^2 \Phi - \sin^2 \varphi} \right);$$
 (C.14)

$$k_0 \triangleq \frac{1}{2}\mu_{n0}C'_{ox}\frac{W}{L}; \tag{C.15}$$

$$\Phi \triangleq \arcsin \sqrt{\frac{I_0}{2k_0V_0^2}}.$$
 (C.16)

where $\varphi \triangleq \omega_0 t$, $\Phi = \omega_0 \Delta t$ is the angle corresponding to time interval Δt , and V_0, I_0 are defined as in Section 3.1. By further processing we find:

$$G_{d}(\varphi) = k_{0}V_{0}\frac{\sin^{2}\Phi - \sin^{2}\varphi}{\sqrt{2\sin^{2}\Phi - \sin^{2}\varphi}}, \quad (C.17)$$

$$\int_{-\Delta t}^{\Delta t} G_{d}(t)dt = \frac{1}{\omega_{0}} \int_{-\Phi}^{\Phi} G_{d}(\varphi)d\varphi$$

$$\approx k_{0}V_{0}\frac{\Phi^{2}}{\omega_{0}}$$

$$\approx \frac{I_{0}}{2\omega_{0}V_{0}}, \quad (C.18)$$

$$\Delta Q_d \approx -\frac{Q_d(-\Delta t)}{C_d} \frac{I_0}{2\omega_0 V_0}, \quad (C.19)$$

$$\beta \triangleq \frac{I_0}{2C_d\omega_0 V_0} \Longrightarrow \Delta Q_d \approx -\beta \cdot Q_d(-\Delta t). \tag{C.20}$$

Since we have determined the charge losses during the transition, we can rewrite (3.14) as follows:

redistribution of the charge and losses during sharing of the transition input charge
$$Q_d(n+1) = \overbrace{(1-\alpha)(1-\beta)Q_d(n)}^{\text{redistribution}} + \overbrace{(1-\alpha)Q_{in}(n)}^{\text{redistribution}}, \qquad (C.21)$$

and procede to determine the total transfer function using the same approximations as above.

$$\mathcal{H}(z) \triangleq \frac{Q_d(z)}{Q_{in}(z)} = \frac{1-\alpha}{z - (1-\alpha)(1-\beta)},$$
 (C.22)

$$\mathcal{H}(s) \approx \frac{1-\alpha}{1+\frac{sT_0}{2}-(1-\alpha)(1-\beta)},$$
 (C.23)

$$\frac{v_{IF}}{i_{RF}} = \frac{2}{\omega_0 C_d} \mathcal{H}(s) \quad \approx \quad \frac{2}{\omega_0 C_d} \cdot \frac{1 - \alpha}{\alpha + \beta - \alpha \beta} \cdot \frac{1}{1 + \frac{sT_0}{2(\alpha + \beta - \alpha \beta)}} \cdot (C.24)$$

We can easily verify that the following limits hold:

$$\begin{array}{lll} \beta \gg \alpha & \Longrightarrow & \text{losses due to } C_p \text{ are negligible} \\ & \Longrightarrow & \frac{v_{IF}}{i_{RF}} \approx \frac{2}{\pi} \frac{R_{IF}}{1 + sC_dR_{IF}} \text{ as in (3.10)}, \\ \alpha \gg \beta & \Longrightarrow & \text{losses due to } C_p \text{ dominate} \\ & \Longrightarrow & \frac{v_{IF}}{i_{RF}} \approx \frac{2}{\omega_0 C_p} \cdot \frac{1}{1 + \frac{sT_0}{2\alpha}} \text{ as in (3.20)}. \end{array}$$

In Figure C.2, numerical results from (C.24) are shown together with simulation results (simulation parameters are the same as in Table 3.1, and reproduced in Table C.1 for convenience). It is possible to refine the figures given by (C.24) by means of numerical analysis: referring to Figure C.3, we consider $Z_{tank} = 0$, and we write the differential equations at nodes IF+ and IF-.

$$(C_p + C_d) \frac{\mathrm{d}v_{IF-}}{\mathrm{d}t} - C_d \frac{\mathrm{d}v_{IF+}}{\mathrm{d}t} + g_{m_4} v_{IF-} = 0;$$
 (C.25)

$$(C_p + C_d) \frac{\mathrm{d}v_{IF+}}{\mathrm{d}t} - C_d \frac{\mathrm{d}v_{IF-}}{\mathrm{d}t} + g_{m_3} v_{IF+} = 0.$$
 (C.26)

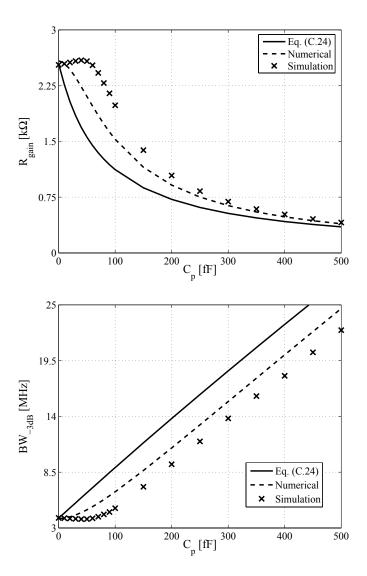


Figure C.2: Calculation results from (C.24) without and with numerical correction (continuous and dashed lines respectively) together with simulation results ("x" markers). The low-frequency gain (on top) and the -3dB bandwidth (bottom) are shown.

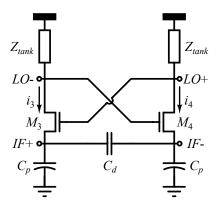


Figure C.3: Equivalent circuit for the numerical analysis based on differential equations at nodes IF+ and IF-.

Table C.1: Main parameters for results in Figure C.2.

f_0	1.6GHz
I_0	$1 \mathrm{mA}$
L_{tank}	$5\mathrm{nH}$
$Q@f_0$	20
C_d	10pF

To solve the system, we use expressions for $g_{m_3} - g_{m_4}$ directly, as given by (C.13)-(C.14). The results are plotted in Figure C.2 (dashed lines). Numerical corrections provide good matching, especially for high values of C_p , but the behaviour in the *transition* region, which is probably going to be more relevant for practical design, is better grasped by the asymptotic limits in (3.10) and (3.20).

 $^{{}^{1}}C_{p} < 100$ fF in the plot.

Appendix D

Layout automation based on SKILL scripting language

The layout of integrated inductors is a tedious procedure that can be greatly sped up using automated tools. Cadence design environment offers a very advanced scripting language named SKILL. A parametric cell can be created using this language: once the cell has been added to the design library, the layout tool will prompt the user a single mask containing all customizable parameters, and then it will design the cell according to the input.

In the case of the toroidal inductor described in Chapter 4, we refer to Figure D.1a to determine the sequence of drawing steps. Given a certain number of turns n, internal radius r_0 , minimum width w_s and cross-section ℓ , before drawing the pattern for the vias we must check that all metal lines respect the minimum clearance, otherwise the parameter r_0 must be adjusted. Once the pattern is drawn, we can complete the structure with metal lines, as shown in Figure D.1b. In the following pages the SKILL code for all the necessary steps is presented. It's preferable to separate the definition of user parameters and the actual design procedure (which in turn can make use of several other procedures).

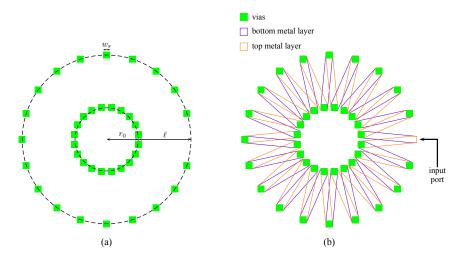


Figure D.1: Top view of the toroid: the pattern for the placement of the vias (a) and the complete structure (b).

Constants

```
Pi=3.141596
sq2=1.414213562
minGrid=0.005; in micrometers
```

This procedure fixes points on the grid

```
procedure(grid(x step)
  step=max( minGrid fix(round(step/minGrid))*minGrid )
  fix(round(x/step))*step
) ;grid
```

A procedure to build multilayer metal connections

```
procedure( rectS( cv w 1 step x0 y0 layerlist vialist)
  let( ( minw, ; minimum dimensions
    layer, ; index
    nx1, ny1, nx2, ny2,; number of contacts for each axis
    wvia1, dist1,; VIAs' geometry parameters
```

```
wvia2, dist2,; VIAs' geometry parameters
   border1, border2,; VIAs' geometry parameters
   POenc ; PADOPEN geometry parameters
 wvia1 = 0.19
 dist1 = 0.29
 wvia2 = 0.36
 dist2 = 0.54
 border1= 0.05
 border2= 0.09
 POenc = 0.7
 minw = 0
 w = max( minw grid( w step ) )
 1 = max( minw grid( 1 step ) )
 x0 = grid(x0 step)
 y0 = grid(y0 step)
 foreach( layer layerlist
    if( layer=="alucap"
      dbCreatePolygon( cv list("alucap" "induct")
       list(x0:y0 x0:(y0+w) (x0+1):(y0+w) (x0+1):y0))
    else
      dbCreatePolygon( cv layer
       list(x0:y0 x0:(y0+w) (x0+1):(y0+w) (x0+1):y0))
    ); if( layer=="alucap" )
  ); foreach( layer layerlist )
 nx1 = fix((1-2*border1+dist1)/(wvia1+dist1))
 ny1 = fix( (w-2*border1+dist1)/(wvia1+dist1) )
 nx2 = fix((1-2*border2+dist2)/(wvia2+dist2))
 ny2 = fix( (w-2*border2+dist2)/(wvia2+dist2) )
 foreach( layer vialist
    if( layer=="M6_M5"
      leCreateContact( cv layer (x0+grid((0.5*1) step)):(y0+grid((0.5*w)
       step)) "RO" wvia2 wvia2 ny2 nx2 (wvia2+dist2) (wvia2+dist2)
"center" "center" )
    else if( layer=="padopen"
      then
        dbCreatePolygon( cv list("padopen" "induct")
          list((x0+P0enc):(y0+P0enc) (x0+P0enc):(y0+w-P0enc)
          (x0+1-P0enc):(y0+w-P0enc) (x0+1-P0enc):(y0+P0enc)))
      else
        leCreateContact( cv layer
```

```
(x0+grid((0.5*1) step)):(y0+grid((0.5*w) step))
          "RO" wvia1 wvia1 ny1 nx1 (wvia1+dist1) (wvia1+dist1)
          "center" "center" ))
    ); multiple if
 ); foreach( layer vialist )
 ) ; let
) ; procedure( rectS )
; User interface
pcDefinePCell(
list( ddGetObj( "LV_PCells" ) "rectStack" "layout" )
( (Width float 11.99)
   (Length float 30.0)
   (Grid float 0.005)
   (XO float 0.0)
   (YO float 0.0)
   (metal1 boolean "TRUE")
   (metal2 boolean "TRUE")
   (metal3 boolean "FALSE")
   (metal4 boolean "FALSE")
   (metal5 boolean "FALSE")
   (metal6 boolean "FALSE")
   (alucap boolean "FALSE")
cv=pcCellView
layerlist = list( )
vialist = list( )
minlayer = 8
maxlayer = 0
when( metal1 == t
 minlayer = 1
 maxlayer = 1)
when( metal2 == t
 minlayer = min( minlayer 2 )
 maxlayer = max( maxlayer 2 ))
when( metal3 == t
 minlayer = min( minlayer 3 )
 maxlayer = max( maxlayer 3 ))
when( metal4 == t
 minlayer = min( minlayer 4 )
 maxlayer = max( maxlayer 4 ))
when( metal5 == t
 minlayer = min( minlayer 5 )
 maxlayer = max( maxlayer 5 ))
when( metal6 == t
```

```
minlayer = min( minlayer 6 )
 maxlayer = max( maxlayer 6 ))
when( alucap == t
 minlayer = min( minlayer 7 )
 maxlayer = max( maxlayer 7 ))
if( minlayer==1
then
 layerlist = cons( "metal1" layerlist)
) ; if( minlayer==1 )
if( minlayer==2
 layerlist = cons( "metal2" layerlist)
); if( minlayer==2)
if( minlayer==3
then
 layerlist = cons( "metal3" layerlist)
); if( minlayer==3)
if( minlayer==4
 layerlist = cons( "metal4" layerlist)
); if( minlayer==4)
if( minlayer==5
 layerlist = cons( "metal5" layerlist)
); if( minlayer==5)
if( minlayer==6
then
 layerlist = cons( "metal6" layerlist)
); if( minlayer==6)
if( minlayer==7
 layerlist = cons( "alucap" layerlist)
); if( minlayer==7)
for( i (minlayer+1) maxlayer
 if( i==2
 then
   layerlist = cons( "metal2" layerlist)
   vialist = cons( "M2_M1" vialist)
 ); if( i==2)
 if( i==3
 then
   layerlist = cons( "metal3" layerlist)
   vialist = cons( "M3_M2" vialist)
 ) ; if( i==3 )
 if( i==4
 then
   layerlist = cons( "metal4" layerlist)
   vialist = cons( "M4_M3" vialist)
 ); if( i==4)
```

```
if( i==5
 then
   layerlist = cons( "metal5" layerlist)
   vialist = cons( "M5_M4" vialist)
 ); if( i==5)
 if( i==6
 then
   layerlist = cons( "metal6" layerlist)
   vialist = cons( "M6_M5" vialist)
 ) ; if( i==6 )
 if( i==7
    layerlist = cons( "alucap" layerlist)
   vialist = cons( "padopen" vialist)
 ) ; if( i==7 )
); for( i (minlayer+1) maxlayer )
layerlist = reverse(layerlist)
vialist = reverse(vialist)
rectS( cv Width Length Grid XO YO layerlist vialist)
) ; pcDefinePCell
```

A procedure to build multilayer metal connections with 45° rotation

```
procedure( rectS45( cv w 1 step x0 y0 layerlist vialist)
let( ( minw, ; minimum dimensions
  weff, leff; effective dimensions
  layer, ; index
 na1, nb1, na2, nb2,; number of contacts for each axis
  wvia1, dist1,; VIAs' geometry parameters
  wvia2, dist2,; VIAs' geometry parameters
  border1, border2; VIAs' geometry parameters
  POenc ; PADOPEN geometry parameters
  wvia1 = 0.19
  dist1 = 0.29
  wvia2 = 0.36
  dist2 = 0.54
  border1= 0.05
  border2= 0.09
  POenc = 0.7
  POenc = grid( POenc*sq2 step )
  minw = 4.4
  w = max( minw w )
```

```
1 = max(w1)
x0 = grid(x0 step)
y0 = grid(y0 step)
weff = max( grid( w/sq2 step ) grid( minw/sq2 step ) )
leff = max( grid( 1/sq2 step ) grid( minw/sq2 step ) )
foreach( layer layerlist
  if( layer == "alucap"
  then
    dbCreatePolygon( cv list("alucap" "induct") list(x0:y0 (x0-weff):
      (y0+weff) (x0+leff-weff):(y0+leff+weff) (x0+leff):(y0+leff)))
    dbCreatePolygon( cv layer list(x0:y0 (x0-weff):(y0+weff)
      (x0+leff-weff):(y0+leff+weff) (x0+leff):(y0+leff)))
  ); if( layer=="alucap")
); foreach( layer layerlist )
na1 = fix((weff-0.5*wvia1+0.5*dist1-2*border1)/(wvia1)
            +dist1))
na2 = fix((weff-0.5*wvia2+0.5*dist2-2*border2)/(wvia2)
            +dist2) )
nb1 = fix( (leff-weff)/(wvia1+dist1) )
nb2 = fix( (leff-weff)/(wvia2+dist2) )
foreach( layer vialist
 if( layer=="M6_M5"
 then
    for( i 1 na2
      leCreateContact( cv layer (x0+leff-weff):(y0+leff+weff-wvia2
        -2*border2-(wvia2+dist2)*(i-1)) "R0" wvia2 wvia2 1 (2*i-1)
        (wvia2+dist2) (wvia2+dist2) "center" "center" )
      leCreateContact( cv layer x0:(y0+wvia2+2*border2+(wvia2
        +dist2)*(i-1)) "R0" wvia2 wvia2 1 (2*i-1) (wvia2+dist2)
        (wvia2+dist2) "center" "center" )
    ); for( i 1 na2 )
    for( i 1 nb2
      leCreateContact( cv layer (x0+leff-weff-(i-1)*(wvia2+dist2)):(y0
        +leff-dist2+2*border2-0.5*wvia2-(i-1)*(wvia2+dist2)) "R0" wvia2
        wvia2 1 2*na2-1) (wvia2+dist2) (wvia2+dist2) "center" "center" )
    ); for( i 1 nb2 )
  else if( layer=="padopen"
  then
    dbCreatePolygon( cv list("padopen" "induct") list(x0:(y0+POenc)
      (x0-weff+P0enc):(y0+weff) (x0+leff-weff):(y0+leff+weff-P0enc)
      (x0+leff-POenc):(y0+leff)))
  else
    for( i 1 na1
      leCreateContact( cv layer (x0+leff-weff):(y0+leff+weff-wvia1
        -2*border1-(wvia1+dist1)*(i-1)) "R0" wvia1 wvia1 1 (2*i-1)
        (wvia1+dist1) (wvia1+dist1) "center" "center" )
      leCreateContact( cv layer x0:(y0+wvia1+2*border1+(wvia1
```

```
+dist1)*(i-1)) "R0" wvia1 wvia1 1 (2*i-1) (wvia1+dist1)
          (wvia1+dist1) "center" "center" )
      ); for( i 1 na1 )
      for( i 1 nb1
        leCreateContact( cv layer (x0+leff-weff-(i-1)*(wvia1+dist1)):(y0
          +leff-dist1+2*border1-0.5*wvia1-(i-1)*(wvia1+dist1)) "R0" wvia1
          wvia1 1 (2*na1-1) (wvia1+dist1) (wvia1+dist1) "center" "center" )
      ); for( i 1 nb1 )
    ) ; multiple if
  ); foreach( layer vialist )
  ) ; let
) ; procedure( rectS45 )
; User interface
pcDefinePCell(
list( ddGetObj( "LV_PCells" ) "rectStack45" "layout" )
( (Width float 11.99)
  (Length float 30.0)
  (Grid float 0.005)
  (X0 float 0.0)
  (Y0 float 0.0)
  (metal1 boolean "TRUE")
  (metal2 boolean "TRUE")
  (metal3 boolean "FALSE")
  (metal4 boolean "FALSE")
  (metal5 boolean "FALSE")
  (metal6 boolean "FALSE")
  (alucap boolean "FALSE")
cv=pcCellView
layerlist = list( )
vialist = list( )
minlayer = 8
maxlayer = 0
when( metal1 == t
 minlayer = 1
  maxlayer = 1)
when( metal2 == t
 minlayer = min( minlayer 2 )
  maxlayer = max( maxlayer 2 ))
when( metal3 == t
 minlayer = min( minlayer 3 )
  maxlayer = max( maxlayer 3 ))
when( metal4 == t
```

```
minlayer = min( minlayer 4 )
 maxlayer = max( maxlayer 4 ))
when( metal5 == t
 minlayer = min( minlayer 5 )
 maxlayer = max( maxlayer 5 ))
when( metal6 == t
 minlayer = min( minlayer 6 )
 maxlayer = max( maxlayer 6 ))
when( alucap == t
 minlayer = min( minlayer 7 )
 maxlayer = max( maxlayer 7 ))
if( minlayer==1
then
 layerlist = cons( "metal1" layerlist)
) ; if( minlayer==1 )
if( minlayer==2
 layerlist = cons( "metal2" layerlist)
); if( minlayer==2)
if( minlayer==3
 layerlist = cons( "metal3" layerlist)
); if( minlayer==3)
if( minlayer==4
then
 layerlist = cons( "metal4" layerlist)
); if( minlayer==4)
if( minlayer==5
then
 layerlist = cons( "metal5" layerlist)
); if( minlayer==5)
if( minlayer==6
then
 layerlist = cons( "metal6" layerlist)
); if( minlayer==6)
if( minlayer==7
then
 layerlist = cons( "alucap" layerlist)
) ; if( minlayer==7 )
for( i (minlayer+1) maxlayer
 if( i==2
 then
   layerlist = cons( "metal2" layerlist)
   vialist = cons( "M2_M1" vialist)
 ) ; if( i==2 )
 if( i==3
    layerlist = cons( "metal3" layerlist)
    vialist = cons( "M3_M2" vialist)
```

```
); if( i==3)
 if( i==4
  then
   layerlist = cons( "metal4" layerlist)
   vialist = cons( "M4_M3" vialist)
 ); if( i==4 )
 if( i==5
 then
    layerlist = cons( "metal5" layerlist)
   vialist = cons( "M5_M4" vialist)
  ); if( i==5)
 if( i==6
   layerlist = cons( "metal6" layerlist)
   vialist = cons( "M6_M5" vialist)
 ); if( i==6)
 if( i==7
 then
    layerlist = cons( "alucap" layerlist)
   vialist = cons( "padopen" vialist)
 ); if( i==7)
); for( i (minlayer+1) maxlayer )
layerlist = reverse(layerlist)
vialist = reverse(vialist)
rectS45( cv Width Length Grid XO YO layerlist vialist)
) ; pcDefinePCell
```

Procedures to determine toroid's dimensions

```
+k)*phi))/(R*sin(k*phi)-(l+R)*sin((k+0.5)*phi)))**2))
    if( ((dist<minDist)&&(dist>0))
     minDist=dist
    ); if((dist<minDist))
   k=k+1
  ); while(((R*sin((1+k)*phi)-R*sin(k*phi))>w))
 minDist
procedure( Rmin( w s n step )
let( (phi,R1,dist1,R2,dist2,distm)
 phi=2*Pi/n
 R1=0.5*(w+s)/sin(0.5*phi)
  dist1=dmin( R1 phi w s l )
 R2=2*R1
  dist2=dmin( R2 phi w s l )
  while( (dist2<s)
   R2=R2+R1
   dist2=dmin( R2 phi w s l )
  ); while((dist2<s))
  distm=dmin((0.5*(R1+R2))) phi w s 1)
  while( ((R2-R1)>2*step)
    if( distm>s
    then
     R2=0.5*(R1+R2)
     dist2=distm
    else
     R1=0.5*(R1+R2)
     dist1=distm
    ) ; if( distm>s )
    distm=dmin((0.5*(R1+R2))) phi w s 1)
  ); while(((R2-R1)>2*step))
  grid( R2 step )
```

Procedures to draw the toroid

```
procedure( toro( cv n R l w step isInt mlayerlist mvialist )
let( ( i, phi,; cycle indexes
 x1, y1, x2, y2,; moving coordinates
 leff,; check for internal or external pins
 blayer, tlayer; layers
 minw = 4.4
 maxw = 11.99
 mins = 0.6
 w = max( minw grid( w 2*step ) )
 w = min(maxw w)
 1 = max( grid( 1 2*step ) (w+mins) )
 if( isInt
 then
   leff = 1
 else
   leff = -1
  ) ; if( isInt )
 R = max(grid(R step) Rmin(w mins n step))
 blayer = list()
 tlayer = list()
 blayer = cons( "pintext" blayer )
 blayer = cons( "metal1" blayer )
 tlayer = cons( "pintext" tlayer )
 tlayer = cons( "metal6" tlayer )
 phi=0
 x1=R+0.5*(1-leff)
 y1=0
 for(i 0 n-1
   x2 = grid((R+0.5*(l+leff))*cos(phi+Pi/n)) step)
   y2 = grid((R+0.5*(l+leff))*sin(phi+Pi/n)) step)
   if( (y1<y2)
    then
      if( (x1<x2)
      then
        dbCreatePolygon(cv list( "metal1" "pintext" ) list( (x1-0.5*w):
          (y1-0.5*w) (x1+0.5*w):(y1-0.5*w) (x2+0.5*w):(y2-0.5*w) (x2+0.5*w)
          +0.5*w):(y2+0.5*w) (x2-0.5*w):(y2+0.5*w) (x1-0.5*w):(y1+0.5*w)))
      else
       dbCreatePolygon(cv list( "metal1" "pintext" ) list( (x1-0.5*w):
          (y1-0.5*w) (x1+0.5*w):(y1-0.5*w) (x1+0.5*w):(y1+0.5*w) (x2+0.5*w)
          +0.5*w):(y2+0.5*w) (x2-0.5*w):(y2+0.5*w) (x2-0.5*w):(y2-0.5*w)))
      )
    else
      if((x1<x2)
```

```
then
        dbCreatePolygon(cv list( "metal1" "pintext" ) list( (x2-0.5*w):
          (y2-0.5*w) (x2+0.5*w):(y2-0.5*w) (x2+0.5*w):(y2+0.5*w) (x1
          +0.5*w):(y1+0.5*w) (x1-0.5*w):(y1+0.5*w) (x1-0.5*w):(y1-0.5*w)))
      else
        dbCreatePolygon(cv list( "metal1" "pintext" ) list( (x2-0.5*w):
          (y2-0.5*w) (x2+0.5*w):(y2-0.5*w) (x1+0.5*w):(y1-0.5*w) (x1
          +0.5*w):(y1+0.5*w) (x1-0.5*w):(y1+0.5*w) (x2-0.5*w):(y2+0.5*w)))
      )
    )
   rectS( cv w w step (x2-0.5*w) (y2-0.5*w) mlayerlist
      mvialist)
    phi = phi+2*Pi/n
    x1 = grid((R+0.5*(1-leff))*cos(phi)) step)
    y1 = grid((R+0.5*(1-leff))*sin(phi)) step)
    if( (y1<y2)
    then
     if( (x1<x2)
      then
        dbCreatePolygon( cv list("metal6" "pintext") list( (x1-0.5*w):
          (y1-0.5*w) (x1+0.5*w):(y1-0.5*w) (x2+0.5*w):(y2-0.5*w) (x2+0.5*w)
          +0.5*w):(y2+0.5*w) (x2-0.5*w):(y2+0.5*w) (x1-0.5*w):(y1+0.5*w)))
        dbCreatePolygon( cv list("metal6" "pintext") list( (x1-0.5*w):
          (y1-0.5*w) (x1+0.5*w):(y1-0.5*w) (x1+0.5*w):(y1+0.5*w) (x2-0.5*w)
          +0.5*w):(y2+0.5*w) (x2-0.5*w):(y2+0.5*w) (x2-0.5*w):(y2-0.5*w)))
      )
    else
      if( (x1<x2)
      then
        dbCreatePolygon( cv list("metal6" "pintext") list( (x2-0.5*w):
          (y2-0.5*w) (x2+0.5*w):(y2-0.5*w) (x2+0.5*w):(y2+0.5*w) (x1
          +0.5*w):(y1+0.5*w) (x1-0.5*w):(y1+0.5*w) (x1-0.5*w):(y1-0.5*w)))
      else
        dbCreatePolygon( cv list("metal6" "pintext") list( (x2-0.5*w):
          (y2-0.5*w)(x2+0.5*w):(y2-0.5*w)(x1+0.5*w):(y1-0.5*w)(x1
          +0.5*w):(y1+0.5*w) (x1-0.5*w):(y1+0.5*w) (x2-0.5*w):(y2+0.5*w)))
      )
    )
   rectS( cv w w step (x1-0.5*w) (y1-0.5*w) mlayerlist
     mvialist)
 ); for(i 0 n-1)
  ) ; let
); toro
; User interface
pcDefinePCell(
```

```
list( ddGetObj( "LV_PCells" ) "toro" "layout" )
( (Width float 4.4)
  (Length float 6.4)
  (N_coils int 44)
  (Radius float 58.84)
  (Grid float 0.005)
  (Internal boolean "TRUE")
  (metal1 boolean "TRUE")
  (metal2 boolean "TRUE")
  (metal3 boolean "FALSE")
  (metal4 boolean "FALSE")
  (metal5 boolean "FALSE")
cv=pcCellView
blayer
mlayerlist = list( "metal1" "metal2" "metal3" "metal4"
  "metal5" )
mvialist = list( "M2_M1" "M3_M2" "M4_M3" "M5_M4" "M6_M5" )
minlayer = 6
maxlayer = 0
when( metal1 == t
 minlayer = 1
 maxlayer = 1)
when( metal2 == t
 minlayer = min( minlayer 2 )
 maxlayer = max( maxlayer 2 ))
when( metal3 == t
 minlayer = min( minlayer 3 )
 maxlayer = max( maxlayer 3 ))
when( metal4 == t
 minlayer = min( minlayer 4 )
 maxlayer = max( maxlayer 4 ))
when( metal5 == t
 minlayer = min( minlayer 5 )
 maxlayer = max( maxlayer 5 ))
if( minlayer==1
 mlayerlist = remd( "metal1" mlayerlist)
); if( minlayer==1)
if( minlayer==2
then
 mlayerlist = remd( "metal2" mlayerlist)
); if( minlayer==2 )
if( minlayer==3
then
 mlayerlist = remd( "metal3" mlayerlist)
); if( minlayer==3)
```

```
if( minlayer==4
then
 mlayerlist = remd( "metal4" mlayerlist)
); if( minlayer==4)
if( minlayer==5
then
 mlayerlist = remd( "metal5" mlayerlist)
); if( minlayer==5)
for( i (minlayer+1) maxlayer
 if( i==2
 then
   mlayerlist = remd( "metal2" mlayerlist)
   mvialist = remd( "M2_M1" mvialist)
 ) ; if( i==2 )
 if( i==3
 then
   mlayerlist = remd( "metal3" mlayerlist)
   mvialist = remd( "M3_M2" mvialist)
 ); if( i==3)
 if( i==4
 then
   mlayerlist = remd( "metal4" mlayerlist)
   mvialist = remd( "M4_M3" mvialist)
 ) ; if( i==4 )
 if( i==5
 then
   mlayerlist = remd( "metal5" mlayerlist)
   mvialist = remd( "M5_M4" mvialist)
 ); if( i==5)
); for( i (minlayer+1) maxlayer )
mlayerlist = reverse(mlayerlist)
mvialist = reverse(mvialist)
toro( cv N_coils Radius Length Width Grid
 Internal mlayerlist mvialist)
) ; pcDefinePCell
```

APPENDIX E

Publications

The following papers resulted from the different research projects related to this study.

Published

- Andreani, P., Wang, X., Vandi, L. and Fard, A., "A Study of Phase Noise in Colpitts and LC-Tank CMOS Oscillators", *IEEE Journal of Solid-State Circuits*, vol. 40, no. 6, 2005. [94]
- Vandi, L., Andreani, P., Temporiti, E., Sacchi, E., Bietti, I., Ghezzi,
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- Liscidini, A., Mazzanti, A., Tonietto, R., Vandi, L, Andreani, P. and Castello, R., "A 5.4mW GPS CMOS Quadrature Front-End Based on a Single-Stage LNA-Mixer-VCO", *Proceedings of IEEE International Solid-State Circuits Conference*, 2006. [1]

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Accepted for publication

Vandi, L., Andreani, P., Temporiti, E., Sacchi, E., Bietti, I., Ghezzi,
 C. and Castello, R., "A toroidal inductor integrated in a standard
 CMOS process", Accepted for publication on Springer Analog Integrated Circuits and Signal Processing International Journal. [95]

Liscidini, A., Mazzanti, A., Tonietto, R., Vandi, L., Andreani, P. and Castello, R., "Single-Stage Low Power Quadrature RF Receiver Front-End: The LMV cell", Accepted for publication on *IEEE Journal of Solid-State Circuits*. [66]

Submitted for publication

Vandi, L., Andreani, P., Tired, T. and Mattisson, S., "A novel approach to negative feedback in RX front-ends", Submitted for publication to IEEE 24th NORCHIP Conference, 2006. [96]

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