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Tanner, D. B.; Cummings, K. D.; Jacobsen, Claus Schelde

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Far-Infrared Study of the Charge Density Wave in Tetrathiafulvalene Tetracyanoquinodimethane (TTF-TCNQ)

D. B. Tanner and K. D. Cummings

Department of Physics, The Ohio State University, Columbus, Ohio 43210

and

C. S. Jacobsen

Physics Laboratory 3, Technical University of Denmark, DK-2800 Lyngby, Denmark

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Detailed far-infrared measurements at temperatures from 25 to 300 K provide strong support for a charge-density-wave mechanism for the dc conductivity and microwave dielectric constant of tetrathiafulvalene tetracyanoquinodimethane (TTF-TCNQ). At low temperatures the charge-density wave is pinned at 40 cm^{-1} while at higher temperatures it appears at zero frequency. Values are obtained for the effective mass ($20m^*$) and lifetime ($1.6 \times 10^{-12} \text{ sec}$ at 60 K) of the charge-density wave.

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Despite considerable experimental and theoretical efforts over the past eight years, the properties of tetrathiafulvalene tetracyanoquinodimethane (TTF-TCNQ) remain controversial. The principal point of dispute is between single-particle models and models incorporating collective effects. The single-particle models explain the temperature dependence of the conductivity by scattering of the conduction electrons from phonons^{1,2} or other electrons.³ The collective picture assumes that the important $2k_F$ phonons are in equilibrium with the electrons rather than with the lattice. This assumption, an extreme case of phonon drag, leads to a sliding charge-density-wave (CDW) picture for the conductivity.⁴⁻⁷

The importance of collective effects is supported by diffuse x-ray⁸ and neutron⁹ scattering measurements [which observe a three-dimensionally (3D) ordered CDW below 38 K and one-dimensional (1D) fluctuations up to 150 K] and by conductivity measurements at high pressure¹⁰ (which show a drop in conductivity when the CDW becomes commensurate to the lattice). The CDW mechanism is, however, certainly not universally accepted at the present time. In this Letter we report the results of far-infrared reflection measurements on TTF-TCNQ at temperatures between 25 and 300 K. Our data indicate that near 60 K the dc conductivity of this material is dominated by a low-frequency mode which we interpret as a sliding charge-density wave. At temperatures below 38 K the CDW becomes pinned at relatively high frequencies ($\sim 40 \text{ cm}^{-1}$) where it is the major contributor to the microwave dielectric constant.

Far infrared (FIR) measurements at low temperatures are central to the single-particle ver-

sus collective-mode issue since the latter model makes rather definite predictions about the far-infrared properties.¹¹ At low temperatures, the CDW is pinned by 3D interactions or impurities and is observed as an FIR absorption peak.⁵ At higher temperatures, the CDW slides through the crystal, contributing a Drude-like term to the frequency-dependent conductivity,

$$\sigma_1(\omega) = \sigma_{1sp} + \Omega_p^2 \tau_c / 4\pi (1 + \omega^2 \tau_c^2), \quad (1)$$

where σ_{1sp} is the single-particle contribution to the conductivity, τ_c is the collective-mode lifetime, and $\Omega_p = (4\pi n e^2 / M^*)^{1/2}$ with n the carrier density (0.59 per molecule for TTF-TCNQ) and M^* the CDW effective mass. Because the lifetime of the collective mode is very long, its contribution to the conductivity is only significant at very low frequencies, $\omega < 1/\tau_c$.

Polarized FIR reflectance measurements were made over 7 to 700 cm^{-1} using a Michelson interferometer in conjunction with a germanium bolometer detector operating at 1 K. The exit aperture of the interferometer was imaged onto the sample by a large (30-cm-diam and 18-cm-focal length) mirror and the reflected light was focused by an identical mirror into the entrance of the light pipe leading to the detector. The polarizer was a gold wire grid on polyethylene.

The sample was a mosaic of nine TTF-TCNQ crystals, each ~ 1.2 by 0.15 cm in size, mounted in an Air Products Helitran cryostat. One crystal had four electrical connections to its back surface to enable resistance measurements; the sample temperature was determined by this resistance. This calibration was necessary because room-temperature radiation heated the sample

and prevented cooling below 25 K. The mosaic intercepted more than 90% of the far-infrared beam and gave a signal-to-noise ratio in the reflectance exceeding 100. The reflectance at a given temperature, $R(T)$, was determined from four measurements: $R(T) = [S_1(T)/S_1(300)][S_2(300)/S_2(\text{Au})]$. Here $S_1(T)$ is the signal from the mosaic at temperature T ; $S_1(300)$ is the signal at room temperature; $S_2(300)$ is the room-temperature signal from a mosaic; and $S_2(\text{Au})$ is the signal from the same mosaic after coating with gold.

Figure 1 shows the \vec{b} -axis polarized reflectance, $R(\omega)$, of TTF-TCNQ at five temperatures between 25 and 300 K. At room temperature,¹² $R(\omega) \rightarrow 1$ at $\omega \rightarrow 0$. The slope of $R(\omega)$ is negative at low frequencies, as expected for a good conductor. At higher frequencies $R(\omega) \sim 0.75$. At 160 and 60 K, $R(\omega)$ is qualitatively similar; its magnitude increases as the temperature is lowered while the negative slope at low frequencies [with $R(\omega) \rightarrow 1$ at $\omega \rightarrow 0$] remains. At 25 K the reflectance is profoundly different: It has positive slope at low frequencies and a maximum around 40–80 cm^{-1} followed by deep minimum at 120 cm^{-1} . The 34-K reflectance is intermediate between the 25-K and

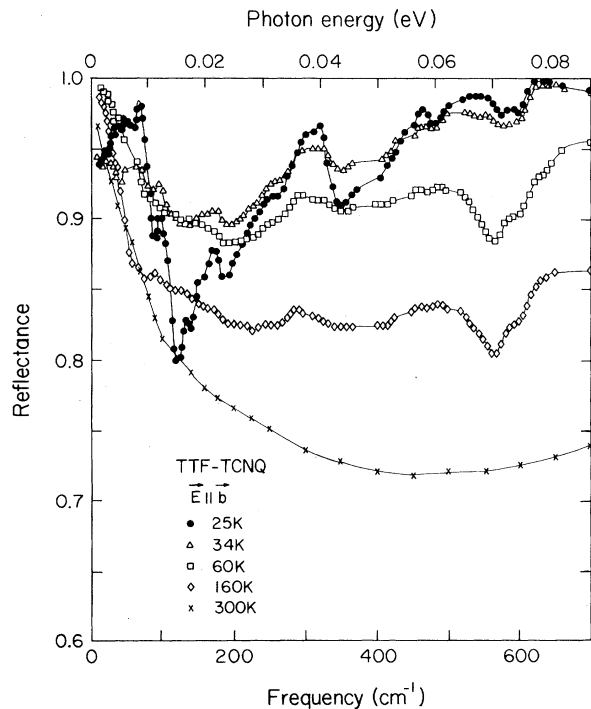


FIG. 1. Reflectance of TTF-TCNQ for electric field along the \vec{b} -axis. Note the suppressed zero. Lines are drawn as a guide to the eye.

the 60-K data.

We have performed a Kramers-Kronig analysis of the reflectance data.¹³ At frequencies about 700 cm^{-1} the reflectance was extrapolated smoothly from the measured data, with a transition¹² at 1600 cm^{-1} to values calculated from the near-infrared reflectance data of Bright, Garito, and Heeger.¹⁴ The reflectance at low frequencies was smoothly extrapolated using the Hagen-Reubens relation in the conducting regime and was assumed constant at lower temperatures.¹³

Figure 2 shows the \vec{b} -axis frequency-dependent conductivity of TTF-TCNQ. The low-frequency conductivity is in complete agreement with the dc conductivity of these "ordinary" crystals.¹⁵ It rises from 500 $\Omega^{-1} \text{cm}^{-1}$ at room temperature to 5000 $\Omega^{-1} \text{cm}^{-1}$ at 60 K and then falls to around 200 $\Omega^{-1} \text{cm}^{-1}$ at 34 K and is small at 25 K. We stress that the low-frequency conductivity is not affected in any significant way by the choice of extrapolation procedure. Assuming constant reflectance below our lowest data point reduced the value of the 60-K conductivity at that point by 20% or so; the narrow structure at low frequency, however, still remained. At 25 K the conductivity has a strong peak at 40 cm^{-1} . This feature is reduced in strength and broadened in the 34-K data and is not present at 60 K and above. There is other structure at higher frequencies, most notably an intense peak at 290 cm^{-1} .

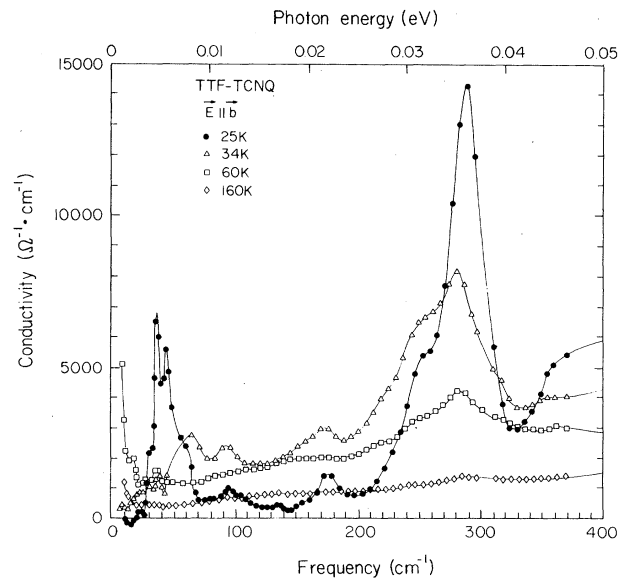


FIG. 2. Frequency-dependent \vec{b} -axis conductivity of TTF-TCNQ at four temperatures: 25, 34, 60, and 160 K.

Figure 3 shows the real part of the \vec{b} -axis dielectric function of TTF-TCNQ. The low-frequency FIR data are in excellent numerical agreement with microwave dielectric constant measurements. We find $\epsilon_1(0) = 3200$ at 25 K, within 10% of the values of Gunning *et al.*¹⁶ at these temperatures. At 34 K, the dielectric constant has increased to 6000 while it is negative and very large in magnitude in the highly conducting state.

The 40-cm⁻¹ feature in the 25-K conductivity has a large oscillator strength. We identify this feature as the phase mode of the pinned CDW in TTF-TCNQ. The effective mass M^* of the pinned charge-density wave can be estimated from the oscillator-strength sum rule,

$$\int_0^{80} \sigma_1(\omega) d\omega = \frac{1}{2}\pi (ne^2/M^*) = \frac{1}{8}\Omega_p^2. \quad (2)$$

The upper limit of the integral is taken as 80 cm⁻¹, where the high-frequency wing of the peak has dropped into the baseline. We obtain $M^* = 60m_e = 20m^*$ (m_e is the free-electron mass and m^* is the ordinary band mass).

While small, this effective mass is consistent with microwave dielectric constant data. According to Lee, Rice, and Anderson there are three contributions to the static dielectric constant in a

CDW system,

$$\epsilon_1 = \epsilon_\infty + \frac{2}{3}\omega_p^2/(2\Delta)^2 + \Omega_p^2/\omega_T^2, \quad (3)$$

where ϵ_∞ is the high-frequency dielectric constant, ω_p is the conduction-electron plasma frequency, 2Δ is the Peierls electronic energy gap, and ω_T is the CDW pinning frequency. Using the above estimate of the CDW effective mass, we obtain $\Omega_p = 2000$ cm⁻¹; our conductivity data indicate that $\omega_T = 40$ cm⁻¹. The CDW contribution to the dielectric constant is then 2700. At frequencies well above the pinning frequency only the first two terms on the right-hand side of Eq. (3) contribute to ϵ_1 . At 25 K, $\epsilon_1 \sim 600$ for 100 cm⁻¹ < ω < 200 cm⁻¹, a reasonable value for the single-particle contribution to ϵ_1 . The total dielectric constant is then predicted to be 3300; experimentally¹⁶ it is between 3000 and 3500.

At 60 K, the CDW is at zero frequency where its finite dc conductivity drives $\epsilon_1(\omega)$ negative at low frequencies. [Note that $\epsilon_1(\omega) > 0$ above 100 cm⁻¹.] The zero crossing occurs at a screened CDW plasma frequency $\Omega_p/\sqrt{\epsilon_H}$ where ϵ_H is the value of $\epsilon_1(\omega)$ above the zero. From the data in Fig. 3 we obtain $\Omega_p = 1000$ cm⁻¹ at 60 K and $\Omega_p = 500$ cm⁻¹ at 160 K. The CDW appears to lose oscillator strength with increasing temperature.

With Ω_p known, the lifetime of the collective mode may be estimated. We use Eq. (1) [but put $\sigma_{1sp}(0) = 0$] and take 5000 Ω^{-1} cm⁻¹ for the 60-K conductivity. We then obtain $\tau_c = 1.6 \times 10^{-12}$ sec ($1/\tau_c = 4$ cm⁻¹) in agreement with the width of the narrow low-frequency peak in $\sigma_1(\omega)$ at 60 K in Fig. 2. In contrast the single-particle lifetime¹⁴ is $\tau_{sp} \sim 6 \times 10^{-15}$ sec.

The 25-K results for the CDW mass and pinning frequency are significantly different from those of earlier low-temperature FIR studies of TTF-TCNQ single crystals by Coleman *et al.*¹⁷ and by Eldridge and Bates.¹⁸ (We note that the data in Fig. 2 resemble early results on sublimed thin films of TTF-TCNQ, although the magnitude of the conductivity was much lower and the frequency of the structure identified as the pinned mode was higher.¹⁹) The polarized reflectance data of Coleman *et al.*¹⁷ did not extend above 100 cm⁻¹ and so did not include the 125-cm⁻¹ reflectance minimum seen in our 25-K data. The magnitude of the reflectance was also lower than ours (from depolarization in the light pipes that they used). In their fit to a model CDW dielectric function they assumed the reflectance minimum to be below rather than above their frequency range and therefore obtained a low value for the pinning frequen-

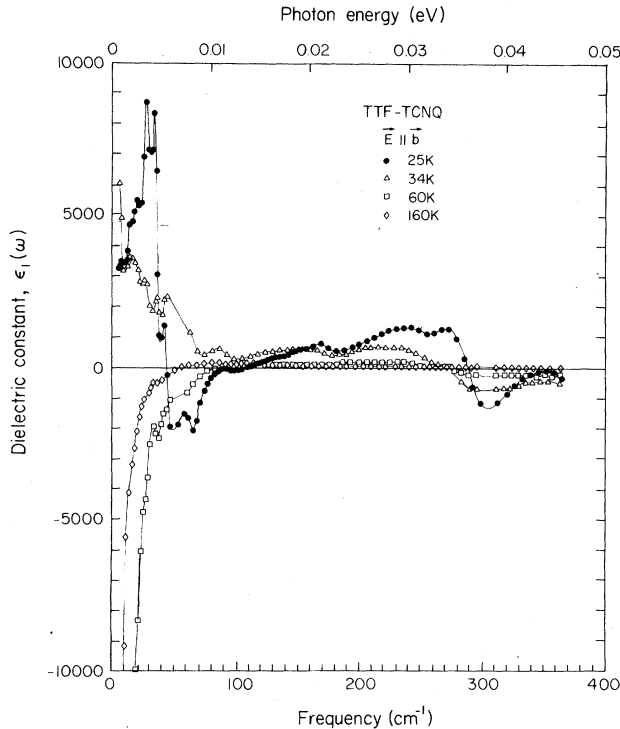


FIG. 3. Real part of the \vec{b} -axis dielectric function of TTF-TCNQ.

cy and, consequently, too large a CDW effective mass.

Eldridge and Bates¹⁸ used a TTF-TCNQ crystal as a far-infrared detector. If $\alpha d \ll 1$ (α is the absorption coefficient; d the thickness) their signal would be proportional to α ; otherwise it would saturate to $1-R$ (R is the reflectance). Our Kramers-Kronig analysis indicates that $\alpha > 10^4 \text{ cm}^{-1}$ at 25 K for frequencies above 30 cm^{-1} so that their samples, with $d > 3 \times 10^{-3} \text{ cm}$, were in the limit $\alpha d \gg 1$; hence the signal was saturated at $1-R$. Their data closely resemble our results for $1-R$ at frequencies below 300 cm^{-1} .

In summary, our data support the picture of TTF-TCNQ that attributes its unusual dc conductivity and microwave dielectric constant to collective effects associated with the charge-density wave. TTF-TCNQ is neither a metal above 58 K nor an ordinary semiconductor below 38 K.

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