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# The Flow Field Upstream Of A Horizontal Axis Wind Turbine

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THE FLOW FIELD UPSTREAM OF  
A HORIZONTAL AXIS WIND TURBINE

Technical Report

by

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## ABSTRACT

A mathematical model is developed for a steady-state axi-symmetric upstream flow of a porous disc, in a uniform flow field. The special case of the upstream flow of a windmill, with and without a nacelle, is treated. First, the windmill is considered as a uniform distribution of sources and then as a linear distribution of sources. Solutions for the blade disc of the wind field upstream are obtained in the form of streamlines and velocity vector components.

Sample flow patterns upstream of the blade disc of the UMass 25 kW wind turbine are presented for several power levels. Documented computer programs applicable to any wind turbine are appended.

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## INTRODUCTION

The problem of the exact solution of the flow through a porous disc has been of interest since the original works of G.I. Taylor [1].

This solution can have two major applications:

- a) Flow through screens (power absorbing devices)
- b) Flow through power developing devices.

The first point of view has to do with porous bodies which have been of practical importance in the past. Most of these bodies, which are like parachutes, fish nets, wind breaks, slotted injection discs in oil combustion chambers, and fire developments in forest, can be modeled by a screen.

The second point of view considered in this work deals with the wind power developing machines. These devices can be modeled by a very porous disc in the wind field.

The first comprehensive analysis of flow through screens was carried out by Taylor and Batchelor (1949), [1], [2]. They were interested in the effect of the screen on the turbulence, and therefore their study was oriented towards the non-uniform channel flow, passing through a flat screen. In 1959, Elder considered the more general case of an irregular-shaped and non-uniform screen in a two dimensional channel flow [3].

The problem of a finite plane screen in an infinite flow field was first considered by Küchemann and Weber (1953) [4]. Later, in 1963, Taylor considered the problem in a two-dimensional case [5]. This was done by replacing the screen with uniformly-distributed sources.

Finally, Koo and James considered the more general case of the two-dimensional flow around a submerged screen [2].

Following the idea of modeling any wind machine by the combination of sources, sinks, or vortices, it was proposed that a windmill can be modeled by a porous screen. This work was oriented towards solving the problem of a three-dimensional porous disc in a steady-state, axisymmetric, uniform flow. The screen was modeled by a distribution of sources and the problem was divided into two cases. First, the simple case of modeling the windmill by a uniformly-distributed disc of sources. Second, a more realistic model was considered. Taking into consideration the fact that the development of power is higher in the outer region of the windmill blades, the blade disc was modeled by a linearly-distributed disc of sources. The effect of a nacelle and its relative orientation to the blade disc was studied in both cases.

The velocity field and the streamlines were constructed for some numerical examples in association with the 25 kW windmill at the University of Massachusetts Solar Habitat I.

## THEORETICAL ANALYSIS

### Introduction to the Modeling Idea

The well-known idea of modeling flow fields through a combination of individual factors is used in order to model a porous disc as a distribution of sources in space within a uniform flow field.

J.K. Koo and D.F. James [2] developed G.I. Taylor's [5] idea of modeling a two-dimensional screen by a distribution of sources, by modeling the screen in a duct. This work is oriented to find a general solution for a three-dimensional screen, using a distribution of sources on a disc, in an axi-symmetric uniform flow.

The development of the basic equations is based on the following procedure: first, the potential of a source located at an arbitrary point in space is determined, second, based on this potential, the cases of a disc with a uniform or linear distribution of sources are considered and the potential on the axis of the disc founded, and third, based on the harmonic and more particularly the symmetric properties of the potential function and the solution on the disc's axis by the use of zonal harmonies, the general solution of the function is constructed.

The model is completed by the superposition of a uniform flow on the potential of the disc.

In the case of modeling a windmill, the effect of the nacelle on the disc's axis can be modeled by a single source, which can result in different body shapes.

The solution is in the form of an infinite series of the Legendre and Associated Legendre polynomials. The velocity field and streamlines

are constructed by a computer program, and will be described later in this report.

### Potential of a Source at an Arbitrary Point

The potential of a point source at the origin can be written as [6]:

$$\Phi = \frac{k}{4\pi} \frac{1}{|\vec{r}|}$$

where  $k$  is the source strength.

The potential at a point  $P$  of a point source at  $S$  is: (Fig. 1)

$$\Phi = \frac{k}{4\pi} \frac{1}{|\vec{r}'|} = \frac{k}{4\pi} \frac{1}{|\vec{r} - \vec{s}|}$$

knowing

$$\vec{r} = \vec{r}(x, y, z)$$

$$\vec{s} = \vec{s}(\xi, \eta, \zeta)$$

the potential is:

$$\Phi = \frac{k}{4\pi} \frac{1}{[(x-\xi)^2 + (y-\eta)^2 + (z-\zeta)^2]^{1/2}} \quad (1)$$

By using a coordinate transformation as shown in Fig. 2, the potential can be transformed to spherical coordinate as:

$$\begin{aligned} \xi &= \int \cos \beta & x &= R \cos \theta \\ \eta &= \int \sin \beta \cos \alpha & \text{and } y &= R \sin \theta \cos \alpha \\ \zeta &= \int \sin \beta \sin \alpha & z &= R \sin \theta \sin \alpha \end{aligned} \quad (1A)$$

then,



$$\Phi = \frac{k}{4\pi} \frac{1}{\left[ (R \cos \theta - \rho \cos \beta)^2 + (R \sin \theta \cos \alpha - \rho \sin \beta \cos \delta)^2 + (R \sin \theta \sin \alpha - \rho \sin \beta \sin \delta)^2 \right]^{1/2}} \quad (2A)$$

If the source is in  $y - z$  plane, where  $\beta = \pi/2$ , then

$$\Phi = \frac{k}{4\pi} \frac{1}{\left[ R^2 \cos^2 \theta + (R \sin \theta \cos \alpha - \rho \cos \delta)^2 + (R \sin \theta \sin \alpha - \rho \sin \delta)^2 \right]^{1/2}} \quad (2)$$

### Model for the Body of a Windmill

The nacelle of a windmill can be approximated as a paraboloid of revolution. This is modeled by a point source in a uniform flow [7]. (See Fig. 3).

Taking  $r_1$  and  $\eta$  from the geometry of the nacelle, and assuming a  $v_0$ , the source strength  $k$  and its position  $r_2$  can be found by:

$$r_0 = \sqrt{\frac{k}{4\pi U_0}} \quad \text{and} \quad \eta = \sqrt{\frac{k}{\pi U_0}}$$

Therefore, the potential for the nacelle in a uniform flow can be written as:

$$\Phi = \frac{k}{4\pi} \frac{1}{|\vec{r}_0|} + U_0 r \cos \theta \quad (3)$$

### Potential of the Distributed Disc of Sources

Introduction. The total potential  $\Phi$  of disc's sources is the solution to the Laplace Equation ( $\nabla^2 \Phi = 0$ ), with the appropriate boundary conditions.

Since the flow is axi-symmetric, that is, independent of  $\alpha$ , (See

Fig. 2) the solution of  $\Delta^2 \Phi = 0$  can be written as [8]:

$$\Phi = \sum_{n=0}^{\infty} \left( A_n r^n + \frac{B_n}{r^{n+1}} \right) P_n(\cos \theta) \quad (4)$$

where:  $P_n(x)$  = Legendre Polynomial of the first kind and  $n$ th order.

Knowing that on the  $x$ -axis  $\theta$  is zero and that  $P_n(0) = 1$  [8]; the solution on the  $x$ -axis can be written as:

$$\Phi_x = \sum_{n=0}^{\infty} \left( A_n r^n + \frac{B_n}{r^{n+1}} \right) \quad (5)$$

Hence, to determine  $A_n$  and  $B_n$ , the solution on the  $x$ -axis should be found, expanded in a power series of  $r$ , and then equated to equation (5). Consequently, the first step is to find the disc's source potential on the  $x$ -axis.

Consider a disc of distributed sources with the radius  $A$ , and source strength per unit area  $k$  as shown in Fig. 4.

The element of area  $dA$  is equal to  $\rho d\gamma d\rho$  for:  $0 < \rho \leq A$ , and  $0 \leq \gamma \leq 2\pi$ . Using equation (2) the differential potential at any field point  $(r, \theta)$  can be written in terms of the differential source distribution, as:

$$d\Phi = \frac{k'}{4\pi} \frac{1}{\left[ R^2 \cos^2 \theta + (R \sin \theta \cos \alpha - \rho \cos \alpha)^2 + (R \sin \theta \sin \alpha - \rho \sin \alpha)^2 \right]^{1/2}} \quad (6)$$

To get the solution on the  $x$ -axis, equation (6) can be restricted to the  $x$ -axis, that is  $\theta = 0$ . The differential potential on the  $x$ -axis is:

$$d\Phi_x = \frac{k'}{4\pi} \frac{1}{\left[ R^2 + \rho^2 \cos^2 \alpha + \rho^2 \sin^2 \alpha \right]^{1/2}}$$

and,

$$d\Phi_x = \frac{k'}{4\pi} \frac{1}{[R^2 + \rho^2]^{1/2}} \quad (7)$$

Integration of equation (7) depends on  $k$ . The problem can be divided into two cases: a)  $k$  is constant or a uniformly-distributed source disc and b)  $k$  is linear in  $\rho$  or a linearly-distributed source disc.

Uniformly distributed disc of sources.

( $k = \text{const}$ )

If  $k$  is constant, equation (7) can be integrated as:

$$\Phi_x = \frac{k'}{2} [(R^2 + A^2)^{1/2} - R] \quad (8)$$

Using a binomial expansion it can be shown that, for  $R$  smaller than  $A$ ,

$$\left(1 + \left(\frac{R}{A}\right)^2\right)^{1/2} = \sum_{n=0}^{\infty} [P_n(0) + P_{n-2}(0)] \left(\frac{R}{A}\right)^n \quad (9)$$

and for  $R$  greater than  $A$ :

$$\left(1 + \left(\frac{A}{R}\right)^2\right)^{1/2} = \sum_{n=0}^{\infty} [P_n(0) + P_{n+2}(0)] \left(\frac{A}{R}\right)^{n+2} \quad (10)$$

Therefore  $\Phi_x$  in equation (8) can be expanded as [9]:

$$\Phi_x = \frac{k'A}{2} \sum_{n=0}^{\infty} [P_n(0) + P_{n-2}(0)] \left(\frac{R}{A}\right)^n - \frac{k'A}{2} \left(\frac{R}{A}\right) \quad R < A \quad (11)$$

and

$$\Phi_x = \frac{k'A}{2} \sum_{n=0}^{\infty} [P_n(0) + P_{n-2}(0)] \left(\frac{A}{R}\right)^{n+1} \quad R > A \quad (12)$$

The equating of equations (11) and (12) to equation (5) allows  $A_n$  and  $B_n$  to be determined for the general solution  $\phi$ . This leads to:

$$A_n = \frac{k'A}{2} [P_n(o) + P_{n-2}(o)] A^{-n} \quad (*)$$

$$A_1 = \frac{k'A}{2}$$

$R < A$

$$B_n = 0$$

and,

$$A_n = 0$$

$$B_n = \frac{k'A}{2} [P_n(o) + P_{n+2}(o)] A^{(n+1)}$$

$R > A$

Therefore, the general solution becomes:

$$\Phi = \frac{k'A}{2} \sum_{n=0}^{\infty} [P_n(o) + P_{n-2}(o)] \left(\frac{R}{A}\right)^n P_n(\cos\theta) - \frac{k'A}{2} \left(\frac{R}{A}\right) P_1(\cos\theta) \quad (13)$$

when  $R < A$ .

and for  $R > A$ .

$$\Phi = \frac{k'A}{2} \sum_{n=0}^{\infty} [P_n(o) + P_{n+2}(o)] \left(\frac{A}{R}\right)^{n+1} P_n(\cos\theta) \quad (14)$$

Linearly distributed disc of sources. For a linearly distributed disc of sources,  $k = m_p$  where  $m$  is a constant. Hence equation (7) can be written as

$$d\Phi_s = \frac{m}{4\pi} \frac{\rho^2 d\Omega d\chi}{[R^2 + \rho^2]^{1/2}}$$

---

\* $P_n(x) = 0$ , if  $n < 0$

and integration of this equation will have the following form

$$\Phi_x = \frac{m}{4\pi} \int_0^{2\pi} \int_0^A \frac{\rho^2 d\rho}{[R^2 + \rho^2]^{1/2}} d\varphi$$

which can be reduced to:

$$\Phi_x = \frac{m}{2} \int_0^{2A} \frac{\rho^2}{[R^2 + \rho^2]^{1/2}} d\rho \quad (15)$$

In order to expand  $\Phi_x$  in a power series of  $R$ , the space is divided into two regions: (See Fig. 5)

a)  $R < A$

b)  $R \geq A$

In region (a), each point  $Q$  (which is inside a sphere of radius  $A$ ) is affected by two kinds of sources: First, those whose distance from the origin ( $\rho$ ) is less than  $R$  (Radius of Point  $Q$ ), that is  $\rho < R$ , and second, those with  $\rho > R$ .

Therefore, for  $R < A$ , the potential on the  $x$ -axis can be written as:

$$\Phi_x = \Phi_1 + \Phi_2$$

where  $\Phi_1$  is the potential of the sources with  $\rho < R$  ( $\rho$  changes from zero to  $R$ ) and  $\Phi_2$  is the potential of the sources with  $\rho > R$  ( $\rho$  changes from  $R$  to  $A$ ).

From equation (15),  $\Phi_1$  may be written as:

$$\begin{aligned}\Phi_1 &= \frac{m}{2} \int_0^R \frac{\rho^2}{R^2} [1 + \rho^2/R^2]^{-1/2} d\rho \\ &= \frac{mR^2}{2} \int_0^1 t^2 (1+t^2)^{-1/2} dt\end{aligned}\quad (15A)$$

where  $t = \rho/R$ .

A binomial expression can be written as:

$$\begin{aligned}(t^2 + 1)^{-1/2} &= 1 + \sum_{n=1}^{\infty} \frac{-1/2(-1/2-1)(-1/2-2)\dots(-1/2-n+1)}{n!} t^{2n} \\ &= 1 + \sum_{n=1}^{\infty} a_n t^{2n} \quad \text{for } t^2 \leq 1\end{aligned}\quad (16)$$

where:

$$a_n = \frac{-1/2(-1/2-1)(-1/2-2)\dots(-1/2-n+1)}{n!}$$

Using equation (16) it is obvious that:

$$t^2(1+t^2)^{-1/2} = t^2 + \sum_{n=1}^{\infty} a_n t^{2n+2}$$

Using this expansion in equation (15A) results in:

$$\Phi_1 = \frac{mR^2}{2} \int_0^1 t^2 dt + \frac{mR^2}{2} \sum_{n=1}^{\infty} a_n \int_0^1 t^{2n+2} dt$$

therefore,

$$\Phi_1 = \frac{mR^2}{2} \left[ \frac{1}{3} + \sum_{n=1}^{\infty} \frac{a_n}{2n+3} \right] \quad (17)$$

Now solving for  $\phi_2$ , equation (15) can be written as:

$$\begin{aligned}
\Phi_2 &= \frac{m}{2} \int_R^A \rho^2 [R^2 + \rho^2]^{-1/2} d\rho \\
&= \frac{mR}{2} \int_R^A \frac{\rho}{R} [R^2/\rho^2 + 1]^{-1/2} d\rho \\
&= \frac{mR^2}{2} \int_{R/A}^1 \frac{1}{t^3} (1+t^2)^{-1/2} dt \quad (15B)
\end{aligned}$$

where,  $t = R/\rho < 1$ .

Using the binomial expansion shown in equation (16), it can be shown that:

$$\begin{aligned}
\frac{1}{t^3} (1+t^2)^{-1/2} &= \frac{1}{t^3} + \sum_{n=1}^{\infty} a_n t^{2n-3} \\
&= \frac{1}{t^3} - \frac{2}{t} + \sum_{n=2}^{\infty} a_n t^{2n-3}
\end{aligned}$$

Considering this result, equation (15B) can be written as:

$$\Phi_2 = \frac{mR^2}{2} \left[ \int_{R/A}^1 \frac{1}{t^3} dt - \frac{1}{2} \int_{R/A}^1 \frac{1}{t} dt + \sum_{n=2}^{\infty} a_n \int_{R/A}^1 t^{2n-3} dt \right]$$

therefore;

$$\begin{aligned}
\Phi_2 &= \frac{mR^2}{2} \left[ -\frac{1}{2} + \sum_{n=2}^{\infty} \frac{a_n}{2n-2} + \sum_{n=2}^{\infty} \frac{a_n}{2n-2} \left(\frac{R}{A}\right)^{2n-2} + \right. \\
&\quad \left. \frac{1}{2} \left(\frac{R}{A}\right)^{-2} + \frac{1}{2} \ln(R/A) \right] \quad (18)
\end{aligned}$$

It was shown that

$$\Phi_x = \Phi_1 + \Phi_2 \quad \text{for } R < A$$

So the combination of equations (17) and (18) will result in  $\Phi_x$  for  $R < A$ , that is:

$$\Phi_x = \frac{mR^2}{2} \left[ \frac{1}{3} - \frac{1}{2} + \sum_{n=1}^{\infty} \frac{a_n}{2n+3} + \sum_{n=2}^{\infty} \frac{a_n}{2n-2} + \frac{1}{2} t_1^{-2} - \sum_{n=2}^{\infty} \frac{a_n}{2n-2} t_1^{2n-2} + \frac{1}{2} \ln t_1 \right]$$

where,  $t_1 = (R/A) < 1$

Rearranging this formula results in:

$$\begin{aligned} \Phi_x &= \frac{mA^2}{2} t_1^2 \left[ \frac{1}{3} - \frac{1}{2} - \frac{1}{10} + \sum_{n=2}^{\infty} \frac{(4n+1)a_n}{(2n+3)(2n-2)} + \frac{1}{2} t_1^{-2} - \sum_{n=2}^{\infty} \frac{a_n}{2n-2} t_1^{2n-2} + \frac{1}{2} \ln t_1 \right] \\ &= \frac{mA^2}{2} \left[ \left( \sum_{n=2}^{\infty} \frac{(4n+1)a_n}{(2n+3)(2n-2)} \right) + \frac{8}{30} \right] t_1^2 + \frac{1}{2} - \sum_{n=2}^{\infty} \frac{a_n}{2n-2} t_1^{2n} + \frac{1}{2} t_1^2 \ln t_1 \end{aligned}$$

The first three terms of the expansion of  $\ln t_1$  for the case of  $t_1 \leq 1$  are

$$\ln t_1 \cong (t_1 - 1) - \frac{1}{2} (t_1 - 1)^2 + \frac{1}{3} (t_1 - 1)^3 = \frac{1}{3} t^3 - \frac{3}{2} t^2 + 3t - \frac{11}{6}$$

Substitution of this in  $\Phi_x$  will give the final result of  $\Phi_x$  for  $R < A$ .



$$\Phi_x = \frac{mA^2}{2} \left[ \frac{1}{2} + \left( \left( \sum_{n=2}^{\infty} \frac{(4n+1)a_n}{(2n+3)(2n-2)} \right) - \frac{76}{60} \right) \left( \frac{R}{A} \right)^2 + \frac{3}{2} \left( \frac{R}{A} \right)^3 - \frac{3}{4} \left( \frac{R}{A} \right)^4 + \frac{1}{6} \left( \frac{R}{A} \right)^5 - \sum_{n=2}^{\infty} \frac{a_n}{2n-3} \left( \frac{R}{A} \right)^{2n} \right] \quad (19)$$

Equation (19) shows the potential of the x-axis for the linearly-distributed disc of sources when  $R < A$ . For the other region ( $R > A$ ), equation (15) can be written as

$$\Phi_x = \frac{mR}{2} \int_0^A \frac{\rho^2}{R^2} \left[ 1 + \frac{\rho^2}{R^2} \right]^{-1/2} d\rho \quad (15C)$$

If  $t = \frac{\rho}{R} \leq 1$  and  $t_2 = \frac{A}{R}$ , equation (15C) can be rewritten as:

$$\Phi_x = \frac{mR^2}{2} \int_0^{t_2} t^2 \left[ 1 + t^2 \right]^{-1/2} dt$$

It was previously shown that

$$t^2 (1 + t^2)^{-1/2} = t^2 + \sum_{n=1}^{\infty} a_n t^{2n+2}$$

Hence,

$$\Phi_x = \frac{mR^2}{2} \int_0^{t_2} t^2 dt + \frac{mR^2}{2} \sum_{n=1}^{\infty} a_n \int_0^{t_2} t^{2n+2} dt$$

A simple integration results in:

$$\Phi_x = \frac{mA^2}{2} \left[ \frac{1}{3} \left( \frac{A}{R} \right) + \sum_{n=1}^{\infty} \frac{a_n}{2n+3} \left( \frac{A}{R} \right)^{2n+1} \right] \quad (20)$$

Equations (19) and (20) are the potential of a linearly-distributed disc of sources of radius  $A$ , on the x-axis, for  $R$  smaller than  $A$ , and  $R$  greater than  $A$ , respectively.

To find the general solution, as was pointed out previously, these equations should be equated to equation (5). The result of the comparison determines  $A_n$  and  $B_n$  for both cases of  $R > A$ , and  $R < A$ .

Comparison of the equation (19) with the equation (5) shows that:

$$\begin{aligned}
 A_0 &= \frac{mA^2}{2} \left( \frac{1}{2} \right) \\
 A_2 &= \frac{mA^2}{2} \left( \left( \sum_{n=2}^{\infty} \frac{(4n+1) a_n}{(2n+3)(2n-2)} \right) - \frac{76}{60} \right) \left( \frac{1}{A} \right)^2 \\
 A_3 &= \frac{mA^2}{2} \left( \frac{3}{2} \right) \left( \frac{1}{A} \right)^3 \\
 A_4 &= \frac{mA^2}{2} \left( \frac{3}{4} \right) \left( \frac{1}{A} \right)^4 \\
 A_5 &= \frac{mA^2}{2} \left( \frac{1}{6} \right) \left( \frac{1}{A} \right)^5 \\
 A_n &= \frac{a_n}{2n-3} \left( \frac{1}{A} \right)^{2n} \frac{mA^2}{2} \quad n=2, \infty \\
 B_n &= 0.
 \end{aligned}
 \left. \vphantom{\begin{aligned} A_0 \\ A_2 \\ A_3 \\ A_4 \\ A_5 \\ A_n \\ B_n \end{aligned}} \right\} R < A \quad (21)$$

Comparison of the equation (20) with the equation (5) shows the following results.

$$\begin{aligned}
 A_n &= 0 \\
 B_0 &= \frac{mA^2}{2} \left( \frac{1}{3} \right) A \\
 B_n &= \frac{mA^2}{2} \frac{a_n}{2n+3} A^{2n+1}
 \end{aligned}
 \left. \vphantom{\begin{aligned} A_n \\ B_0 \\ B_n \end{aligned}} \right\} R > A \quad (22)$$

where:

$$a_n = \frac{-1/2 (-1/2 - 1) (-1/2 - 2) \dots (-1/2 - n + 1)}{n!}$$

The general solution can be determined by substitution of equations (21) and (22) in equation (4). The following results can thus be derived:

For  $R < A$ :

$$\begin{aligned} \Phi = \frac{mA^2}{2} & \left[ \frac{1}{2} P_0(\cos \theta) + \left( \sum_{n=2}^{\infty} \frac{(4n+1)a_n}{(2n+3)(2n-2)} - \frac{76}{60} \right) \left(\frac{R}{A}\right)^2 P_2(\cos \theta) \right. \\ & + \frac{3}{2} \left(\frac{R}{A}\right)^3 P_3(\cos \theta) - \frac{3}{4} \left(\frac{R}{A}\right)^4 P_4(\cos \theta) + \frac{1}{6} \left(\frac{R}{A}\right)^5 P_5(\cos \theta) \\ & \left. - \sum_{n=2}^{\infty} \frac{a_n}{2n-2} \left(\frac{R}{A}\right)^{2n} P_{2n}(\cos \theta) \right] \quad (21) \end{aligned}$$

And, for  $R > A$ :

$$\Phi = \frac{mA^2}{2} \left[ \frac{1}{3} \left(\frac{A}{R}\right) P_0(\cos \theta) + \sum_{n=1}^{\infty} \frac{a_n}{2n+3} \left(\frac{A}{R}\right)^{2n+1} P_{2n}(\cos \theta) \right] \quad (22)$$

where:

$a_n$  = binomial coefficients

$$a_n = \frac{-1/2 (-1/2 - 1) (-1/2 - 2) \dots (-1/2 - n + 1)}{n!}$$

Summary. As a conclusion to part four, it is important to make the following summary. The general solution to the potential of a disc of distributed sources of radius  $A$  and source density per unit area  $k$  has been found. The solution has been determined for two cases. The solution to the uniform distribution of sources ( $k = \text{const.}$ ) is shown in the equations (13) and (14). For the linear distribution of sources ( $k = m\rho$ ), the potential is established in equations (21) and (22).

### Superposition of the Potentials

To complete the flow model for a porous disc and windmill in uniform flow, the necessary potentials should be superimposed.

For the operating ease of the porous disc or a windmill without a nacelle, the potential is:

$$\phi = \phi_1 + \phi_2$$

where:

$\phi_1$  = potential of the disc

$\phi_2$  = potential of the uniform flow

In the case of a windmill, the potential can be written as:

$$\phi = \phi_1 + \phi_2 + \phi_3$$

where:

$\phi_1$  = potential of the disc

$\phi_2$  = potential of the single source at the position  $r_2$  on the x-axis as shown in Fig. 3

$\phi_3$  = potential of the uniform flow

$\phi_1$  was established in equations, (13), (14), (21) and (22).  $\phi_2$  can be determined from equations (1) and (3) as follows:

$$\Phi_2 = \frac{k}{4\pi} \frac{1}{[(x-\eta)^2 + (y-\eta)^2 + (z-\xi)^2]^{1/2}} \quad (1)$$

In this case,  $\eta = \epsilon = 0$ .  $\xi = r_2$

so  $\phi_2$  can be written as

$$\Phi_2 = \frac{k}{4\pi} \frac{1}{[(x-r_2)^2 + y^2 + z^2]^{1/2}}$$

Using relations (1A),  $\Phi_2$  can be transformed to a spherical coordinate.

$$x = R \cos \theta$$

$$y = R \sin \theta \cos \alpha$$

$$z = R \sin \theta \sin \alpha$$

and

$$\xi = r_2 \cos \beta, \quad \beta = 0$$

Hence:

$$\Phi_2 = \frac{k}{4\pi} \frac{1}{[(R \cos \theta) - r_2]^2 + R^2 \sin^2 \theta}^{1/2}$$

$\Phi_1$  is the potential of a uniform flow, and it can be represented by:

$$\Phi_3 = v R \cos \theta \quad (24)$$

### Velocity Field

General Note. The velocity field can be determined by superposition of the velocities. The task of this section is to find the components of the velocity vector for each potential.

The relation between the potential function and the velocity vector is known to be:

$$\vec{v} = \nabla \Phi \quad (25)$$

The gradient in spherical coordinates can be shown as:

$$\vec{v} = \frac{\partial \Phi}{\partial r} \hat{a}_r + \frac{1}{r} \frac{\partial \Phi}{\partial \theta} \hat{a}_\theta + \frac{1}{r \sin \theta} \frac{\partial \Phi}{\partial \psi} \hat{a}_\psi$$

In the case of axi-symmetry the gradient reduces to:

$$\vec{U} = \frac{\partial \Phi}{\partial r} \hat{a}_r + \frac{1}{r} \frac{\partial \Phi}{\partial \theta} \hat{a}_\theta \quad (26)$$

Hence, the components of the velocity vector can be shown as

$$U_R = - \frac{\partial \Phi}{\partial r} \quad , \quad \text{and} \quad U_T = - \frac{1}{r} \frac{\partial \Phi}{\partial \theta} \quad (27)$$

where:

$U_R$  = Radial Component of the velocity vector

$U_T$  = Tangential Component of the velocity vector

Now one can supply this generate note to any special case.

Uniform distribution of Sources. For this case the potential was found and shown in equations (13) and (14) as follows:

$$\Phi_1 = \frac{k'A}{2} \sum_{n=0}^{\infty} [P_n(0) + P_{n-2}(0)] \left(\frac{R}{A}\right)^n P_n(\cos \theta) - \quad (13)$$

$$\frac{k'A}{2} \left(\frac{R}{A}\right) P_1(\cos \theta) \quad \text{for } R < A$$

$$\Phi_1 = \frac{k'A}{2} \sum_{n=2}^{\infty} [P_n(0) + P_{n+2}(0)] \left(\frac{A}{R}\right)^{n+1} P_n(\cos \theta) \quad \text{for } R > A \quad (14)$$

a) in the case of  $R < A$ , the components of the velocity vector can be derived as follows:

$$U_{R1} = - \frac{\partial \Phi_1}{\partial R}$$

Therefore, the differentiation of equation (13) yields:

$$U_{R1} = - \frac{k'A}{2} \left[ \sum_{n=0}^{\infty} \frac{n}{A} \left(\frac{R}{A}\right)^{n-1} [P_n(0) + P_{n-2}(0)] P_n(\cos \theta) \right]$$

$$+ \frac{k'A}{2} \left(\frac{1}{A}\right) P_1(\cos \theta)$$

Hence:

$$UR1 = \frac{k'}{2} P_1(\cos\theta) - \frac{k'}{2} \sum_{n=0}^{\infty} n \left(\frac{R}{A}\right)^{n-1} [P_n(0) + P_{n-2}(0)] P_n(\cos\theta) \quad (28)$$

The tangential component can be written as:

$$UT1 = - \frac{1}{R} \frac{\partial \bar{\Phi}_1}{\partial \theta}$$

Using  $\bar{\Phi}_1$  from equation (13), it can be written that

$$UT1 = - \frac{1}{R} \frac{k'A}{2} \left( \sum_{n=0}^{\infty} \left(\frac{R}{A}\right)^n [P_n(0) + P_{n-2}(0)] \frac{\partial P_n(\cos\theta)}{\partial \theta} \right) + \frac{1}{R} \left(\frac{R}{A}\right) \frac{k'A}{2} \frac{\partial P_1(\cos\theta)}{\partial \theta} \quad (29)$$

Using the chain rule, the differentials can be changed to

$$\frac{\partial P_n(\cos\theta)}{\partial \theta} = \frac{\partial P_n(\cos\theta)}{\partial (\cos\theta)} \cdot \frac{d \cos\theta}{d\theta} \quad (30)$$

From the definition of the Associated Legendre polynomials:

$$P_n^m(x) = (1-x^2)^{m/2} \frac{d^m P_n(x)}{d x^m}$$

and by rearrangement:

$$\frac{d^m P_n(x)}{d x^m} = (1-x^2)^{-m/2} P_n^m(x) \quad (31)$$

where:

$P_n^m(x)$  = Associated Legendre polynomial of the first kind, the  $n$  th order, and the  $m$  degree.

Using expression (31) in the equation (30) it can easily be shown that

$$\frac{\partial P_n(\cos\theta)}{\partial\theta} = (1 - \cos^2\theta)^{-1/2} P_n^1(\cos\theta) \cdot (-\sin\theta)$$

and

$$\frac{\partial P_n(\cos\theta)}{\partial\theta} = - P_n^1(\cos\theta)$$

Using the notation  $AS_n(\cos\theta)$  for  $P_n^1(\cos\theta)$ , the above equation can be written as:

$$\frac{\partial P_n(\cos\theta)}{\partial\theta} = - AS_n(\cos\theta) \quad (32)$$

Substitution of the equation (32) into the equation (29) yields:

$$UT_1 = \frac{k'}{2} \left\{ \left[ \sum_{n=0}^{\infty} \left(\frac{R}{A}\right)^{n-1} [P_n(0) + P_{n-2}(0)] AS_n(\cos\theta) \right] - AS_1(\cos\theta) \right\} \quad (33)$$

b) In the case of  $R > A$ , following the same procedure for  $\phi_1$  from the equation (14) it can be written that

$$UR_1 = \frac{k'}{2} \sum_{n=0}^{\infty} (n+1) \left(\frac{A}{R}\right)^{n+2} [P_n(0) + P_{n-2}(0)] P_n(\cos\theta) \quad (34)$$

and

$$UT_1 = \frac{k'}{2} \sum_{n=0}^{\infty} \left(\frac{A}{R}\right)^{n+2} [P_n(0) + P_{n-2}(0)] AS_n(\cos\theta) \quad (35)$$

Accordingly, the velocity components of the uniformly-distributed disc of sources can be written as equations (29) and (33) for the case



of  $R < A$ , and as equations (34) and (35) for the case of  $R > A$ .

Linearly distributed case. For the case of the linearly-distributed disc of sources, it can be shown that the components of the velocity field are as follows: (See Appendix 1 for the details).

For  $R < A$

$$\begin{aligned}
 U_{R1} = & -\frac{mA}{2} \left[ 2 \left( \sum_{n=2}^{\infty} \frac{(4n+1) a_n}{(2n+3)(2n-2)} \right) - \frac{76}{60} \right] \left( \frac{R}{A} \right) P_2(\cos \theta) \\
 & + \frac{9}{2} \left( \frac{R}{A} \right)^2 P_3(\cos \theta) - 3 \left( \frac{R}{A} \right)^3 P_4(\cos \theta) \\
 & + \frac{5}{6} \left( \frac{R}{A} \right)^4 P_5(\cos \theta) - \sum_{n=2}^{\infty} \frac{n a_n}{(n-1)} \left( \frac{R}{A} \right)^{2n-1} P_{2n}(\cos \theta) \quad (36)
 \end{aligned}$$

$$\begin{aligned}
 U_{T1} = & \frac{mA}{2} \left( \frac{A}{R} \right) \left[ \frac{1}{2} A S_0(\cos \theta) + \left( \sum_{n=2}^{\infty} \frac{(4n+1) a_n}{(2n+3)(2n-2)} \right) - \frac{76}{60} \right] \\
 & \left( \frac{R}{A} \right)^2 A S_2(\cos \theta) + \frac{3}{2} \left( \frac{R}{A} \right)^3 A S_3(\cos \theta) - \\
 & \frac{3}{4} \left( \frac{R}{A} \right)^4 A S_4(\cos \theta) + \frac{1}{6} \left( \frac{R}{A} \right)^5 A S_5(\cos \theta) - \\
 & \sum_{n=2}^{\infty} \frac{a_n}{(2n-2)} \left( \frac{R}{A} \right)^{2n} A S_{2n}(\cos \theta) \quad (37)
 \end{aligned}$$

And for the case of  $R > A$ , it can be written:

$$UR1 = \frac{mA}{2} \left[ \frac{1}{3} \left( \frac{A}{R} \right)^2 P_0(\cos\theta) + \sum_{n=1}^{\infty} \frac{(2n+1) a_n}{2n+3} \left( \frac{A}{R} \right)^{2n+2} P_{2n}(\cos\theta) \right] \quad (38)$$

and

$$UT1 = \frac{mA}{2} \left[ \frac{1}{3} \left( \frac{A}{R} \right)^2 AS_0(\cos\theta) + \sum_{n=1}^{\infty} \frac{a_n}{2n+3} \left( \frac{A}{R} \right)^{2n+2} AS_{2n}(\cos\theta) \right] \quad (39)$$

Uniform flow. It is quite simple to show that the components of the velocity field for the uniform flow are as follows:

$$UR3 = -U \cos\theta \quad (40)$$

and

$$UT3 = U \sin\theta \quad (41)$$

Single source at the position  $r_2$  (representing the nacelle). Differentiation of the equation (23) according to equation (27) yields the velocity components due to the body shape of the nacelle as:

$$UR2 = \frac{k}{4\pi} \left( \frac{R - r_2 \cos\theta}{[(R \cos\theta - r_2)^2 + R^2 \sin^2\theta]^{3/2}} \right) \quad (42)$$

and

$$UT2 = \frac{k}{4\pi} \left( \frac{r_2 \sin\theta}{[(R \cos\theta - r_2)^2 + R^2 \sin^2\theta]^{3/2}} \right) \quad (43)$$

Remarks. In order to utilize the previously derived formulas for the velocity field, two computer programs were written. Both of these programs were written for the general case of the presence of all three potentials.

The first program, Program Project 1 in Appendix 2A, was designed for the uniformly distributed disc of source in general, and its application to the University of Massachusetts windmill in particular.

The second program, Program Project 2 in Appendix 2B, deals with the linearly distributed disc of sources in general, and its application to the University of Massachusetts windmill in particular.

The output of both the programs is the velocity field in cartesian coordinates at each point. The output has the form of:  $(x, y)$ , the coordinate of the point,  $(u_x, u_y)$ ,  $x$ , and  $y$  components of the velocity vector at the point  $(x, y)$  and  $(\theta)$ , the angle between  $u_x$  and  $u_y$ .

Streamline construction. The construction of streamlines is based on the fact that no flow crosses a specific stream tube. Once a selection of the streamlines starting point has been made, other points of the same streamline can be found by an application of the conservation of mass.

In reference to Fig. 6, the velocity through the stream tubes is as follows:  $A_0$  is  $u(0)$ ,  $A_1$  is  $u(1)$ ,  $A_2$  is  $u(2)$ , and so on, where  $u(0)$  is the  $x$ -component of the velocity at an  $x$ -station and  $y=0$ , and the same applies for  $(1)$ ,  $(2)$ .....,  $u(a)$ .

Fig. 7 shows the cross section of the stream tubes.

Accordingly, the area corresponding to any velocity  $u(n)$  can be found as follows:

$$\begin{aligned}
 A_0 &= (Dy/2)^2 \pi \\
 A_1 &= (Dy/2 + Dy)^2 \pi - (Dy/2)^2 \pi \\
 A_2 &= (Dy/2 + 2Dy)^2 \pi - A_1 \\
 &\vdots \\
 A_n &= \left(\frac{Dy}{2} + nDy\right)^2 \pi - A_{n-1}
 \end{aligned}
 \quad \left. \vphantom{\begin{aligned} A_0 \\ A_1 \\ A_2 \\ \vdots \\ A_n \end{aligned}} \right\} (44)$$

and

Having the point (a), the starting point of the streamline, the mass flow rate or the volumetric flow rate can be found by summing up the flow rate through tubes  $A_0$  to  $A_a$  as shown in Fig. 8.

When the flow rate through the stream tube (a) is determined at one  $x$  station, other points of the same stream tube can subsequently be calculated.

At each station  $x$ , the flow rate through the tubes with cross-section  $A_0, A_1, A_2 \dots$  and  $A_a$  (shown in Fig. 7) should be easily found. The desired stream tube can be determined by summing up these calculated flow rates,  $F_1, F_2 \dots$  up to  $F_n$  where  $n$  indicates the point where this summation is equal to the reference flow rate. If  $n$  is known, the  $y$ -coordinate can be determined.

The task of the streamline construction is performed by two computer programs. The first one constructs the streamline for a uniformly distributed disc of sources in general, and for the University of Massachusetts windmill, in particular (see Project S1, Appendix 2C). The second program deals with construction of the streamlines for a linearly distributed disc of sources in general, and for the University of Massachusetts windmill in particular, (see Project S2, Appendix 2D).

The input of these programs can be the starting point of the streamline or the flow rate through the stream tube.

The output indicates the form of flow rate, and y coordinate indicates the streamline at each x station.

Applying the one-dimensional momentum theory to the windmills. The source strength density per unit area  $k$  can be found in the case of the uniformly distributed source disc and the slope of  $k$ , (that is  $m$ ) for the linearly distributed source disc, by the one dimensional momentum theory and Lagally's theorem.

According to the Lagally's theorem [10], the force exerted upon a point source in a uniform flow is  $\delta \lambda u$ , where  $\delta$  is the density of the fluid,  $\lambda$  is the source strength, and  $u$  is the free stream velocity of the uniform flow.

a) Linearly distributed disc.

In this case, the source-strength density is equal to  $mr$ , where  $m$  is the slope of  $k$  as described in section four. According to Lagally's theorem, the force on the disc is:

$$F = \delta k u \quad (45)$$

where  $k$  is the total strength of the disc. However,  $k$  can be determined as follows:

$$k = \int_0^A mr (2\pi r dr) = \frac{2}{3} \pi m A^2 \quad (46)$$

Substituting equation (46) into equation (45);

$$F = \frac{2}{3} \pi m A^3 \delta U \quad (47)$$

From one dimensional momentum theory it can be shown [11] that the total force on the disc is (See Fig. 9)

$$F = \delta \pi A^2 u (U - U_1) \quad (48)$$

where  $u$  is the velocity through the disc and  $U_1$  is the velocity down stream of the disc.

Using the Bernoulli's equation and the fact that  $F = \pi A^2 \Delta P$ , it can be shown that [8]:

$$F = \delta \frac{S}{2} (U^2 - U_1^2) \quad (49)$$

and

$$u = \frac{U + U_1}{2} \quad (50)$$

where  $S$  is the area of the disc.

From the definition of the axial interference coefficient  $a$ , it can be written that

$$u = (1-a)U \quad (51)$$

Using equation (50)

$$U_1 = (1-2a)U$$

and

$$U_1 - U = -2aU$$

or

$$U - U_1 = 2aU \quad (52)$$

which can be written as

$$U_1 + U = 2aU + 2U = 2(a+1)U \quad (53)$$

Multiplication of equation (52) by (53) will result in

$$U^2 - U_1^2 = 4a(1+a)U^2$$

substitution into equation (49) leads to;

$$F = 2\delta\pi A^2 a(1+a)U^2$$

The slope  $m$  can be determined by equating equations (47) and (54), that is

$$m = \frac{3a(1+a)}{A} U \quad (55)$$

b) Uniformly distributed disc.

Following the same procedure, the total disc strength in this case is

$$k = \hat{k}\pi r^2 \quad (56)$$

and the total force on the disc from the Lagally's theorem is:

$$F = U\hat{k}\pi A^2\delta \quad (57)$$

Equating this formula with the equation (54) will provide  $k$ , hence

$$\hat{k} = 2a(1+a)U \quad (58)$$

It can be shown [12] that the relation between the axial interference factor ( $a$ ) and the power coefficient  $C_p$  is as follows:

$$C_p = 4a(1-a)^2 \quad (59)$$

The maximum power is developed when  $a = 1/3$  [8].

Thus it has been shown through Lagally's theorem and the simple one-dimensional momentum theory that there is a unique relation between the strength of the distributed source disc  $k$  and the power coefficient  $C_p$  of the windmill.

## RESULTS

The velocity field and some characteristic streamlines have been calculated for the numerical values appropriate to the 25 kW windmill at the University of Massachusetts Solar Habitat I (See Fig. 10).

A characteristic free stream velocity of 38 ft/sec (11.58 m/sec) have been considered. The body of this windmill can be modeled by a single source of  $K=647.74\text{lbm/sec}$  (294 kg/sec) in a free stream of  $U = 38$  ft/sec.

A sample result of the velocity field upstream the windmill is shown at the end of Project 1 and Project 2 [see Appendix 2A and 2B].

The velocity profile of the uniformly-distributed disc and the linearly-distributed disc model is shown in figs. (11) and (12).

The streamlines constructed for different cases are shown in Fig. 13 through 20.

It is obvious from Figs. 13 through 20 that the effective change in the free stream velocity is almost negligible for more than two radii upstream of the blade disc.

Fig. 13 represents the uniformly distributed disc model without the body for the case of  $C_p = C_{p_{\max}} = 0.5$ . It is shown that the disc samples almost 57 percent of the volume of the far upstream wind. For the same case, if the body is located at the center of the blade disc (Fig. 15), 51 percent of the flow is sampled. By moving the single source, which forms the body to the position ( $R_2 = 5.02$  ft), to model the University of Massachusetts windmill, the percentage of the flow sampled reduces to 35.

The same analysis for the linearly-distributed disc model (Figs. 14, 16, and 18) indicates similar results. The percentage of the wind sampled



changes from 39 to 38 and then to 30.

Comparison of the uniform model and linear model shows that for the same conditions, the linear model samples less flow than the uniform disc. To see this, Figs. 17 and 18 must be compared.

The uniformly-distributed disc samples 35 percent of the flow, while the linearly distributed disc samples 30 percent.

The effect of the body is not restricted to the sampling problem. Although this is the case for the uniformly-distributed model, in the case of the more realistic linearly-distributed model, a significant difference is observed - the problem of leaking.

Figs. 16 and 18 show that due to the higher resistance at the outer region of the blades, some of the sampled flow appears to leak through the central region near the body. This is quite obvious in Fig. 12B, where at the station  $x=1^{ft}$  ( $0.3^m$ ), the  $y$  component of the velocity is down towards the center between  $y=4^{ft}$  ( $1.2^m$ ) and  $y=10^{ft}$  ( $3^m$ ).

Using the velocity field the pressure increase in front of the disc can quite easily be calculated.

In the case of the linearly-distributed source disc the Bernoulli equation on the streamline  $\psi_0$  (see Fig. 18) can be written as

$$\frac{1}{2} U^2 + P/S = \text{const.} \quad (60)$$

At  $x = \infty$ , the velocity is 38 ft/sec and the pressure is  $P_\infty$ , with density  $\delta$ . At  $x = 1$ , and  $y = 16$ , the velocity is:

$$U = \sqrt{28^2 + 19^2} = 33.8 \text{ ft/sec}$$

then, it is easy to show that

$$\text{or } P = P_{\infty} + 149.58 \quad (61)$$

$$P = 14.801 \text{ psi } (101.84 \text{ kN/m}^2)$$

Following the same procedure for the uniformly-distributed disc of sources, the pressure would be:

$$P = P_{\infty} + 168.9 \quad (62)$$

$$\text{or } P = 14.802 \text{ psi } (101.85 \text{ kN/m}^2)$$

Comparing equation (60) and (62), it is quite obvious that the pressure on the streamline  $\psi_0$  increases more in front of the linear distributed source disc.

Fig. 20 shows the effect of a change in the power coefficient,  $C_p$ . The percentage of the flow sampled changes from 57 to 80 percent while the  $C_p$  is changed from  $C_{p_{\max}}$  to  $C_{p_{\max}}/2$ . Also, in the limit as  $C_p \rightarrow 0$ , all the flow would pass through the blade disc unaffected.

The variation of the velocity on the stagnation streamline has been represented in Fig. 21. The diagram shows the velocity variation for the body and the linearly-distributed disc with the body. It also shows that the stagnation point has shifted forwards, and that the velocity decreases more rapidly with the blade disc than for the body alone.

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APPENDIX 1

## APPENDIX 1

1)  $R < A$ ,

The potential  $\Phi$  was shown in the equation (21).

then 
$$U_R = - \frac{\partial \Phi}{\partial R}$$

$$\begin{aligned} U_R = & - \frac{mA^2}{2} \left[ \left( \sum_{n=2}^{\infty} \frac{(4n+1)a_n}{(2n+3)(2n-2)} \right) - \frac{76}{60} \right] \left( \frac{2R}{A^2} \right) P_2(\cos\theta) \\ & + \frac{3}{2} \left( \frac{3R^2}{A^3} \right) P_3(\cos\theta) - \frac{3}{4} \left( \frac{4R^3}{A^4} \right) P_4(\cos\theta) \\ & - \frac{1}{6} \left( \frac{5R^4}{A^5} \right) P_5(\cos\theta) - \\ & \left. \sum_{n=2}^{\infty} \frac{a_n}{2n-2} \frac{2n R^{2n-1}}{A^{2n}} P_{2n}(\cos\theta) \right] \end{aligned}$$

This can be arranged as:

$$\begin{aligned} U_R = & - \frac{mA^2}{2} \left[ \frac{1}{A} \left( \sum_{n=2}^{\infty} \frac{(4n+1)a_n}{(2n+3)(2n-2)} \right) - \frac{76}{60} \right] \left( \frac{R}{A} \right) P_2(\cos\theta) \\ & + \frac{9}{2A} \left( \frac{R}{A} \right)^2 P_3(\cos\theta) - \frac{3}{A} \left( \frac{R}{A} \right)^3 P_4(\cos\theta) + \end{aligned}$$

$$\frac{5}{6A} \left(\frac{R}{A}\right)^4 P_3(\cos\theta) - \frac{1}{A} \sum_{n=2}^{\infty} \frac{n a_n}{n-1} \left(\frac{R}{A}\right)^{2n-1} P_{2n}(\cos\theta) ]$$

which can be reduced to:

$$\begin{aligned} U_R = & -\frac{mA}{2} \left[ 2 \left( \sum_{n=2}^{\infty} \frac{(4n+1) a_n}{(2n+3)(2n-2)} \right) - \frac{76}{60} \right] \left(\frac{R}{A}\right) P_2(\cos\theta) \\ & + \frac{9}{2} \left(\frac{R}{A}\right)^2 P_3(\cos\theta) - 3 \left(\frac{R}{A}\right)^3 P_4(\cos\theta) \\ & + \left(\frac{5}{6}\right) \left(\frac{R}{A}\right)^4 P_5(\cos\theta) - \\ & \sum_{n=2}^{\infty} \frac{n a_n}{(n-1)} \left(\frac{R}{A}\right)^{2n-1} P_{2n}(\cos\theta) ] \end{aligned}$$

which is the equation (36).

Equation (37) can be derived as follows, using the same potential (equation 21) and knowing that

$$U_T = -\frac{1}{R} \frac{\partial \Phi}{\partial \theta}$$

it is easy to show that:

$$\begin{aligned} U_T = & -\frac{1}{R} \frac{mA^2}{2} \left[ -\frac{1}{2} A S_0(\cos\theta) + \right. \\ & \left. \left( \sum_{n=2}^{\infty} \frac{(4n+1) a_n}{(2n+3)(2n-2)} \right) - \frac{76}{60} \right] \left(\frac{R}{A}\right)^2 (-A S_2(\cos\theta)) \end{aligned}$$

$$\begin{aligned}
& + \frac{3}{2} \left( \frac{R}{A} \right)^3 (-AS_3(\cos\theta)) - \frac{3}{4} \left( \frac{R}{A} \right)^4 (-AS_4(\cos\theta)) \\
& + \frac{1}{6} \left( \frac{R}{A} \right)^5 (-AS_5(\cos\theta)) - \\
& \left[ \sum_{n=2}^{\infty} \frac{a_n}{2n-2} \left( \frac{R}{A} \right)^{2n} (-AS_{2n}(\cos\theta)) \right]
\end{aligned}$$

By a little algebraic multiplication equation (37) can easily be derived.

2)  $R > A$ .

In this case the equation 22 should be differentiated to produce UR, and UT.

APPENDIX 2: A & B



```

00100 PROGRAM PROJ(INPUT,OUTPUT)
00110 REAL K,L,K1
00120C*****
00130C
00140C
00150C PROGRAM FOR UNIFORMLY DISRIBUTED DISC OF SOURCES
00160C WITH THE BODY
00170C
00180C
00190C THIS PROGRAM FINDS THE VELOCITY FIELD(UX,UY,ETA)
00200C FOR A UNIFORMLY DISTRIBUTED DISC OF SOURCES AND
00210C A STRONG SOURCE AT POSITION R2 ON THE X-AXIS.
00220C
00230C
00240C TO RUN THIS PROGRAM:
00250C ONE SHODD INPUT THE FOLLOWINGS
00260C IN RESPONSE TO THE ASK FOR DATA:
00270C NX,NY,DELTX,DELTY,U,A,A1,R2,K1
00280C ACCORDING TO THE DEFINITION OF
00290C THE PARAMETERS GIVEN BELOW:
00340C
00350C
00360C Y-AXIS
00370C :
00380C +-----+ (NY-1)*DELTY
00390C :
00400C :
00410C : VELOCITY SPACE
00420C :
00430C :
00440C +-----+ X-AXIS
00450C ORIGIN NX*DELTX
00460C
00470C PARAMETERS:
00500C NX: NUMBER OF THE INCREMENTS IN X-DIREC.
00510C NY:NUMBER OF THE INCREMENTS IN Y-DIREC.
00520C DELTX: INCREMENT IN X-DIRECTION
00530C DELTY: INCREMENT IN Y-DIRECTION
00540C A:RADIOUS OF THE DISC
00550C XX: VALUE OF THE X-AXIS
00560C YY: VALUE OF THE Y-AXIS
00570C K1: SRENGTH OF THE BODY SOURCE
00580C R2:POSITION OF THE BODY SOURCE
00590C U: FREE STREAM VELOCITY
00600C A1: AXIAL INTERFERENCE FACTOR
00610C K:SOURCE STRENGTH DENSITY ON THE DISC
00620C R AND T: RADICUS AND ANGLE OF POLAR
00630C COORDINATE IN FT.AND RAD.

```

```

00640C      UT1: TANGANTIAL VELOCITY AT POIN (R,T)      *
00650C      DUE TO THE DISC.                            *
00660C      UR1: RADIAL VELOCITY AT POINT (R,T)        *
00670C      DUE TO THE DISC                            *
00680C      UT2: TANGANTIAL VELOCITY AT POIN (R,T)    *
00690C      DUE TO THE BODY                            *
00700C      UR2: RADIAL VELOCITY AT POIN (R,T)        *
00710C      DUE TO THE BODY                            *
00720C      UT3: TANGANTIAL VELOCITY AT POINT (R,T)   *
00730C      DUE TO THE FREE STREAM.                   *
00740C      UR3: RADIAL VELOCITY AT POINT (R,T)        *
00750C      DUE TO THE FREE STREAM.                   *
00760C      P(N,X): LEGENDRE POLYNOMIAL OF THE         *
00770C      FIRST KIND AND NTH ORDER DEFINED           *
00780C      BY THE FUNCTION P(N,X).                    *
00790C      AS(N,X): ASSOCIATED LEGENDRE OF THE        *
00800C      FIRST KIND ,FIRST DEGREE.                  *
00810C      AND,NTH ORDER DEFINED BY THE              *
00820C      AS(N,X) FUNCTION                            *
00830C      UX: X COMPONENT OF THE VELOCITY AT         *
00840C      ANY POINT (R,T)OR ITS EQUIVALENT(XX,YY)    *
00850C      UY: Y COMPONENT OF THE VELOCITY AT         *
00860C      ANY POINT(R,T) OR ITS EQUIVALENT(XX,YY)    *
00870C      ETA: ANGL BETWEEN UX AND, UY IN DEG.      *
00880C      *
00890C      *
00900C-----THE PROGRAM WAS RUN FOR THE FOLLOWING    *
00910C      VALUES:NX=20,NY=21,DELTX=1,DELTY=1,U=38     *
00920C      ,A=16,A1=1/3(FOR MAX. POWER),R2=0.          *
00930C      ,K1=647.78                                    *
00940C      *
00950C      ALSO SEE THE PROJECT REPORT                  *
00960C*****
00970 READ,NX,NY,DELTX,DELTY,U,A,A1,R2,K1
00980 IN=50
00981 PRINT 210,NX,NY,DELTX,DELTY,U,A,A1,R2,K1
00982 210 FORMAT(///,*NX=*,I5,/*NY=*,I5,/*DELTX=*,F5.3,/*DELTY=*,F5.3,
00983+/*U=*,F10.4,/*A=*,F6.2,/*A1=*,F10.6,/*R2=*,F10.5,/*K1=*,F10.4)
00990 PRINT 21
01000 21 FORMAT(//,25X,*VELOCITY FIELD FOR UNIFORM DISC AND THE BODY
01010+*,/,25X,*-----*)
01020 PRINT 10
01030 10 FORMAT(5X,*X IN FT*,9X,*Y IN FT*,6X,*UX IN FT/SEC*,5X,*UY IN FT/SEC
01040+5X,*ETA IN DEG.*,/,5X,*-----*,9X,*-----*,6X,*-----*,5X
01041+,*-----*,5X,*-----*)
01050 DO 600 III=1,NX
01060 DO 500 II=1,NY
01070 YY=(II-1)*DELTY
01080 XX=FLOAT(III)*DELTX
01090 K=2*A1*(1.+A1)*U
01100 R=((XX**2)+(YY**2))*0.5)

```

```

01110 N0=0
01120 N1=1
01130 T=ATAN((YY/XX))
01140 T1=COS(T)
01150 T2=SIN(T)
01160 T0=0.
01170 IF(R.GE.A) GO TO 5
01180C
01190C
01200C          DISC FOR THE CASE R<A
01210C
01220C
01230 AK=K/2.
01240 X=R/A
01250 SUM=0.
01260 DO 1 I=1,IN
01270 N=I-1
01280 M=N-2
01290 SUM=SUM+(N*(X**(N-1))*(P(N,T0)+P(M,T0))*P(N,T1))
01300 1 CONTINUE
01310 UR1=AK*(P(N1,T1)-SUM)
01320 SUM=0.
01330 DO 2 I=1,IN
01340 N=I-1
01350 M=N-2
01360 SUM=SUM+((X**(N-1))*(P(N,T0)+P(M,T0))*AS(N,T1))
01370 2 CONTINUE
01380 UT1=AK*(SUM-AS(N1,T1))
01390 GO TO 200
01400 5 CONTINUE
01410C
01420C
01430C          DISC FOR THE CASE R>A
01440C
01450C
01460 X=A/R
01470 SUM=0.
01480 DO 3 I=1,IN
01490 N=I-1
01500 M=N+2
01510 SUM=SUM+((N+1)*(X**M)*(P(N,T0)+P(M,T0))*P(N,T1))
01520 3 CONTINUE
01530 UR1=(K/2)*SUM
01540 SUM=0.
01550 DO 4 I=1,IN
01560 N=I-1
01570 M=N+2
01580 SUM=SUM+((X**M)*(P(N,T0)+P(M,T0))*AS(N,T1))
01590 4 CONTINUE
01600 UT1=(K/2)*SUM
01610 200 CONTINUE

```

```
01620C
01630C
01640C          BODY CALCULATIONS
01650C
01660C
01670 D15=(R-(R2*T1))
01680 D16=((R*T1)-R2)**2
01690 D17=(R**2)*(T2**2)
01700 D18=((D16+D17)**(3/2))
01710 AK1=(K1/(4*3.14))
01720 UR2=AK1*D15/D18
01730 D19=R2*T2
01740 D110=((((R*T1)-R2)**2)+((R**2)*(T2**2)))*(3./2.)
01750 AK3=K1/(4.*3.14)
01760 UT2=AK3*D19/D110
01770C
01780C
01790C          UNIFORM FLOW
01800C
01810C
01820 UR3=-U*T1
01830 UT3=U*T2
01840C
01850C
01860C          SUPERPOSITION
01870C
01880C
01890 UR=UR1+UR2+UR3
01900 UT=UT1+UT2+UT3
01910C
01920C
01930C          EVALUATION OF UX , UY , AND ETA
01940C
01950C
01960 UX=(UR*T1)-(UT*T2)
01970 UY=(UR*T2)+(UT*T1)
01980 XY=UY/UX
01990 ET=ATAN(XY)*360./(2.*3.14)
02000 PRINT 300,XX,YY,UX,UY,ET
02010 300 FORMAT(5(5X,F10.4))
02020 500 CONTINUE
02030 PRINT 700
02040 700 FORMAT(/)
02050 600 CONTINUE
02060 END
02070C
02080C
02090C
02100C
02110C
02120C
```

```
02130C
02140C
02150C
02160C
02170C
02180C
02190C-----*
02200 FUNCTION P(N,X)
02210C
02220C
02230C THIS IS A FUNCTION TO CALCULATE LEGENDER POLYNOMIALS OF THE
02240C FIRST KIND
02250C
02260C
02270 Q0=0.
02280 P=Q0
02290 IF(N.LT.0) GO TO 1
02300 Q0=1.
02310 P=Q0
02320 IF(N.EQ.0) GO TO 1
02330 Q1=X
02340 P=Q1
02350 IF(N.EQ.1) GO TO 1
02360 Q2=((3*(X**2))-1.)/2.
02370 P=Q2
02380 IF(N.EQ.2) GO TO 1
02390 F1=((3*(X**2))-1.)/2.
02400 F2=X
02410 I=3
02420 2 CONTINUE
02430 P=(((2.*I)-1.)/I)*X*F1-((I-1.)/I)*F2
02440 IF(I.EQ.N) GO TO 1
02450 I=I+1
02460 F2=F1
02470 F1=P
02480 GO TO 2
02490 1 CONTINUE
02500 RETURN
02510 END
02520C
02530C
02540C
02550C
02560C
02570C
02580C
02590C
02600C
02610C
02620C
02630C
```

```
02640C
02650C-----*
02660 FUNCTION AS(N,X)
02670 DIMENSION R(1000)
02680C
02690C
02700C THIS FUNCTION CALCULATES ASSOCIATED LEGENDER POLYNOMIALS
02710C OF THE FIRST KIND AND FIRST POWER
02720C
02730C
02740 R0=0.
02750 RA=R0
02760 IF(N.LE.0) GO TO 3
02770 R1=(1-(X**2))**0.5
02780 RA=R1
02790 IF(N.EQ.1) GO TO 3
02800 R2=3.*X*((1-(X**2))**0.5)
02810 RA=R2
02820 IF(N.EQ.2) GO TO 3
02830 P1=X
02840 P2=((3*(X**2))-1.)/2.
02850 DO 10 I=3,N
02860 R(I)=(((2*I)-1)*((1-(X**2))**0.5)*P2)+R1
02870 P=(((2.*I)-1.)/I)*X*P2-((I-1.)/I)*P1
02880 P1=P2
02890 P2=P
02900 R1=R2
02910 R2=R(I)
02920 RA=R(I)
02930 10 CONTINUE
02940 3 CONTINUE
02950 AS=RA
02960 RETURN
02970 END
02980C
READY.
```

RNN

78/07/04. 22.15.48.  
FILE PROJ1

? 16,16,4.,2.,38.,16.,0.33,5.02,647.78

NX= 16  
NY= 16  
DELTX=4.000  
DELY=2.000  
U= 38.0000  
A= 16.00  
A1= .330000  
R2= 5.02000  
K1= 647.7800

VELOCITY FIELD FOR UNIFORM DISC AND THE BODY

| X IN FT | Y IN FT | UX IN FT/SEC | UY IN FT/SEC. | ETA IN I |
|---------|---------|--------------|---------------|----------|
| 4.0000  | 0.      | -75.9304     | -.0000        | .0000    |
| 4.0000  | 2.0000  | -30.1489     | 10.0254       | -18.4028 |
| 4.0000  | 4.0000  | -22.8738     | 8.2835        | -19.9175 |
| 4.0000  | 6.0000  | -23.1659     | 8.6163        | -20.4124 |
| 4.0000  | 8.0000  | -24.2123     | 9.0017        | -20.4046 |
| 4.0000  | 10.0000 | -25.3713     | 9.5443        | -20.6261 |
| 4.0000  | 12.0000 | -26.7800     | 10.2562       | -20.9665 |
| 4.0000  | 14.0000 | -28.7235     | 10.9299       | -20.8435 |
| 4.0000  | 16.0000 | -31.3540     | 10.9288       | -19.2263 |
| 4.0000  | 18.0000 | -33.6426     | 9.9691        | -16.5141 |
| 4.0000  | 20.0000 | -35.1916     | 8.5545        | -13.6697 |
| 4.0000  | 22.0000 | -36.0882     | 7.2744        | -11.4023 |
| 4.0000  | 24.0000 | -36.6225     | 6.2429        | -9.6790  |
| 4.0000  | 26.0000 | -36.9612     | 5.4270        | -8.3573  |
| 4.0000  | 28.0000 | -37.1885     | 4.7762        | -7.3223  |
| 4.0000  | 30.0000 | -37.3484     | 4.2495        | -6.4945  |
| 8.0000  | 0.      | -11.4735     | -.0000        | .0000    |
| 8.0000  | 2.0000  | -16.0516     | 5.3424        | -18.4181 |
| 8.0000  | 4.0000  | -21.1687     | 6.4546        | -16.9657 |
| 8.0000  | 6.0000  | -24.1212     | 6.7699        | -15.6853 |
| 8.0000  | 8.0000  | -26.0884     | 7.0450        | -15.1196 |
| 8.0000  | 10.0000 | -27.6714     | 7.3130        | -14.8112 |
| 8.0000  | 12.0000 | -29.1343     | 7.5240        | -14.4878 |

|         |         |          |        |          |
|---------|---------|----------|--------|----------|
| 8.0000  | 14.0000 | -30.6546 | 7.7273 | -14.1553 |
| 8.0000  | 16.0000 | -31.9906 | 7.4480 | -13.1128 |
| 8.0000  | 18.0000 | -33.2753 | 7.0658 | -11.9945 |
| 8.0000  | 20.0000 | -34.3391 | 6.5193 | -10.7552 |
| 8.0000  | 22.0000 | -35.1599 | 5.9136 | -9.5521  |
| 8.0000  | 24.0000 | -35.7708 | 5.3255 | -8.4722  |
| 8.0000  | 26.0000 | -36.2219 | 4.7920 | -7.5400  |
| 8.0000  | 28.0000 | -36.5576 | 4.3234 | -6.7481  |
| 8.0000  | 30.0000 | -36.8110 | 3.9175 | -6.0777  |
| 12.0000 | 0.      | -23.9398 | -.0000 | .0000    |
| 12.0000 | 2.0000  | -24.4340 | 1.8019 | -4.2199  |
| 12.0000 | 4.0000  | -25.6093 | 3.1907 | -7.1055  |
| 12.0000 | 6.0000  | -26.9719 | 4.1169 | -8.6829  |
| 12.0000 | 8.0000  | -28.2769 | 4.7155 | -9.4724  |
| 12.0000 | 10.0000 | -29.4728 | 5.0455 | -9.7192  |
| 12.0000 | 12.0000 | -30.5842 | 5.3377 | -9.9048  |
| 12.0000 | 14.0000 | -31.6260 | 5.4193 | -9.7285  |
| 12.0000 | 16.0000 | -32.5898 | 5.3807 | -9.3800  |
| 12.0000 | 18.0000 | -33.4611 | 5.2281 | -8.8848  |
| 12.0000 | 20.0000 | -34.2222 | 4.9879 | -8.2967  |
| 12.0000 | 22.0000 | -34.8652 | 4.6934 | -7.6707  |
| 12.0000 | 24.0000 | -35.3948 | 4.3755 | -7.0507  |
| 12.0000 | 26.0000 | -35.8241 | 4.0575 | -6.4651  |
| 12.0000 | 28.0000 | -36.1698 | 3.7539 | -5.9283  |
| 12.0000 | 30.0000 | -36.4481 | 3.4723 | -5.4447  |
| 16.0000 | 0.      | -27.4909 | .0000  | -.0000   |
| 16.0000 | 2.0000  | -28.4354 | .8270  | -1.6668  |
| 16.0000 | 4.0000  | -28.9778 | 1.7890 | -3.5345  |
| 16.0000 | 6.0000  | -29.6169 | 2.5308 | -4.8866  |
| 16.0000 | 8.0000  | -30.3266 | 3.0826 | -5.8069  |
| 16.0000 | 10.0000 | -31.0677 | 3.4780 | -6.3908  |
| 16.0000 | 12.0000 | -31.8054 | 3.7426 | -6.7146  |
| 16.0000 | 14.0000 | -32.5177 | 3.8954 | -6.8346  |
| 16.0000 | 16.0000 | -33.1900 | 3.9526 | -6.7949  |
| 16.0000 | 18.0000 | -33.8105 | 3.9300 | -6.6335  |
| 16.0000 | 20.0000 | -34.3707 | 3.8442 | -6.3850  |
| 16.0000 | 22.0000 | -34.8659 | 3.7123 | -6.0807  |
| 16.0000 | 24.0000 | -35.2961 | 3.5504 | -5.7469  |
| 16.0000 | 26.0000 | -35.6647 | 3.3721 | -5.4040  |
| 16.0000 | 28.0000 | -35.9776 | 3.1881 | -5.0665  |
| 16.0000 | 30.0000 | -36.2419 | 3.0058 | -4.7435  |
| 20.0000 | 0.      | -30.9024 | -.0000 | .0000    |
| 20.0000 | 2.0000  | -30.9754 | .5997  | -1.1098  |
| 20.0000 | 4.0000  | -31.1848 | 1.1631 | -2.1371  |
| 20.0000 | 6.0000  | -31.5059 | 1.6621 | -3.0213  |



|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 20.0000 | 8.0000  | -31.9068 | 2.0805 | -3.7326 |
| 20.0000 | 10.0000 | -32.3569 | 2.4133 | -4.2676 |
| 20.0000 | 12.0000 | -32.8305 | 2.6629 | -4.6394 |
| 20.0000 | 14.0000 | -33.3077 | 2.8356 | -4.8685 |
| 20.0000 | 16.0000 | -33.7734 | 2.9401 | -4.9777 |
| 20.0000 | 18.0000 | -34.2167 | 2.9859 | -4.9897 |
| 20.0000 | 20.0000 | -34.6297 | 2.9831 | -4.9260 |
| 20.0000 | 22.0000 | -35.0076 | 2.9420 | -4.8062 |
| 20.0000 | 24.0000 | -35.3482 | 2.8720 | -4.6473 |
| 20.0000 | 26.0000 | -35.6514 | 2.7817 | -4.4638 |
| 20.0000 | 28.0000 | -35.9187 | 2.6785 | -4.2669 |
| 20.0000 | 30.0000 | -36.1529 | 2.5682 | -4.0653 |

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 24.0000 | 0.      | -32.4816 | -.0000 | .0000   |
| 24.0000 | 2.0000  | -32.5221 | .4046  | -.7131  |
| 24.0000 | 4.0000  | -32.6404 | .7920  | -1.3907 |
| 24.0000 | 6.0000  | -32.8271 | 1.1477 | -2.0034 |
| 24.0000 | 8.0000  | -33.0689 | 1.4609 | -2.5308 |
| 24.0000 | 10.0000 | -33.3511 | 1.7255 | -2.9632 |
| 24.0000 | 12.0000 | -33.6592 | 1.9393 | -3.2991 |
| 24.0000 | 14.0000 | -33.9805 | 2.1033 | -3.5438 |
| 24.0000 | 16.0000 | -34.3040 | 2.2208 | -3.7060 |
| 24.0000 | 18.0000 | -34.6211 | 2.2964 | -3.7967 |
| 24.0000 | 20.0000 | -34.9253 | 2.3354 | -3.8276 |
| 24.0000 | 22.0000 | -35.2119 | 2.3439 | -3.8102 |
| 24.0000 | 24.0000 | -35.4780 | 2.3275 | -3.7553 |
| 24.0000 | 26.0000 | -35.7220 | 2.2916 | -3.6724 |
| 24.0000 | 28.0000 | -35.9436 | 2.2412 | -3.5698 |
| 24.0000 | 30.0000 | -36.1434 | 2.1804 | -3.4541 |

|         |        |          |        |         |
|---------|--------|----------|--------|---------|
| 28.0000 | 0.     | -33.5582 | -.0000 | .0000   |
| 28.0000 | 2.0000 | -33.5830 | .2876  | -.4909  |
| 28.0000 | 4.0000 | -33.6558 | .5660  | -.9639  |
| 28.0000 | 6.0000 | -33.7724 | .8269  | -1.4034 |

## \*TIME LIMIT\*

SRU 21.302 UNTS.

T,500

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 28.0000 | 8.0000  | -33.9266 | 1.0638 | -1.7970 |
| 28.0000 | 10.0000 | -34.1108 | 1.2720 | -2.1367 |
| 28.0000 | 12.0000 | -34.3169 | 1.4489 | -2.4189 |
| 28.0000 | 14.0000 | -34.5373 | 1.5937 | -2.6433 |
| 28.0000 | 16.0000 | -34.7649 | 1.7072 | -2.8127 |
| 28.0000 | 18.0000 | -34.9936 | 1.7912 | -2.9317 |
| 28.0000 | 20.0000 | -35.2185 | 1.8485 | -3.0061 |
| 28.0000 | 22.0000 | -35.4357 | 1.8823 | -3.0421 |

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 28.0000 | 24.0000 | -35.6424 | 1.8958 | -3.0463 |
| 28.0000 | 26.0000 | -35.8366 | 1.8926 | -3.0245 |
| 28.0000 | 28.0000 | -36.0174 | 1.8756 | -2.9825 |
| 28.0000 | 30.0000 | -36.1842 | 1.8479 | -2.9250 |

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 32.0000 | 0.      | -34.3276 | -.0000 | .0000   |
| 32.0000 | 2.0000  | -34.3437 | .2126  | -.3549  |
| 32.0000 | 4.0000  | -34.3913 | .4199  | -.6999  |
| 32.0000 | 6.0000  | -34.4681 | .6169  | -1.0259 |
| 32.0000 | 8.0000  | -34.5710 | .7995  | -1.3254 |
| 32.0000 | 10.0000 | -34.6958 | .9643  | -1.5928 |
| 32.0000 | 12.0000 | -34.8379 | 1.1093 | -1.8247 |
| 32.0000 | 14.0000 | -34.9927 | 1.2333 | -2.0195 |
| 32.0000 | 16.0000 | -35.1558 | 1.3361 | -2.1776 |
| 32.0000 | 18.0000 | -35.3229 | 1.4183 | -2.3005 |
| 32.0000 | 20.0000 | -35.4906 | 1.4812 | -2.3910 |
| 32.0000 | 22.0000 | -35.6559 | 1.5262 | -2.4523 |
| 32.0000 | 24.0000 | -35.8165 | 1.5555 | -2.4880 |
| 32.0000 | 26.0000 | -35.9705 | 1.5708 | -2.5018 |
| 32.0000 | 28.0000 | -36.1168 | 1.5743 | -2.4972 |
| 32.0000 | 30.0000 | -36.2545 | 1.5680 | -2.4777 |

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 36.0000 | 0.      | -34.8978 | -.0000 | .0000   |
| 36.0000 | 2.0000  | -34.9087 | .1622  | -.2664  |
| 36.0000 | 4.0000  | -34.9412 | .3212  | -.5269  |
| 36.0000 | 6.0000  | -34.9940 | .4738  | -.7760  |
| 36.0000 | 8.0000  | -35.0653 | .6172  | -1.0089 |
| 36.0000 | 10.0000 | -35.1527 | .7493  | -1.2217 |
| 36.0000 | 12.0000 | -35.2535 | .8684  | -1.4117 |
| 36.0000 | 14.0000 | -35.3649 | .9734  | -1.5775 |
| 36.0000 | 16.0000 | -35.4839 | 1.0639 | -1.7183 |
| 36.0000 | 18.0000 | -35.6080 | 1.1399 | -1.8345 |
| 36.0000 | 20.0000 | -35.7345 | 1.2019 | -1.9273 |
| 36.0000 | 22.0000 | -35.8613 | 1.2506 | -1.9982 |
| 36.0000 | 24.0000 | -35.9866 | 1.2870 | -2.0492 |
| 36.0000 | 26.0000 | -36.1089 | 1.3123 | -2.0825 |
| 36.0000 | 28.0000 | -36.2270 | 1.3278 | -2.1002 |
| 36.0000 | 30.0000 | -36.3401 | 1.3348 | -2.1046 |

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 40.0000 | 0.      | -35.3327 | -.0000 | .0000   |
| 40.0000 | 2.0000  | -35.3404 | .1271  | -.2061  |
| 40.0000 | 4.0000  | -35.3634 | .2520  | -.4085  |
| 40.0000 | 6.0000  | -35.4010 | .3728  | -.6036  |
| 40.0000 | 8.0000  | -35.4520 | .4876  | -.7883  |
| 40.0000 | 10.0000 | -35.5150 | .5948  | -.9600  |
| 40.0000 | 12.0000 | -35.5884 | .6933  | -1.1167 |
| 40.0000 | 14.0000 | -35.6703 | .7822  | -1.2569 |
| 40.0000 | 16.0000 | -35.7590 | .8609  | -1.3799 |

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 40.0000 | 18.0000 | -35.8525 | .9293  | -1.4856 |
| 40.0000 | 20.0000 | -35.9493 | .9874  | -1.5742 |
| 40.0000 | 22.0000 | -36.0475 | 1.0356 | -1.6465 |
| 40.0000 | 24.0000 | -36.1460 | 1.0744 | -1.7034 |
| 40.0000 | 26.0000 | -36.2435 | 1.1045 | -1.7463 |
| 40.0000 | 28.0000 | -36.3390 | 1.1265 | -1.7765 |
| 40.0000 | 30.0000 | -36.4317 | 1.1414 | -1.7954 |
| 44.0000 | 0.      | -35.6728 | -.0000 | .0000   |
| 44.0000 | 2.0000  | -35.6784 | .1017  | -.1634  |
| 44.0000 | 4.0000  | -35.6951 | .2020  | -.3244  |
| 44.0000 | 6.0000  | -35.7226 | .2995  | -.4805  |
| 44.0000 | 8.0000  | -35.7601 | .3929  | -.6298  |
| 44.0000 | 10.0000 | -35.8067 | .4811  | -.7702  |
| 44.0000 | 12.0000 | -35.8614 | .5633  | -.9004  |
| 44.0000 | 14.0000 | -35.9230 | .6388  | -1.0193 |
| 44.0000 | 16.0000 | -35.9902 | .7070  | -1.1260 |
| 44.0000 | 18.0000 | -36.0619 | .7678  | -1.2203 |
| 44.0000 | 20.0000 | -36.1368 | .8209  | -1.3020 |
| 44.0000 | 22.0000 | -36.2138 | .8666  | -1.3716 |
| 44.0000 | 24.0000 | -36.2919 | .9051  | -1.4294 |
| 44.0000 | 26.0000 | -36.3701 | .9367  | -1.4761 |
| 44.0000 | 28.0000 | -36.4476 | .9619  | -1.5125 |
| 44.0000 | 30.0000 | -36.5238 | .9812  | -1.5396 |
| 48.0000 | 0.      | -35.9442 | -.0000 | .0000   |
| 48.0000 | 2.0000  | -35.9484 | .0829  | -.1322  |
| 48.0000 | 4.0000  | -35.9609 | .1649  | -.2628  |
| 48.0000 | 6.0000  | -35.9815 | .2449  | -.3901  |
| 48.0000 | 8.0000  | -36.0097 | .3220  | -.5127  |
| 48.0000 | 10.0000 | -36.0449 | .3956  | -.6291  |
| 48.0000 | 12.0000 | -36.0865 | .4648  | -.7384  |
| 48.0000 | 14.0000 | -36.1337 | .5292  | -.8396  |
| 48.0000 | 16.0000 | -36.1856 | .5884  | -.9321  |
| 48.0000 | 18.0000 | -36.2413 | .6421  | -1.0155 |
| 48.0000 | 20.0000 | -36.3001 | .6901  | -1.0896 |
| 48.0000 | 22.0000 | -36.3611 | .7324  | -1.1545 |
| 48.0000 | 24.0000 | -36.4236 | .7692  | -1.2104 |
| 48.0000 | 26.0000 | -36.4868 | .8006  | -1.2576 |
| 48.0000 | 28.0000 | -36.5501 | .8269  | -1.2966 |
| 48.0000 | 30.0000 | -36.6129 | .8483  | -1.3279 |
| 52.0000 | 0.      | -36.1647 | -.0000 | .0000   |
| 52.0000 | 2.0000  | -36.1679 | .0687  | -.1089  |
| 52.0000 | 4.0000  | -36.1775 | .1367  | -.2166  |
| 52.0000 | 6.0000  | -36.1932 | .2033  | -.3221  |
| 52.0000 | 8.0000  | -36.2149 | .2679  | -.4241  |
| 52.0000 | 10.0000 | -36.2421 | .3299  | -.5219  |

|         |         |          |       |         |
|---------|---------|----------|-------|---------|
| 52.0000 | 12.0000 | -36.2743 | .3888 | -.6144  |
| 52.0000 | 14.0000 | -36.3110 | .4442 | -.7012  |
| 52.0000 | 16.0000 | -36.3517 | .4956 | -.7815  |
| 52.0000 | 18.0000 | -36.3957 | .5430 | -.8551  |
| 52.0000 | 20.0000 | -36.4424 | .5861 | -.9218  |
| 52.0000 | 22.0000 | -36.4913 | .6248 | -.9814  |
| 52.0000 | 24.0000 | -36.5417 | .6593 | -1.0341 |
| 52.0000 | 26.0000 | -36.5932 | .6894 | -1.0799 |
| 52.0000 | 28.0000 | -36.6452 | .7155 | -1.1192 |
| 52.0000 | 30.0000 | -36.6973 | .7377 | -1.1522 |

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 56.0000 | 0.      | -36.3466 | -.0000 | .0000   |
| 56.0000 | 2.0000  | -36.3491 | .0577  | -.0910  |
| 56.0000 | 4.0000  | -36.3566 | .1149  | -.1812  |
| 56.0000 | 6.0000  | -36.3689 | .1711  | -.2698  |
| 56.0000 | 8.0000  | -36.3858 | .2259  | -.3558  |
| 56.0000 | 10.0000 | -36.4071 | .2787  | -.4388  |
| 56.0000 | 12.0000 | -36.4325 | .3292  | -.5179  |
| 56.0000 | 14.0000 | -36.4615 | .3770  | -.5928  |
| 56.0000 | 16.0000 | -36.4939 | .4220  | -.6629  |
| 56.0000 | 18.0000 | -36.5291 | .4639  | -.7279  |
| 56.0000 | 20.0000 | -36.5667 | .5024  | -.7876  |
| 56.0000 | 22.0000 | -36.6063 | .5377  | -.8419  |
| 56.0000 | 24.0000 | -36.6474 | .5695  | -.8908  |
| 56.0000 | 26.0000 | -36.6896 | .5980  | -.9343  |
| 56.0000 | 28.0000 | -36.7327 | .6232  | -.9725  |
| 56.0000 | 30.0000 | -36.7761 | .6452  | -1.0056 |

|         |         |          |        |        |
|---------|---------|----------|--------|--------|
| 60.0000 | 0.      | -36.4988 | -.0000 | .0000  |
| 60.0000 | 2.0000  | -36.5008 | .0491  | -.0771 |
| 60.0000 | 4.0000  | -36.5067 | .0978  | -.1536 |
| 60.0000 | 6.0000  | -36.5164 | .1458  | -.2288 |
| 60.0000 | 8.0000  | -36.5299 | .1926  | -.3022 |
| 60.0000 | 10.0000 | -36.5469 | .2380  | -.3733 |
| 60.0000 | 12.0000 | -36.5671 | .2817  | -.4416 |
| 60.0000 | 14.0000 | -36.5904 | .3234  | -.5066 |
| 60.0000 | 16.0000 | -36.6165 | .3629  | -.5680 |
| 60.0000 | 18.0000 | -36.6450 | .3999  | -.6256 |
| 60.0000 | 20.0000 | -36.6755 | .4345  | -.6791 |
| 60.0000 | 22.0000 | -36.7079 | .4664  | -.7284 |
| 60.0000 | 24.0000 | -36.7418 | .4957  | -.7733 |
| 60.0000 | 26.0000 | -36.7767 | .5223  | -.8140 |
| 60.0000 | 28.0000 | -36.8126 | .5462  | -.8505 |
| 60.0000 | 30.0000 | -36.8490 | .5675  | -.8828 |

|         |        |          |        |        |
|---------|--------|----------|--------|--------|
| 64.0000 | 0.     | -36.6276 | -.0000 | .0000  |
| 64.0000 | 2.0000 | -36.6292 | .0422  | -.0660 |
| 64.0000 | 4.0000 | -36.6339 | .0841  | -.1316 |

|         |         |          |       |        |
|---------|---------|----------|-------|--------|
| 64.0000 | 6.0000  | -36.6418 | .1254 | -.1962 |
| 64.0000 | 8.0000  | -36.6526 | .1659 | -.2595 |
| 64.0000 | 10.0000 | -36.6663 | .2053 | -.3210 |
| 64.0000 | 12.0000 | -36.6827 | .2434 | -.3803 |
| 64.0000 | 14.0000 | -36.7017 | .2799 | -.4372 |
| 64.0000 | 16.0000 | -36.7229 | .3147 | -.4913 |
| 64.0000 | 18.0000 | -36.7462 | .3477 | -.5424 |
| 64.0000 | 20.0000 | -36.7713 | .3787 | -.5904 |
| 64.0000 | 22.0000 | -36.7981 | .4076 | -.6350 |
| 64.0000 | 24.0000 | -36.8261 | .4344 | -.6762 |
| 64.0000 | 26.0000 | -36.8553 | .4591 | -.7140 |
| 64.0000 | 28.0000 | -36.8854 | .4816 | -.7484 |
| 64.0000 | 30.0000 | -36.9161 | .5019 | -.7794 |

END.

SRU 52.309 UNTS.

RUN COMPLETE.

LNH

```

00100 PROGRAM PROJ2(INPUT,OUTPUT)
00110 IN=50
00120*****
00130*
00140*
00150*   PROGRAM FOR LINEARLY DISTRIBUTED DISC OF SOURCES
00160*           WITH THE BODY
00170*
00180*
00190*   THIS PROGRAM CALCULATES THE VELOCITY FIELD
00200*   (UX,UY,ETA) FOR A LINEARLY DISTRIBUTED DISC
00210*   OF SOURCES AND A STRONG SOURCE AT THE POSITION
00220*   R2 ON THE X-AXIS.
00230*
00240*
00250*
00260* TO RUN THIS PROGRAM:
00270*           ONE SHOULD INPUT THE FOLLOWINGS
00280*           IN RESponce TO THE ASK FOR DATA
00290*           NX,NY,DELTX,DELTY,U,A,A1,R2,K1
00300*           (ACORDING TO THE DEFINITION OF
00310*           THE PARAMETERS GIVEN BELOW:
00320*
00330*
00340*
00350*   PARAMETERS:
00360*           SEE PROJ1 PROGRAM DUCUMENTS.
00370*           ALSO SEE THE PROJECT REPORT
00380*
00390*
00400*
00410*
00420*
00430*****
00440 REAL M,K1,L
00450 READ,NX,NY,DELTX,DELTY,U,A,A1,R2,K1
00460 PRINT 210
00470 210 FORMAT(////,10X,*VELOCITY FIELD FOR LINEARLY DISTRIBUTED DISC+BODY*
00480 PRINT 220,NX,NY,DELTX,DELTY,U,A,A1,R2,K1
00490 220 FORMAT(10X,*-----*,//,
00500+*NX=*,I5,/,*NY=*,I5,/,*DELTX=*,F5.2,/,*DELTY=*,F5.2,/,*U=*,F10.4,/,
00510+*A=*,F10.4,/,*A1=*,F6.4,/,*R2=*,F10.4,/,*K1=*,F10.3,///,
00520 PRINT 200
00530 200 FORMAT(5X,*X IN FT.*,10X,*Y IN FT.*,7X,*UX IN FT/SEC.*
00531+,7X,*UY IN FT./SEC.*,5X,*ETA IN DEG.*,/,5X,*-----*,10X,
00532+*-----*,7X,*-----*,7X,*-----*,5X,
00533+*-----*)
00550 DO 500 III=1,NX
00560 DO 600 II=1,NY

```

```

00570 YY=(II-1)*DELTY
00580 XX=FLOAT(III)*DELTX
00590 R=((XX**2)+(YY**2)**0.5)
00600 6 CONTINUE
00610C          RADIAL COMPONENT OF THE DISC VELOCITY FIELD
00620 N0=0
00630 N1=1
00640 N2=2
00650 N3=3
00660 N4=4
00670 N5=5
00680 X=R/A
00690 PROD=-0.5
00700 M=3*A1*(1.+A1)*U/A
00710 T2=ATAN((YY/XX))
00720 T1=COS(T2)
00730 T3=SIN(T2)
00740 IF(R.GE.A) GO TO 5
00750 SUM1=0.
00760 DO 1 I=2,IN
00770 PROD=PROD*(-0.5-(I-1))/I
00780 SUM1=SUM1+(((4*I)+1.)/(((2*I)+3.)*((2*I)-2.))*PROD)
00790 1 CONTINUE
00800 S1=SUM1
00810 PROD1=-0.5
00820 SUM2=0.
00830 DO 2 I=2,IN
00840 N=2*I
00850 PROD1=PROD1*(-0.5-(I-1))/I
00860 SUM2=SUM2+(((I*PROD1)/(I-1))*X**((2*I)-1))*F(N,T1))
00870 2 CONTINUE
00880 S2=SUM2
00890 AK1=-M*A/2.
00900 D0=2*(S1-(76./60.))*X*F(N2,T1)
00910 D1=(9./2.)*(X**2)*F(N3,T1)
00920 D2=3.*(X**3)*F(N4,T1)
00930 D3=(5./6.)*(X**4)*F(N5,T1)
00940 D4=S2
00950 UR1=AK1*(D0+D1-D2+D3-D4)
00960C          RADIAL COMPONENT OF THE BODY VELOCITY FIELD
00970 D5=(R-(R2*T1))
00980 D6=((R*T1)-R2)**2)
00990 D7=(R**2)*(T3**2)
01000 D8=((D6+D7)**(3/2))
01010 AK=(K1/(4*3.14))
01020 UR2=AK*D5/D8
01030 UR3=-U*T1
01040C          TANGENTIAL VELOCITY OF THE FREE STREAM
01050 UT3=U*T3
01060C          TANGENTIAL VELOCITY OF THE BODY
01070 D9=R2*T3

```

```

01080 D10=((((R*TI)-R2)**2)+((R**2)*(T3**2)))**(3./2.)
01090 AK3=K1/(4.*3.14)
01100 UT2=AK3*D9/D10
0110C      TANGENTIAL VELOCITY OF THE DISC
01120 PRDD2=-0.5
01130 SUM3=0.
01140 DO 3 I=2,IN
01150 N=2*I
01160 PRDD2=PRDD2*(-0.5-(I-1))/I
01170 SUM3=SUM3+((PRDD2/((2*I)-2))*((X**N)*KAS(N,T1)))
01180 3 CONTINUE
01190 S3=SUM3
01200 AK4=(MKA/2.)**(1./X)
01210 D11=(1./2.)*KAS(N0,T1)
01220 D12=(S1-(76./60.))*((X**2)*KAS(N2,T1))
01230 D13=(3./2.)*((X**3)*KAS(N3,T1))
01240 D14=(3./4.)*((X**4)*KAS(N4,T1))
01250 D15=(1./6.)*((X**5)*KAS(N5,T1))
01260 D16=S3
01270 UT1=AK4*(D11+D12+D13-D14+D15-D16)
01280 UTR=(UT2+UT3)
01290 UTD=(UT1+UT3)
01300 UT=(UT1+UT2+UT3)
01310 UR=UR1+UR2+UR3
01320 UC=((UT**2)+(UR**2))*0.5)
01330 UX=(UR*TI)-(UT*T3)
01340 UY=(UR*T3)+(UT*TI)
01350 XY=UY/UX
01360 ET=ATAN(XY)*360./((2*3.14)
01370 URB=(UR2+UR3)/U
01380 URD=(UR1+UR3)/U
01390 PRINT 100,XX,YY,UX,UY,ET
01400 100 FORMAT(5X,F10.4,5X,F10.4,5X,F14.4,5X,F10.4)
01410 GO TO 600
01420C   FOR R GT THAN A
01430 5 X=A/R
01440C   RADIAL COMPONENTS
01450 D5=(R-(R2*TI))
01460 D6=((R*TI)-R2)**2)
01470 D7=(R**2)*(T3**2)
01480 D8=((D6+D7)**(3/2))
01490 AK=(K1/(4.*3.14))
01500 UR2=AK*D5/D8
01510 UR3=-U*TI
01520 PRDD=1.
01530 SUM1=0.
01540 DO 10 I=1,IN
01550 PRDD=PRDD*(-0.5-(I-1))/I
01560 N=2*I
01570 SUM1=SUM1+(((2*I)+1)*PRDD)/(((2*I)+3))*((X**((2*I)+2))*P(N,T1))
01580 10 CONTINUE

```



```

01590 S1=SUM1
01600 S2=(1./3.)*(X**2)*P(N0,T1)
01610 AK1=M*A/2.
01620 UR1=AK1*(S2+S1)
01630C      TANGANTIAL COMPONENTS
01640 UT3=U*T3
01650 D9=R2*T3
01660 D10=(((R*T1)-R2)**2)+((R**2)*(T3**2))**(3./2.)
01670 AK3=K1/(4.*3.14)
01680 UT2=AK3*D9/D10
01690 PROD=1.
01700 SUM=0.
01710 DO 30 I=1,IN
01720 N=I*2
01730 PROD=PROD*(-0.5-(I-1))/I
01740 SUM=SUM+(PROD/((2*I)+3))*(X**((2*I)+2))*AS(N,T1)
01750 30 CONTINUE
01760 S3=SUM
01770 S4=(1./3.)*(X**2)*AS(N0,T1)
01780 AK4=M*A/2.
01790 UT1=AK4*(S4+S3)
01800 UTR=UT2+UT3
01810 UTD=UT1+UT3
01820 UT=UT1+UT2+UT3
01830 URB=(UR2+UR3)/U
01840 URD=(UR1+UR3)/U
01850 UR=UR1+UR2+UR3
01860 UC=(((UT**2)+(UR**2))**0.5)
01870 UX=(UR*T1)-(UT*T3)
01880 UY=(UR*T3)+(UT*T1)
01890 XY=UY/UX
01900 ET=ATAN(XY)*360./(2.*3.14)
01910 PRINT 100,XX,YY,UX,UY,ET
01920 600 CONTINUE
01930 800 FORMAT(//)
01940 PRINT 800
01950 500 CONTINUE
01960 END
01970*****
01980 FUNCTION F(N,X)
01990C THIS IS A FUNCTION TO CALCULATE LEGENDER POLYNOMIALS OF THE
02000C FIRST KIND AND NTH ORDER
02010 Q0=1.
02020 F=Q0
02030 IF(N.EQ.0) GO TO 1
02040 Q1=X
02050 F=Q1
02060 IF(N.EQ.1) GO TO 1
02070 Q2=((3*(X**2))-1.)/2.
02080 F=Q2
02090 IF(N.EQ.2) GO TO 1

```

```

02100 P1=((3*(X**2))-1.)/2.
02110 P2=X
02120 I=3
02130 2 CONTINUE
02140 P=((((2.*I)-1.)/I)*X*P1)-(((I-1.)/I)*P2
02150 IF(I.EQ.N) GO TO 1
02160 I=I+1
02170 P2=P1
02180 P1=P
02190 GO TO 2
02200 1 CONTINUE
02210 RETURN
02220 END
02230C*****
02240 FUNCTION AS(N,X)
02250 DIMENSION R(1000)
02260C THIS IS A FUNCTION TO CALCULATE ASSOCIATE LEGENDER POLYNOMIAL
02270C OF THE FIRST KIND AND POWER(M).
02280 R0=0.
02290 RA=R0
02300 IF(N.EQ.0) GO TO 3
02310 R1=(1-(X**2))**0.5
02320 RA=R1
02330 IF(N.EQ.1) GO TO 3
02340 R2=3.*X*((1-(X**2))**0.5)
02350 RA=R2
02360 IF(N.EQ.2) GO TO 3
02370 P1=X
02380 P2=((3*(X**2))-1.)/2.
02390 DO 10 I=3,N
02400 R(I)=(((2*I)-1)*((1-(X**2))**0.5)*P2)+R1
02410 P=(((2.*I)-1.)/I)*X*P2-(((I-1.)/I)*P1
02420 P1=P2
02430 P2=P
02440 R1=R2
02450 R2=R(I)
02460 RA=R(I)
02470 10 CONTINUE
02480 3 CONTINUE
02490 AS=RA
02500 RETURN
02510 END
READY.

```

RUN

78/07/04. 22.37.42.

FILE PROJ2

? 8,5,4,,2,,38,,16,,0.33,5.02,647.78

-----  
 VELOCITY FIELD FOR LINEARLY DISTRIBUTED DISC+BODY  
 -----

NX= 8  
 NY= 5  
 DELTX= 4.00  
 DELTY= 4.00  
 U= 38.0000  
 A= 16.0000  
 A1= .3300  
 R2= 5.0200  
 K1= 647.780

| X IN FT. | Y IN FT. | UX IN FT./SEC. | UY IN FT./SEC. | ETA IN DEG. |
|----------|----------|----------------|----------------|-------------|
| 4.0000   | 0.       | -80.5193       | -.0000         | .0000       |
| 4.0000   | 4.0000   | -25.7175       | 4.9155         | -10.8261    |
| 4.0000   | 8.0000   | -22.2953       | 4.4470         | -11.2858    |
| 4.0000   | 12.0000  | -18.6772       | 8.8719         | -25.4212    |
| 4.0000   | 16.0000  | -30.2627       | 10.7436        | -19.5553    |
| 8.0000   | 0.       | -11.3074       | -.0000         | .0000       |
| 8.0000   | 4.0000   | -20.3965       | 5.2585         | -14.4641    |
| 8.0000   | 8.0000   | -23.9431       | 5.8885         | -13.8240    |
| 8.0000   | 12.0000  | -26.1749       | 8.7563         | -18.5060    |
| 8.0000   | 16.0000  | -31.6941       | 7.0805         | -12.5996    |
| 12.0000  | 0.       | -21.7867       | -.0000         | .0000       |
| 12.0000  | 4.0000   | -23.4000       | 2.4939         | -6.0864     |
| 12.0000  | 8.0000   | -26.2924       | 3.7251         | -8.0681     |
| 12.0000  | 12.0000  | -30.7332       | 4.8705         | -9.0097     |
| 12.0000  | 16.0000  | -32.5288       | 5.1147         | -8.9404     |

|         |         |          |        |         |
|---------|---------|----------|--------|---------|
| 16.0000 | 0.      | -27.9716 | .0000  | -.0000  |
| 16.0000 | 4.0000  | -29.4585 | 1.6599 | -3.2266 |
| 16.0000 | 8.0000  | -30.6262 | 2.8337 | -5.2889 |
| 16.0000 | 12.0000 | -31.9344 | 3.5027 | -6.2626 |
| 16.0000 | 16.0000 | -33.2036 | 3.7821 | -6.5017 |
| 20.0000 | 0.      | -31.1806 | -.0000 | .0000   |
| 20.0000 | 4.0000  | -31.4346 | 1.0854 | -1.9786 |
| 20.0000 | 8.0000  | -32.0854 | 1.9553 | -3.4891 |
| 20.0000 | 12.0000 | -32.9272 | 2.5318 | -4.3991 |
| 20.0000 | 16.0000 | -33.8059 | 2.8335 | -4.7936 |
| 24.0000 | 0.      | -32.6392 | -.0000 | .0000   |
| 24.0000 | 4.0000  | -32.7848 | .7521  | -1.3149 |
| 24.0000 | 8.0000  | -33.1794 | 1.3941 | -2.4073 |
| 24.0000 | 12.0000 | -33.7283 | 1.8646 | -3.1659 |
| 24.0000 | 16.0000 | -34.3369 | 2.1540 | -3.5914 |
| 28.0000 | 0.      | -33.6528 | -.0000 | .0000   |
| 28.0000 | 4.0000  | -33.7439 | .5442  | -.9244  |
| 28.0000 | 8.0000  | -33.9976 | 1.0263 | -1.7299 |
| 28.0000 | 12.0000 | -34.3659 | 1.4047 | -2.3418 |
| 28.0000 | 16.0000 | -34.7930 | 1.6647 | -2.7406 |
| 32.0000 | 0.      | -34.3873 | -.0000 | .0000   |
| 32.0000 | 4.0000  | -34.4476 | .4074  | -.6779  |
| 32.0000 | 8.0000  | -34.6182 | .7773  | -1.2869 |
| 32.0000 | 12.0000 | -34.8728 | 1.0821 | -1.7782 |
| 32.0000 | 16.0000 | -35.1783 | 1.3085 | -2.1313 |

END.

SRU 8.155 UNTS.

RUN COMPLETE.

LNH

```

00100 PROGRAM PROJS1(INPUT,OUTPUT)
00110 DIMENSION UF(50,100),AR(50),YOUT(100)
00120 REAL K,L,K1
00130 IN=50
00140C
00150 READ,NS,NA,NE,DELTX,DELTY,U,A,A1,R2,K1
00160*****
00170*
00180*
00190*
00200*   PROGRAM FOR UNIFORMLY DISTRIBUTED DISC OF
00210*   SOURCES AND THE BODY.
00220*
00230*
00240*   THIS PROGRAM CONSTRUCTS THE STREAM LINES OF THE
00250*   FLOW FIELD MENTIONED ABOVE, IN TE FORM OFTHE
00260*   COORDINATES OF THE POINTS ON A SPCIFIC STREAM
00270*   LINE(SPECIFIED BY ITS STARTING POINT OR BY ITS
00280*   MASS FLOW RATE AS WILL BE DESCRIBED LATER.)AS ITS
00290*   Y AND X COORDINATES; THAT IS FOR EACH Y ,THE
00300*   CORESPONDING X VALUE OF THE STREAM LINE IS DETERMINED
00310*
00320*
00330*
00340*   ! Y
00350*   !
00360*   !
00370* (NS-1)*DIY!.....\   STREAM LINE: SAI=X
00380*   !   . \-----
00390*   !   .
00400*   !   .
00410*   !   .
00420*   !   .
00430*   !   .
00440*   !   .
00450*   !   .
00460*   !   .
00470* ORIGIN +----- X
00480*           NA*DX                               NE*DX
00490*
00500*
00510*
00520* TO RUN THIS PROGRAM:
00530*           A)HAVING THE STARTING POINT OF THE
00540*           STREAM LINE; ONE SOULD INPUT THE
00550*           FOLLOWINGS IN RESPONCE TO THE ASK
00560*           FOR DATA:NS,NA,NE,DELTX,DELTY,U,A
00570*           ,A1,R2,K1,ACORDING TO THE DEFINITION

```

```

00580*           OF THE PARAMETERS GIVEN BELOW.
00590*           (NOTE: IN THE PRESENCE OF THE BODY
00600*           NA*UX SOULD BE GREATER THAN
00610*           THE STAGNATION POINT OF THE
00620*           BODY.)
00630*           B)HAVING THE MASS FLOW RATE THROUGH
00640*           THE STREAM TUBE,ONE SOULD FIRST
00650*           CHANGE THE STATEMENT REFM=SUMZ TO
00660*           REFM=MASS FLOW INTERESTED IN,(NOTE:
00670*           SINCE FLOW IS IN MINUS X DIRECTION,
00680*           THE MASS FLOW WILL HAVE A MINUS
00690*           SIGN)THEN THE FOLOWING IS INPUT
00700*           IN RESPONCE TO THE ASK FOR DATA:
00710*           NS,NA,NE,DELTX,DELTY,U,A,A1,R2,K1.
00720*           (NOTE: IN THIS CASE NS SOULD BE
00730*           SELECTED ON AN ESTIMATE BASE,THAT
00740*           IS A VALUE THAT ONE IS SUREIS GREATER
00750*           THAN THE REAL VALUE.ALSO THE SAME
00760*           RESTRICTION IS VALIED FOR NA AS
00770*           WAS POINTED OUT IN PART A )
00780*
00790*
00800*
00810*           OUTPUT OF THE PROGRAM:
00820*           VOLUMETRIC FLOW RATE THROUGH THE
00830*           STREAM TUBE,AND X-Y COORDINATE OF
00840*           THE STREAM LINE.
00850*
00860*
00870*           PARAMETERS:
00880*           SEE PROJ1 PROGRAM DUCUMENTS
00890*           SEE THE PROJECT REPORT
00900*           SEE THE ABOVE DIAGRAM,
00910*           NS: STARTING POINT OF THE STREAM LINE ON
00920*           THE Y-AXIS.
00930*           NA: STARTING POINT ON THE X-AXIS.
00940*           NE: ENDING POINT OF THE STREAM LINE ON THE
00950*           X-AXIS.
00960*
00970*
00980*-----THE PROGRAM WAS RUN FOR:
00990*           NS=12,NA=3,NE=16,DELTX=1.,
01000*           DELTY=1.,U=38.,A=16.,A1=0.33,
01010*           R2=0.,K1=647.78
01020*
01030*
01040******
01050*
01060 PRINT 201,NS,NA,NE,DELTX,DELTY,U,A,A1,R2,K1
01070 201 FORMAT(///,20X,*STREAM LINE CONSTRUCTION FOR UNIFORMLY DISTRIBUTED*
01080+/,30X,*DISC OF SOURCES AND THE BODY*,/,20X,*-----*

```

```

01090+,*-----*,//,*NS=*,IS,/*NA=*,IS,/*NE=*,IS,/*DELTX=*,
01100+F6.4,/*DELT Y=*,F6.4,/*U=*,F10.5,/*A=*,F4.1,/*A1=*,F6.4,
01110+/,*R2=*,F10.6,/*K1=*,F10.4)
01120C
01130C
01140C
01150C
01160C
01170 DO 600 III=1,NE
01180 DO 500 II=1,NS
01190 YY=(II-1)*1.
01200 XX=FLD(III)
01210 K=2.*A1*(1.+A1)*U
01220 R=((XX**2)+(YY**2))*0.5)
01230 NO=0
01240 N1=1
01250 T=ATAN((YY/XX))
01260 T1=COS(T)
01270 T2=SIN(T)
01280 TO=0.
01290 IF(R.GE.A) GO TO 5
01300C*****          DISC FOR R<A          *****
01310 AK=K/2.
01320 X=R/A
01330 SUM=0.
01340 DO 1 I=1,IN
01350 N=I-1
01360 M=N-2
01370 SUM=SUM+(N*(X**(N-1))*(P(N,TO)+P(M,TO))*P(N,T1))
01380 1 CONTINUE
01390 UR1=AK*(P(N1,T1)-SUM)
01400 SUM=0.
01410 DO 2 I=1,IN
01420 N=I-1
01430 M=N-2
01440 SUM=SUM+((X**(N-1))*(P(N,TO)+P(M,TO))*AS(N,T1))
01450 2 CONTINUE
01460 UT1=AK*(SUM-AS(N1,T1))
01470 GO TO 200
01480 5 CONTINUE
01490C*****          DISC FOR R>A          *****
01500 X=A/R
01510 SUM=0.
01520 DO 3 I=1,IN
01530 N=I-1
01540 M=N+2
01550 SUM=SUM+((N+1)*(X**M)*(P(N,TO)+P(M,TO))*P(N,T1))
01560 3 CONTINUE
01570 UR1=(K/2)*SUM
01580 SUM=0.
01590 DO 4 I=1,IN

```

```

01600 N=I-1
01610 M=N+2
01620 SUM=SUM+(((X**M)*(P(N,T0)+P(M,T0))*AS(N,T1))
01630 4 CONTINUE
01640 UT1=(K/2)*SUM
01650 200 CONTINUE
01660C***** THE BODY *****
01670 D15=(R-(R2*T1))
01680 D16=((R*T1)-R2)**2)
01690 D17=(R**2)*(T2**2)
01700 D18=((D16+D17)**(3/2))
01710 AK1=(K1/(4*3.14))
01720 UR2=AK1*D15/D18
01730 D19=R2*T2
01740 D110=((((R*T1)-R2)**2)+((R**2)*(T2**2))**3/2)
01750 AK3=K1/(4*3.14)
01760 UT2=AK3*D19/D110
01770C***** UNIFORM FLOW *****
01780 UR3=-U*T1
01790 UT3=U*T2
01800C***** SUPERPOSITION *****
01810 UR=UR1+UR2+UR3
01820 UT=UT1+UT2+UT3
01830 UX=(UR*T1)-(UT*T2)
01840 UY=(UR*T2)+(UT*T1)
01850 XY=UY/UX
01860 ET=ATAN(XY)*360./(2.*3.14)
01870C
01880C VELOCITY FIELD AND STREAM TUBE AREA STORAGE
01890C
01900 UF(III,II)=UX
01910 500 CONTINUE
01920 600 CONTINUE
01930 AT=0.
01940 DO 601 I=2,NS
01950 AR(1)=3.14*((DELTY/2)**2)
01960 IO=I-1
01970 AT=AT+AR(IO)
01980 AR(I)=(3.14*(((I-1)*DELTY)+(DELTY/2))**2))-AT
01990 601 CONTINUE
02000C REFERENCE MASS FLOW CONSTRUCTION
02010C
02020 SUMZ=0.
02030 DO 6 I=1,NS
02040 SUMZ=SUMZ+(AR(I)*UF(NA,I))
02050 6 CONTINUE
02060 REFM=SUMZ
02070 PRINT 501,REFM
02080 501 FORMAT(//,20X,*VOLUMETRIC FLOW RATE THROUGH THE STREAM TUBE=*,F14.3)
02090 PRINT 504
02100 504 FORMAT(////,20X,*X IN FT.*,12X,*Y IN FT.*,/,20X,*-----*,10X,*-----

```



```

02110+-----*)
02120C
02130C     STREAM LINES CONSTRUCTION
02140C
02150 AL=0.1
02160 DO 7 I=NA,NE
02170 FLOW=0.
02180 J=1
02190 10 CONTINUE
02200 FLOW1=FLOW
02210 FLOW=FLOW+(AR(J)*UF(I,J))
02220 ALLOW=(FLOW-REFM)
02230 IF(ABS(ALLOW).LE.AL) GO TO 8
02240 IF(ALLOW.GT.0) GO TO 9
02250 DELTY=(FLOW1/(FLOW1+FLOW))
02260 YOUT(I)=((J-2)*DELTY)+(DELTY*DELTY)
02270 GO TO 11
02280 9 J=J+1
02290 GO TO 10
02300 8 YOUT(I)=(J-1)*DELTY
02310 11 CONTINUE
02320 XXX=I*DELTX
02330 YOUT(I)=YOUT(I)+(DELTY/2)
02340 PRINT 505,XXX,YOUT(I)
02350 505 FORMAT(20X,F10.3,10X,F10.3)
02360 7 CONTINUE
02370 END
02380C
02390C
02400C
02410C
02420C
02430C
02440C
02450C
02460C
02470C
02480C
02490C
02500C
02510*****
02520 FUNCTION P(N,X)
02530C
02540C
02550C THIS IS A FUNCTION TO CALCULATE LEGENDER POLYNOMIALS OF THE
02560C FIRST KIND
02570C
02580C
02590 Q0=0.
02600 P=Q0
02610 IF(N.LT.0) GO TO 1

```

```

02620 Q0=1.
02630 P=Q0
02640 IF(N.EQ.0) GO TO 1
02650 Q1=X
02660 P=Q1
02670 IF(N.EQ.1) GO TO 1
02680 Q2=((3*(X**2))-1.)/2.
02690 P=Q2
02700 IF(N.EQ.2) GO TO 1
02710 P1=((3*(X**2))-1.)/2.
02720 P2=X
02730 I=3
02740 2 CONTINUE
02750 F=((((2.*I)-1.)/I)*X*P1)-(((I-1.)/I)*P2
02760 IF(I.EQ.N) GO TO 1
02770 I=I+1
02780 P2=P1
02790 P1=P
02800 GO TO 2
02810 1 CONTINUE
02820 RETURN
02830 END
02840 FUNCTION AS(N,X)
02850C*****
02860 DIMENSION R(1000)
02870C
02880C
02890C THIS FUNCTION CALCULATES ASSOCIATED LEGENDER POLYNOMIALS
02900C OF THE FIRST KIND AND FIRST POWER
02910C
02920C
02930 R0=0.
02940 RA=R0
02950 IF(N.LE.0) GO TO 3
02960 R1=(1-(X**2))*0.5
02970 RA=R1
02980 IF(N.EQ.1) GO TO 3
02990 R2=3.*X*((1-(X**2))*0.5)
03000 RA=R2
03010 IF(N.EQ.2) GO TO 3
03020 P1=X
03030 P2=((3*(X**2))-1.)/2.
03040 DO 10 I=3,N
03050 R(I)=(((2*I)-1)*((1-(X**2))*0.5)*P2)+R1
03060 P=(((2.*I)-1.)/I)*X*P2-((I-1.)/I)*P1
03070 P1=P2
03080 P2=P
03090 R1=R2
03100 R2=R(I)
03110 RA=R(I)
03120 10 CONTINUE

```

03130 3 CONTINUE  
03140 AS=RA  
03150 RETURN  
03160 END  
READY.

RUN

78/07/04. 22.52.15.  
FILE PROJS1

? 14,7,20,1.,1.,38.,16.,0.33,5.02,647.78

STREAM LINE CONSTRUCTION FOR UNIFORMLY DISTRIBUTED  
DISC OF SOURCES AND THE BODY

---

NS= 14  
NA= 7  
NE= 20  
DELTX=1.0000  
DELY=1.0000  
U= 38.00000  
A=16.0  
A1= .3300  
R2= 5.020000  
K1= 647.7800

\*TIME LIMIT\*

SRU 25.287 UNTS.

T,2000

VOLUMETRIC FLOW RATE THROUGH THE STREAM TUBE= -15675.736

| X IN FT. | Y IN FT. |
|----------|----------|
| 7.000    | 13.500   |
| 8.000    | 12.957   |
| 9.000    | 12.957   |
| 10.000   | 12.957   |
| 11.000   | 12.958   |
| 12.000   | 12.958   |
| 13.000   | 12.959   |
| 14.000   | 12.959   |
| 15.000   | 12.959   |
| 16.000   | 12.960   |

|        |        |
|--------|--------|
| 17.000 | 12.960 |
| 18.000 | 12.960 |
| 19.000 | 12.960 |
| 20.000 | 11.957 |

END.

SRU 57.256 UNTS.

RUN COMPLETE.

## APPENDIX 2D

## LIST

78/07/04. 22.56.16.

FILE PROJS2

```

00100 PROGRAM PROJS2(INPUT,OUTPUT)
00110 DIMENSION UF(50,100),AR(50),YOUT(100)
00120 IN=50
00130C*****
00131C
00132C
00133C   PROGRAM FOR LINEARLY DISTRIBUTED DISC OF
00134C   SOURCES AND THE BODY.
00135C
00136C
00137C
00138C           THIS PROGRAM CONSTRUCTS THE STREAM LINES
00139C           FOR THE LINEARLY DISTRIBUTED DISC OF SOURCES
00140C           IN COMPLETE ACCORDANCE TO PROGRAM PROJS1.
00141C
00142C
00143C*****
00200 REAL K1,M,K,L
00220C
00230 READ,NS,NA,NE,DELTX,DELTY,U,A,A1,R2,K1
00231 PRINT 201,NS,NA,NE,DELTX,DELTY,U,A,A1,R2,K1
00232 201 FORMAT(///,20X,*STREAM LINE CONSTRUCTION FOR LINEARLY DISTRIBUTED*
00233+,/,30X,*DISC OF SOURCES AND THE BODY*,/,20X,*-----*,
00234+,*-----*,///,*NS=*,I5,/,*NA=*,I5,/,*NE=*,I5,/,*DELTX=*,
00235+F6.4,/,*DELTY=*,F6.4,/,*U=*,F10.5,/,*A=*,F4.1,/,*A1=*,F6.4,
00236+/,*R2=*,F10.6,/,*K1=*,F10.4)
00240 DO 500 III=1,NE
00250 DO 600 II=1,NS
00260 YY=(II-1)*1.
00290 XX=FLOAT(III)
00340 R=(((XX**2)+(YY**2))**0.5)
00360 6 CONTINUE
00370C           RADIAL COMPONENT OF THE DISC VELOCITY FIELD
00380 N0=0
00390 N1=1
00400 N2=2
00410 N3=3
00420 N4=4
00430 N5=5
00440 X=R/A
00450 PROD=-0.5
00460 M=3*A1*(1.+A1)*U/A
00470 T2=ATAN((YY/XX))
00480 T1=COS(T2)
00490 T3=SIN(T2)

```

```

00500 IF(R.GE.A) GO TO 5
00510 SUM1=0.
00520 DO 1 I=2,IN
00530 PROD=PROD*(-0.5-(I-1))/I
00540 SUM1=SUM1+((((4*I)+1.)/(((2*I)+3.)*((2*I)-2.)))*PROD)
00550 1 CONTINUE
00560 S1=SUM1
00570 PROD1=-0.5
00580 SUM2=0.
00590 DO 2 I=2,IN
00600 N=2*I
00610 PROD1=PROD1*(-0.5-(I-1))/I
00620 SUM2=SUM2+(((I*PROD1)/(I-1))*(X**((2*I)-1))*P(N,T1))
00630 2 CONTINUE
00640 S2=SUM2.
00650 AK1=-M*A/2.
00660 D0=2*(S1-(76./60.))*X*P(N2,T1)
00670 D1=(9./2.)*(X**2)*P(N3,T1)
00680 D2=3.*(X**3)*P(N4,T1)
00690 D3=(5./6.)*(X**4)*P(N5,T1)
00700 D4=S2
00710 UR1=AK1*(D0+D1-D2+D3-D4)
00720C          RADIAL COMPONENT OF THE BODY VELOCITY FIELD
00730 D5=(R-(R2*T1))
00740 D6=(((R*T1)-R2)**2)
00750 D7=(R**2)*(T3**2)
00760 D8=((D6+D7)**(3/2))
00770 AK=(K/(4*3.14))
00780 UR2=AK*D5/D8
00790 UR3=-U*T1
00800C          TANGENTIAL VELOCITY OF THE FREE STREAM
00810 UT3=U*T3
00820C          TANGENTIAL VELOCITY OF THE BODY
00830 D9=R2*T3
00840 D10=(((R*T1)-R2)**2)+((R**2)*(T3**2))**((3./2.))
00850 AK3=K/(4.*3.14)
00860 UT2=AK3*D9/D10
00870C          TANGENTIAL VELOCITY OF THE DISC
00880 PROD2=-0.5
00890 SUM3=0.
00900 DO 3 I=2,IN
00910 N=2*I
00920 PROD2=PROD2*(-0.5-(I-1))/I
00930 SUM3=SUM3+((PROD2/((2*I)-2))*(X**N)*AS(N,T1))
00940 3 CONTINUE
00950 S3=SUM3
00960 AK4=(M*A/2.)*(1./X)
00970 D11=(1./2.)*AS(N0,T1)
00980 D12=(S1-(76./60.))*X**2)*AS(N2,T1)
00990 D13=(3./2.)*(X**3)*AS(N3,T1)
01000 D14=(3./4.)*(X**4)*AS(N4,T1)

```

```

01010 D15=(1./6.)*(X**5)*AS(N5,T1)
01020 D16=S3
01030 UT1=AK4*(D11+D12+D13-D14+D15-D16)
01040 UTB=(UT2+UT3)
01050 UTD=(UT1+UT3)
01060 UT=(UT1+UT2+UT3)
01070 UR=UR1+UR2+UR3
01080 UC=((UT**2)+(UR**2)**0.5)
01090 UX=(UR*T1)-(UT*T3)
01100 UY=(UR*T3)+(UT*T1)
01110 XY=UY/UX
01120 ET=ATAN(XY)*360./(2*3.14)
01130 URB=(UR2+UR3)/U
01140 URD=(UR1+UR3)/U
01150 GO TO 603
01160C   FOR R GT THAN A
01170 S X=A/R
01180C   RADIAL COMPONENTS
01190 D5=(R-(R2*T1))
01200 D6=((R*T1)-R2)**2)
01210 D7=(R**2)*(T3**2)
01220 D8=((D6+D7)**(3/2))
01230 AK=(K/(4.*3.14))
01240 UR2=AK*D5/D8
01250 UR3=-U*T1
01260 PROD=1.
01270 SUM1=0.
01280 DO 10 I=1,IN
01290 PROD=PROD*(-0.5-(I-1))/I
01300 N=2*I
01310 SUM1=SUM1+((((2*I)+1)*PROD)/((2*I)+3))*(X**((2*I)+2))*P(N,T1)
01320 10 CONTINUE
01330 S1=SUM1
01340 S2=(1./3.)*(X**2)*P(N0,T1)
01350 AK1=M*A/2.
01360 UR1=AK1*(S2+S1)
01370C   TANGENTIAL COMPONENTS
01380 UT3=U*T3
01390 D9=R2*T3
01400 D10=((((R*T1)-R2)**2)+((R**2)*(T3**2))**3./2.)
01410 AK3=K/(4.*3.14)
01420 UT2=AK3*D9/D10
01430 PROD=1.
01440 SUM=0.
01450 DO 30 I=1,IN
01460 N=I*2
01470 PROD=PROD*(-0.5-(I-1))/I
01480 SUM=SUM+(PROD/((2*I)+3))*(X**((2*I)+2))*AS(N,T1)
01490 30 CONTINUE
01500 S3=SUM
01510 S4=(1./3.)*(X**2)*AS(N0,T1)

```



```

01520 AK4=M*A/2.
01530 UT1=AK4*(S4+S3)
01540 UTB=UT2+UT3
01550 UTD=UT1+UT3
01560 UT=UT1+UT2+UT3
01570 URB=(UR2+UR3)/U
01580 URD=(UR1+UR3)/U
01590 UR=UR1+UR2+UR3
01600 UC=((UT**2)+(UR**2))*0.5)
01610 UX=(UR*T1)-(UT*T3)
01620 UY=(UR*T3)+(UT*T1)
01630 XY=UY/UX
01640 ET=ATAN(XY)*360./(2.*3.14)
01650 603 CONTINUE
01660C VELOCITY FIELD AND STREAM TUBE AREA STORAGE
01670 UF(III,II)=UX
01680 600 CONTINUE
01690 500 CONTINUE
01691 AT=0.
01700 DO 601 I=2,NS
01710 AR(I)=3.14*(((DELTY/2)**2))
01720 IO=I-1
01721 AT=AT+AR(IO)
01730 AR(I)=(3.14*(((I-1)*DELTY)+(DELTY/2)**2))-AT
01740 601 CONTINUE
01750C
01760C REFERENCE MASS FLOW RATE
01770 SUMZ=0.
01780 DO 65 I=1,NS
01790 SUMZ=SUMZ+(AR(I)*UF(NA,I))
01800 65 CONTINUE
01810 REFM=SUMZ
01820 PRINT 562,REFM
01821 562 FORMAT(30X,*REFERENCE FLOW RATE=*,F10.4)
01830C
01840C STREAM LINE CONSTRUCTION
01850C
01860 PRINT 19
01870 19 FORMAT(/,5X,2X,*X IN FT.*,6X,2X,*Y IN FT.*,/,5X,*-----*,
01871+5X,*-----*)
01880 AL=0.1
01890 DO 71 I=NA,NE
01900 FLOW=0.
01910 J=1
01920 1011 CONTINUE
01930 FLOW1=FLOW
01940 FLOW=FLOW+(AR(J)*UF(I,J))
01950 ALLOW=(FLOW-REFM)
01960 IF(ABS(ALLOW).LE.AL) GO TO 8
01970 IF(ALLOW.GT.0) GO TO 9
01980 DELTY=(FLOW1/(FLOW1+FLOW))

```

```

01990 YOUT(I)=((J-2)*DELTY)+(DELT1Y*DELTY)
02000 GO TO 11
02010 9 J=J+1
02020 GO TO 1011
02030 8 YOUT(I)=(J-1)*DELTY
02040 11 CONTINUE
02050 YOUT(I)=YOUT(I)+(DELTY/2)
02060 XXX=I*DELTX
02070 PRINT 73,XXX,YOUT(I)
02071 73 FORMAT(5X,F10.3,5X,F10.3)
02080 71 CONTINUE
02090 END
02100C
02110C
02120C
02130C*****
02140 FUNCTION F(N,X)
02150C THIS IS A FUNCTION TO CALCULATE LEGENDER POLYNOMIALS OF THE
02160C FIRST KIND AND NTH ORDER
02170 Q0=1.
02180 F=Q0
02190 IF(N.EQ.0) GO TO 1
02200 Q1=X
02210 F=Q1
02220 IF(N.EQ.1) GO TO 1
02230 Q2=((3*(X**2))-1.)/2.
02240 F=Q2
02250 IF(N.EQ.2) GO TO 1
02260 F1=((3*(X**2))-1.)/2.
02270 F2=X
02280 I=3
02290 2 CONTINUE
02300 F=(((2.*I)-1.)/I)*X*F1-((I-1.)/I)*F2
02310 IF(I.EQ.N) GO TO 1
02320 I=I+1
02330 F2=F1
02340 F1=F
02350 GO TO 2
02360 1 CONTINUE
02370 RETURN
02380 END
02390C*****
02400 FUNCTION AS(N,X)
02410 DIMENSION R(1000)
02420C THIS IS A FUNCTION TO CALCULATE ASSOCIATE LEGENDER POLYNOMIAL
02430C OF THE FIRST KIND AND POWER(M).
02440 R0=0.
02450 RA=R0
02460 IF(N.EQ.0) GO TO 3
02470 R1=(1-(X**2))**0.5
02480 RA=R1

```

```
02490 IF(N.EQ.1) GO TO 3
02500 R2=3.*X*((1-(X**2))**0.5)
02510 RA=R2
02520 IF(N.EQ.2) GO TO 3
02530 P1=X
02540 P2=((3*(X**2))-1.)/2.
02550 DO 10 I=3,N
02560 R(I)=(((2*I)-1)*((1-(X**2))**0.5)*P2)+R1
02570 P=(((2.*I)-1.)/I)*X*P2-((I-1.)/I)*P1
02580 P1=P2
02590 P2=P
02600 R1=R2
02610 R2=R(I)
02620 RA=R(I)
02630 10 CONTINUE
02640 3 CONTINUE
02650 AS=RA
02660 RETURN
02670 END
READY.
```

RUN

78/07/04. 23.01.56.  
FILE PROJ2

? 14,7,20,1.,1.,38.,16.,0.33,5.02,647.78

STREAM LINE CONSTRUCTION FOR LINEARLY DISTRIBUTED  
DISC OF SOURCES AND THE BODY

---

NS= 14  
NA= 7  
NE= 20  
DELTX=1.0000  
DELY=1.0000  
U= 38.00000  
A=16.0  
A1= .3300  
R2= 5.020000  
K1= 647.7800

\*TIME LIMIT\*

SRU 24.620 UNTS.

T,5000

REFERENCE FLOW RATE=\*5390.1555

| X IN FT. | Y IN FT. |
|----------|----------|
| 7.000    | 13.500   |
| 8.000    | 12.960   |
| 9.000    | 12.959   |
| 10.000   | 12.958   |
| 11.000   | 12.958   |
| 12.000   | 12.959   |
| 13.000   | 11.955   |
| 14.000   | 11.956   |
| 15.000   | 11.957   |
| 16.000   | 11.958   |

|        |        |
|--------|--------|
| 17.000 | 11.958 |
| 18.000 | 11.958 |
| 19.000 | 11.958 |
| 20.000 | 11.958 |

END.

SRU 46.583 UNTS.

RUN COMPLETE.

## APPENDIX 3

## Future Work Recommendations

A more realistic case can be considered in which the strength of the sources on the disc, which falls inside the body of the windmill is put to be zero.

Fig. (22) shows the profile of the source density  $\hat{k}$ , in the case of linearly distributed source as calculated in this work. Fig. (23) shows the improvement which should be made. This improvement, although physically more realistic, is very cumbersome mathematically.

The potential for the improved model on the x-axis can be found as was described in Chapter II. It can be easily shown that:

$$\Phi = \frac{m}{2} \left[ \left( \frac{A}{2} - b \right) \sqrt{A^2 + R^2} - \frac{R^2}{2} \ln \frac{A + \sqrt{A^2 + R^2}}{b^2 + \sqrt{b^2 + R^2}} + \frac{b}{2} \sqrt{b^2 + R^2} \right]$$

Using Lagallay's theorem and the one dimensional momentum theory it can be shown that

$$m = \frac{3A^2 a(1+a)U}{(A^3 - b^3)}$$

Then what remains to be done is the expansion in zonal harmonics

FIGURES

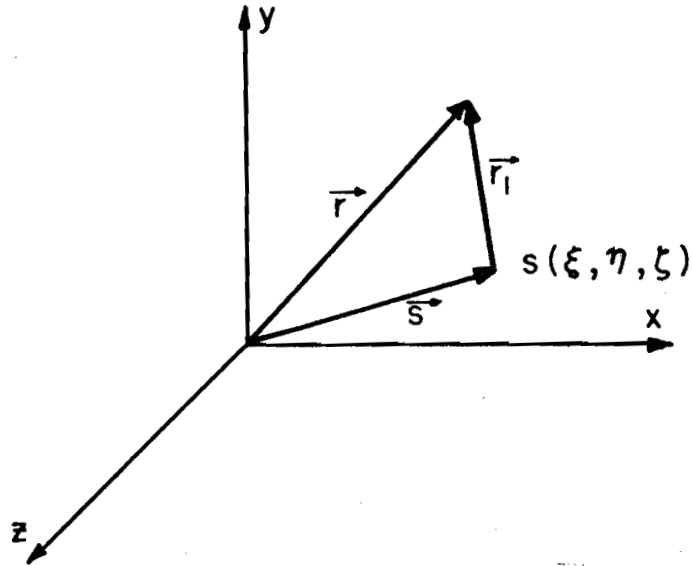


FIGURE 1.

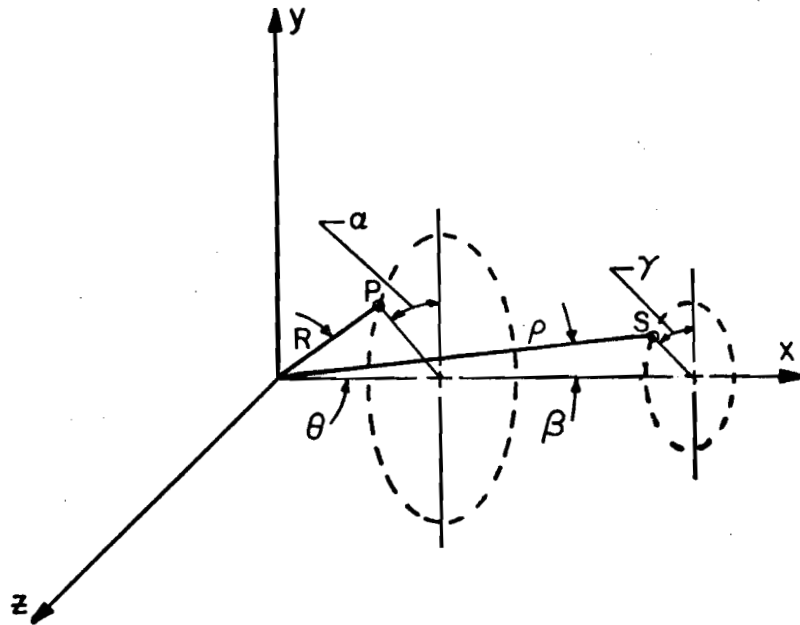


FIGURE 2.



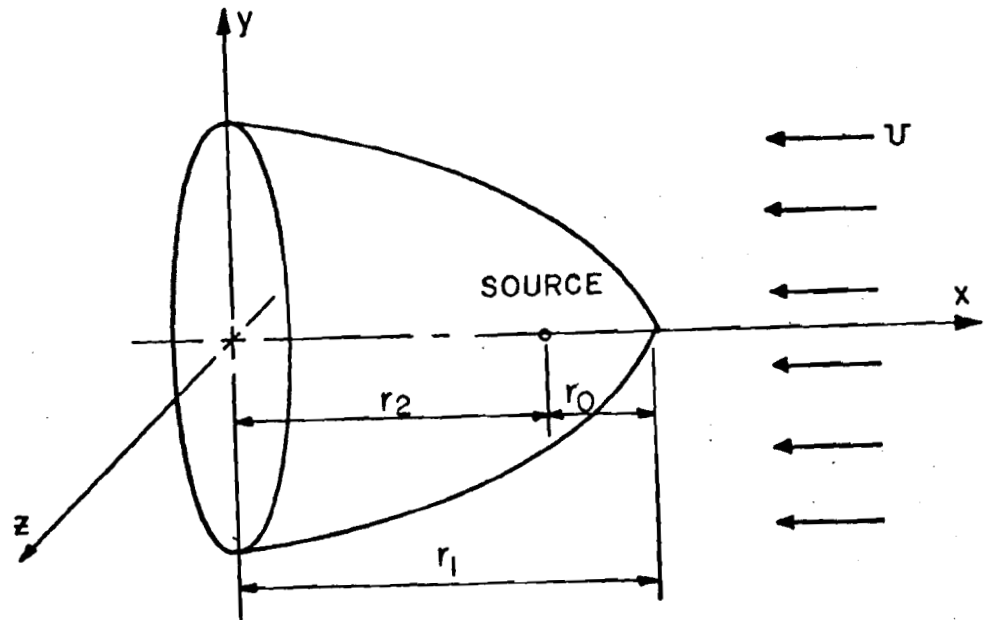


FIGURE 3.

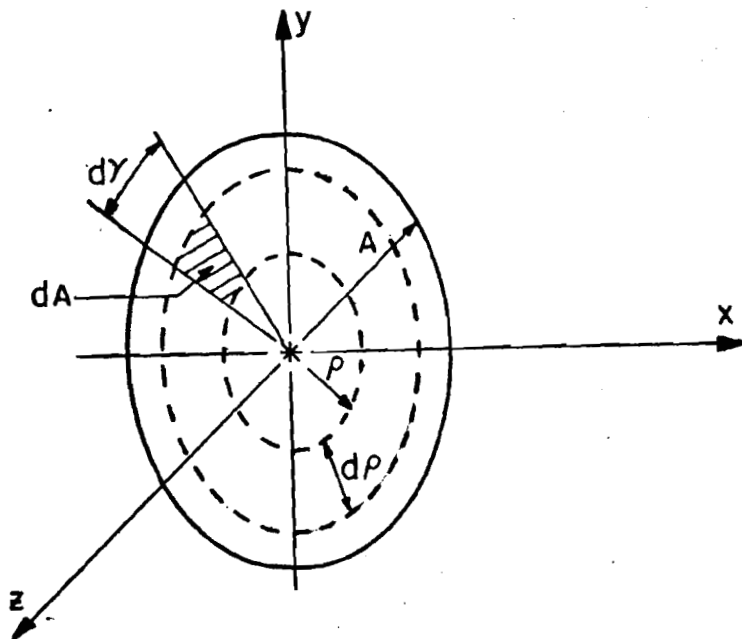


FIGURE 4.

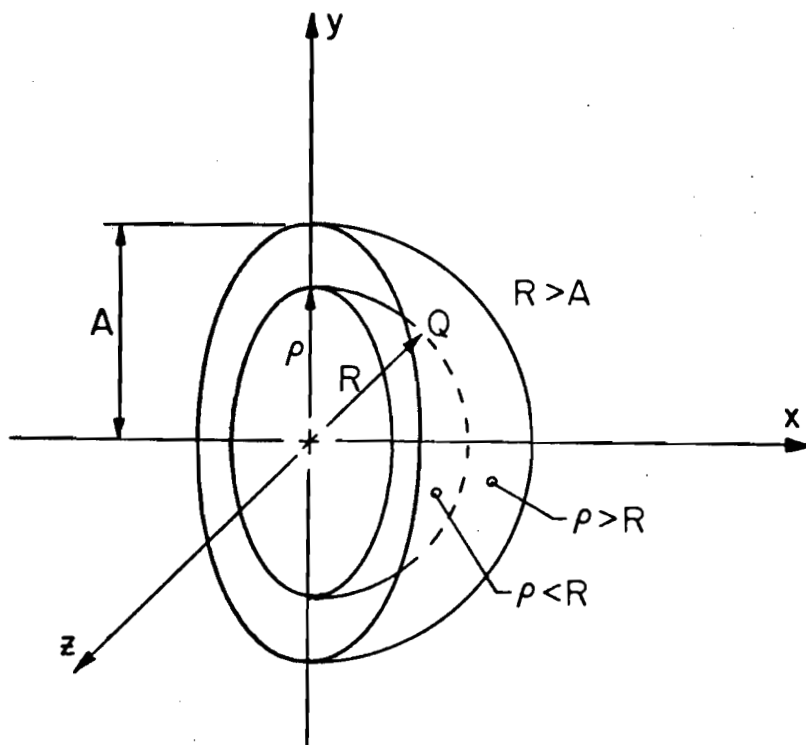


FIGURE 5.

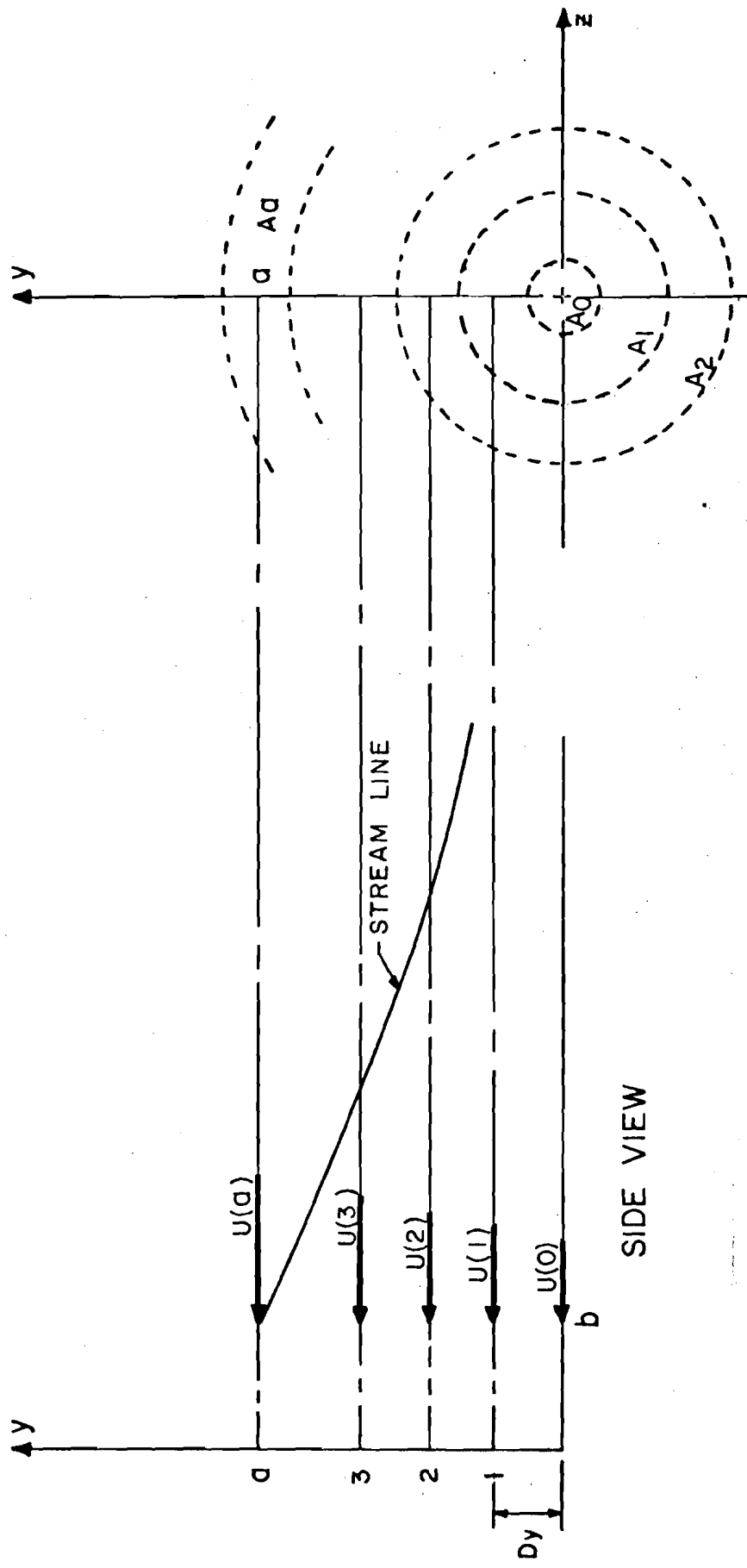


FIGURE 6.

FIGURE 7.

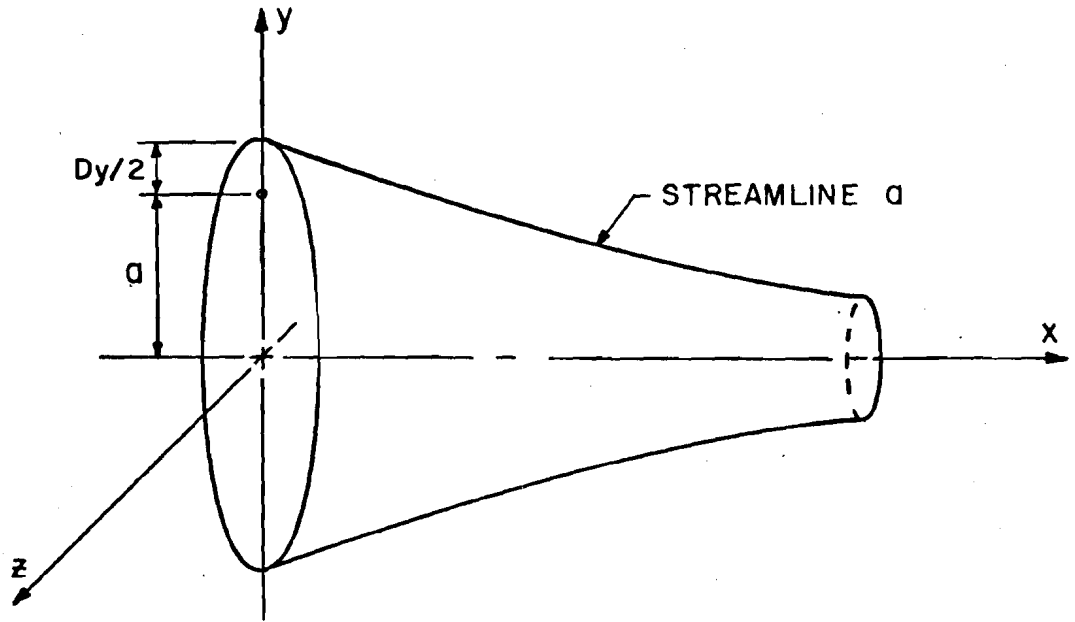


FIGURE 8.

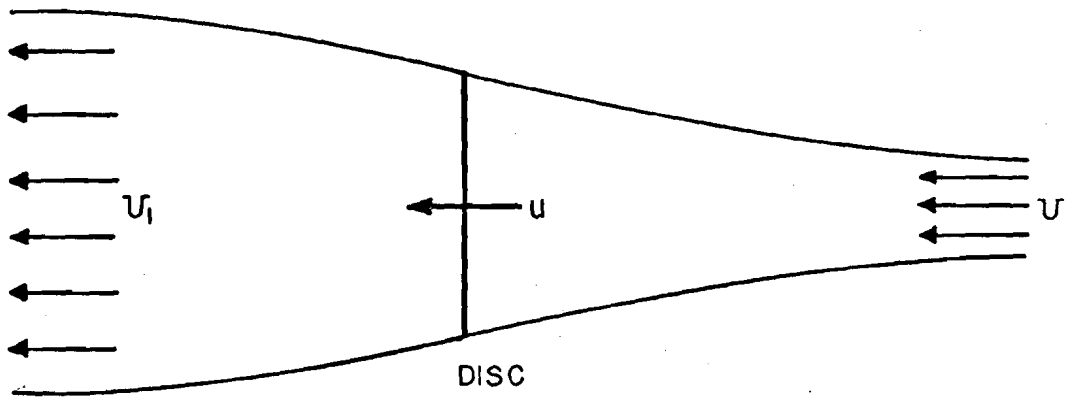


FIGURE 9.

ENERGY ALTERNATIVE PROGRAM  
UNIVERSITY of MASSACHUSETTS  
ON-SITE ARRANGEMENT DRAWING  
NEW ENGLAND WIND FURNACE PROJECT  
Drawing No. 01-1000-1-1  
Scale: 1/4" = 1'-0"  
Date: 1/24/81

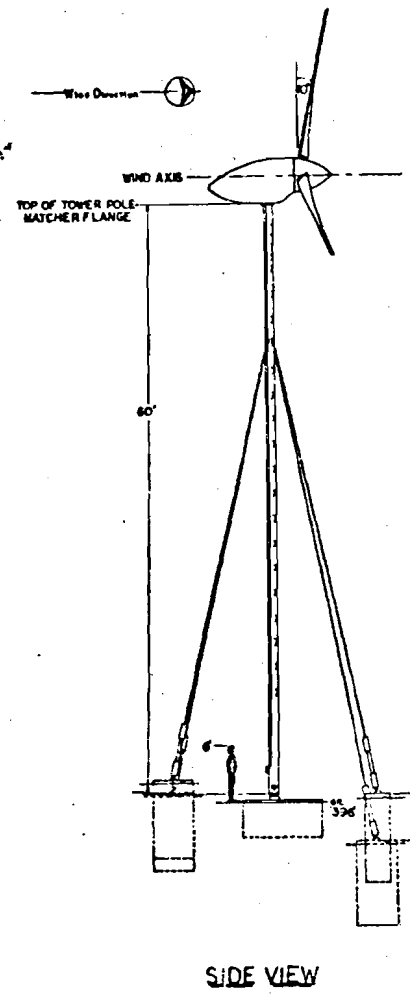
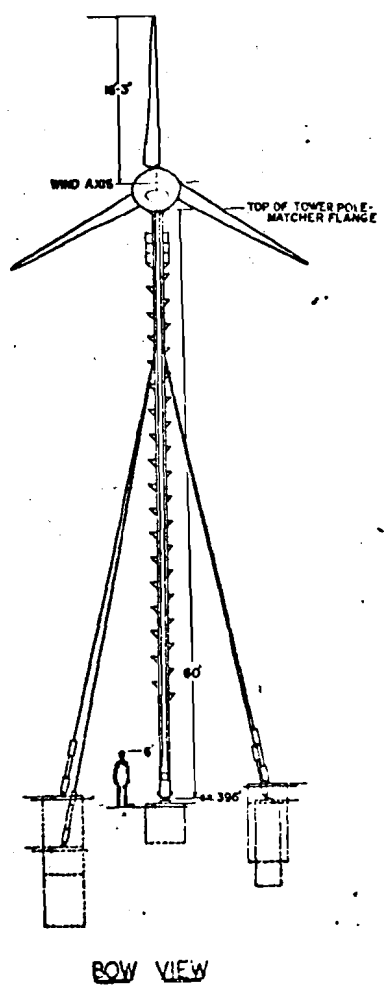
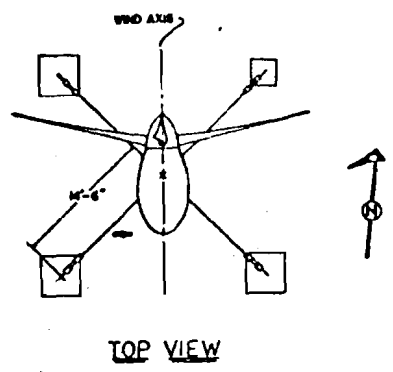


FIGURE 10.

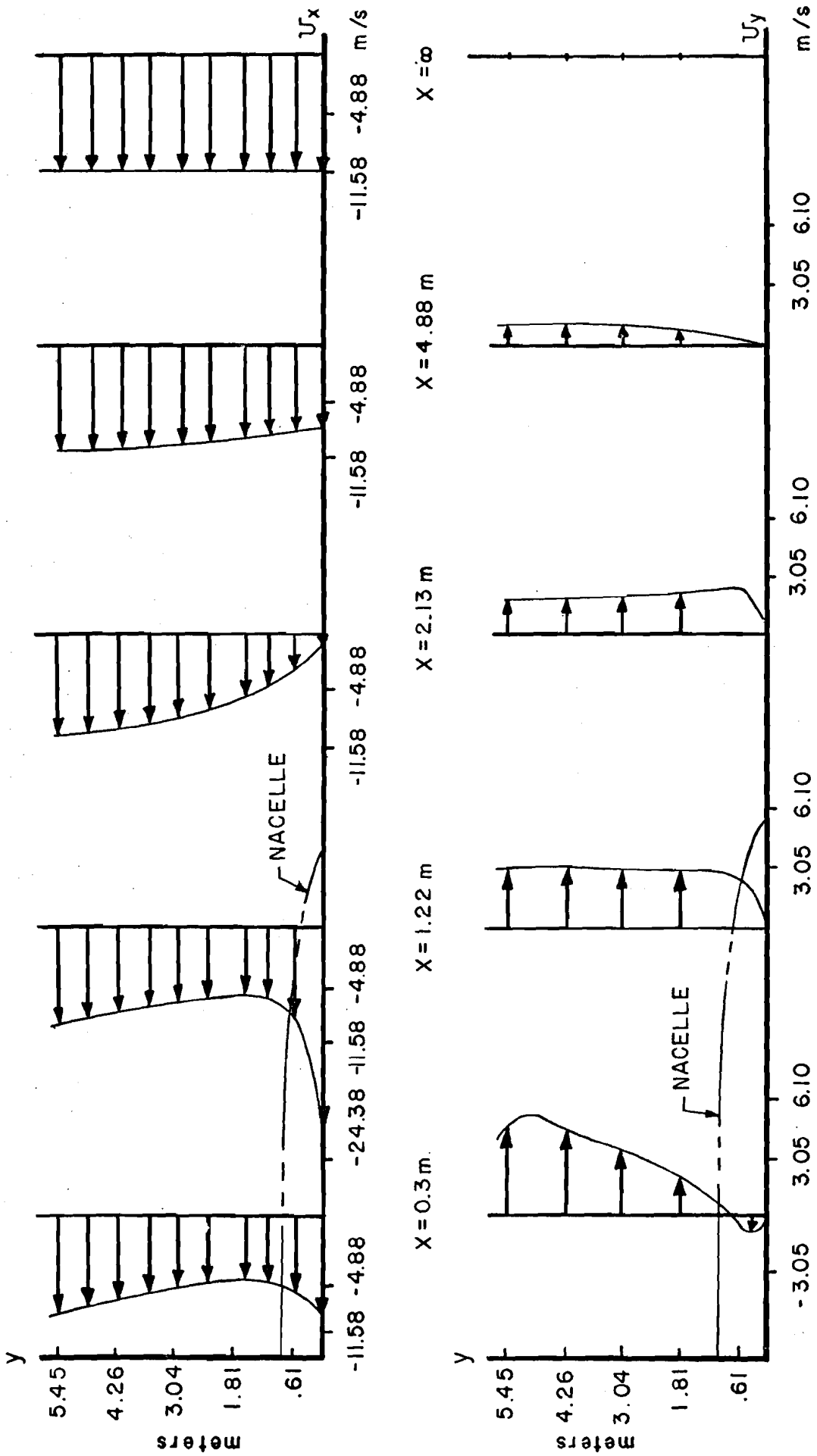


FIGURE 11.

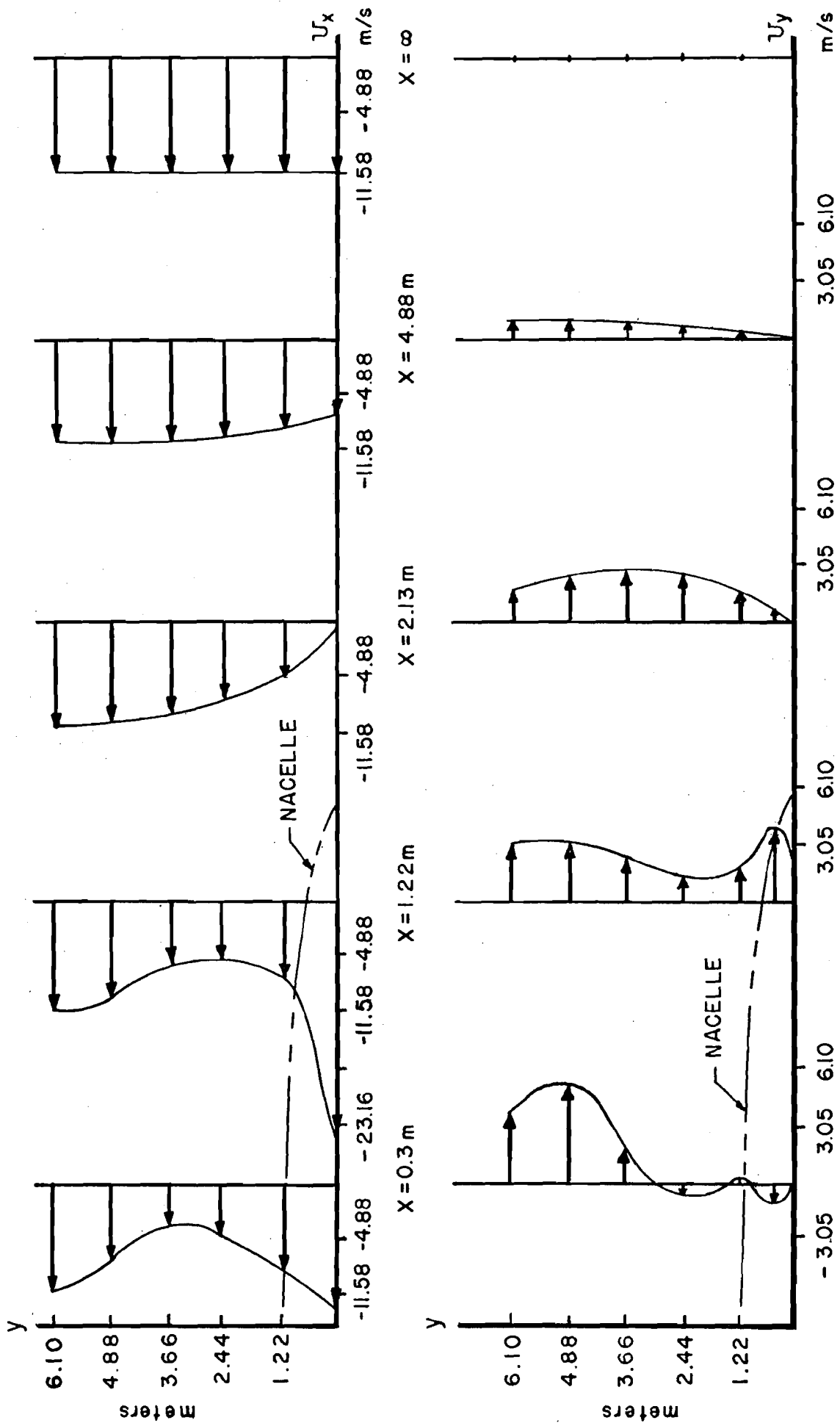


FIGURE 12.

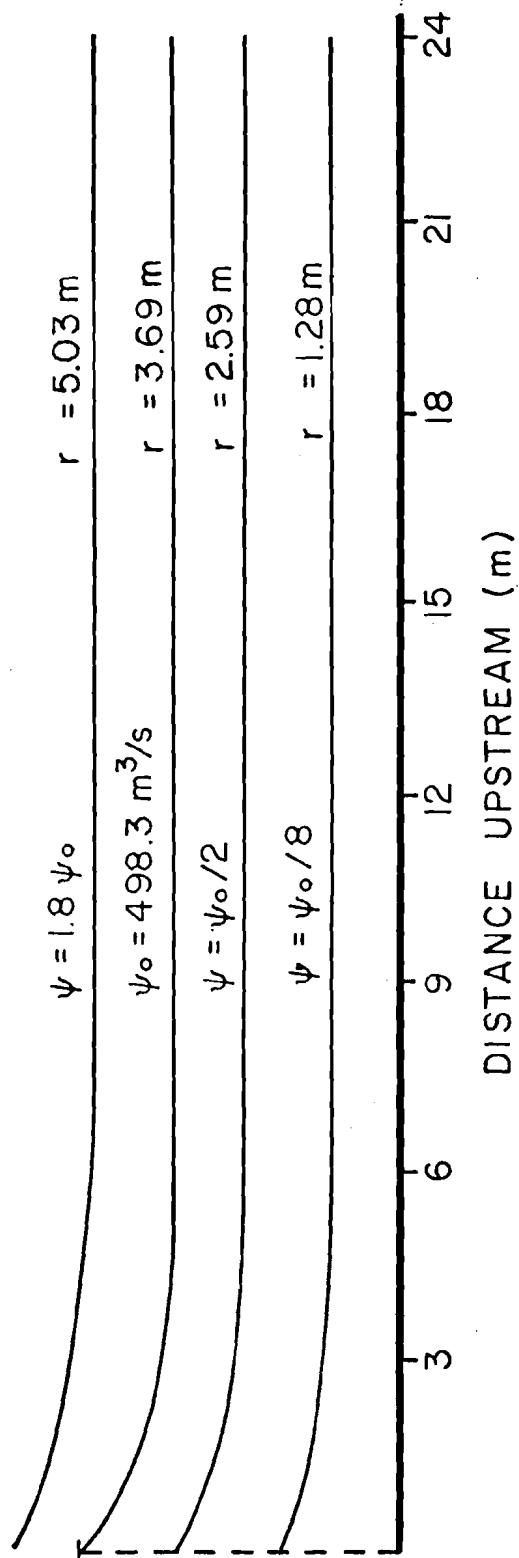


FIGURE 13.



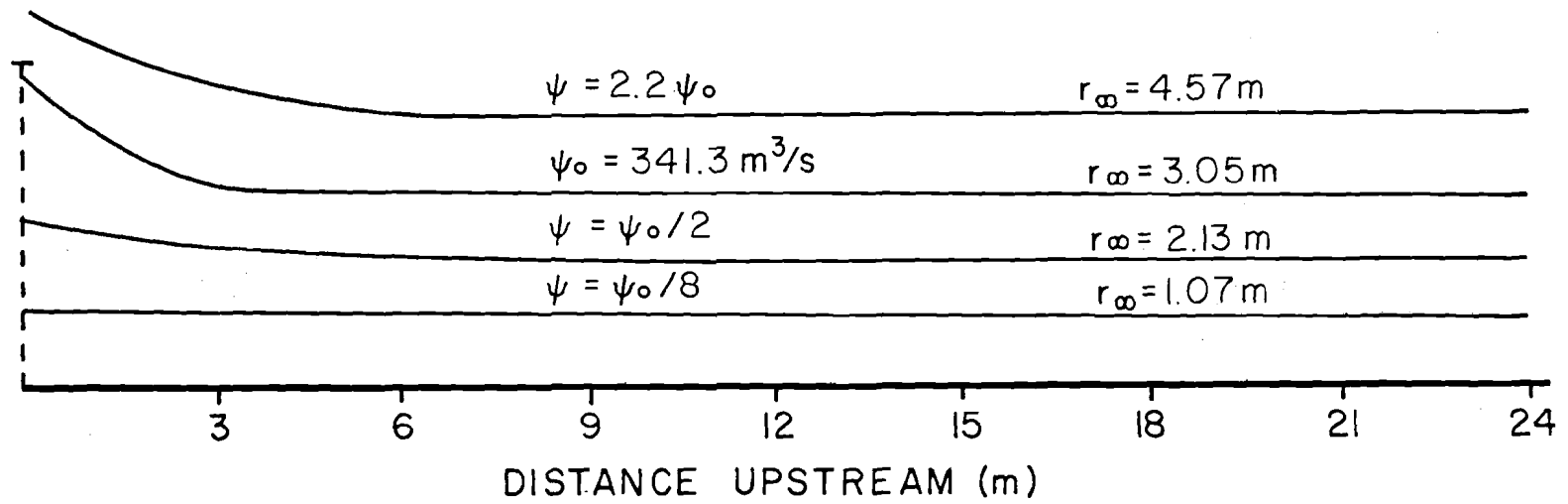


FIGURE 14.

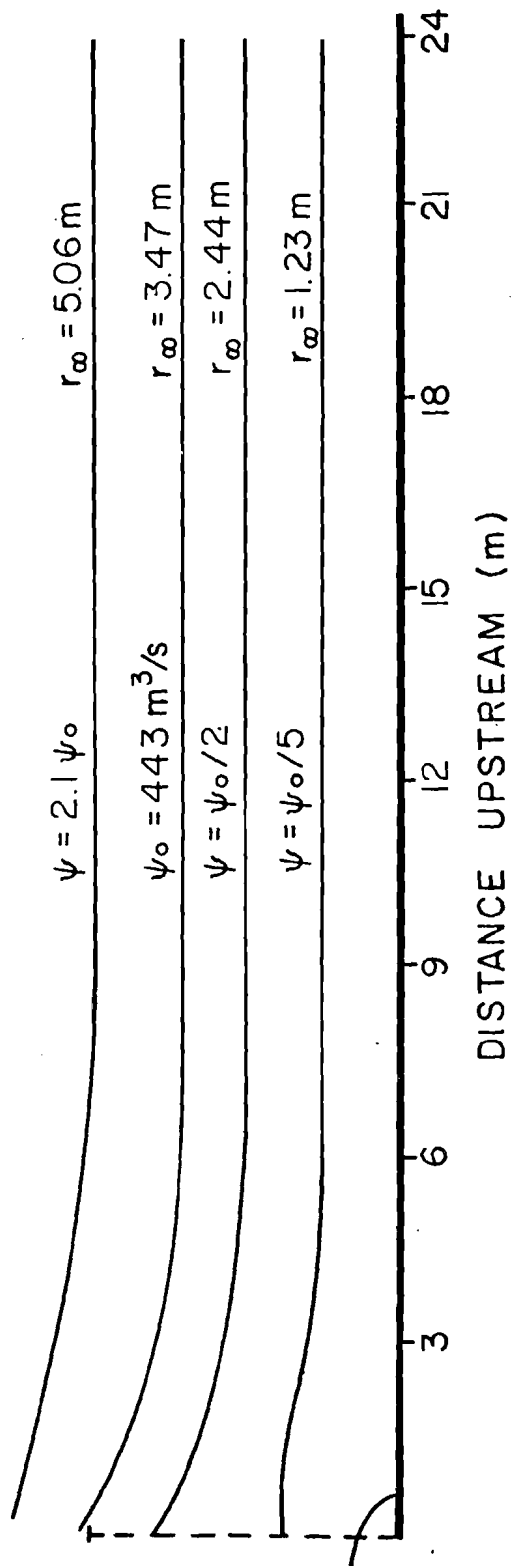


FIGURE 15.

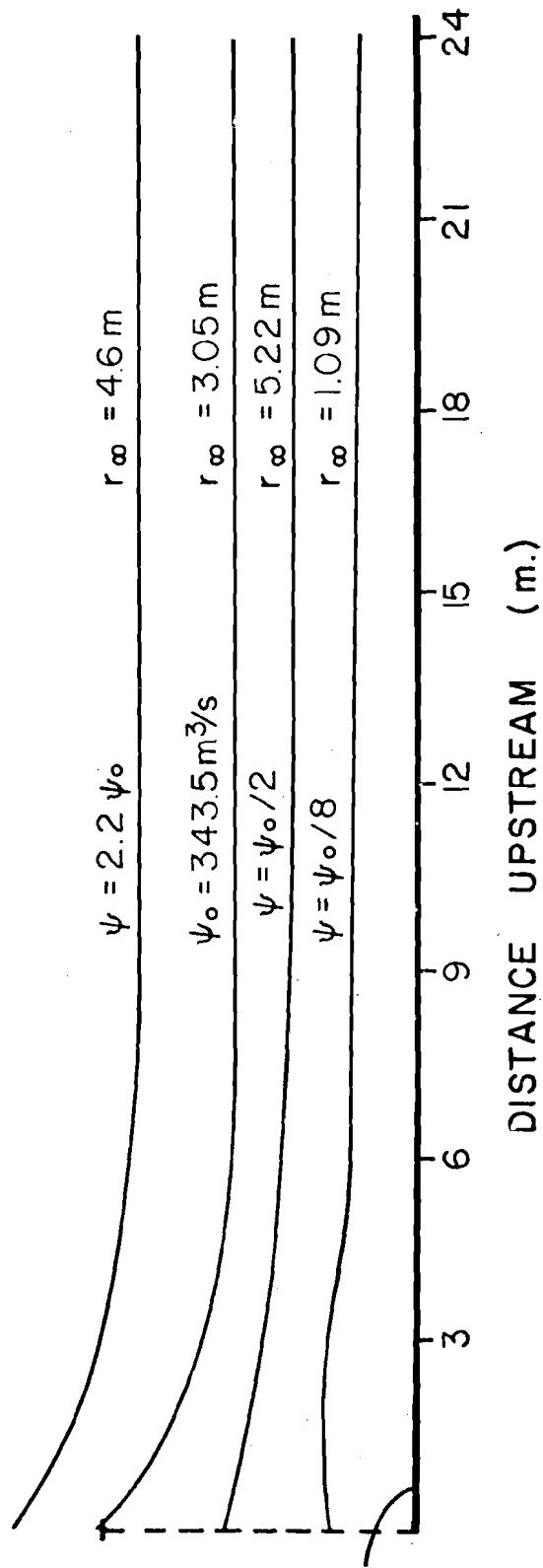


FIGURE 16.

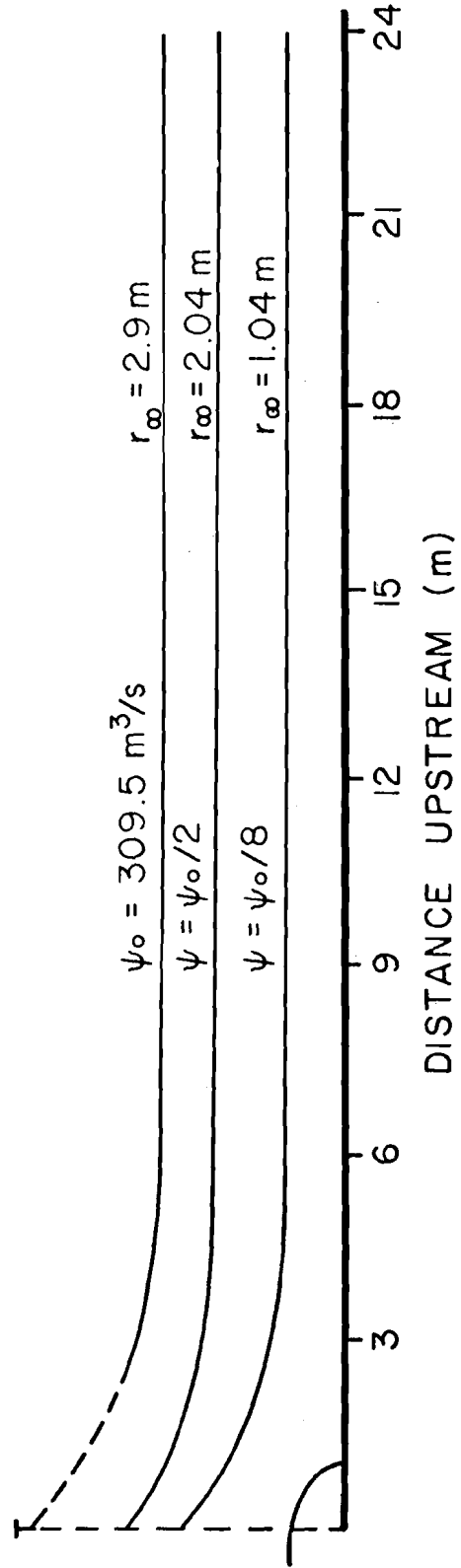


FIGURE 17.

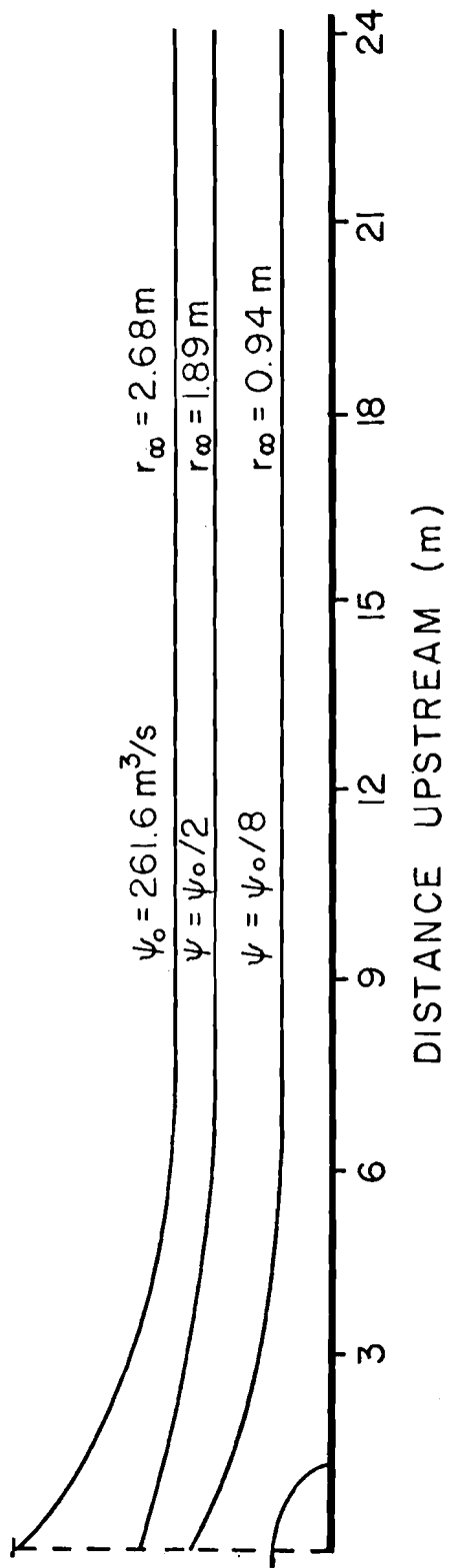


FIGURE 18.

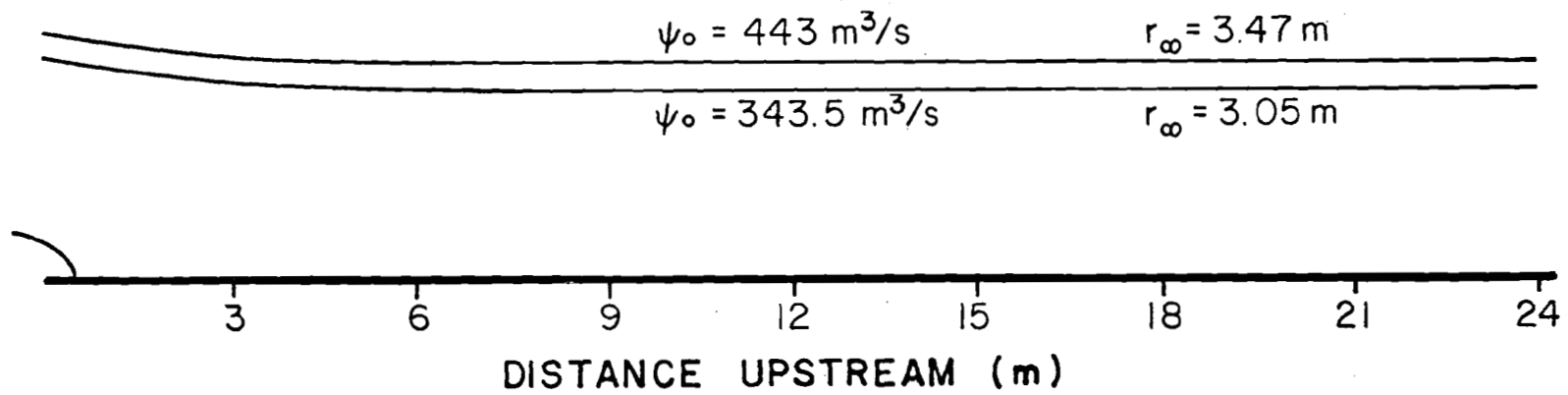


FIGURE 19.

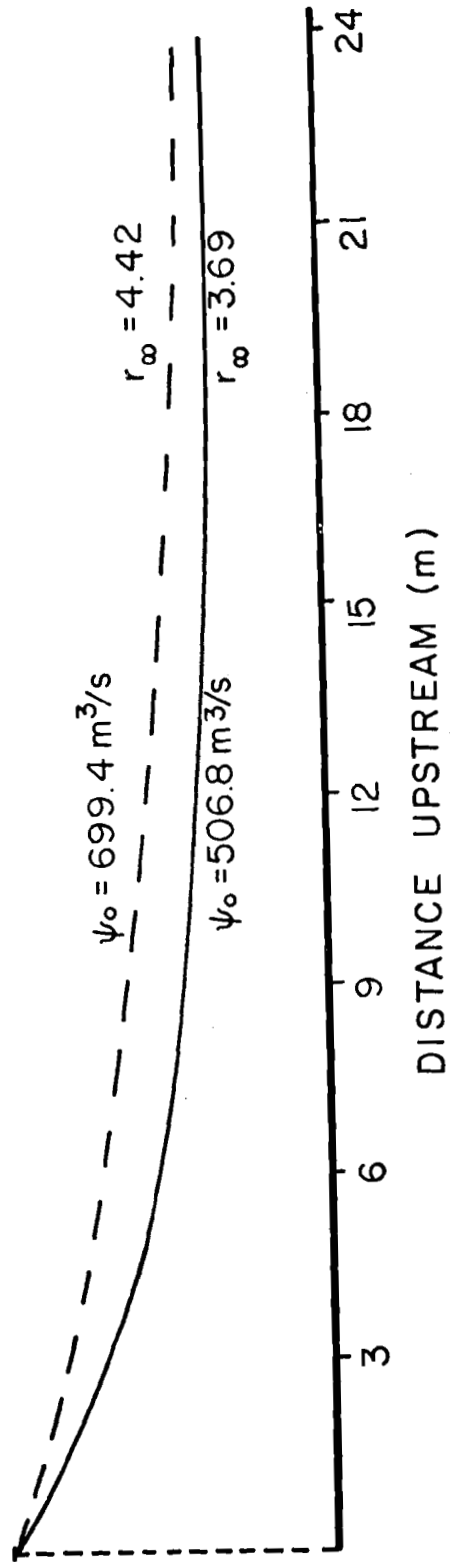
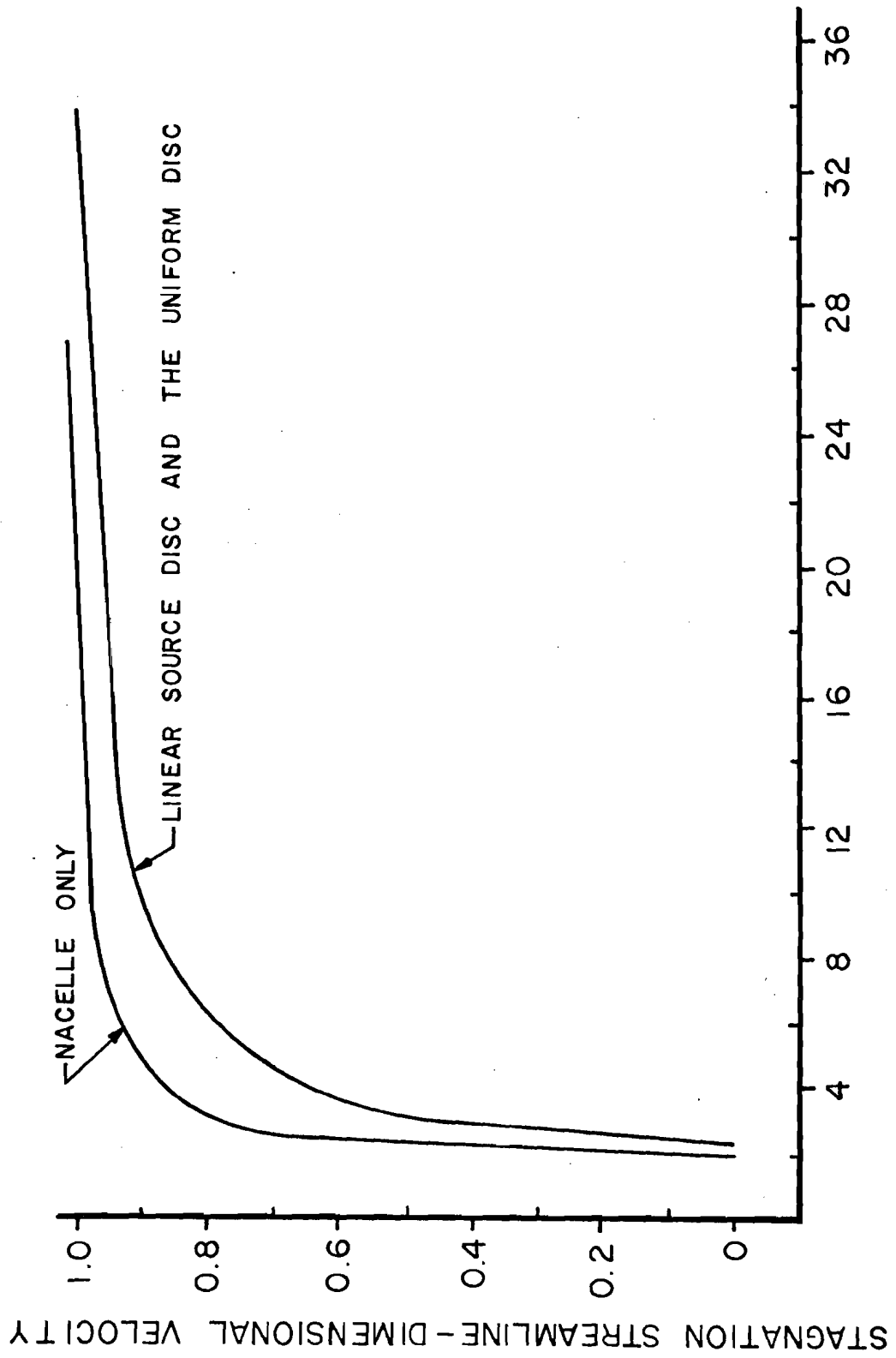


FIGURE 20.



x - METER UPSTREAM OF BLADES

FIGURE 21.



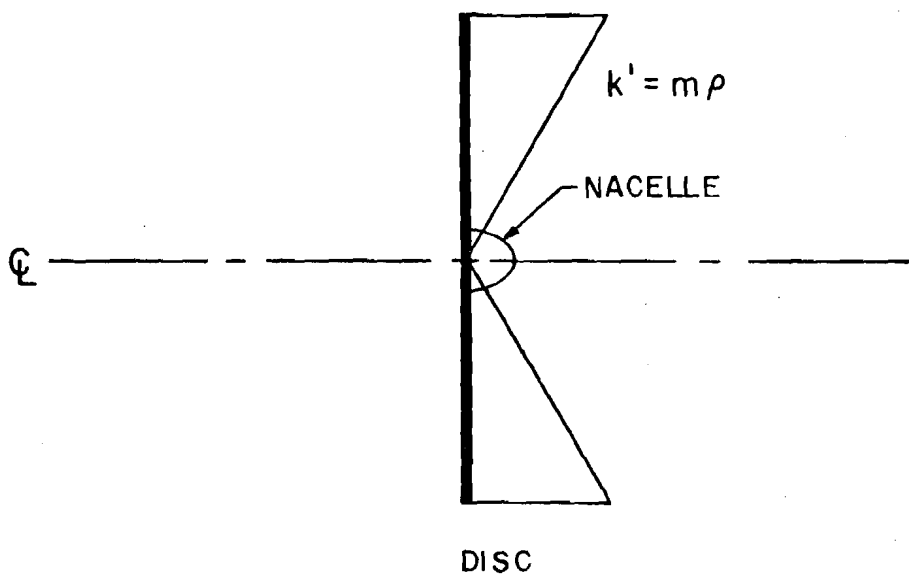


FIGURE 22.

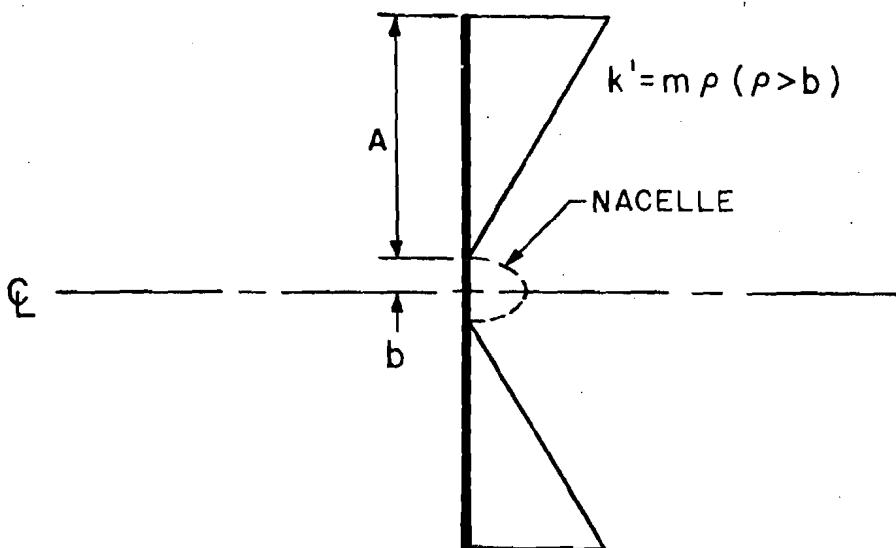


FIGURE 23.