

Implementation and calibration of finite-length plastic hinge elements for use in seismic structural collapse analysis

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ABSTRACT

Finite-length plastic hinge (FLPH) models have shown advantages over the concentrated plasticity hinge (CPH) models. However, empirical phenomenological relationships, such as Modified Ibarra-Medina-Krawinkler (ModIMK) deterioration model, were mainly calibrated for use in CPH models. ModIMK relationships are versatile and have been applied to steel, reinforced concrete, and timber structures. Herein, a calibration procedure of FLPH models and a unified algorithm for use of ModIMK relationships in CPH and FLPH models are presented. Results from included examples validate the proposed algorithms, which were implemented in OpenSees. Additionally, results highlight that FPLH models avoid errors and convergence pitfalls of CPH models.

Keywords: calibration, collapse, deterioration, finite elements, finite-length plastic hinge, concentrated plasticity, frame models, seismic analysis.

INTRODUCTION

Accurately modeling the behavior of structural members under large cyclic deformations is paramount for the quantification of the seismic performance of structures with some degree of confidence. The behavior of structural elements under these extreme loading conditions is

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20 extremely complex. Thus, several simulation approaches have been proposed which include
21 models of varying complexity and computational cost. Continuum models are generally con-
22 sidered as the most reliable approach for estimating the seismic demands of structures to
23 localized and global collapse, but they are typically complex and lead to extensive compu-
24 tational effort.

25 Concentrated plasticity hinge (CPH) elements are used herein as a reference modelling
26 approach, considering the vast experience on the use of these in the modeling of buildings
27 under seismic loads (Ibarra and Krawinkler, 2005; Medina and Krawinkler, 2005; Haselton
28 and Deierlein, 2007; PEER/ATC, 2010). In these models, each structural element is modeled
29 as the association of a linear elastic beam element and a nonlinear spring at each member
30 modified. The correct linear-elastic solution for the entire member is only obtained if the
31 end rotational springs are approximated as rigid-plastic. This is usually achieved using an
32 ad-hoc stiffness modifier parameter, n_{Factor} , for the zero-length springs. However, the defi-
33 nition of the ideal value n_{Factor} is not trivial, as a low value leads to erroneous results and
34 a high value results in numerical instability and convergence issues. As discussed in detail
35 in this work, the use of n_{Factor} also increases significantly the complexity of the implemen-
36 tation of nonlinear constitutive models. If a CPH model is used in the development of a
37 structural model, moment-rotation relationships directly obtained from experimental tests
38 can be employed to define the nonlinear zero-length springs that control element flexural
39 response. Distributed plasticity models (Spacone et al., 1996; Neuenhofer and Filippou,
40 1997) have also been widely used in the development of numerical models and have been
41 implemented in finite element softwares. Based on this formulation of distributed plasticity,
42 alternative approaches have been proposed by different authors to limit the effects of local-
43 ization phenomena related to non-objective strain-softening response of distributed plasticity
44 force-based beam-column elements (Coleman and Spacone, 2001). Force-based finite-length
45 plastic hinge (FLPH) beam-column elements were developed by Scott and Fenves (2006)
46 and Addessi and Ciampi (2007) as an alternate formulation to address localization issues.

47 The FLPH elements include two discrete length plastic hinge zones at element ends and a
48 linear elastic region in between the hinge zones, all of which are incorporated through an
49 appropriate element numerical integration scheme. When compared to the CPH approach,
50 this model has been shown to be advantageous, namely in what concerns to modeling effort,
51 computational cost, clear separation between member and connection nonlinearity, and more
52 realistic modelling of yielding progression and hinge rotations. When empirically calibrated
53 moment-rotation models are used to define the inelastic FLPH elements, a calibration pro-
54 cedure is needed. This was first identified by Scott and Ryan (2013) and a solution was
55 proposed by Ribeiro et al. (2014) for sections exhibiting softening response under monotonic
56 loading. Combined with empirically calibrated constitutive relationships, these models al-
57 low for reliable estimation of the seismic structural demands up to the onset of collapse with
58 limited computational cost.

59 Many hysteric laws have been proposed in the last decades to model the performance
60 of different structural elements and structural materials subjected to large cyclic displace-
61 ments. The main observed nonlinear phenomena include cyclic deterioration in stiffness
62 (Takeda et al., 1970) and strength (Pincheira et al., 1999; Sivaselvan and Reinhorn, 2000),
63 and pinching under load reversal (Roufaiel and Meyer, 1987). Among these models, the Mod-
64 ified Ibarra-Medina-Krawinkler (Lignos, 2008), denoted ModIMK, was selected herein for its
65 versatility. The ModIMK model has been applied to RC (Haselton and Deierlein, 2007),
66 steel (Lignos and Krawinkler, 2011), and timber structures (Ibarra and Krawinkler, 2005).
67 Since these models were mainly developed to describe force-displacement (e.g. moment-
68 rotation) relations for use in concentrated hinges, their use in FLPH elements requires mod-
69 ifications, alternative implementation, and special calibration considerations. In fact, as
70 shown in Scott and Ryan (2013) and Ribeiro et al. (2014), the use of simple scaling of the
71 constitutive law by the plastic hinge length to define a moment-curvature relation in FLPH
72 models leads to inconsistent pushover results.

73 The main objective of this paper is to present a unified implementation algorithm of

74 the ModIMK deterioration models for use in CPH and FLPH models. For the CPH model,
75 new implementations are provided for updating the unloading stiffness and the post-yield
76 hardening ratio, as well as, the computation of the committed member displacements and the
77 updated spring displacements. For the FLPH models, an extended calibration procedure is
78 proposed, which updates the flexural stiffness of the interior sections of the member to provide
79 objective and consistent element responses when empirically calibrated moment rotations
80 rules are employed for cyclic analysis. The formulation and implementation proposed was
81 included in a modified version of the Open System for Earthquake Engineering Simulation
82 (OpenSees, Mazzoni et al. (2009) 2.4.3, r5695) framework. Results from included examples
83 validate the proposed algorithms. Additionally, results highlight that FPLH models avoid
84 errors and convergence pitfalls of CPH models.

85 BACKGROUND

86 Concentrated Plasticity Hinge Models

87 In CPH models, two discrete zero-length hinges are defined at member ends and a linear
88 elastic region is defined in-between the two zero-length hinges. These three components are
89 associated in series to define a CPH member. The flexibility matrix of this member, \mathbf{f}_{mem} ,
90 is given by:

$$91 \quad \mathbf{f}_{\text{mem}} = \mathbf{f}_{\text{sI}} + \mathbf{f}_{\text{int}} + \mathbf{f}_{\text{sJ}} \quad (1)$$

92 where \mathbf{f}_{int} is the flexibility of the linear-elastic interior element and \mathbf{f}_{sI} and \mathbf{f}_{sJ} are the flexibil-
93 ities of the springs at ends I and J , respectively. The elastic stiffness of the member is given
94 as the inverse of the flexibility matrix shown in Equation 1. The model stiffness is therefore
95 related to the elastic stiffness of the rotational springs and the beam-column element, which
96 are connected in series. When a CPH model is used to model an elastic beam-column mem-
97 ber, the correct linear-elastic solution for the entire model is only obtained if rigid-plastic
98 zero-length end springs are defined. Thus, the linear elastic stiffness of the springs at both

99 ends are amplified by a constant factor n_{Factor} such that the initial stiffness of the springs
 100 is large, but not so large as to pose numerical instability. During the nonlinear analysis,
 101 the n_{Factor} must be considered when the stiffness of the element is computed (e.g., when
 102 updating the post-yield stiffness).

103 **Finite-length Plastic-hinge Models**

104 The use of finite-length plastic-hinge (FLPH) models in nonlinear analysis requires the
 105 definition of one single element in which inelastic hinge zones with discrete length are defined
 106 at element ends. The FLPH elements (Scott and Fenves, 2006; Addressi and Ciampi, 2007)
 107 are based on force-based distributed plasticity formulations in which the element integration
 108 is performed using methods that allow for the definition of a user defined hinge length at
 109 element ends.

110 In this model both end sections are assigned a nonlinear behavior, whereas the element
 111 interior is typically assumed to have an elastic behavior, however this assumption is not
 112 necessary. The flexibility of the FLPH element is computed as:

$$113 \quad \mathbf{f} = \int_{L_{pI}} \mathbf{b}(x)^T \mathbf{f}_S(x) \mathbf{b}(x) dx + \int_{L_{int}} \mathbf{b}(x)^T \mathbf{f}_S(x) \mathbf{b}(x) dx + \int_{L_{pJ}} \mathbf{b}(x)^T \mathbf{f}_S(x) \mathbf{b}(x) dx \quad (2)$$

114 where L_{pI} and L_{pJ} are the length of the plastic hinges at element ends, L_{int} is the length
 115 of the linear-elastic element interior, $\mathbf{b}(x)$ is the interpolation function matrix, and \mathbf{f}_S is the
 116 section flexibility, nonlinear for the first and third term, and typically linear for the second
 117 term. For other formulation details see Scott and Fenves (2006), for example.

118 **CONSTITUTIVE LAWS FOR CYCLIC LOADING**

119 In this paper, the modified Ibarra-Medina-Krawinkler deterioration model (Lignos, 2008),
 120 ModIMK model in short, was chosen for its versatility in modeling degrading hysteretic
 121 response of structural elements. This model was empirically calibrated for reproducing the
 122 moment-rotation relation of reinforced concrete (Haselton and Deierlein, 2007) and steel
 123 structural components (Lignos and Krawinkler, 2011). The ModIMK model is based on: (i)

124 a backbone curve defining the reference monotonic behavior, (ii) a set of rules defining the
 125 hysteretic behavior between the positive and negative backbone curves; and (iii) a set of
 126 rules that define up to six modes of deterioration of the hysteretic behavior.

127 Figure 1(a) illustrates the parameters that define the backbone curve. This curve is
 128 defined by three strength parameters: effective yield strength (or basic strength), F_y , cap-
 129 ping strength, F_C (or post-yield strength hardening ratio F_C/F_y), and residual strength, F_r ;
 130 and four deformation parameters: yield deformation, d_y , pre-capping plastic deformation
 131 for monotonic loading, d_p , post-capping plastic deformation, d_{pc} , and ultimate deformation
 132 capacity, d_u . The ModIMK model defines six modes of cyclic strength and stiffness deteri-
 133 oration: (i) basic strength, (ii) post-yield hardening ratio, (iii) post-capping strength, (iv)
 134 unloading stiffness, (v) reloading stiffness, and (vi) pinching behavior. Figures 1(b) to 1(d)
 135 illustrate three models that have been proposed in the literature based on different combina-
 136 tions of these six modes of deterioration. All three models share the same backbone curve.

137 The models are:

- 138 • Bilinear hysteretic response (Bilin) model with strength deterioration (Figure 1b);
- 139 • Peak-oriented model with strength and stiffness deterioration (Figure 1c);
- 140 • Pinching model with strength and stiffness deterioration (Figure 1d).

141 In the ModIMK models, the rates of cyclic deterioration are controlled by a characteristic
 142 total hysteretic energy dissipation capacity E_t and an energy based rule developed in Rah-
 143 nama and Krawinkler (1993). The characteristic total hysteretic energy dissipation capacity
 144 E_t is obtained from experimental results.

145 The energy based rule developed by Rahnama and Krawinkler (1993) expresses the cyclic
 146 deterioration in excursion i , β_i :

$$147 \quad \beta_i = \left(\frac{E_i}{E_t - \sum_{j=1}^i E_j} \right)^c \leq 1 \quad (3)$$

148 where E_i is the hysteretic energy dissipated in excursion i , and $\sum E_j \leq E_t$ is the hysteretic

149 energy dissipated in all previous excursions in both positive and negative directions. The
150 exponent c defines the rate of deterioration. According to Rahnama and Krawinkler (1993),
151 a reasonable range of values for c is between 1.0 and 2.0. β_i ranges between 0 and 1.

152 The generalized stiffness or strength parameter, X , can be updated through:

$$153 \quad X_i = (1 - \beta_k) \times X_{i-1} \quad (4)$$

154 where X_i is the value of the parameter in excursion i and β_k is the value of deterioration
155 parameter.

156 The ModIMK is used herein to model the behavior of plastic hinges. However, the
157 implementation of this model within a finite element framework is complex and dependent
158 on the type of finite element used. In the following sections the details regarding a consistent
159 and unified implementation of these models is provided for CPH and FLPH models.

160 **IMPLEMENTATION OF MODIMK MODELS IN HINGE ELEMENTS**

161 Figure 2 shows the general procedure used to update the ModIMK model parameters.
162 This procedure is a direct application of the proposal by Ibarra et al. (2005), and it is
163 detailed here for completeness of the discussion on new implementation that follows in the
164 next subsections. At the beginning of the analysis, the model parameters are initialized. In
165 the elastic range, no change in these parameters occurs and no update of the constitutive
166 law is required. The unloading stiffness is the only parameter which is updated when a load
167 reversal takes place in the inelastic range. In a finite element implementation, the stiffness
168 must be known before the reversal, requiring the updating of the unloading stiffness in all
169 steps in the inelastic range. Furthermore, this is the only deterioration mode for which a
170 common deterioration parameter is used in both loading directions.

171 The remaining parameters are updated at the end of the unloading branch ($F_{n-1} \times F_n <$
172 0), denoted by point Y in Figure 1. At this point, dissipated energy in the previous excursion
173 is computed. This allows for updating of the reloading stiffness, the basic strength, the strain

174 hardening ratio, the capping point, and the pinching parameters for the current excursion.
 175 The procedure is then repeated for each excursion reaching the nonlinear range.

176 **Implementation in Concentrated Plastic Hinge Models**

177 In the CPH model, to guarantee the rigid plastic behavior of the springs, their initial
 178 stiffness is given by:

$$179 \quad k_{s,m} = (n_{Factor} + 1) \times K_{mem}, \quad m = I, J \quad (5)$$

180 where K_{mem} is the elastic stiffness of the member. In the case of double curvature, $K_{mem} =$
 181 $6EI/L$, where EI the is cross-section flexural stiffness, and L is the member length. Since
 182 the elastic stiffness of the member is related to the elastic stiffness of the rotational springs
 183 and the interior elastic element, which are connected in series, the stiffness of the interior
 184 element, k_{int} , is also affected by n_{Factor} , as:

$$185 \quad k_{int} = \frac{n_{Factor} + 1}{n_{Factor}} \times K_{mem} = \frac{6EI_{mod}}{L} \quad (6)$$

186 where $EI_{mod} = \frac{n_{Factor} + 1}{n_{Factor}} EI$ is the modified elastic stiffness of the element interior.

187 In the post-yielding region, member stiffness is computed by multiplying the elastic stiff-
 188 ness by the post-yielding ratio, α . Since the elastic stiffness of the zero-length spring is
 189 affected by the n_{Factor} , an adjusted post-yielding ratio of the spring, α' (ratio of the tangent
 190 stiffness, k_{Ts} , to the linear elastic stiffness, k_s) is given by:

$$191 \quad \alpha' = \frac{k_{Ts}}{k_s} = \frac{\alpha}{1 + n_{Factor} \times (1 - \alpha)} \quad (7)$$

192 The introduction of an n_{Factor} in the definition of the zero-length springs requires that
 193 several modifications are considered in the ModIMK implementation and general deteri-
 194 oration model given in Equation 4. The adjusted implementation details when defining
 195 moment-rotation empirical relations in CPH models are presented next for each of the six
 196 deterioration modes. For comparison purposes, a simplified implementation, where the effect

197 of n_{Factor} is not considered in the updating of model parameters, is denoted as CPH-original.
 198 In general, two main adjustments are made. First, the stiffness of the nonlinear spring is
 199 updated so that the stiffness of the entire element is equal to the objective stiffness. Second,
 200 the displacements of the springs need to be updated so that the correct target displace-
 201 ments (rotations) of the element are achieved. In what regards strengths, since the force
 202 (moment) in the spring is equal to the force (moment) in the element ends, no adjustment
 203 is required. Therefore, the basic and post-capping strength deterioration follows the general
 204 form of Equation 4.

205 The zero-length spring stiffness is affected by the n_{Factor} and the post-yielding ratio of the
 206 spring defined in Equation 7 is used. When computing the deterioration of the post-yielding
 207 hardening ratio the general model described in Equation 4 is not applicable. Instead, the
 208 deterioration of the post-yielding hardening ratio is computed using the new procedure shown
 209 in Figure 3. In this procedure, first, the member hardening ratio of the previous excursion
 210 is computed using the inverse of Equation 7. Second, given β_i , the member post-yielding
 211 stiffness is updated. Lastly, Equation 7 is used to compute the updated hardening ratio of
 212 the nonlinear springs.

213 Since the unloading stiffness deterioration depends on the energy dissipated up to the
 214 beginning of the unloading branch rather than that dissipated in a complete excursion, an
 215 implementation different than the one proposed by Ibarra and Krawinkler (2005) is used
 216 herein for this parameter. Equation 4 is thus replaced by:

$$217 \quad K_{u,n}^{member} = \left[\prod_{j=1}^{i-1} (1 - \beta_{k,j}) \right] \times (1 - \beta_{k,n}) \times K_0 = \gamma_k \times K_0 \quad (8)$$

218 where i is the total number of inelastic excursions up to load step n , $\beta_{k,j}$ is the deterioration
 219 parameter associated with completed inelastic excursion j , $\beta_{k,n}$ is the deterioration param-
 220 eter computed considering the energy dissipated in excursion i up to load step n , γ_k is the
 221 cumulative deterioration of the unloading stiffness and K_0 is the member initial elastic stiff-
 222 ness. The procedure starts by computing the residual energy dissipation capacity, $E_t - \sum E_j$

223 and the damage parameter β_k . Equation 8 is then used to update the unloading stiffness of
 224 the element based on its elastic stiffness. The unloading spring stiffness is thus given by:

$$225 \quad K_{u,n}^{spring} = \left(\frac{\gamma_k}{1 + n_{Factor} \times (1 - \gamma_k)} \right) \times K_0 \quad (9)$$

226 where K_0 and $K_{u,n}^{spring}$ are the original member elastic stiffness and updated unloading stiffness
 227 of the zero-length spring in loading step n .

228 The reloading stiffness deterioration is modeled by increasing the absolute value of the
 229 target displacement of the member, d_i , corresponding to the horizontal coordinate of point
 230 Y in Figure 1c, in each direction as:

$$231 \quad d_i = (1 + \beta_i) \times d_{i-1}^{max} \quad (10)$$

232 where d_{i-1}^{max} is the maximum displacement observed up to the $i - 1$ excursion in the same
 233 direction.

234 The implemented algorithm for computing the reloading stiffness deterioration in CPH
 235 models is presented in Figure 4. Firstly, the maximum displacement of the member in
 236 previous excursions, $d_{i-1}^{max,member}$, is computed using the general relation between spring and
 237 member rotations:

$$238 \quad d^{spring} = d^{member} - d^{elastic} = d^{member} - F(d^{member}) \times K_{member} \quad (11)$$

239 where $F(d^{member})$ is the force associated with the displacement d^{member} , obtained with the
 240 backbone curve computed for the current step of the analysis. $F(d^{member}) \times K_{member}$ is
 241 thus the elastic deformation of the member, associated with the force $F(d^{member})$ under the
 242 assumption of double curvature.

243 The updated member maximum displacement is then updated using Equation 10. Then,
 244 the updated backbone curve for this excursion is defined, based on the updated basic
 245 strength, post-yielding ratio, and post-capping strength. This is then used to compute

246 the force $F(d_i^{max,member})$. The maximum deformation of the zero-length spring can then be
247 calculated using Equation 11.

248 Finally, the reloading stiffness is defined using point Y in Figure 1 and the new maximum
249 deformation point ($d_i^{max,spring}$; $F^i(d_i^{max,member})$). The maximum deformation is monitored in
250 each load step.

251 The implementation of updates of the pinching parameters is similar to that described
252 for the reloading stiffness. The additional notable point in reloading (see point P in Fig-
253 ure 1d) is computed by multiplying the yielding displacement and the corresponding force by
254 parameters A_{pinch} and F_p^\pm , respectively. Firstly, the maximum deformation in the member
255 is calculated, using the relationship presented in Equation 11. Then, the intermediate point
256 for pinching response is computed for the member by multiplying factors A_{pinch} and F_p^\pm (for
257 positive loading direction) to the maximum deformation and associated force, respectively.
258 Once this intermediate point is found, the corresponding intermediate point for the zero-
259 length spring is computed using Equation 11. Finally, the stiffness associated with the two
260 branches that characterize pinching response can be computed for the CPH member.

261 **Implementation and Calibration in Finite-length Plastic-hinge Models**

262 If the deteriorating models described herein are applied to *FPLH* elements, the imple-
263 mentations developed by Ibarra and Krawinkler (2005) do not require modifications, as the
264 objective stiffness and displacements can be directly assigned to the member. This results
265 in a much simpler implementation based on the general algorithm presented in Figure 2 and
266 the general updating Equation (Eq. 4). This is one of the main advantages of using FLPH
267 models, i.e. that the original hysteretic laws do not need adjustments as is the case when
268 CPH models are used.

269 As shown by Scott and Ryan (2013), employing a moment-rotation constitutive law divid-
270 ing the rotations by the plastic hinge length to obtain a moment-curvature relation produces
271 inconsistent results and the objective moment-rotation response is not recovered. Thus,
272 Scott and Ryan (2013) proposed a calibration procedure to address this issue. However, the

273 calibration procedure was developed for hardening responses only. An alternate calibration
 274 procedure developed by Ribeiro et al. (2014) was proposed for both hardening and softening
 275 responses under monotonic loading. This procedure is extended here for cyclic loading.

276 The detailed formulation of the FLPH elements is presented in Scott and Fenves (2006).
 277 In the interest of brevity, only a description of key aspects is presented here. The member
 278 flexibility using the modified Gauss-Radau integration scheme is given by:

$$279 \quad \mathbf{f} = \sum_{i=1}^{N_{pI}} (\mathbf{b}^T \mathbf{f}_s \mathbf{b}|_{x=\xi_i}) w_i + \int_{L_{int}} \mathbf{b}(x)^T \mathbf{f}_s(x) \mathbf{b}(x) dx + \sum_{i=N_{pI}+1}^{N_{pI}+N_{pJ}} (\mathbf{b}^T \mathbf{f}_s \mathbf{b}|_{x=\xi_i}) w_i \quad (12)$$

280 where N_{pI} and N_{pJ} are the number of integration points associated with the plastic hinges
 281 at the element ends, and $\mathbf{f}_s(x)$ is the section flexibility. For the modified Gauss-Radau
 282 integration $N_{pI} = N_{pJ} = 2$. The interior element term (middle term in the right hand side
 283 of Equation 12) can be computed analytically or numerically. In the latter case, the Gauss-
 284 Legendre integration scheme can be used. If two integration points are placed in this region,
 285 a total of six integration points are defined along the member length. The location ξ_i of
 286 the integration points associated with the modified Gauss-Radau plastic hinge integration is
 287 given by:

$$288 \quad \xi = \{\xi_I, \xi_{int}, \xi_J\} \quad (13)$$

289 where:

$$290 \quad \begin{aligned} \xi_I &= \left\{0; \frac{8L_{pI}}{3}\right\} \\ \xi_{int} &= \left\{4L_p + \frac{L_{int}}{2} \times \left(1 - \frac{1}{\sqrt{3}}\right); 4L_p + \frac{L_{int}}{2} \times \left(1 + \frac{1}{\sqrt{3}}\right)\right\} \\ \xi_J &= \left\{L - \frac{8L_{pJ}}{3}; L\right\} \end{aligned} \quad (14)$$

291 The corresponding weights w_i are given by:

$$292 \quad \mathbf{w} = \{\mathbf{w}_I, \mathbf{w}_{int}, \mathbf{w}_J\} \quad (15)$$

293 where:

$$294 \quad \mathbf{w}_I = \{L_{pI}; 3L_{pI}\} \quad \mathbf{w}_{int} = \left\{ \frac{L_{int}}{2}; \frac{L_{int}}{2} \right\} \quad \mathbf{w}_J = \{3L_{pJ}; L_{pJ}\} \quad (16)$$

295 In this case, the element flexibility is then given by:

$$296 \quad \mathbf{f} = \sum_{i=1}^6 (\mathbf{b}^T \mathbf{f}_s \mathbf{b}|_{x=\xi_i}) w_i \quad (17)$$

297 The inclusion of experimentally calibrated moment-rotation relations to define the behav-
 298 ior of nonlinear regions can be implemented by modifying the flexural stiffness at integration
 299 points in the elastic region of the FLPH member (see Figure 5), so that the flexibility matrix
 300 of the calibrated FLPH member is equal to a reference flexibility, which is considered as
 301 that of the CPH model with $n_{Factor} \rightarrow \infty$. In a 2D beam-column element, a system of three
 302 integral equations corresponding to each of the unique flexural coefficients of the element
 303 flexibility matrix is defined. The flexibility matrix of the FLPH element is computed using
 304 Equation 17, where

$$305 \quad \mathbf{b} = \begin{bmatrix} x/L - 1 & x/L \end{bmatrix} \quad (18)$$

306 and

$$307 \quad \{f_s(\xi_1); \dots; f_s(\xi_6)\}^T = 1/EI \cdot \{\alpha_1 \times 6 \times L_{pI}/L; \beta_1; \beta_2; \beta_2; \beta_3; \alpha_2 \times 6 \times L_{pJ}/L\}^T \quad (19)$$

308 where α_1 and α_2 are the ratio between the nonlinear stiffness and the elastic stiffness at end
 309 I and J , respectively, and β_1 , β_2 and β_3 are the flexural modification parameters.

310 The equivalent flexibility matrix, considering the CPH model is given by:

$$311 \quad \mathbf{f}_b = \lim_{n_{Factor} \rightarrow \infty} \left\{ \begin{bmatrix} 1/k_{TI} & 0 \\ 0 & 0 \end{bmatrix} + \frac{L}{6EI_{mod}} \times \begin{bmatrix} 2 & -1 \\ -1 & 2 \end{bmatrix} + \begin{bmatrix} 0 & 0 \\ 0 & 1/k_{TJ} \end{bmatrix} \right\} \quad (20)$$

312 where $k_{TI} = (n_{Factor} + 1) K_{mem}$ and $k_{TJ} = (n_{Factor} + 1) K_{mem}$ are the tangent stiffness of
 313 the springs at ends I and J , respectively, and $EI_{mod} = \frac{n_{Factor} + 1}{n_{Factor}} EI$.

314 From this system of equations, the three elastic stiffness modification parameters, β_1 , β_2 ,
315 and β_3 , are computed as a function of L_{pI} , L_{pJ} , L and n_{Factor} . When the n_{Factor} tends to
316 infinity, β_1 , β_2 and β_3 are given by:

$$\begin{aligned}
317 \beta_1 &= -\frac{54L_{pI}L^3 - 6L_{pI}(60L_{pI} + 60L_{pJ})L^2 + 6L_{pI}(96L_{pI}^2 + 288L_{pI}L_{pJ} + 96L_{pJ}^2)L - 6L_{pI}(256L_{pI}^2L_{pJ} + 256L_{pI}L_{pJ}^2)}{L(3L - 16L_{pJ})(L^2 - 20LL_{pI} + 4L_{pJ}L + 64L_{pI}^2)} \\
318 \beta_2 &= -\frac{3(4L_{pI} - L + 4L_{pJ})(3L^2 - 12LL_{pI} - 12LL_{pJ} + 32L_{pI}L_{pJ})}{L(3L - 16L_{pI})(3L - 16L_{pJ})} \\
319 \beta_3 &= -\frac{54L_{pJ}L^3 - 6L_{pJ}(60L_{pI} + 60L_{pJ})L^2 + 6L_{pJ}(96L_{pI}^2 + 288L_{pI}L_{pJ} + 96L_{pJ}^2)L - 6L_{pJ}(256L_{pI}^2L_{pJ} + 256L_{pI}L_{pJ}^2)}{L(3L - 16L_{pI})(L^2 - 20LL_{pJ} + 4L_{pI}L + 64L_{pJ}^2)}
\end{aligned} \tag{21}$$

320 Assuming both plastic hinges at member ends have similar lengths L_p , the stiffness modifying
321 factors (β_1 , β_2 and β_3 , see Figure 5) are given by:

$$\begin{aligned}
322 \beta_1 &= \beta_3 = -\frac{6(3L^2L_p - 24LL_p^2 + 32L_p^3)}{L(L - 8L_p)^2} \\
323 \beta_2 &= \frac{3(3L^3 - 48L^2L_p + 224LL_p^2 - 256L_p^3)}{L(3L - 16L_p)^2}
\end{aligned} \tag{22}$$

324 As shown, these factors do not depend on the stiffness terms α_i , $\{i = 1, 2\}$ and therefore
325 are constant during the analysis. Therefore Equation 22 only needs to be applied once
326 at the beginning the analysis, implying a very limited computational cost. Moreover, this
327 independence between the stiffness terms and the flexural modification factors makes this
328 implementation independent of the constitutive law employed.

329 For the *FPLH* model, in terms of calibration, the only other parameter that needs
330 adjusting is the the total energy dissipation capacity E_t . This term is defined empirically
331 for the moment-rotation relation, and can be defined, for moment-curvature, as:

$$332 E_t^{M-\chi} = E_t^{M-\theta} / L_p \tag{23}$$

333 All other parameters follow the general models developed and implemented by Ibarra

334 and Krawinkler (2005).

335 NUMERICAL EXAMPLES

336 In this section a simple structure subjected to a set of cyclic pushover analyses is used to
337 evaluate the accuracy and stability of the proposed implementations. The algorithms and
338 procedures discussed were implemented in a modified version of the Open System for Earth-
339 quake Engineering Simulation (OpenSees, Mazzoni et al. (2009), 2.4.3, r5695) framework.
340 In the examples, a simply supported beam subjected to different end moments (see Figures
341 6 to 10) is analyzed under cyclic displacement control considering the three material models
342 discussed. The beam has a 24 feet (7.33m) span and the model parameters for all material
343 models are presented in Table 1. The ultimate rotation, θ_u and the plastic hinge length, L_p ,
344 were taken equal to 0.4 rad and $L/16$, respectively, for all cases. For the Pinching model,
345 three additional parameters that define the mid-point in the reloading branch are assumed
346 to be equal to 0.4.

347 Figures 6, 7, and 8 shows results for analyses performed using the pinching model
348 for moment gradients defined with one end moment, two anti-symmetric end moments, and
349 two symmetrical end moments. The first set of results compares the results obtained using
350 the CPH model, both considering direct application of Equation 4 (CPH-original) and using
351 the proposed implementation (CPH-updated), with those obtained with the finite length
352 plastic hinge model (FLPH) and an analytical solution. For the CPH-original, n_{Factor} was
353 taken equal to 10 to reduce numerical instabilities, following recommendations in Ibarra and
354 Krawinkler (2005) and Zareian and Medina (2010). For CPH-updated n_{Factor} was taken equal
355 to 10 and 1000. Results show that all implementations lead to acceptable results. However,
356 the CPH-original and CPH-updated, considering a n_{Factor} equal to 10, lead to a noticeable
357 over-estimation of the elastic stiffness. This error propagates to the entire analysis, as can be
358 seen at the end of the unloading branch. Moreover, as a result of not updating the stiffness of
359 the elastic element interior during analysis, the CPH-original also leads to significant errors
360 in the unloading and reloading stiffnesses. The analysis using the FLPH elements provide

361 the results closest to the theoretical results, being clearly the most accurate model. Figures
362 6, 7, and 8 show that the amplitude of observed errors decrease with increase in the moment
363 gradient along the element length, being smaller for the anti-symmetric loading and larger
364 for the symmetric loadings. In addition, it is clear that the use of the CPH-original model
365 does not allow for obtaining accurate results as the direct application of Equation 4, i.e.
366 not considering the implementation procedures proposed herein, is not enough for correctly
367 updating model parameters during the analysis. Figures 9 and 10, which show the results
368 obtained for the Peak-oriented and Bilin models indicate that the conclusions drawn for the
369 pinching model hold for the other material models.

370 In Figure 11 the errors in the elastic stiffness are plotted for the FLPH model and for
371 the CPH-updated implementation with n_{Factor} values between 10 and 1000. Results show
372 convergence of the error when the CPH-updated implementation is used. However, even
373 for large $n_{factors}$ the CPH-updated produces the largest errors when estimating the elastic
374 stiffness. It is clear that the FLPH model results in very small errors, only comparable with
375 those obtained for the CPH-updated with an n_{Factor} equal to 1000. The results presented
376 refer to the Bilin model, but conclusions hold for all constitutive models implemented.

377 To compare the numerical stability of different implementations, results of an elementary
378 assessment are shown in Figure 12. The models were analysed considering the Krylov-Newton
379 algorithm (Scott and Fenves, 2010) under displacement control analyses. Pseudo-time steps
380 between 1×10^{-7} and 1×10^{-3} are used in the analyses. The norm of the displacement
381 increment convergence test is used with a threshold of 1×10^{-8} . Figure 12 shows that FLPH
382 and the CPH-original with n_{Factor} equal to 10 converged for all time steps. However, n_{Factor}
383 values between 100 and 500 required a pseudo-time step smaller than 1×10^{-5} for achieving
384 convergence. For a n_{Factor} equal to 1000, a pseudo-time step of 1×10^{-7} was necessary to
385 achieve convergence. Although this is not an exhaustive convergence stability analysis, the
386 results indicate that the FLPH is significantly more stable. Similar stability is obtained for
387 the CPH-original model only if n_{Factor} is taken equal to 10 which, as shown above, leads to

388 significant overestimation of the elastic stiffness.

389 CONCLUSIONS

390 Within the wide range of member models available in the literature, concentrated plas-
391 ticity hinge (CPH) models have been the reference model for earthquake engineering studies
392 during the last decade. However, finite-length plastic hinge (FLPH) models have been re-
393 cently shown to be advantageous over the CPH models. A significant reduction in modelling
394 effort, as well as in computational cost, a clear distinction between member and connection
395 nonlinearities, and more realistic modelling of yielding progression and hinge rotations are
396 the most important advantages of the FLPH model.

397 In this work, results obtained for cyclic analysis using implementation and calibration of
398 the FLPH models are discussed and compared to those resulting from two implementations
399 used for updating parameters of the unloading stiffness and other deterioration modes in
400 the CPH models. All implementations were performed in the Open System for Earthquake
401 Engineering Simulation (OpenSees) making use of the ModIMK material models, which have
402 been widely used for simulating steel, RC, and timber frame structures.

403 In terms of the implementation, the main conclusions of this paper are:

- 404 1. a new unified implementation was developed in the OpenSees framework, where the
405 ModIMK material models can now be used in both CPH and FLPH models;
- 406 2. the implementation of the ModIMK in the CPH models proved to be significantly more
407 complex than that done for FLPH models. This results from the use, in this case,
408 of three separate components, two zero-length springs and an elastic beam-column
409 interior element. In addition, the elastic stiffness of the zero-length springs needs to
410 be amplified in order to obtain the correct member flexibility matrix, which requires
411 further adjustments in the updating procedure of all parameters of the springs;
- 412 3. in FLPH models, the main difficulty lies, not on the implementation of the Mod-
413 IMK material models, but in the need to calibrate the element to consider empirical

414 moment-rotation relationships;

- 415 4. although a calibration procedure is required for the FLPH elements, this procedure
416 can be used independently of the constitutive law. For the CPH models, custom
417 implementations are required if different constitutive laws are to be used;
- 418 5. for FPLH models, once the formulation of the calibration is defined, the implementa-
419 tion procedure is significantly simpler and applicable to a wide range of constitutive
420 deterioration models, thus not restricted to the ModIMK relationships;
- 421 6. the FLPH calibration proposed was validated for nonlinear cyclic analysis.

422 Based on the numerical results shown:

- 423 1. in general, CPH and FLPH models can provide reasonable results for nonlinear cyclic
424 analysis;
- 425 2. for a beam element with anti-symmetric end moments, CPH models provide accurate
426 results independently of the n_{Factor} that is used to amplify the elastic stiffness of the
427 zero-length springs;
- 428 3. for a beam element with other moment gradients, non-negligible errors are obtained
429 for the elastic stiffness if the n_{Factor} in CPH models is not large enough (e.g., approx-
430 imately 5% error is obtained for symmetric bending moments for $n_{Factor} = 10$); these
431 errors propagate throughout the analysis;
- 432 4. CPH models with large n_{Factor} values give rise to numerical instabilities;
- 433 5. calibrated FLPH models provided the most accurate results.

434 In summary, even though the use of FLPH models in large numerical studies requires
435 more investigation, the results presented in this work indicate that these models are suitable
436 for being used in large numerical simulations, being more stable, accurate, and versatile.

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502 **List of Tables**

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TABLE 1. ModIMK model parameters used in the numerical examples

Model	EI (kN.m ²)	M_y^+ and M_y^- (kN.m)	M_c/M_y	θ_p (rad)	θ_{pc} (rad)	κ	$E_t^{M-\theta}$ (kN.m)
All models	2.33×10^6	1911	1.05	0.233	0.156	0.4	2255

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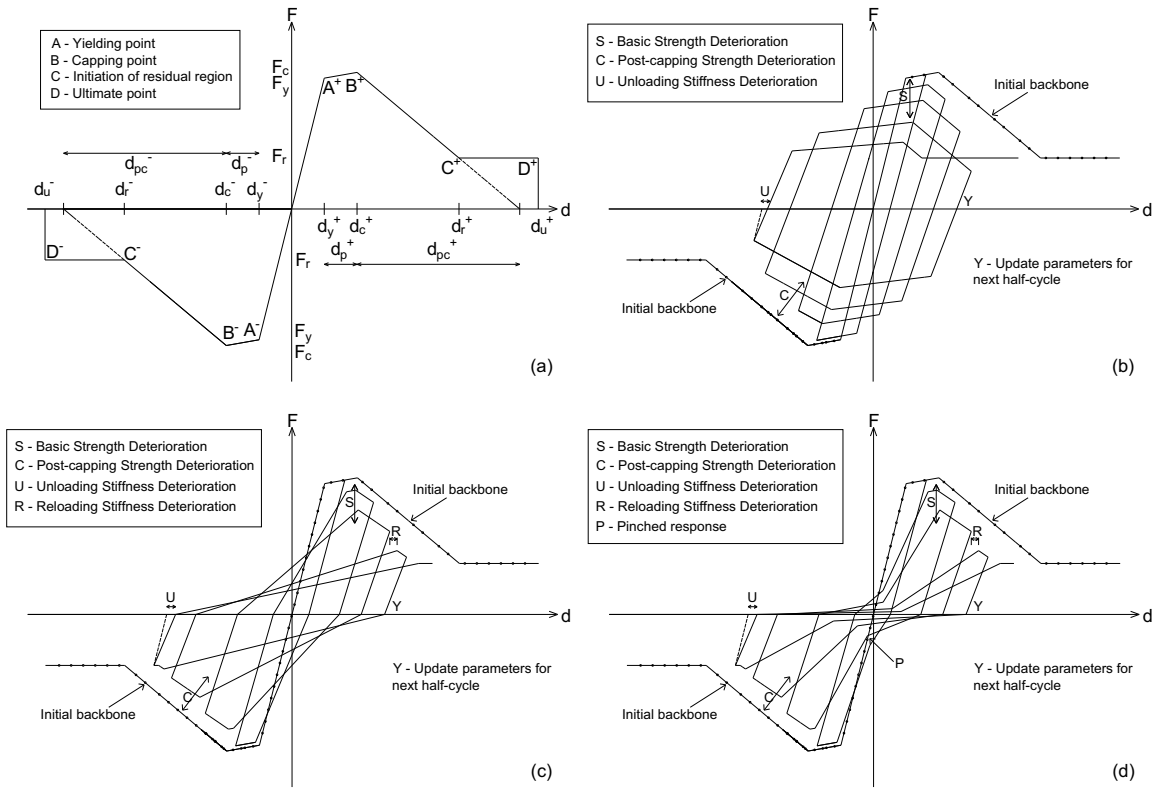


FIG. 1. Modified Ibarra-Medina-Krawinkler deterioration models: (a) backbone curve, (b) Bilin model, (c) Peak-oriented model, and (d) Pinching model.

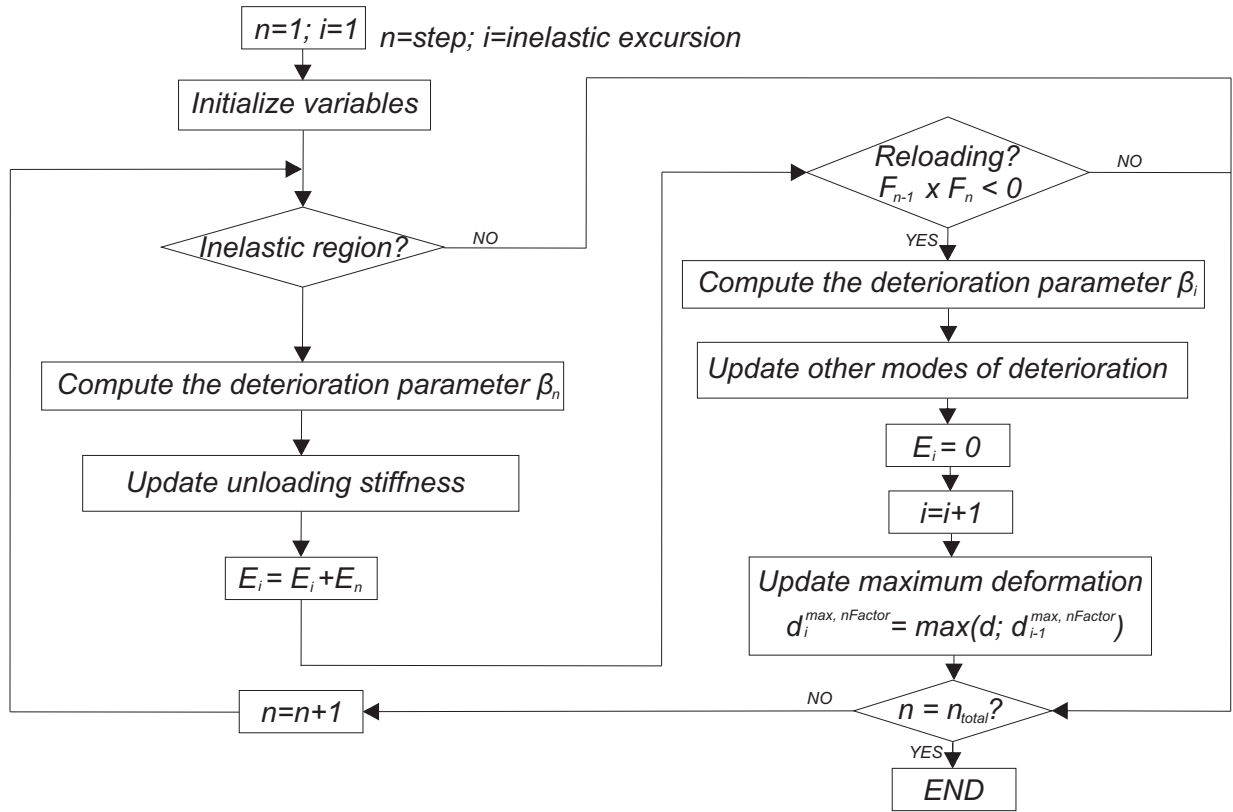


FIG. 2. General procedure for updating model parameters during cyclic analysis

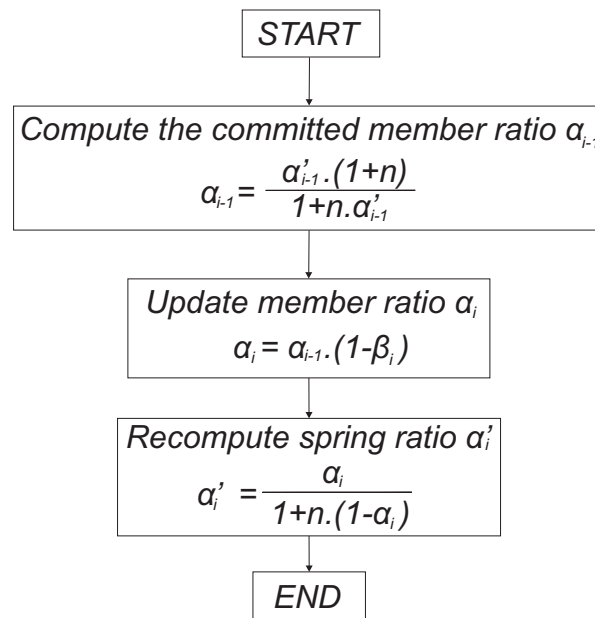


FIG. 3. Procedure for updating post-yielding ratio during cyclic analysis for Concentrated Plasticity Hinge model

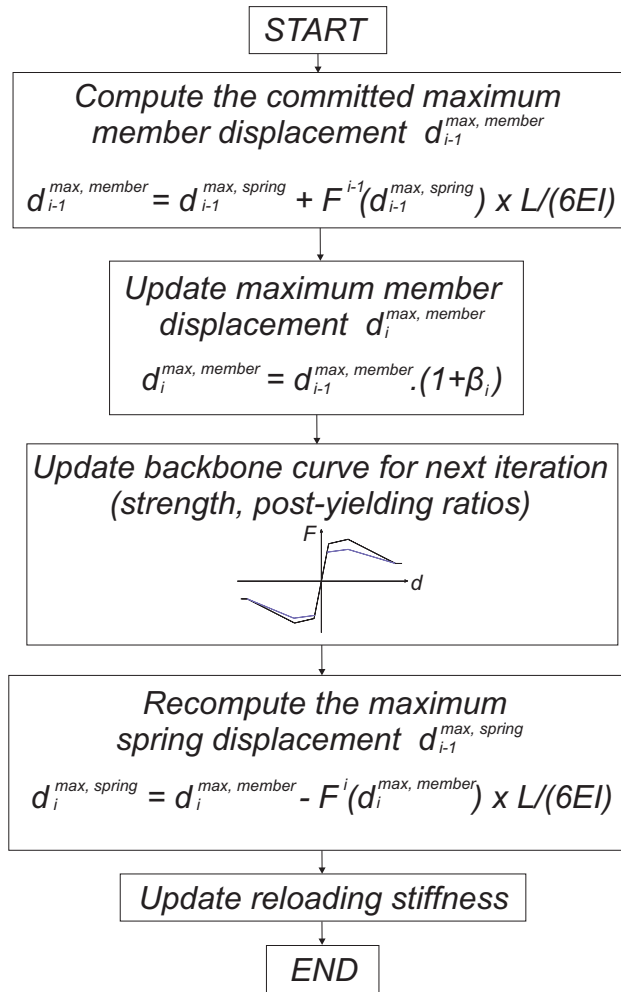


FIG. 4. Procedure for updating reloading stiffness during cyclic analysis for Concentrated Plasticity Hinge model

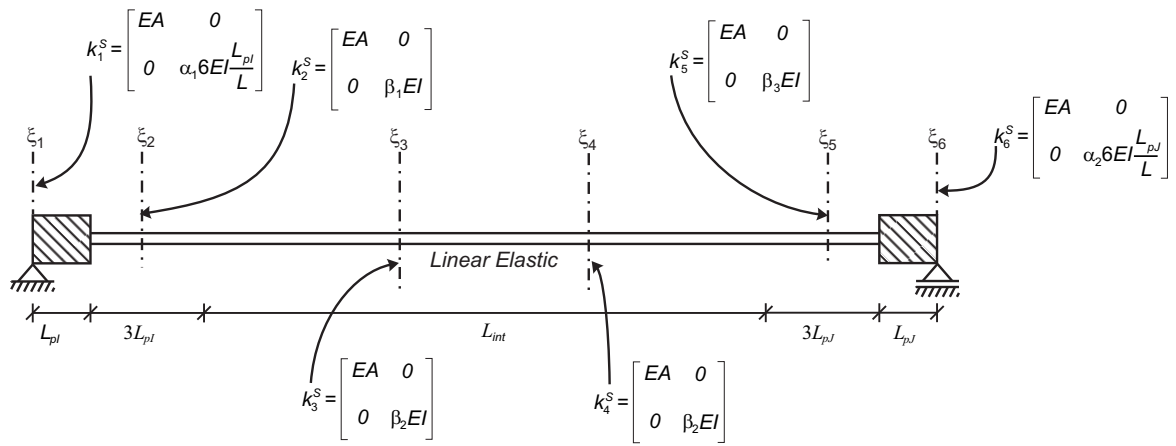


FIG. 5. Modified Gauss-Radau integration scheme considering flexural stiffness modification parameters (β_1 , β_2 and β_3)

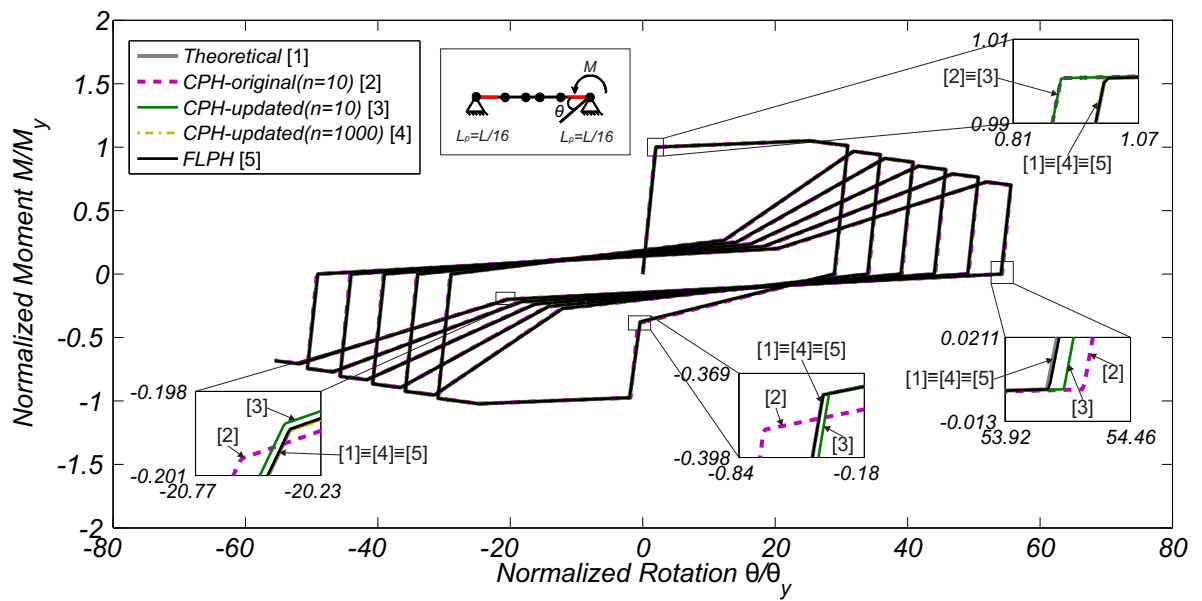


FIG. 6. Pinching model - cyclic analysis considering a single end moment

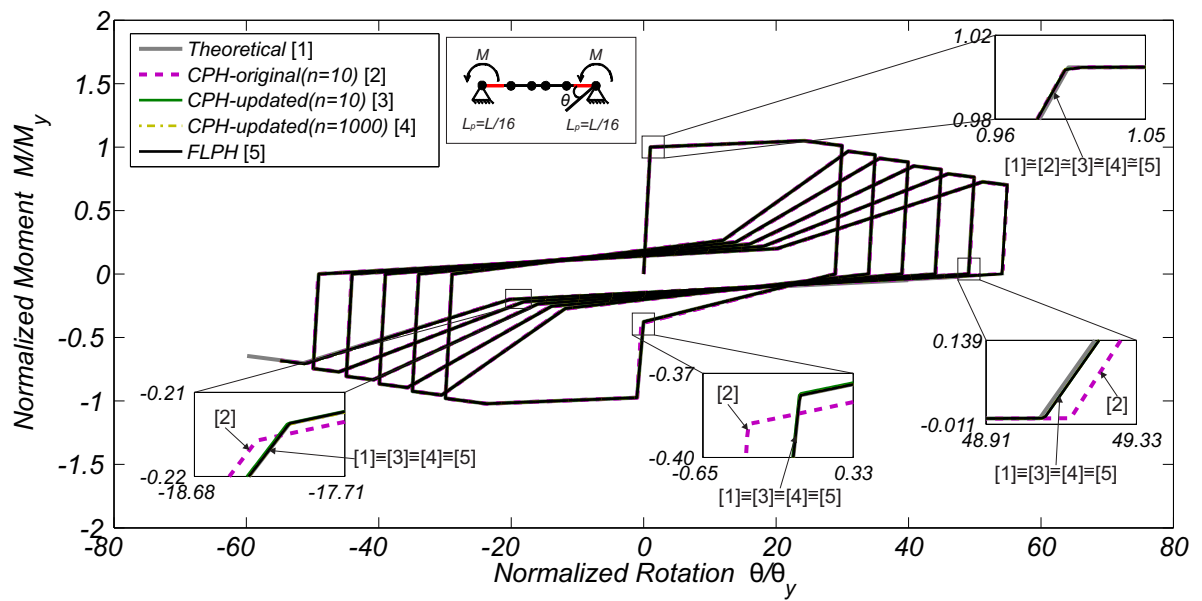


FIG. 7. Pinching model - cyclic analysis considering anti-symmetric end moments

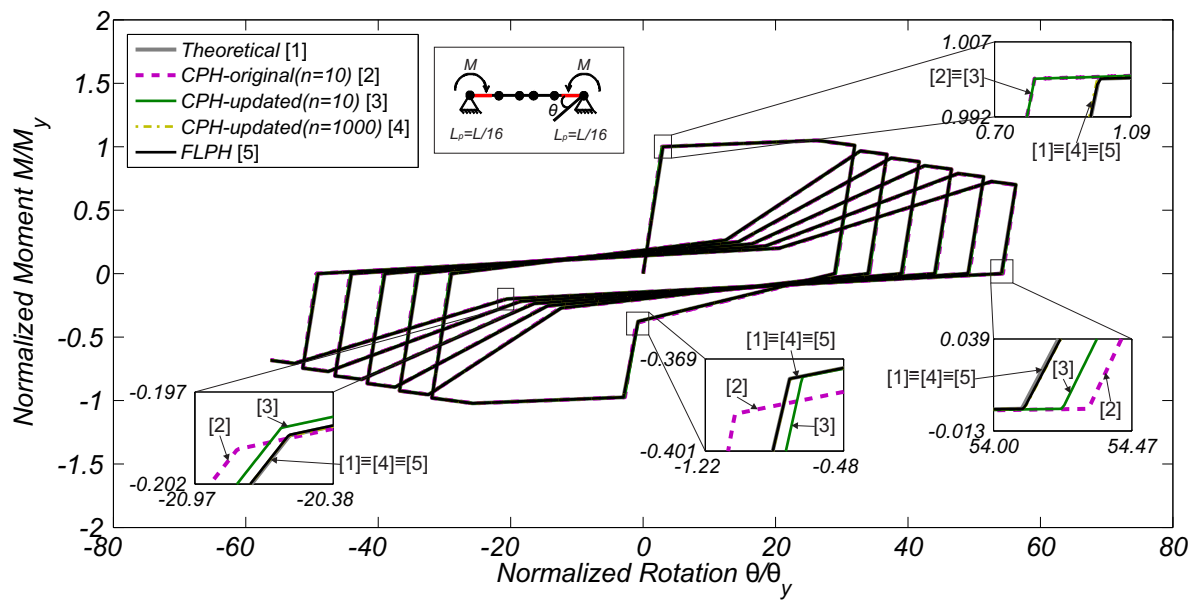


FIG. 8. Pinching model - cyclic analysis considering symmetric end moments

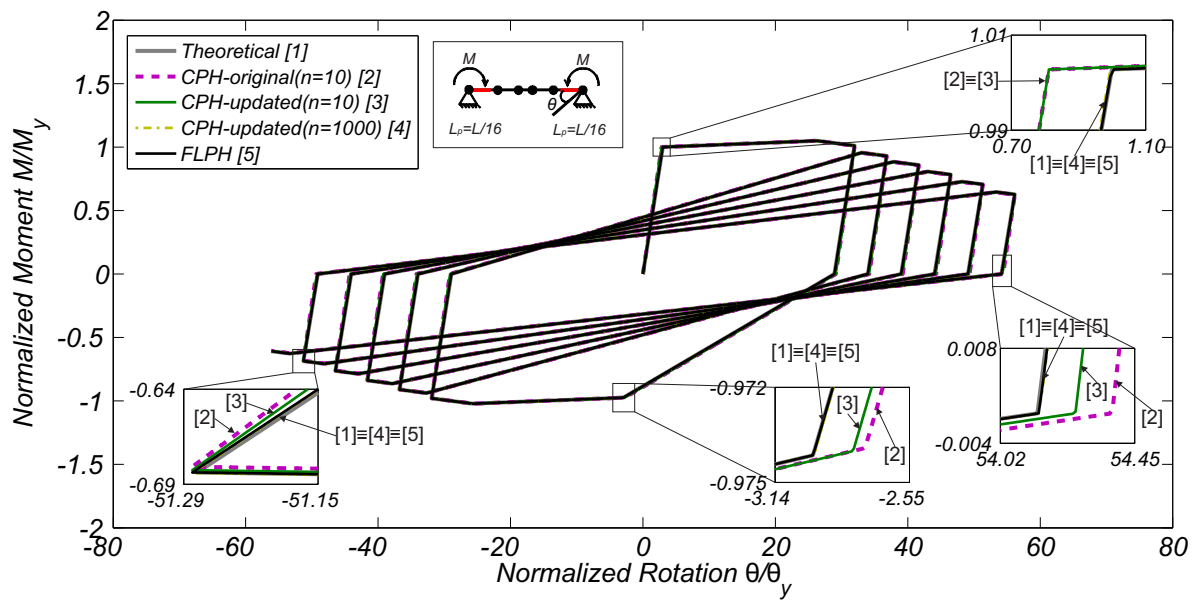


FIG. 9. Peak-oriented model - cyclic analysis considering symmetric end moments

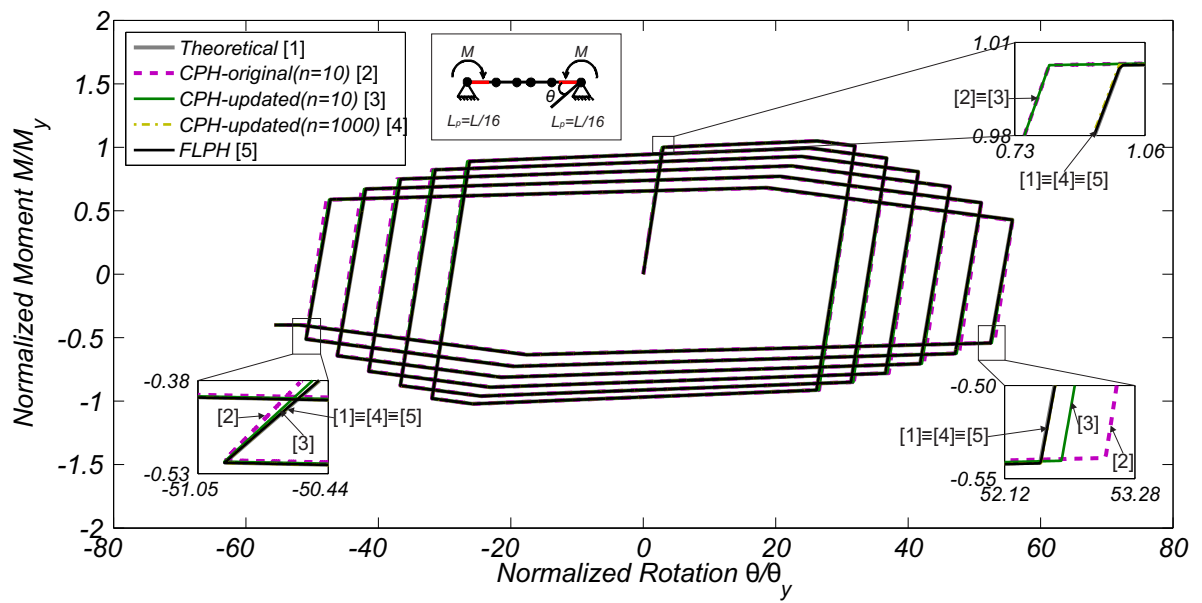


FIG. 10. Bilin model - cyclic analysis considering symmetric end moments

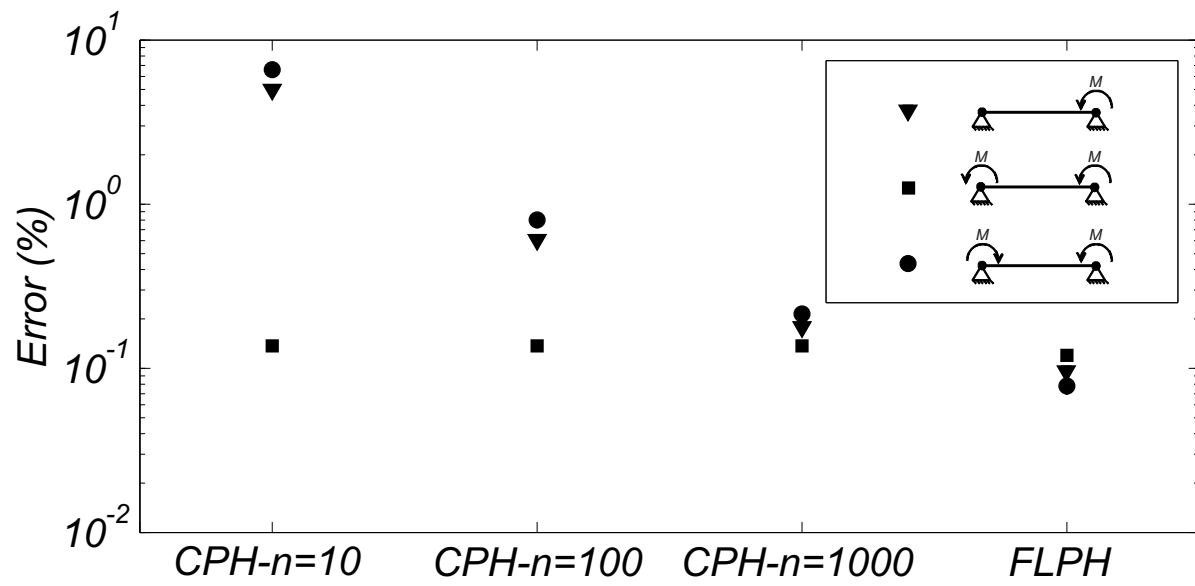


FIG. 11. Comparison of error in the elastic stiffness for CPH-updated with different values of n_{Factor} and FLPH

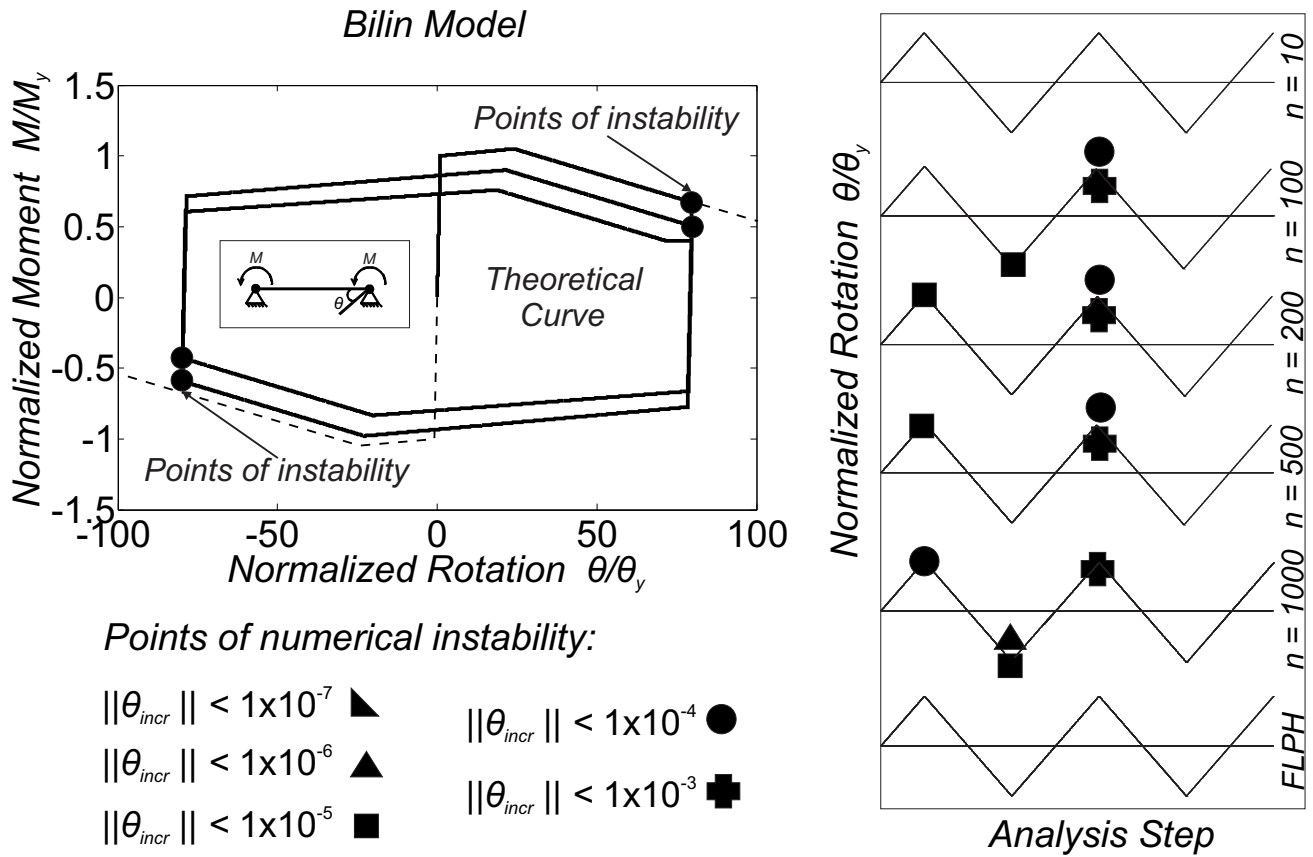


FIG. 12. Convergence stability analysis using the Bilin model