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FOURIER RESTRICTION TO CONVEX SURFACES OF REVOLUTION IN \mathbb{R}^3

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Abstract

If Γ is a C^3 hypersurface in \mathbb{R}^n and $d\sigma$ is induced Lebesgue measure on Γ , then it is well known that a Tomas-Stein Fourier restriction estimate on Γ implies that Γ has a nowhere vanishing Gaussian curvature. In a recent paper, Carbery and Ziesler observed that if induced Lebesgue measure is replaced by affine surface area, then a Tomas-Stein restriction estimate on Γ implies that Γ satisfies the affine isoperimetric inequality. Since the only property needed for a hypersurface to satisfy the affine isoperimetric inequality is convexity, this raised the question of whether a Tomas-Stein restriction estimate can be obtained for flat but convex hypersurfaces in \mathbb{R}^n such as $\Gamma(x) = (x, e^{-1/|x|^m})$, $m = 1, 2, \dots$. We prove that this is indeed the case in dimension $n = 3$.

1. Introduction

Let Γ be a C^3 hypersurface in \mathbb{R}^n and $d\sigma$ a measure on Γ . A Tomas-Stein Fourier restriction estimate for the pair $(\Gamma, d\sigma)$ is an inequality of the form

$$(1) \quad \|\widehat{f}\|_{L^2(d\sigma)} \lesssim \|f\|_{L^{\frac{2n+2}{n+3}}(\mathbb{R}^n)}$$

for $f \in C_0(\mathbb{R}^n)$.

The existence of restriction estimates such as (1), as well as their connection with the geometry of Γ , or with the decay of the Fourier transform of $d\sigma$, has been a subject of great interest. See [9, pp. 368–373] for some important applications of these estimates.

The choice of the measure $d\sigma$ is not completely arbitrary. It usually reflects some aspect of the geometry of Γ . Two important choices of $d\sigma$ are induced Lebesgue measure and affine surface area. In the former case, if Γ is assumed to have non-vanishing Gaussian curvature, (1) is a

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classical result of Tomas and Stein (see [10] and [9]). Conversely, if (1) holds with induced Lebesgue measure, then a result of Iosevich and Lu [3] (see also [2]), implies that Γ has non-vanishing Gaussian curvature. The proof of this converse uses, among other things, a Knapp-type scaling argument. To see how this argument goes, consider the special case where Γ is a surface of revolution given by $\Gamma(x) = (x, \phi(x))$, where $\phi(x) = \gamma(|x|)$, and $\gamma: [0, b) \rightarrow \mathbb{R}$ is increasing and satisfies $\gamma(0) = \gamma'(0) = 0$. For $0 < \delta < b$, let $S_\delta = \{(x, \gamma(|x|)) : |x| \leq \delta\}$ and let f_δ be a smoothed-out characteristic function of S_δ . It is then easy to see that $\|f_\delta\|_{L^2(d\sigma)} \lesssim \delta^{(n-1)/2}$, and that $|\widehat{f_\delta d\sigma}| \gtrsim \delta^{n-1}$ on a $(C/\delta) \times \cdots \times (C/\delta) \times (C/\gamma(\delta))$ box in \mathbb{R}^n (for a suitable constant C). Now if (1) holds then, by duality, the equivalent adjoint restriction estimate

$$(2) \quad \|f \widehat{d\sigma}\|_{L^{\frac{2n+2}{n-1}}(\mathbb{R}^n)} \lesssim \|f\|_{L^2(d\sigma)}$$

also holds. Applying (2) to f_δ we obtain

$$(3) \quad \delta^2 \lesssim \gamma(\delta)$$

and this implies that $\gamma''(0) \neq 0$. In particular γ cannot have vanishing Gaussian curvature at the origin. A more elaborate argument shows that the same conclusion holds in general.

In the latter case, say when $\Gamma(x) = (x, \phi(x))$, the affine surface area on Γ is given as the pushforward under Γ of the $(n-1)$ -dimensional measure $|K_\phi(x)|^{1/(n+1)} dx$, where $K_\phi(x) = \det(\text{Hess } \phi(x))$ is the affine curvature of Γ . To see what kind of geometry on Γ may be expected, take the case of a surface of revolution considered above. The radial assumption on ϕ , e.g. $\phi(x) = \gamma(|x|)$, simplifies matters and one computes that

$$K_\phi(x) = \gamma''(|x|) \left(\frac{\gamma'(|x|)}{|x|} \right)^{n-2}.$$

If we then take $d\sigma$ in the adjoint restriction estimate (2), which is equivalent to (1), to be affine surface area and use the function f_δ in it, we arrive [1] at the inequality

$$\int_0^\delta \left| \gamma''(r) \left(\frac{\gamma'(r)}{r} \right)^{n-2} \right|^{1/(n+1)} r^{n-2} dr \lesssim (\delta^{n-1} \gamma(\delta))^{(n-1)/(n+1)}.$$

But now this inequality does not imply non-vanishing curvature. Rather, it is satisfied by any convex γ , regardless of how flat it is at the origin, e.g. it is satisfied by $\gamma(t) = e^{-1/t^m}$, m any positive integer. In fact, even if ϕ is not radial, there is a similar scaling argument that can be applied, and it leads to the conclusion that ϕ satisfies the affine isoperimetric

inequality of affine differential geometry, which is certainly true whenever ϕ is convex. For more details we refer the reader to [1, pp. 409–410], [5, Chapter 5], and [6].

An earlier result of Sjölin [8] had already established that, if the dimension $n = 2$, and ϕ is convex, then the restriction inequality holds true for affine surface area. The strength of this result, along with the above considerations, suggested that, perhaps, the geometric condition of convexity of ϕ could imply a restriction result for affine surface area in higher dimensions. But if only convexity is to be used, functions such as $\phi(x) = e^{-1/|x|^m}$ have to be admitted. In attempting to prove this result, i.e. to show that convexity implies restriction, Carbery and Ziesler [1] considered the implications of a decay assumption on the Fourier transform of $d\sigma$.

Kenig, Ponce and Vega [4] proved that if the decay assumption

$$(4) \quad \left| \int_{B(0,b)} e^{-2\pi i\xi \cdot \Gamma(x)} |K_\phi(x)|^{\frac{1}{2}+i\alpha} dx \right| \lesssim \frac{(1+|\alpha|)^N}{|\xi_n|}$$

was true for all real α and some integer N , then (2) holds¹. When testing (4) on $\phi(x) = e^{-1/|x|^m}$, Carbery and Ziesler [1] found that it did not hold true in dimension $n = 3$. This, of course, did not mean that there was no restriction result for $\phi(x) = e^{-1/|x|^m}$. More recently, the same restriction question was addressed in [7]. A consequence of the results there implies that if $\phi(\cdot) = \gamma(|\cdot|)$, where γ is convex, $\gamma(0) = \gamma'(0) = 0$, $\gamma^{(3)}(t)$ non-negative, and if

$$\sup_{0 < t < b} \frac{t\gamma''(t)}{\gamma'(t)} \leq C < \infty,$$

then the restriction estimate (1) holds for affine surface area in dimension $n = 3$. Testing this last condition on $\gamma(t) = e^{-1/t^m}$, where $0 < t < b_m$, $b_m = m/(3m + 3)$, one finds that

$$\sup_{0 < t < b_m} \frac{t\gamma''(t)}{\gamma'(t)} = \sup_{0 < t < b_m} \left(\frac{m}{t^m} - m - 1 \right) = \infty.$$

Once again, the function e^{-1/t^m} was precluded from the result.

It turns out that, at least for surfaces of revolution $\Gamma(x) = (x, \phi(x))$, $\phi(x) = \gamma(|x|)$, a Tomas-Stein restriction estimate for affine surface area

¹This connection between decay and restriction is valid in dimensions $n = 2, 3$. In dimensions $n \geq 4$, one has to modify things slightly by inserting a smooth cut-off function into both (2) and (4), see [1] for further details.

does hold in the presence of convexity, if we add the condition that

$$(5) \quad \sup_{0 < t < b} \frac{\gamma(t)\gamma''(t)}{\gamma'(t)^2} \leq C < \infty.$$

Now testing this condition on $\gamma(t) = e^{-1/t^m}$ one finds that

$$(6) \quad \sup_{0 < t < b_m} \frac{\gamma(t)\gamma''(t)}{\gamma'(t)^2} = \sup_{0 < t < b_m} \left(1 - \frac{m+1}{m} t^m \right) \leq 1.$$

We thus have a Tomas-Stein restriction result that includes the surfaces $\Gamma(x) = (x, e^{-1/|x|^m})$ in \mathbb{R}^3 .

The purpose of this paper is to obtain restriction estimates for convex surfaces of revolution in \mathbb{R}^3 . A major role is played by the function

$$\frac{\gamma(t)\gamma''(t)}{\gamma'(t)^2}$$

and our results only require the boundedness of certain L^{p_0} norms of this function. In particular, we obtain a Tomas-Stein restriction estimate for surfaces of revolution in \mathbb{R}^3 satisfying (5). We find it useful to prove our results in a little more general setting. In Section 2 we introduce a family of measures $d\sigma_\gamma$, state a general (L^p, L^q) restriction result for such measures, and obtain as a corollary the result on $\Gamma(x) = (x, e^{-1/|x|^m})$. In Section 3 we present the main component of our proof. In Section 4 we prove our results.

2. Statement of results

Let $0 < b \leq \infty$, and denote by $B(0, b)$ the ball in \mathbb{R}^2 of center 0 and radius b . Let $\mathcal{C}([0, b])$ be the set of all real-valued functions $\gamma \in C^3([0, b])$ such that $\gamma(0) = \gamma'(0) = 0$, $\gamma''(t) > 0$ for $0 < t < b$, and $\gamma^{(3)}(t) \geq 0$ for $0 \leq t < b$.

Suppose $0 \leq \lambda \leq 1$, $1 \leq p$, $p_0 \leq \infty$, $4 \leq q \leq \infty$, and $1/p + 2/q \leq 1$. For $\gamma \in \mathcal{C}([0, b])$, let $d\sigma_\gamma$ be the pushforward under the map $x \rightarrow (x, \gamma(|x|))$ of the two-dimensional measure

$$(7) \quad \left(\frac{\gamma'(|x|)^{3-2\lambda} \gamma''(|x|)^\lambda}{|x| \gamma(|x|)^{1-\lambda}} \right)^{\frac{p'}{2q}} dx$$

with the understanding that when $p' = q = \infty$, $p'/(2q)$ is set to be equal to $1/4$; so that $p'/(2q) = 1/4$ on the sharp line $1/p + 2/q = 1$ including the point $(1/p, 1/q) = (1, 0)$.

Theorem 1. *If $1/p + 2/q = 1 - 1/p_0$, then*

$$(8) \quad \|\widehat{f d\sigma_\gamma}\|_{L^q(\mathbb{R}^3)} \leq C_q \left\| \left(\frac{\gamma(|\cdot|)\gamma''(|\cdot|)}{\gamma'(|\cdot|)^2} \right)^{\frac{\lambda}{2q}} \right\|_{L^{p_0}(B(0,b))} \|f\|_{L^p(d\sigma_\gamma)}$$

for all $(f, \gamma) \in C_0(\mathbb{R}^3) \times \mathcal{C}([0, b])$, where $C_q = 4(2^{7/6}\pi)^{3/(2q)}$.

Notice that if $\lambda = 1$, then the density of the measure (7) is $|K_{\gamma(|\cdot|)}(x)|^{p'/(2q)}$, so if in addition $1/p + 2/q = 1$, then $d\sigma_\gamma$ is the same affine surface area measure we described in Section 1.

Corollary 1. *Suppose $\gamma \in \mathcal{C}([0, b])$ is such that*

$$\left\| \left(\frac{\gamma(|\cdot|)\gamma''(|\cdot|)}{\gamma'(|\cdot|)^2} \right)^{\frac{1}{2q}} \right\|_{L^{p_0}(B(0,b))} < \infty.$$

Let $\lambda = 1$ and $d\sigma = d\sigma_\gamma$. If $1/p + 2/q = 1 - 1/p_0$, then

$$\|\widehat{f d\sigma}\|_{L^q(\mathbb{R}^3)} \lesssim \|f\|_{L^p(d\sigma)}$$

for all $f \in L^p(d\sigma)$.

For example if $\gamma(t) = e^{-1/t^m}$, then by (6),

$$\left\| \left(\frac{\gamma(|\cdot|)\gamma''(|\cdot|)}{\gamma'(|\cdot|)^2} \right)^{\frac{1}{2q}} \right\|_{L^{p_0}(B(0,b_m))} \leq (\pi b_m^2)^{1/p_0} < \infty$$

for $1 \leq p_0 \leq \infty$, and so the adjoint restriction estimate in Corollary 1 holds for $\gamma(t) = e^{-1/t^m}$ whenever $4 \leq q \leq \infty$ and $1/p + 2/q \leq 1$.

If, as another example, we take $\gamma(t) = -t \log(1-t)$, which is in $\mathcal{C}([0, 1])$, then

$$\left\| \left(\frac{\gamma(|\cdot|)\gamma''(|\cdot|)}{\gamma'(|\cdot|)^2} \right)^{\frac{1}{2q}} \right\|_{L^{p_0}(B(0,1))} \approx \left\| \left(-\frac{\log(1-|\cdot|)}{|\cdot|} \right)^{\frac{1}{2q}} \right\|_{L^{p_0}(B(0,1))}$$

is finite for $1 \leq p_0 < \infty$ but not for $p_0 = \infty$ (except if $q = \infty$), and so the adjoint restriction estimate in Corollary 1 holds for $\gamma(t) = -t \log(1-t)$ whenever $4 \leq q \leq \infty$ and $1/p + 2/q < 1$.

3. Main estimate

Let $\tilde{B} = B(0, b) \cap \{x = (x_1, x_2) \in \mathbb{R}^2: x_1, x_2 > 0\}$. The purpose of this section is to prove the following proposition.

Proposition 1. *Suppose $0 < b \leq \infty$ and $\gamma \in \mathcal{C}([0, b])$. Then*

$$\begin{aligned} \int_{\tilde{B}} \int_{\tilde{B}} h(u+v, \gamma(|u|) + \gamma(|v|)) \left(\frac{\gamma'(|u|)^3}{|u|\gamma(|u|)} \frac{\gamma'(|v|)^3}{|v|\gamma(|v|)} \right)^{\frac{1}{4}} du dv \\ \leq (2^{7/6}\pi)^{3/2} \|h\|_{L^1(\mathbb{R}^3)} \end{aligned}$$

for all Lebesgue measurable $h: \mathbb{R}^3 \rightarrow [0, \infty]$.

Proof: Denoting the integral on the left-hand side of the inequality by I, and changing into polar coordinates, we have

$$I = \int_0^b \int_0^b \int_0^{\frac{\pi}{2}} \int_0^{\frac{\pi}{2}} h(re^{i\theta} + se^{i\phi}, \gamma(r) + \gamma(s)) d\theta d\phi \left(\frac{r^3 \gamma'(r)^3 s^3 \gamma'(s)^3}{\gamma(r)\gamma(s)} \right)^{\frac{1}{4}} dr ds.$$

The change of variable $x = re^{i\theta} + se^{i\phi}$ (cf [7]) shows that

$$\begin{aligned} \int_0^{\frac{\pi}{2}} \int_0^{\theta} h(re^{i\theta} + se^{i\phi}, \gamma(r) + \gamma(s)) d\phi d\theta \\ \leq \int_{\sqrt{r^2+s^2} < |x| < r+s} \frac{2h(x, \gamma(r) + \gamma(s))}{\sqrt{(|x|^2 - (r-s)^2)((r+s)^2 - |x|^2)}} dx. \end{aligned}$$

So

$$\begin{aligned} \int_0^{\frac{\pi}{2}} \int_0^{\frac{\pi}{2}} h(re^{i\theta} + se^{i\phi}, \gamma(r) + \gamma(s)) d\phi d\theta \\ \leq \int_{\sqrt{r^2+s^2} < |x| < r+s} \frac{4h(x, \gamma(r) + \gamma(s))}{\sqrt{(|x|^2 - (r-s)^2)((r+s)^2 - |x|^2)}} dx \\ \leq \int_{\sqrt{r^2+s^2} < |x| < r+s} \frac{4h(x, \gamma(r) + \gamma(s))}{\sqrt{(2rs)((r+s)^2 - |x|^2)}} dx \\ \leq \int_{|x| < r+s} \frac{2h(x, \gamma(r) + \gamma(s))}{(rs)^{\frac{3}{4}} \sqrt{r+s-|x|}} dx, \end{aligned}$$

where we have used the inequality $r + s \geq 2\sqrt{rs}$. It follows that

$$\begin{aligned}
 \text{I} &\leq 2 \int_0^b \int_0^b \int_{|x| < r+s} \frac{h(x, \gamma(r) + \gamma(s))}{\sqrt{r+s-|x|}} dx \left(\frac{\gamma'(r)^3 \gamma'(s)^3}{\gamma(r)\gamma(s)} \right)^{\frac{1}{4}} dr ds \\
 &= 2 \int_{B(0,2b)} \int_0^b \int_0^b h(x, \gamma(r) + \gamma(s)) \frac{\chi_E(r, s)}{\sqrt{r+s-x}} \left(\frac{\gamma'(r)^3 \gamma'(s)^3}{\gamma(r)\gamma(s)} \right)^{\frac{1}{4}} dr ds dx \\
 &= 4 \int_{B(0,2b)} \int_0^b \int_0^b h(x, \gamma(r) + \gamma(s)) \frac{\chi_F(r, s)}{\sqrt{r+s-x}} \left(\frac{\gamma'(r)^3 \gamma'(s)^3}{\gamma(r)\gamma(s)} \right)^{\frac{1}{4}} dr ds dx \\
 &= 4 \int_{B(0,2b)} \text{II} dx,
 \end{aligned}$$

where $E = \{(r, s) \in (0, b) \times (0, b) : r + s > |x|\}$, $F = \{(r, s) \in E : s < r\}$, and

$$\text{II} = \int_0^b \int_0^b h(x, \gamma(r) + \gamma(s)) \frac{\chi_F(r, s)}{\sqrt{r+s-x}} \left(\frac{\gamma'(r)^3 \gamma'(s)^3}{\gamma(r)\gamma(s)} \right)^{\frac{1}{4}} dr ds.$$

To estimate II, we shall first apply the change of variable

$$\begin{aligned}
 r &= r(t, y) = \gamma^{-1}(y \sin^2 t) \\
 s &= s(t, y) = \gamma^{-1}(y \cos^2 t),
 \end{aligned}$$

which is defined on the open set

$$\Omega = \left\{ (t, y) \in \mathbb{R}^2 : \frac{\pi}{4} < t < \frac{\pi}{2}, y > 0 \right\};$$

so, with a slight abuse of notation, (r, s) is now a mapping from Ω to \mathbb{R}^2 . The Jacobian of this mapping is

$$J_{(r,s)}(t, y) = \frac{2y \sin t \cos^3 t + 2y \sin^3 t \cos t}{\gamma'(\gamma^{-1}(y \sin^2 t))\gamma'(\gamma^{-1}(y \cos^2 t))} = \frac{y \sin 2t}{\gamma'(r)\gamma'(s)}.$$

But²

$$(9) \quad \gamma(r) = y \sin^2 t \quad \text{and} \quad \gamma(s) = y \cos^2 t,$$

so

$$(10) \quad \gamma'(r) \frac{\partial r}{\partial t} = y \sin 2t \quad \text{and} \quad \gamma'(s) \frac{\partial s}{\partial t} = -y \sin 2t,$$

²To simplify the notation, we are writing $r, s, \partial r/\partial t$, and $\partial s/\partial t$ for $r(t, y), s(t, y), \partial r/\partial t(t, y)$, and $\partial s/\partial t(t, y)$ respectively.

and so

$$y \sin 2t = \sqrt{\gamma'(r)\gamma'(s)} \sqrt{\left| \frac{\partial r}{\partial t} \right| \left| \frac{\partial s}{\partial t} \right|}.$$

Thus

$$J_{(r,s)}(t, y) = \frac{1}{\sqrt{\gamma'(r)\gamma'(s)}} \sqrt{\left| \frac{\partial r}{\partial t} \right| \left| \frac{\partial s}{\partial t} \right|}.$$

But also

$$\frac{\gamma'(r)\gamma'(s)}{\gamma(r)\gamma(s)} \frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| = \frac{y^2 \sin^2 2t}{(y \sin^2 t)(y \cos^2 t)} = 4,$$

so

$$\left(\frac{\gamma'(r)^3 \gamma'(s)^3}{\gamma(r)\gamma(s)} \right)^{\frac{1}{4}} J_{(r,s)}(t, y) = \left(4 \frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \right)^{\frac{1}{4}}.$$

Next, to determine the domain of integration in the ty -plane, we make the following observations. By the convexity of γ , $\gamma(r) + \gamma(|x| - r)$, as a function of r , increases on the interval $(|x|/2, |x|)$. So

$$2\gamma\left(\frac{|x|}{2}\right) \leq \gamma(r) + \gamma(|x| - r) < \gamma(r) + \gamma(s)$$

whenever $|x|/2 < r < |x|$ and $|x| - r < s$, which are in turn satisfied whenever $s < r < |x| < r + s$. Also by the convexity of γ ,

$$2\gamma\left(\frac{|x|}{2}\right) \leq \gamma(|x|) \leq \gamma(r) < \gamma(r) + \gamma(s)$$

whenever $r \geq |x|$ and $s > 0$. Thus

$$2\gamma\left(\frac{|x|}{2}\right) < \gamma(r) + \gamma(s) < 2\gamma(b)$$

whenever $0 < s < r < b$ and $|x| < r + s$. But, by the definition of the mapping (r, s) ,

$$y = \gamma(r) + \gamma(s)$$

for all $(t, y) \in \Omega$, so

$$2\gamma\left(\frac{|x|}{2}\right) < y < 2\gamma(b)$$

whenever $0 < s < r < b$ and $|x| < r + s$. For any such (fixed) y , the range of (r, s) is a curve in \mathbb{R}^2 that "enters" the closure of the domain

of integration of Π when $t = \pi/4$ (i.e. when $s = r$) and “leaves” when $t = \tau(y)$ for some $\tau(y) \in (\pi/4, \pi/2]$. Thus

$$\begin{aligned} \Pi &= \int_{2\gamma(\frac{|x|}{2})}^{2\gamma(b)} \int_{\frac{\pi}{4}}^{\tau(y)} h(x, y) \frac{1}{\sqrt{r+s-|x|}} \left(4 \frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \right)^{\frac{1}{4}} dt dy \\ &= \int_{2\gamma(\frac{|x|}{2})}^{2\gamma(b)} h(x, y) \int_{\frac{\pi}{4}}^{\tau(y)} \frac{\sqrt{2}}{\sqrt{r+s-|x|}} \left(\frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \right)^{\frac{1}{4}} dt dy. \end{aligned}$$

Now, by the definition of $\tau(y)$,

$$r+s = r(t, y) + s(t, y) \geq |x| \quad \text{for} \quad \frac{\pi}{4} \leq t \leq \tau(y),$$

so, in particular,

$$r(\tau(y), y) + s(\tau(y), y) \geq |x|,$$

and hence

$$r+s-|x| \geq r+s-(r(\tau(y), y) + s(\tau(y), y)) \quad \text{for} \quad \frac{\pi}{4} < t < \tau(y).$$

Thus

$$\Pi \leq \int_{2\gamma(\frac{|x|}{2})}^{2\gamma(b)} h(x, y) \int_{\frac{\pi}{4}}^{\tau(y)} \frac{\sqrt{2}}{\sqrt{r+s-r(\tau(y), y) - s(\tau(y), y)}} \left(\frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \right)^{\frac{1}{4}} dt dy.$$

The rest of the proof will be devoted to estimating

$$\frac{1}{\sqrt{r+s-r(\tau(y), y) - s(\tau(y), y)}} \left(\frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \right)^{\frac{1}{4}}$$

for $2\gamma(|x|/2) < y < 2\gamma(b)$ and $\pi/4 < t < \tau(y)$.

We start by examining the function $\partial r/\partial t + \partial s/\partial t$. By (10),

$$\frac{\partial r}{\partial t} + \frac{\partial s}{\partial t} = \frac{y \sin 2t}{\gamma'(r)} - \frac{y \sin 2t}{\gamma'(s)}$$

is negative for $\pi/4 < t < \pi/2$ (since $\gamma'(s) < \gamma'(r)$), so

$$\begin{aligned} \left| \frac{\partial r}{\partial t} + \frac{\partial s}{\partial t} \right| &= \frac{y \sin 2t}{\gamma'(s)} - \frac{y \sin 2t}{\gamma'(r)} \\ &= 2y \left(\frac{\cos t}{\gamma'(s)} \sin t - \frac{\sin t}{\gamma'(r)} \cos t \right) \\ &= 2y \sqrt{\frac{\cos^2 t}{\gamma'(s)^2} + \frac{\sin^2 t}{\gamma'(r)^2}} \sin(t - \phi), \end{aligned}$$

where $\phi = \phi(t)$ is defined by

$$\sin \phi = \frac{(\sin t)/\gamma'(r)}{\sqrt{\frac{\cos^2 t}{\gamma'(s)^2} + \frac{\sin^2 t}{\gamma'(r)^2}}}, \quad \cos \phi = \frac{(\cos t)/\gamma'(s)}{\sqrt{\frac{\cos^2 t}{\gamma'(s)^2} + \frac{\sin^2 t}{\gamma'(r)^2}}}.$$

We shall need precise information about ϕ and $\partial^2 r/\partial t^2 + \partial^2 s/\partial t^2$. For this we need the following easy, but important, observation. By integration by parts,

$$\int_0^\rho 2\gamma'(\alpha)\gamma''(\alpha) d\alpha = 2\gamma(\rho)\gamma''(\rho) - 2 \int_0^\rho \gamma(\alpha)\gamma^{(3)}(\alpha) d\alpha$$

for $0 < \rho < b$, and since $\gamma^{(3)}$ is nonnegative, we get

$$(11) \quad \gamma'(\rho)^2 \leq 2\gamma(\rho)\gamma''(\rho) \quad \text{for } 0 < \rho < b.$$

(This is the only place where we use the assumptions that γ is C^3 and $\gamma^{(3)}$ is nonnegative; everywhere else we need only require of γ to be C^2 and convex.)

Differentiating both sides of (10) with respect to t , we have

$$\gamma''(r) \left(\frac{\partial r}{\partial t} \right)^2 + \gamma'(r) \frac{\partial^2 r}{\partial t^2} = 2y \cos 2t$$

and

$$\gamma''(s) \left(\frac{\partial s}{\partial t} \right)^2 + \gamma'(s) \frac{\partial^2 s}{\partial t^2} = -2y \cos 2t.$$

This combined with (11) gives

$$\frac{\gamma'(r)^2}{2\gamma(r)} \left(\frac{\partial r}{\partial t} \right)^2 + \gamma'(r) \frac{\partial^2 r}{\partial t^2} \leq 2y \cos 2t$$

and

$$\frac{\gamma'(s)^2}{2\gamma(s)} \left(\frac{\partial s}{\partial t} \right)^2 + \gamma'(s) \frac{\partial^2 s}{\partial t^2} \leq -2y \cos 2t.$$

But by (9) and (10),

$$\frac{\gamma'(r)^2}{2\gamma(r)} \left(\frac{\partial r}{\partial t} \right)^2 = 2y \cos^2 t \quad \text{and} \quad \frac{\gamma'(s)^2}{2\gamma(s)} \left(\frac{\partial s}{\partial t} \right)^2 = 2y \sin^2 t,$$

so

$$\gamma'(r) \frac{\partial^2 r}{\partial t^2} \leq -2y \sin^2 t \quad \text{and} \quad \gamma'(s) \frac{\partial^2 s}{\partial t^2} \leq -2y \cos^2 t,$$

and it follows that

$$(12) \quad \gamma'(r) \left(\frac{\partial^2 r}{\partial t^2} + \frac{\partial^2 s}{\partial t^2} \right) \leq -2y < 0$$

for $\pi/4 < t < \pi/2$.

Going back to ϕ , we have

$$\tan \phi = \frac{\gamma'(s)}{\gamma'(r)} \tan t.$$

Now if we let $m(\rho) = \gamma'(\rho)^2/\gamma(\rho)$, $0 < \rho < b$, then by (11), $m'(\rho) \geq 0$ and it follows that $m(s) \leq m(r)$. Hence

$$\tan \phi = \frac{\sqrt{m(s)\gamma(s)}}{\sqrt{m(r)\gamma(r)}} \tan t \leq \frac{\sqrt{\gamma(s)}}{\sqrt{\gamma(r)}} \tan t = \cot t \tan t = 1,$$

and hence $0 < \phi \leq \pi/4$. Thus

$$\begin{aligned} \left| \frac{\partial r}{\partial t} + \frac{\partial s}{\partial t} \right| &\geq 2y \sqrt{\frac{\cos^2 t}{\gamma'(s)^2} + \frac{\sin^2 t}{\gamma'(r)^2}} \sin\left(t - \frac{\pi}{4}\right) \\ &= 2y \sqrt{\frac{\cos^2 t}{(y \sin 2t)^2} \left(\frac{\partial s}{\partial t}\right)^2 + \frac{\sin^2 t}{(y \sin 2t)^2} \left(\frac{\partial r}{\partial t}\right)^2} \sin\left(t - \frac{\pi}{4}\right) \\ &= \sqrt{\frac{1}{\sin^2 t} \left(\frac{\partial s}{\partial t}\right)^2 + \frac{1}{\cos^2 t} \left(\frac{\partial r}{\partial t}\right)^2} \sin\left(t - \frac{\pi}{4}\right) \\ &> \sqrt{\left(\frac{\partial r}{\partial t}\right)^2 + \left(\frac{\partial s}{\partial t}\right)^2} \sin\left(t - \frac{\pi}{4}\right) \\ &\geq \sqrt{2} \frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \left(\frac{2}{\pi}\right) \left(t - \frac{\pi}{4}\right) \end{aligned}$$

for $\pi/4 < t < \pi/2$. Thus

$$\frac{\sqrt{\frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right|}}{\left| \frac{\partial r}{\partial t} + \frac{\partial s}{\partial t} \right|} < \frac{\pi}{2\sqrt{2}} \frac{1}{t - \frac{\pi}{4}}$$

for $\pi/4 < t < \pi/2$.

As we saw above, $\partial r/\partial t + \partial s/\partial t$ is negative on the interval $(\pi/4, \pi/2)$. Also by (12),

$$\frac{\partial^2 r}{\partial t^2} + \frac{\partial^2 s}{\partial t^2} < 0,$$

so $\partial r/\partial t + \partial s/\partial t$, as a function of t , is decreasing on $(\pi/4, \pi/2)$, and so $|\partial r/\partial t + \partial s/\partial t|$ is increasing there. Now applying the mean value theorem, we obtain

$$r + s - r(\tau(y), y) - s(\tau(y), y) \geq (\tau(y) - t) \left| \frac{\partial r}{\partial t} + \frac{\partial s}{\partial t} \right|$$

for $\pi/4 < t < \tau(y)$. Thus

$$\begin{aligned} \frac{\left(\frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \right)^{\frac{1}{4}}}{\sqrt{r + s - r(\tau(y), y) - s(\tau(y), y)}} &\leq \frac{1}{\sqrt{\tau(y) - t}} \frac{\left(\frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \right)^{\frac{1}{4}}}{\left| \frac{\partial r}{\partial t} + \frac{\partial s}{\partial t} \right|^{\frac{1}{2}}} \\ &< \sqrt{\frac{\pi}{2\sqrt{2}}} \frac{1}{\sqrt{\tau(y) - t}} \frac{1}{\sqrt{t - \frac{\pi}{4}}} \end{aligned}$$

for $\pi/4 < t < \tau(y)$. Thus

$$\begin{aligned} &\int_{\frac{\pi}{4}}^{\tau(y)} \frac{\sqrt{2}}{\sqrt{r + s - r(\tau(y), y) - s(\tau(y), y)}} \left(\frac{\partial r}{\partial t} \left| \frac{\partial s}{\partial t} \right| \right)^{\frac{1}{4}} dt \\ &\leq \sqrt{\frac{\pi}{2}} \int_{\frac{\pi}{4}}^{\tau(y)} \frac{dt}{\sqrt{(\tau(y) - t)(t - \frac{\pi}{4})}} \\ &= \frac{\pi^{3/2}}{2^{1/4}}. \end{aligned}$$

Thus

$$\text{II} \leq \frac{\pi^{3/2}}{2^{1/4}} \int_{2^{\gamma(\frac{1-x}{2})}}^{2^{2\gamma(b)}} h(x, y) dy$$

and consequently

$$\text{I} \leq 2^{7/4} \pi^{3/2} \int_{B(0, 2b)} \int_{2^{\gamma(\frac{1-x}{2})}}^{2^{2\gamma(b)}} h(x, y) dy dx \leq (2^{7/6} \pi)^{3/2} \|h\|_{L^1(\mathbb{R}^3)}. \quad \square$$

4. Proof of Theorem 1

Let f be a continuous function on \mathbb{R}^3 which is compactly supported in the third variable, and let $\gamma \in \mathcal{C}([0, b])$. It is enough to show that

$$\|\widehat{f d\sigma}\|_{L^q(\mathbb{R}^3)} \leq (2^{7/6} \pi)^{3/(2q)} \left\| \left(\frac{\gamma(|\cdot|) \gamma''(|\cdot|)}{\gamma'(|\cdot|)^2} \right)^{\frac{\lambda}{2q}} \right\|_{L^{p_0}(B(0, b))} \|f\|_{L^p(d\sigma)},$$

where $d\sigma = \chi_E d\sigma_\gamma$ and $E = \{(x_1, x_2, x_3) \in \mathbb{R}^3 : x_1, x_2 \geq 0\}$. If $q = \infty$, then this follows easily from Hölder's inequality. So we may assume

$q < \infty$. Then the relation $1/p + 2/q = 1 - 1/p_0$ tells us that $p, p_0 > 1$. Also, since

$$\|\widehat{f d\sigma}\|_{L^q(\mathbb{R}^3)} = \|\widehat{f d\sigma} \widehat{f d\sigma}\|_{L^{q/2}(\mathbb{R}^3)}^{1/2} = \|f d\sigma * f d\sigma\|_{L^{q/2}(\mathbb{R}^3)}^{1/2},$$

and since $q/2 \geq 2$, it is enough by the Hausdorff-Young inequality to establish that

$$\begin{aligned} & \int h|f|d\sigma * |f|d\sigma \\ & \leq (2^{7/6}\pi)^{3/q} \left\| \left(\frac{\gamma(|\cdot|)\gamma''(|\cdot|)}{\gamma'(|\cdot|)^2} \right)^{\frac{\lambda}{2q}} \right\|_{L^{p_0}(B(0,b))}^2 \|f\|_{L^p(d\sigma)}^2 \|h\|_{L^{q/2}(\mathbb{R}^3)} \end{aligned}$$

for any nonnegative Lebesgue measurable function h on \mathbb{R}^3 . But by Hölder's inequality,

$$\int h|f|d\sigma * |f|d\sigma \leq \|f\|_{L^p(d\sigma)}^2 \|h\|_{L^{p'}(d\sigma * d\sigma)},$$

so we need to have

$$\|h\|_{L^{p'}(d\sigma * d\sigma)} \leq (2^{7/6}\pi)^{3/q} \left\| \left(\frac{\gamma(|\cdot|)\gamma''(|\cdot|)}{\gamma'(|\cdot|)^2} \right)^{\frac{\lambda}{2q}} \right\|_{L^{p_0}(B(0,b))}^2 \|h\|_{L^{q/2}(\mathbb{R}^3)}.$$

Now this follows from Proposition 1 by writing

$$\begin{aligned} & \|h\|_{L^{p'}(d\sigma * d\sigma)}^{p'} \\ & = \int_{\bar{B}} \int_{\bar{B}} h^{p'}(x+y, \gamma(|x|)+\gamma(|y|)) M(x)^{\frac{p'}{2q}} M(y)^{\frac{p'}{2q}} \frac{N(x)^{\frac{\lambda p'}{2q}} N(y)^{\frac{\lambda p'}{2q}}}{N(x)^{\frac{\lambda p'}{2q}} N(y)^{\frac{\lambda p'}{2q}}} dx dy, \end{aligned}$$

where

$$M(\cdot) = \frac{\gamma'(|\cdot|)^{3-2\lambda} \gamma''(|\cdot|)^\lambda}{|\cdot| \gamma(|\cdot|)^{1-\lambda}} \quad \text{and} \quad N(\cdot) = \frac{\gamma(|\cdot|)\gamma''(|\cdot|)}{\gamma'(|\cdot|)^2},$$

and applying Hölder's inequality to get

$$\begin{aligned}
& \|h\|_{L^{p'}(d\sigma*d\sigma)}^{p'} \\
& \leq \left(\int_{\tilde{B}} \int_{\tilde{B}} N(x)^{\frac{\lambda r p'}{2q}} N(y)^{\frac{\lambda r p'}{2q}} dx dy \right)^{\frac{1}{r}} \\
& \quad \times \left(\int_{\tilde{B}} \int_{\tilde{B}} h^{q/2}(x+y, \gamma(|x|) + \gamma(|y|)) \frac{M(x)^{\frac{1}{4}} M(y)^{\frac{1}{4}}}{N(x)^{\frac{1}{4}} N(y)^{\frac{1}{4}}} dx dy \right)^{\frac{2p'}{q}} \\
& = \left(\int_{\tilde{B}} N(x)^{\frac{\lambda p_0}{2q}} dx \right)^{\frac{2}{r}} \\
& \quad \times \left(\int_{\tilde{B}} \int_{\tilde{B}} h^{q/2}(x+y, \gamma(|x|) + \gamma(|y|)) \left(\frac{\gamma'(|x|)^3}{|x|\gamma(|x|)} \frac{\gamma'(|y|)^3}{|y|\gamma(|y|)} \right)^{\frac{1}{4}} dx dy \right)^{\frac{2p'}{q}} \\
& \leq \left(\int_{\tilde{B}} N(x)^{\frac{\lambda p_0}{2q}} dx \right)^{\frac{2p'}{p_0}} \left((2^{7/6}\pi)^{3/2} \|h^{q/2}\|_{L^1(\mathbb{R}^3)} \right)^{\frac{2p'}{q}} \\
& \leq (2^{7/6}\pi)^{3p'/q} \left\| \left(\frac{\gamma(|\cdot|)\gamma''(|\cdot|)}{\gamma'(|\cdot|)^2} \right)^{\frac{\lambda}{2q}} \right\|_{L^{p_0}(B(0,b))}^{2p'} \|h\|_{L^{q/2}(\mathbb{R}^3)}^{p'},
\end{aligned}$$

where r is the dual exponent to $q/(2p')$ (so that $rp' = p_0$).

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