

Restricted independence in displacement function for better estimation of cyclicity ^{*}

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Abstract

Since the independence of focal values is a sufficient condition to give a number of limit cycles arising from a center-focus equilibrium, in this paper we consider a restricted independence to a parametric curve, which gives a method not only to increase the lower bound for the cyclicity of the center-focus equilibrium but also to be available when those focal values are not independent. We apply the method to a nondegenerate cubic center-focus variety and prove that the cyclicity reaches its an upper bound.

Keywords: center-focus variety; cyclicity; focal value; independence; power sequence.

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1 Introduction

As it appears in the books [3, 5, 19] and articles [4, 6, 13, 14, 18, 20], the discussion on the center-focus equilibria is one of the most important problems in ordinary differential equations. A center-focus equilibrium is an equilibrium at which the linear part of the differential system has a pair of nonvanished pure imaginary eigenvalues. The main interest on the research of these equilibria is the determination of the kind

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of their stability and on their *cyclicity* at this equilibrium, i.e., the greatest number of limit cycles which may arise from a Hopf bifurcation at these equilibria.

We generally consider the family of analytic systems

$$\dot{x} = \alpha x - y + P(x, y, \boldsymbol{\lambda}), \quad \dot{y} = x + \alpha y + Q(x, y, \boldsymbol{\lambda}) \quad (1.1)$$

with a standardized linear part, where $\alpha \in \mathbb{R}$ and $\boldsymbol{\lambda} := (\lambda_1, \dots, \lambda_m) \in \mathbb{R}^m$ are parameters, $P(x, y, \boldsymbol{\lambda})$ and $Q(x, y, \boldsymbol{\lambda})$ are real analytic functions of x and y starting at least with terms of degree two and depending polynomially on $\boldsymbol{\lambda}$. This family, denoted by $LC(\alpha, \boldsymbol{\lambda})$ for short, is referred as a family of linear centers, which has either a center or a focus at the origin $O : (0, 0)$ clearly. Its *center-focus variety* is

$$\mathcal{CF} := \{(\alpha, \boldsymbol{\lambda}) \in \mathbb{R} \times \mathbb{R}^m : \alpha = 0\} \cong \mathbb{R}^m.$$

In this variety the so-called center-focus equilibrium O needs to be identified between focus (called a *weak focus*) and center, which is decided if using finitely many focal values (see [3] for more details). Focal values come from the coefficients of the displacement function $\Pi(\rho) := h(\rho) - \rho$, where h is the Poincaré return map

$$h(\rho) := e^{2\pi\alpha}\rho + \sum_{i=2}^{+\infty} g_i(\alpha, \boldsymbol{\lambda})\rho^i \quad (1.2)$$

and g_i 's are analytic functions of α and $\boldsymbol{\lambda}$ such that $g_2(0, \boldsymbol{\lambda}) \equiv 0$. Let $g_i(\boldsymbol{\lambda}) := g_i(0, \boldsymbol{\lambda})$ for all $i = 2, 3, \dots$. Since P, Q are assumed to be polynomially dependent on $\boldsymbol{\lambda}$, $g_i \in \mathbb{R}[\boldsymbol{\lambda}]$ (the ring of real polynomials in the variable $\boldsymbol{\lambda}$) for all $i \geq 2$ and $g_{2k} \in \langle g_3, \dots, g_{2k-1} \rangle$ (the ideal generated by g_3, \dots, g_{2k-1} over the ring $\mathbb{R}[\boldsymbol{\lambda}]$) for all $k \geq 2$, as indicated in [1, 3]. Those g_{2i+1} 's are called the *focal values*, which are algebraically equivalent to the Lyapunov quantities ([15, 16]). Note that \mathcal{CF} contains the subset

$$\mathcal{C} := V(g_3, \dots, g_{2i+1}, \dots),$$

where $V(g_3, \dots, g_{2i+1}, \dots)$ denotes the algebraic variety of $\langle g_3, \dots, g_{2i+1}, \dots \rangle$ and, by [7, p. 3], is actually the set of all common zeros of all g_{2i+1} 's ($i \geq 1$). As in [3, p. 11], \mathcal{C} is called the *center variety* of the family $LC(\alpha, \boldsymbol{\lambda})$ because the center-focus O of system $LC(0, \boldsymbol{\lambda}')$ is a *center* if and only if $\boldsymbol{\lambda}' \in \mathcal{C}$. For any $\boldsymbol{\lambda}' \in \mathbb{R}^m \setminus \mathcal{C}$ there exists an integer $k \geq 1$ such that $g_{2k+1}(\boldsymbol{\lambda}') \neq 0$ and $g_i(\boldsymbol{\lambda}') = 0$ for all $i = 2, \dots, 2k$, for which we call O a *weak focus of multiplicity k* in system $LC(0, \boldsymbol{\lambda}')$.

As usual the *cyclicity* of a center-focus is the maximal number of limit cycles emerging from it in the phase portrait when we change slightly the parameters of the system (see [3, 8, 10, 17] and references therein). More precisely, the greatest number of limit cycles bifurcated from O is called the *cyclicity* of system $LC(0, \boldsymbol{\lambda}')$ at O (perturbed within the family $LC(\alpha, \boldsymbol{\lambda})$) and denoted by $\mathcal{N}(\boldsymbol{\lambda}')$. In particular, $\mathcal{N}(\boldsymbol{\lambda}')$ is denoted by $\mathcal{N}_c(\boldsymbol{\lambda}')$ (resp. $\mathcal{N}_f(\boldsymbol{\lambda}')$) and called *center cyclicity* (resp. *focus cyclicity*) of system $LC(0, \boldsymbol{\lambda}')$ if $\boldsymbol{\lambda}' \in \mathcal{C}$ (resp. $\boldsymbol{\lambda}' \in \mathbb{R}^m \setminus \mathcal{C}$).

The cyclicity $\mathcal{N}(\boldsymbol{\lambda}')$ is decided not only by the multiplicity of the center-focus equilibrium, but also by the greatest number of independent sign changes in the displacement function near $(0, \boldsymbol{\lambda}')$. For multi-parametric families there are more difficulties in finding the greatest number of nonvanished focal values, but one usually gives its a lower bound. For such an independence, a well-known method is to check the following conditions:

(ID_k-1) every neighborhood of $\boldsymbol{\lambda}'$ contains a $\boldsymbol{\mu}' \in V(g_3, \dots, g_{2k-1})$ such that $g_{2k+1}(\boldsymbol{\mu}') \neq 0$, and

(ID_k-2) for each positive integer $\ell \leq k - 1$ and each $\boldsymbol{\mu}' \in V(g_3, \dots, g_{2\ell+1})$ satisfying that $g_{2\ell+3}(\boldsymbol{\mu}') \neq 0$, every neighborhood of $\boldsymbol{\mu}'$ contains a $\boldsymbol{\mu}'' \in V(g_3, \dots, g_{2\ell-1})$ such that $g_{2\ell+1}(\boldsymbol{\mu}'')g_{2\ell+3}(\boldsymbol{\mu}') < 0$.

Note that there is only **(ID₁-1)** if $k = 1$. As indicated in [2], **(ID_k-1)** and **(ID_k-2)** are known as conditions for the first k focal values g_{2j+1} , $j = 1, \dots, k$, to be independent, under which $\mathcal{N}(\boldsymbol{\lambda}') \geq k$ as shown in [12]. Another method ([3]) is to determine the rank $r(\boldsymbol{\lambda}')$ of the Jacobian matrix

$$\left. \frac{\partial(g_3, \dots, g_{2j+1})}{\partial \boldsymbol{\lambda}} \right|_{\boldsymbol{\lambda}=\boldsymbol{\lambda}'},$$

where all g_3, \dots, g_{2j+1} vanish at $\boldsymbol{\lambda}'$, which asserts $\mathcal{N}(\boldsymbol{\lambda}') \geq r(\boldsymbol{\lambda}')$ because it implies the existence of an independent subsequence of $r(\boldsymbol{\lambda}')$ members in $\{g_3, \dots, g_{2j+1}\}$ as indicated in ([2]). However, it is not easy to verify the independence of focal values or compute the rank of the Jacobian with many parameters. Besides, conditions **(ID_k-1)** and **(ID_k-2)** are strong sufficient conditions for the independence, which remind us to find weaker ones. The rank of the Jacobian gives a lower bound for $\mathcal{N}(\boldsymbol{\lambda}')$, but this bound may not be the best.

In this paper we give a method to increase the lower bound for the cyclicity $\mathcal{N}(\boldsymbol{\lambda}')$ of system (1.1). The method is to find an appropriate curve passing through $(0, \boldsymbol{\lambda}')$ in the space \mathbb{R}^{m+1} of parameters $(\alpha, \boldsymbol{\lambda})$, on which those focal values depend on a single variable in such a way that we can determine easily the number of independent sign changes in the displacement function on the curve, called an *restricted independence* to the curve. Such a restriction may give a larger number of independent sign changes than the rank of the Jacobian, from which we can find more limit cycles bifurcating from the origin O . The result about this method is given in section 2. In section 3 we give some corollaries for easier applications, and practical application of our method with an example. This example has a nonvanished 4th order focal value but does not satisfy the independence condition of focal values, from which one cannot assert that the cyclicity of O is 4. However, using our method, we prove that the cyclicity is exactly 4. Finally, in section 4 we apply the method to a five-parametric family of cubic systems for finding its $\mathcal{N}(\boldsymbol{\lambda}')$.

2 Restricted independence

For $\boldsymbol{\lambda}' \in \mathbb{R}^m \setminus \mathcal{C}$ the equilibrium O is a weak focus of system $LC(0, \boldsymbol{\lambda}')$. Let $\zeta(\boldsymbol{\lambda}')$ be the multiplicity of the weak focus. Then $\zeta(\boldsymbol{\lambda}')$ gives an upper estimate for $\mathcal{N}_f(\boldsymbol{\lambda}')$ because

$$g_{2\zeta(\boldsymbol{\lambda}')+1}(\boldsymbol{\lambda}') \neq 0, \quad \text{and} \quad g_i(\boldsymbol{\lambda}') = 0 \quad \forall i < 2\zeta(\boldsymbol{\lambda}') + 1, \quad (2.1)$$

and the sequence of g_{2k+1} 's ($k = 1, \dots, \zeta(\boldsymbol{\lambda}')$) may not be independent. In contrast, for $\boldsymbol{\lambda}' \in \mathcal{C}$, equilibrium O is a center of system $LC(0, \boldsymbol{\lambda}')$. Since $\mathbb{R}\{\boldsymbol{\lambda}\}_{\boldsymbol{\lambda}'}$, the ring of convergent power series at $\boldsymbol{\lambda}'$, is a Noetherian ring (see [9, p. 147]), every ideal in this ring is finitely generated, which implies the existence of a least integer $\iota(\boldsymbol{\lambda}') > 0$ satisfying that

$$\langle g_3, g_5, \dots, g_{2\iota(\boldsymbol{\lambda}')+1} \rangle_{\boldsymbol{\lambda}'} = \langle g_3, \dots, g_{2i+1}, \dots \rangle_{\boldsymbol{\lambda}'} \quad \text{in } \mathbb{R}\{\boldsymbol{\lambda}\}_{\boldsymbol{\lambda}'}. \quad (2.2)$$

The integer $\iota(\boldsymbol{\lambda}')$, called the *multiplicity of the center* O , gives an upper estimate for $\mathcal{N}_c(\boldsymbol{\lambda}')$. In general, we see that the field \mathbb{R} is a commutative Noetherian ring, which implies by the Hilbert Basis Theorem ([9, p. 144]) that $\mathbb{R}[\boldsymbol{\lambda}]$ is a Noetherian ring. Therefore, there exists the least integer $\iota_p \geq 1$ such that

$$\langle g_3, g_5, \dots, g_{2\iota_p+1} \rangle = \langle g_3, \dots, g_{2i+1}, \dots \rangle \quad \text{in } \mathbb{R}[\boldsymbol{\lambda}]. \quad (2.3)$$

It follows that $V(g_3, \dots, g_{2\iota_p+1}) = V(g_3, \dots, g_{2i+1}, \dots)$ and $\max\{\zeta(\boldsymbol{\lambda}'), \iota(\boldsymbol{\lambda}')\} \leq \iota_p$.

For $\boldsymbol{\lambda}' \in \mathbb{R}^m$ we need to discuss the sign changes among those focal values g_{2i+1} and the real part α of the eigenvalues near $(\alpha, \boldsymbol{\lambda}) = (0, \boldsymbol{\lambda}')$. For convenience, define

$$g_1(\alpha) := 2\pi\alpha$$

complementarily. Our strategy is to restrict those g_{2i+1} 's ($i = 0, 1, \dots, \kappa(\boldsymbol{\lambda}')$), where $\kappa(\boldsymbol{\lambda}') = \zeta(\boldsymbol{\lambda}')$ (or $\iota(\boldsymbol{\lambda}')$) if $\boldsymbol{\lambda}' \in \mathbb{R}^m \setminus \mathcal{C}$ (or \mathcal{C}), to a curve in the space \mathbb{R}^{m+1} of parameters $(\alpha, \boldsymbol{\lambda})$ to see independent sign changes in the displacement function. Consider a continuous curve Υ of the form

$$(\alpha(\eta), \boldsymbol{\lambda}(\eta)) := (d_0\eta^{\alpha_0}, \lambda'_1 + d_1\eta^{\alpha_1}, \dots, \lambda'_m + d_m\eta^{\alpha_m}) \quad (2.4)$$

in the parameter space \mathbb{R}^{m+1} , where $d_i \in \mathbb{R}$ and $\alpha_i > 0$ are indeterminate constants, $i = 0, \dots, m$, and $(\lambda'_1, \dots, \lambda'_m) = \boldsymbol{\lambda}'$. Clearly, $(\alpha(0), \boldsymbol{\lambda}(0)) = (0, \boldsymbol{\lambda}')$, i.e., the curve passes through the point $(0, \boldsymbol{\lambda}')$. The curve is of polynomial form if all α_i 's are positive integers. Restricted to the curve Υ given in (2.4), the focal values are of the form

$$g_1(\alpha(\eta)) = c_0\eta^{w_0} + o(\eta^{w_0}), \quad g_{2i+1}(\boldsymbol{\lambda}(\eta)) = c_i\eta^{w_i} + o(\eta^{w_i}), \quad i = 1, \dots, \kappa(\boldsymbol{\lambda}'), \quad (2.5)$$

where w_i 's are positive constants depending on the α_i 's and the c_i 's are real constants depending on the d_i 's such that, for each i , $c_i \neq 0$ if and only if $g_{2i+1}(\boldsymbol{\lambda}(\eta)) \neq 0$. In particular,

$$c_{\kappa(\boldsymbol{\lambda}')} = g_{2\kappa(\boldsymbol{\lambda}')+1}(\boldsymbol{\lambda}') \neq 0 \quad \text{and} \quad w_{\kappa(\boldsymbol{\lambda}')} = 0$$

if $\boldsymbol{\lambda}' \in \mathbb{R}^m \setminus \mathcal{C}$.

Our method of restricted independence highly depends on the evolution in the power sequence $w := \{w_i\}$, where the w_i 's are reals unless that all α_i 's are chosen as integers, because then the w_i 's are integers. It is worthy mentioning that the sequence $\{w_i\}$ may not be increasing although the corresponding g_{2i+1} is given by the coefficient of the term ρ^{2i+1} in the return map (1.2). Let

$$\Delta(i, j)w := \frac{w_i - w_j}{i - j},$$

called the *difference quotient of w between i and j* . For $i_0 < \dots < i_k$ in $\{0, 1, \dots, \kappa(\boldsymbol{\lambda}')\}$, the power sequence $\{w_i\}$ is said to be *ladder-likely degressive* on the scale (i_0, \dots, i_k) if there are constants $h_{i_0} > h_{i_1} > \dots > h_{i_k} > 0$, called the *degressive rates*, such that

(LD) for each $\nu = 0, \dots, k$,

$$\Delta(i_\nu, j)w \begin{cases} \leq -h_{i_\nu} & \forall j = 0, \dots, i_\nu - 1, \\ \geq -h_{i_\nu} & \forall j = i_\nu + 1, \dots, \kappa(\boldsymbol{\lambda}'). \end{cases}$$

Obviously, the sequence $\{7, 4, 2, 1\}$ is ladder-likely degressive on the scale $(0, 1, 2, 3)$, where we note that $\{3, 2, 1\}$, the sequence of differences between two consecutive terms, is strictly decreasing and we can choose the sequence $\{6, 5/2, 3/2, 1/2\}$ for degressive rates. Note that the concept of ladder-like degressiveness does not require the sequence $\{w_i\}$ to be decreasing but needs the existence of a decreasing subsequence of $\{w_i\}$ with weaker and weaker degressive rates correspondingly. For example, the sequence $\{7, 4, 2, 4, 1\}$ does not decrease but has a ladder-likely degressive scale $(0, 1, 4)$ with the sequence $(14, 5/2, 1/4)$ of degressive rates.

Considering the “=” in (LD), for each i_ν define

$$\Xi(i_\nu) := \{i_\nu\} \cup \{j \in \{0, \dots, \kappa(\boldsymbol{\lambda}')\} : \Delta(i_\nu, j)w = -h_{i_\nu}\},$$

the set of all j 's having the same slope $-h_{i_\nu}$ with respect to i_ν . Let

$$\begin{aligned} \hat{c}_{i_\nu} &:= \sum_{j \in \Xi(i_\nu)} c_j, \\ \mathcal{V} &:= \{i_\nu \in \{i_0, \dots, i_k\} : \exists j \in \{\nu + 1, \dots, k\} \text{ such that } \hat{c}_{i_\nu} \hat{c}_{i_j} < 0 \\ &\quad \text{and } \hat{c}_{i_l} = 0 \forall l = \nu + 1, \dots, j - 1\}, \end{aligned} \tag{2.6}$$

where c_i is the leading coefficient of g_{2i+1} as given in (2.5). Clearly \mathcal{V} is a set of indices for independent sign changes, i.e. an independence restricted to the parameterized curve Υ .

Theorem 2.1. *Suppose that the power sequence of g_{2i+1} 's, the focal values of family $LC(\alpha, \boldsymbol{\lambda})$ given in (1.1) near $(0, \boldsymbol{\lambda}')$, restricted to the parameterized curve (2.4) is ladder-likely degressive on (i_0, \dots, i_k) , i.e., condition (LD) holds. Then $\mathcal{N}(\boldsymbol{\lambda}') \geq \#\mathcal{V}$, the cardinality of the set \mathcal{V} .*

Proof. Let i_s be the greatest member of \mathcal{V} and $s' \in \{s+1, \dots, k\}$ be the corresponding j given in the definition of \mathcal{V} . Clearly,

$$i_s < i_{s'} \quad \text{and} \quad \hat{c}_{i_s} \hat{c}_{i_{s'}} < 0.$$

From (1.2) we get

$$\begin{aligned} \Pi(\rho) &= g_1(\alpha)\rho(1 + \Psi_0(\alpha)) + \sum_{i=2}^{+\infty} g_i(\alpha, \boldsymbol{\lambda})\rho^i \\ &= \sum_{i=0}^{+\infty} g_{2i+1}\rho^{2i+1} (1 + \Psi_{2i+1}(\alpha, \boldsymbol{\lambda}, \rho)), \end{aligned} \quad (2.7)$$

where $\Psi_0(\alpha)$ is an analytic function such that $\Psi_0(0) = 0$, and the functions Ψ_{2i+1} 's are analytic in $(\alpha, \boldsymbol{\lambda}, \rho)$ and vanish at $(0, \boldsymbol{\lambda}, 0)$. On the other hand, by **(LD)** we get $w_{i_\nu} - w_j \leq -(i_\nu - j)h_{i_\nu}$, i.e. $w_{i_\nu} + i_\nu h_{i_\nu} \leq w_j + j h_{i_\nu}$ for all $j = 0, \dots, \kappa(\boldsymbol{\lambda}')$. This implies that for each $\nu \in \{0, \dots, k\}$ we have

$$w_{i_\nu} + (i_\nu + 1)h_{i_\nu} \leq w_j + (j + 1)h_{i_\nu} \quad \forall j = 0, \dots, \kappa(\boldsymbol{\lambda}'). \quad (2.8)$$

In what follows we use (2.7) and (2.8) to discuss in $\mathbb{R}^m \setminus \mathcal{C}$ and \mathcal{C} separately.

For $\boldsymbol{\lambda}' \in \mathbb{R}^m \setminus \mathcal{C}$, by (2.3) and (2.7) we get

$$\Pi(\rho) = \sum_{i=0}^{\iota_p} g_{2i+1}\rho^{2i+1} (1 + \Phi_{2i+1}(\alpha, \boldsymbol{\lambda}, \rho)), \quad (2.9)$$

where Φ_{2i+1} 's are analytic at $(\alpha, \boldsymbol{\lambda}, \rho)$ and vanish when $\alpha = \rho = 0$. Restricted to the curve $(\alpha, \boldsymbol{\lambda}) = (\alpha(\eta), \boldsymbol{\lambda}(\eta))$, from (2.9) we obtain

$$\begin{aligned} \Pi(\rho) &= \sum_{i=0}^{\zeta(\boldsymbol{\lambda}')} c_i \eta^{w_i} \rho^{2i+1} (1 + H_{2i+1}(\eta, \rho)) \\ &+ \begin{cases} \sum_{i=\zeta(\boldsymbol{\lambda}')+1}^{\iota_p} g_{2i+1}(\alpha(\eta), \boldsymbol{\lambda}(\eta)) \rho^{2i+1} (1 + \Phi_{2i+1}(\alpha(\eta), \boldsymbol{\lambda}(\eta), \rho)) & \text{if } \zeta(\boldsymbol{\lambda}') < \iota_p, \\ 0 & \text{if } \zeta(\boldsymbol{\lambda}') = \iota_p, \end{cases} \end{aligned}$$

where $H_{2i+1}(\eta, \rho) \rightarrow 0$ as $(\eta, \rho) \rightarrow (0, 0)$. Then, for each $\nu = 0, 1, \dots, k$, we have

$$\begin{aligned} \Pi(\eta^{h_{i_\nu}/2}) &= \sum_{i=0}^{\zeta(\boldsymbol{\lambda}')} c_i \eta^{w_i + (i+1/2)h_{i_\nu}} (1 + \tilde{H}_{2i+1}(\eta)) \\ &+ \begin{cases} \sum_{i=\zeta(\boldsymbol{\lambda}')+1}^{\iota_p} g_{2i+1}(\alpha(\eta), \boldsymbol{\lambda}(\eta)) \eta^{(i+1/2)h_{i_\nu}} (1 + \tilde{H}_{2i+1}(\eta)) & \text{if } \zeta(\boldsymbol{\lambda}') < \iota_p, \\ 0 & \text{if } \zeta(\boldsymbol{\lambda}') = \iota_p, \end{cases} \end{aligned} \quad (2.10)$$

where

$$\tilde{H}_{2i+1}(\eta) := H_{2i+1}(\eta, \eta^{h_{i_\nu}/2}) \rightarrow 0 \quad \text{as } \eta \rightarrow 0. \quad (2.11)$$

Applying (2.8) to (2.10) we get

$$\begin{aligned}
 \Pi(\eta^{h_{i_\nu}/2}) &= \eta^{w_{i_\nu}+(i_\nu+1/2)h_{i_\nu}} \sum_{i \in \Xi(i_\nu)} c_i \left(1 + \tilde{H}_{2i+1}(\eta)\right) \\
 &+ \sum_{i \in \{0, \dots, \zeta(\boldsymbol{\lambda}')\} \setminus \Xi(i_\nu)} c_i \eta^{w_i+(i+1/2)h_{i_\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) \\
 &+ \begin{cases} \sum_{i=\zeta(\boldsymbol{\lambda}')+1}^{\iota_p} g_{2i+1}(\alpha(\eta), \boldsymbol{\lambda}(\eta)) \eta^{(i+1/2)h_{i_\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) & \text{if } \zeta(\boldsymbol{\lambda}') < \iota_p, \\ 0 & \text{if } \zeta(\boldsymbol{\lambda}') = \iota_p, \end{cases} \\
 &= \eta^{w_{i_\nu}+(i_\nu+1/2)h_{i_\nu}} \left\{ \hat{c}_{i_\nu} + \sum_{i \in \Xi(i_\nu)} c_i \tilde{H}_{2i+1}(\eta) \right\} \\
 &+ \sum_{i \in \{0, \dots, \zeta(\boldsymbol{\lambda}')\} \setminus \Xi(i_\nu)} c_i \eta^{w_i+(i+1/2)h_{i_\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) \\
 &+ \begin{cases} \sum_{i=\zeta(\boldsymbol{\lambda}')+1}^{\iota_p} g_{2i+1}(\alpha(\eta), \boldsymbol{\lambda}(\eta)) \eta^{(i+1/2)h_{i_\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) & \text{if } \zeta(\boldsymbol{\lambda}') < \iota_p, \\ 0 & \text{if } \zeta(\boldsymbol{\lambda}') = \iota_p, \end{cases} \\
 &= \hat{c}_{i_\nu} \eta^{w_{i_\nu}+(i_\nu+1/2)h_{i_\nu}} (1 + \Psi_f(\eta)) \quad \forall i_\nu \in \mathcal{V} \cup \{i_{s'}\}, \tag{2.12}
 \end{aligned}$$

where \hat{c}_{i_ν} is defined in (2.6) and

$$\begin{aligned}
 \Psi_f(\eta) &:= \sum_{i \in \Xi(i_\nu)} \frac{c_i}{\hat{c}_{i_\nu}} \tilde{H}_{2i+1}(\eta) + \sum_{i \in \{0, \dots, \zeta(\boldsymbol{\lambda}')\} \setminus \Xi(i_\nu)} \frac{c_i}{\hat{c}_{i_\nu}} \eta^{w_i-w_{i_\nu}+(i-i_\nu)h_{i_\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) \\
 &+ \begin{cases} \sum_{i=\zeta(\boldsymbol{\lambda}')+1}^{\iota_p} \frac{g_{2i+1}(\alpha(\eta), \boldsymbol{\lambda}(\eta))}{\hat{c}_{i_\nu}} \eta^{(i-i_\nu)h_{i_\nu}-w_{i_\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) & \text{if } \zeta(\boldsymbol{\lambda}') < \iota_p, \\ 0 & \text{if } \zeta(\boldsymbol{\lambda}') = \iota_p. \end{cases} \tag{2.13}
 \end{aligned}$$

By (2.11) we have that $\Psi_f(\eta) \rightarrow 0$ as $\eta \rightarrow 0$.

For $\boldsymbol{\lambda}' \in \mathcal{C}$, from (2.2) and (2.7) we get

$$\Pi(\rho) = \sum_{i=0}^{\iota(\boldsymbol{\lambda}')} g_{2i+1} \rho^{2i+1} (1 + \Phi_{2i+1}(\alpha, \boldsymbol{\lambda}, \rho)), \tag{2.14}$$

where Φ_{2i+1} 's are analytic at $(\alpha, \boldsymbol{\lambda}', \rho)$ and vanish when $\alpha = \rho = 0$. Restricted to the curve $(\alpha, \boldsymbol{\lambda}) = (\alpha(\eta), \boldsymbol{\lambda}(\eta))$, from (2.14) we obtain

$$\Pi(\rho) = \sum_{i=0}^{\iota(\boldsymbol{\lambda}')} c_i \eta^{w_i} \rho^{2i+1} (1 + H_{2i+1}(\eta, \rho)),$$

where $H_{2i+1}(\eta, \rho) \rightarrow 0$ as $(\eta, \rho) \rightarrow (0, 0)$. Then

$$\Pi(\eta^{h_{i_\nu}/2}) = \sum_{i=0}^{\iota(\boldsymbol{\lambda}')} c_i \eta^{w_i+(i+1/2)h_{i_\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) \quad \text{for } \nu = 0, 1, \dots, k, \tag{2.15}$$

where

$$\tilde{H}_{2i+1}(\eta) := H_{2i+1}(\eta, \eta^{h_{i\nu}/2}) \rightarrow 0 \quad \text{as } \eta \rightarrow 0. \quad (2.16)$$

Applying (2.8) to (2.15) we get

$$\begin{aligned} \Pi(\eta^{h_{i\nu}/2}) &= \eta^{w_{i\nu}+(i\nu+1/2)h_{i\nu}} \sum_{i \in \Xi(i\nu)} c_i \left(1 + \tilde{H}_{2i+1}(\eta)\right) \\ &\quad + \sum_{i \in \{0, \dots, \iota(\boldsymbol{\lambda}')\} \setminus \Xi(i\nu)} c_i \eta^{w_i+(i+1/2)h_{i\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) \\ &= \eta^{w_{i\nu}+(i\nu+1/2)h_{i\nu}} \left(\hat{c}_{i\nu} + \sum_{i \in \Xi(i\nu)} c_i \tilde{H}_{2i+1}(\eta)\right) \\ &\quad + \sum_{i \in \{0, \dots, \iota(\boldsymbol{\lambda}')\} \setminus \Xi(i\nu)} c_i \eta^{w_i+(i+1/2)h_{i\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right) \\ &= \hat{c}_{i\nu} \eta^{w_{i\nu}+(i\nu+1/2)h_{i\nu}} (1 + \Psi_c(\eta)) \quad \forall i\nu \in \mathcal{V} \cup \{i_{s'}\}, \end{aligned} \quad (2.17)$$

where $\hat{c}_{i\nu}$ is defined in (2.6) and

$$\Psi_c(\eta) := \sum_{i \in \Xi(i\nu)} \frac{c_i}{\hat{c}_{i\nu}} \tilde{H}_{2i+1}(\eta) + \sum_{i \in \{0, \dots, \iota(\boldsymbol{\lambda}')\} \setminus \Xi(i\nu)} \frac{c_i}{\hat{c}_{i\nu}} \eta^{w_i-w_{i\nu}+(i-i\nu)h_{i\nu}} \left(1 + \tilde{H}_{2i+1}(\eta)\right). \quad (2.18)$$

By (2.16) we have that $\Psi_c(\eta) \rightarrow 0$ as $\eta \rightarrow 0$.

Finally, for $\boldsymbol{\lambda}' \in \mathbb{R}^m$ we can choose $\varepsilon > 0$ small enough such that the point $(\alpha(\hat{\eta}), \boldsymbol{\lambda}(\hat{\eta}))$ on the parameterized curve, where $\hat{\eta} = \varepsilon^\beta$ and $\beta > 2/h_{i_{s'}}$ is a constant, lies in an arbitrarily small neighborhood of $(0, \boldsymbol{\lambda}')$. The monotonicity of $\{h_{i\nu}\}$, given just before condition **(LD)**, implies that

$$\hat{\eta}^{h_{i\nu}/2} < \hat{\eta}^{h_{i_{s'}/2}} = \varepsilon^{\beta h_{i_{s'}/2}} < \varepsilon,$$

for all $i\nu \in \mathcal{V}$. Let $\hat{\iota} := \#\mathcal{V}$, the cardinality of \mathcal{V} . It follows that the $\hat{\iota} + 1$ points $\hat{\eta}^{h_{i\nu}/2}$, $i\nu \in \mathcal{V} \cup \{i_{s'}\}$, all lie in $(0, \varepsilon)$ and increase as ν increases. Thus, from (2.12) when $\boldsymbol{\lambda}' \in \mathbb{R}^m \setminus \mathcal{C}$, or from (2.17) when $\boldsymbol{\lambda}' \in \mathcal{C}$, we get

$$\Pi(\hat{\eta}^{h_{i\nu}/2}) = \hat{c}_{i\nu} \varepsilon^{\beta w_{i\nu}+(i\nu+1/2)\beta h_{i\nu}} (1 + \Phi(\varepsilon)) \quad \text{for each } i\nu \in \mathcal{V} \cup \{i_{s'}\}, \quad (2.19)$$

where $\Phi(\varepsilon) := \Psi_f(\varepsilon^\beta)$ or $\Psi_c(\varepsilon^\beta)$, which tends to 0 as $\varepsilon \rightarrow 0$ as defined in (2.13) and (2.18). The formula (2.19) shows that, for sufficiently small $\varepsilon > 0$, $\Pi(\hat{\eta}^{h_{i\nu}/2})$ has the same sign as $\hat{c}_{i\nu}$ for each $i\nu \in \mathcal{V} \cup \{i_{s'}\}$. Therefore, the sign of Π alters once at each of these $\hat{\iota} + 1$ points, implying that the equation $\Pi(\rho) = 0$ has at least $\hat{\iota}$ roots in $(0, \varepsilon)$, i.e., $\mathcal{N}(\boldsymbol{\lambda}') \geq \hat{\iota}$. The proof is completed. \square

In roughly speaking, $\#\mathcal{V}$ is a number of sign changes caused by those terms whose powers compose a ladder-likely degressive sequence. In order to apply Theorem 2.1, we need to find a curve $(\alpha, \boldsymbol{\lambda}) = (\alpha(\eta), \boldsymbol{\lambda}(\eta))$ passing through $(0, \boldsymbol{\lambda}')$ in the parameter space \mathbb{R}^{m+1} such that the condition **(LD)** holds. Then, one can choose an appropriate $(\alpha, \boldsymbol{\lambda})$ near $(0, \boldsymbol{\lambda}')$ on the curve to obtain the number $\#\mathcal{V}$ of limit cycles of system $LC(\alpha, \boldsymbol{\lambda})$ near the center O . The curve Υ is found by solving the indeterminate α_i 's (in the w_i 's) and h_j 's from the inequalities given in **(LD)**, which will be illustrated in section 4.

3 Some corollaries

In this section we give two corollaries of Theorem 2.1 for easier applications in some cases. Those cases come from special cases of the condition **(LD)**.

Suppose that the power sequence of g_{2i+1} 's at $(0, \boldsymbol{\lambda}')$ satisfies

$$\text{(D)} \quad w_0 - w_1 \geq \dots \geq w_{i-1} - w_i \geq \dots \geq w_{\kappa(\boldsymbol{\lambda}')-1} - w_{\kappa(\boldsymbol{\lambda}')} > 0.$$

This means that the sequence $\{w_i\}$ is decreasing and the gaps between two consecutive terms become smaller and smaller.

For those “=” in condition **(D)**, we consider $w_{\varsigma-1} - w_{\varsigma} = w_{\varsigma} - w_{\varsigma+1}$ for some ς in $\{0, 1, \dots, \kappa(\boldsymbol{\lambda}')\}$ and let integers $\varsigma_1, \varsigma_2 \in \{0, \dots, \kappa(\boldsymbol{\lambda}')\}$ denote the indices such that $\varsigma_1 < \varsigma < \varsigma_2$ and

$$w_{\varsigma_1-1} - w_{\varsigma_1} > w_{\varsigma_1} - w_{\varsigma_1+1} = \dots = w_{\varsigma_2-1} - w_{\varsigma_2} > w_{\varsigma_2} - w_{\varsigma_2+1}, \quad (3.1)$$

i.e. ς_1 and ς_2 are respectively the first left index and the first right index near ς which destroy the equality “=” in **(D)**. Define $w_{-1} = w_{\kappa(\boldsymbol{\lambda}')+1} := +\infty$ complementarily.

Then, for each $\varsigma \in \{0, 1, \dots, \kappa(\boldsymbol{\lambda}')\}$, we define

$$\widehat{\Xi}(\varsigma) := \begin{cases} \{\varsigma_1, \dots, \varsigma, \dots, \varsigma_2\} & \text{if (3.1) holds,} \\ \{\varsigma\} & \text{if } w_{\varsigma-1} - w_{\varsigma} > w_{\varsigma} - w_{\varsigma+1}, \end{cases}$$

which can be used to find a scale (i_0, i_1, \dots, i_k) , where $i_0 := 0$, $i_k := \kappa(\boldsymbol{\lambda}')$ and $i_1, \dots, i_{k-1} \in \{1, \dots, \kappa(\boldsymbol{\lambda}') - 1\}$ such that

$$\begin{aligned} \widehat{\Xi}(i_j) &\neq \widehat{\Xi}(i_l) && \text{for } 0 \leq j \neq l \leq k, \\ \widehat{\Xi}(i_j) &\neq \widehat{\Xi}(i_j - 1) && \text{for } 0 < j \leq k, \\ \cup_{\nu=0}^k \widehat{\Xi}(i_\nu) &= \{0, 1, \dots, \kappa(\boldsymbol{\lambda}')\}. \end{aligned}$$

On the scale we define the set \mathcal{V} as in (2.6), where

$$\hat{c}_{i_\nu} := \sum_{j \in \widehat{\Xi}(i_\nu)} c_j \quad \text{for } \nu = 0, 1, \dots, k,$$

and c_i is the leading coefficient of g_{2i+1} as given in (2.5).

Corollary 3.1. $\mathcal{N}(\boldsymbol{\lambda}') \geq \#\mathcal{V}$ if the power sequence of g_{2i+1} 's at $(0, \boldsymbol{\lambda}')$ satisfies condition **(D)**.

Proof. Under condition **(D)**, define

$$\begin{aligned} h_i &:= (w_{i-1} - w_{i+1})/2, && \forall i = 1, \dots, \kappa(\boldsymbol{\lambda}') - 1, \\ h_0 &:= 2(w_0 - w_1), && h_{\kappa(\boldsymbol{\lambda}')} := (w_{\kappa(\boldsymbol{\lambda}')-1} - w_{\kappa(\boldsymbol{\lambda}')})/2, \end{aligned} \quad (3.2)$$

i.e. h_i is the average of the differences $w_{i-1} - w_i$ and $w_i - w_{i+1}$. From **(D)** and (3.2), it is easy to see that $h_0 > h_i \geq h_j > h_{\kappa(\boldsymbol{\lambda}')} > 0$ for $i < j$ in $\{1, \dots, \kappa(\boldsymbol{\lambda}') - 1\}$. Moreover, $h_i = h_j$ if and only if $w_{i-1} - w_i = w_i - w_{i+1} = \dots = w_j - w_{j+1}$. By the choice of the scale (i_0, i_1, \dots, i_k) given before this corollary, we have that $h_{i_0} > h_{i_1} > \dots > h_{i_k} > 0$. We claim that this sequence $\{w_i\}$ is ladder-likely degressive on the scale (i_0, i_1, \dots, i_k) . In fact, from **(D)** we see that for $j < i_\nu$

$$\begin{aligned} (i_\nu - j)h_{i_\nu} &= (i_\nu - j) \left(\frac{w_{i_\nu-1} - w_{i_\nu}}{2} + \frac{w_{i_\nu} - w_{i_\nu+1}}{2} \right) \leq (i_\nu - j)(w_{i_\nu-1} - w_{i_\nu}) \\ &\leq \sum_{i=j+1}^{i=i_\nu} (w_{i-1} - w_i) = w_j - w_{i_\nu}, \end{aligned}$$

where “=” holds if and only if $w_j - w_{j+1} = \dots = w_{i_\nu-1} - w_{i_\nu} = w_{i_\nu} - w_{i_\nu+1}$ because of **(D)**. Similarly, for $j > i_\nu$ we obtain

$$\begin{aligned} (i_\nu - j)h_{i_\nu} &= (i_\nu - j) \left(\frac{w_{i_\nu-1} - w_{i_\nu}}{2} + \frac{w_{i_\nu} - w_{i_\nu+1}}{2} \right) \leq (i_\nu - j)(w_{i_\nu} - w_{i_\nu+1}) \\ &= (j - i_\nu)(w_{i_\nu+1} - w_{i_\nu}) \leq \sum_{i=i_\nu+1}^{i=j} (w_i - w_{i-1}) = w_j - w_{i_\nu}, \end{aligned}$$

where “=” holds if and only if $w_{i_\nu-1} - w_{i_\nu} = w_{i_\nu} - w_{i_\nu+1} = \dots = w_{j-1} - w_j$. Thus, (2.8) holds in our case, which implies that the h_{i_ν} 's ($\nu = 0, 1, \dots, k$) satisfy condition **(LD)** and therefore the claim is proved. Therefore, $\mathcal{N}(\boldsymbol{\lambda}') \geq \#\mathcal{V}$ by Theorem 2.1. \square

Since condition **(D)** requires that the sequence of power-differences $\{w_0 - w_1, \dots, w_{\kappa(\boldsymbol{\lambda}')-1} - w_{\kappa(\boldsymbol{\lambda}')}\}$ is non-increasing, we have the following three cases:

(D0) there is no “=” in **(D)**;

(D1) “=” appears in **(D)** in discontinuous manner;

(D2) “=” appears continuously in **(D)**, i.e. there exists i such that $\dots \geq w_{i-1} - w_i = w_i - w_{i+1} = w_{i+1} - w_{i+2} \geq \dots$

Corollary 3.2. $\mathcal{N}(\boldsymbol{\lambda}') = \kappa(\boldsymbol{\lambda}')$ if either $c_i c_{i+1} < 0$ for all $i = 0, \dots, \kappa(\boldsymbol{\lambda}') - 1$ in the case **(D0)**, or $\hat{c}_{i_\nu} \hat{c}_{i_\nu+1} < 0$ for all $\nu = 0, \dots, \kappa(\boldsymbol{\lambda}') - 1$ in the case **(D1)**.

Proof. In case **(D0)** choosing the scale (i_0, i_1, \dots, i_k) as $(0, 1, \dots, \kappa(\boldsymbol{\lambda}'))$ and defining the h_{i_ν} 's as in (3.2), we compute $\hat{c}_{i_\nu} = c_{i_\nu}$ for all $\nu = 0, 1, \dots, \kappa(\boldsymbol{\lambda}')$. In case **(D1)** choosing the scale (i_0, i_1, \dots, i_k) and defining the h_{i_ν} 's as above, we see that $k = \kappa(\boldsymbol{\lambda}')$ and $h_{i_0} > h_{i_1} > \dots > h_{i_k} > 0$. Then the result follows from Corollary 3.1. \square

Although the cyclicity $\mathcal{N}(\boldsymbol{\lambda}')$ may reach the upper estimate $\kappa(\boldsymbol{\lambda}')$ in cases **(D0)** and **(D1)**, for which Corollary 3.2 gives sufficient conditions, we do not have such a

result yet in case **(D2)** because the equality $\widehat{\Xi}(i) = \widehat{\Xi}(i+1)$ known by the definition of $\widehat{\Xi}$ implies that $\#\mathcal{V} \leq k < \kappa(\boldsymbol{\lambda}')$.

Remark that Theorem 6.6 of [11] can also be employed to case **(D0)** but does not work for case **(D1)**. On the other hand, $\hat{c}_{i\nu} = c_{i\nu}$ for all $\nu = 0, 1, \dots, \kappa(\boldsymbol{\lambda}')$ in case **(D0)** but, there are some i such that $\hat{c}_{i\nu} = c_{i\nu-1} + c_{i\nu} + c_{i\nu+1}$ in case **(D1)**. For example, $\hat{c}_0 = c_0$, $\hat{c}_1 = c_1$, $\hat{c}_2 = c_1 + c_2 + c_3$ and $\hat{c}_3 = c_3$ when $w_0 - w_1 > w_1 - w_2 = w_2 - w_3 > w_3 - w_4 \dots$

We consider the following family of polynomial differential systems

$$\dot{x} = \alpha x - y + \sum_{i=1}^4 a_{2i+1} x(x^2 + y^2)^i, \quad \dot{y} = x + \alpha y + \sum_{i=1}^4 a_{2i+1} y(x^2 + y^2)^i, \quad (3.3)$$

parameterized by $(\alpha, \boldsymbol{\lambda}) := (\alpha, \lambda_1, \lambda_2) \in \mathbb{R}^3$, where $a_3 := \lambda_1$, $a_5 := -\lambda_2^2$, $a_7 := 3\lambda_2$ and $a_9 := -1$. One can compute its focal values

$$g_1 = 2\pi\alpha, \quad g_3 = 2\pi\lambda_1, \quad g_5 = -2\pi\lambda_2^2, \quad g_7 = 6\pi\lambda_2, \quad g_9 = -2\pi, \quad (3.4)$$

where each g_{2i+1} is the remainder of the original g_{2i+1} divided by the Gröbner basis of ideal $\langle g_3, \dots, g_{2i-1} \rangle$ in the order $\lambda_1 \prec \lambda_2$. Thus, $\kappa(\boldsymbol{\lambda}') = \zeta(\boldsymbol{\lambda}') = 4$, where $\boldsymbol{\lambda}' = (0, 0)$. Note that the independence condition of focal values, i.e. **(ID_k-1)** and **(ID_k-2)**, do not hold for g_1, g_3, \dots, g_7 because $g_7 = 0$ if $g_5 = 0$, which implies that we cannot obtain 4 limit cycles by verifying the classical independence of focal values. Using our above mentioned method, we choose the curve

$$\Upsilon : \alpha = -\eta^9, \quad \lambda_1 = \eta^5, \quad \lambda_2 = \eta,$$

in the $(\alpha, \lambda_1, \lambda_2)$ -space. Restricted to Υ those focal values given in (3.4) can be written in the form (2.5) taking

$$\begin{aligned} w_0 = 9, \quad w_1 = 5, \quad w_2 = 2, \quad w_3 = 1, \quad w_4 = 0, \\ c_0 = -2\pi, \quad c_1 = 2\pi, \quad c_2 = -2\pi, \quad c_3 = 6\pi, \quad c_4 = -2\pi. \end{aligned}$$

One can check that

$$w_0 - w_1 > w_1 - w_2 > w_2 - w_3 = w_3 > 0,$$

and compute that

$$\hat{c}_0 = c_0 = -2\pi, \quad \hat{c}_1 = c_1 = 2\pi, \quad \hat{c}_2 = c_2 = -2\pi, \quad \hat{c}_3 = c_2 + c_3 + c_4 = 2\pi, \quad \hat{c}_4 = c_4 = -2\pi,$$

which implies by Corollary 3.2 in case **(D1)** that $\mathcal{N}(\boldsymbol{\lambda}') = 4$.

4 Application to cubic systems

In this section we apply our method to a family of cubic polynomial differential systems with 5 parameters. Consider

$$\begin{aligned} \dot{x} &= \alpha x - y + (\lambda_1 - \lambda_3)x^2 + \lambda_2 xy + \lambda_3 y^2 - (9 + \lambda_2^2 + \lambda_3 \lambda_4)x^2 y + 2y^3, \\ \dot{y} &= x + \alpha y - x^3 - (12 + \lambda_2^2 + \lambda_3^2 + \lambda_3 \lambda_4)xy^2, \end{aligned} \quad (4.1)$$

parameterized by $(\alpha, \boldsymbol{\lambda}) := (\alpha, \lambda_1, \dots, \lambda_4) \in \mathbb{R}^5$. It is easy to compute the first five nonzero focal values

$$g_1 = 2\pi\alpha, \quad g_3 = \frac{\pi}{4}\lambda_1\lambda_2, \quad g_5 = -\frac{\pi}{24}\lambda_2\lambda_3^3, \quad g_7 = \frac{\pi}{96}\lambda_2\lambda_3^2\lambda_4, \quad g_9 = -\frac{\Omega(\lambda_2)\pi}{960}\lambda_2\lambda_3, \quad (4.2)$$

where $\Omega(\lambda) := 6720 + 1265\lambda^2 + 61\lambda^4$ and, for a short statement, each g_{2i+1} is the remainder of the original g_{2i+1} divided by the Gröbner basis of ideal $\langle g_3, \dots, g_{2i-1} \rangle$ in the order $\lambda_1 \prec \lambda_2 \prec \lambda_3 \prec \lambda_4$. Family (4.1) has the center variety $\mathcal{C} = \Gamma_1 \cup \Gamma_2$, where

$$\Gamma_1 := \{\boldsymbol{\lambda} \in \mathbb{R}^4 : \lambda_1 = \lambda_3 = 0\} \quad \text{and} \quad \Gamma_2 := \{\boldsymbol{\lambda} \in \mathbb{R}^4 : \lambda_2 = 0\}.$$

In fact, family (4.1)| $_{\alpha=0}$ is time-reversible for $\boldsymbol{\lambda} \in \Gamma_1 \cup \Gamma_2$, but on the contrary

$$V(g_3) \cap \mathbb{R}^4 \supsetneq V(g_3, g_5) \cap \mathbb{R}^4 = \Gamma_1 \cup \Gamma_2 = V(g_3, \dots, g_{2i+1}, \dots) \cap \mathbb{R}^4, \quad (4.3)$$

by the expressions of g_3 and g_5 .

Proposition 4.1. *For $\boldsymbol{\lambda} \in \mathbb{R}^4 \setminus \mathcal{C}$ the cyclicity $\mathcal{N}_f(\boldsymbol{\lambda})$ and the multiplicity $\zeta(\boldsymbol{\lambda})$ of O in the family (4.1) satisfy that either $\mathcal{N}_f(\boldsymbol{\lambda}) = \zeta(\boldsymbol{\lambda}) = 1$ or $\mathcal{N}_f(\boldsymbol{\lambda}) = \zeta(\boldsymbol{\lambda}) = 2$, which holds if either $\lambda_1\lambda_2 \neq 0$ or $\lambda_1 = 0 \neq \lambda_2\lambda_3$ correspondingly.*

Proof. The results can be proved by checking the independence condition of focal values, i.e. **(ID_k-1)** and **(ID_k-2)** for $k = 2$. Actually, by (4.3), the origin is a weak focus of multiplicity at most 2 when $\boldsymbol{\lambda} \in \mathbb{R}^4 \setminus \mathcal{C}$. For such a $\boldsymbol{\lambda}$, by the definitions of Γ_1 and Γ_2 , there are only two cases: either $\lambda_1\lambda_2 \neq 0$ or $\lambda_1 = 0 \neq \lambda_2\lambda_3$. It is easy to check that g_1, g_3 and g_5 are independent at $(0, \boldsymbol{\lambda})$ when $\boldsymbol{\lambda}$ satisfies that $\lambda_1 = 0 \neq \lambda_2\lambda_3$, and that g_1 and g_3 are independent at $(0, \boldsymbol{\lambda})$ when $\boldsymbol{\lambda}$ satisfies that $\lambda_1\lambda_2 \neq 0$.

Meanwhile, the results of this proposition can also be proved by using our main theorem or corollaries. In fact, in the case that $\lambda_1 = 0 \neq \lambda_2\lambda_3$, consider the parametric curve

$$\alpha(\eta) := -\text{sgn}(\lambda_2\lambda_3)\eta^{10}, \quad \lambda_1(\eta) := \text{sgn}(\lambda_3)\eta^3, \quad \lambda_i(\eta) := \lambda_i + \eta, \quad i = 2, 3, 4.$$

With the restriction to the curve, we can compute $w_0 = 10, w_1 = 3, w_2 = 0$ and $c_0 = -2\pi\text{sgn}(\lambda_2\lambda_3), c_1 = \pi\lambda_2\text{sgn}(\lambda_3)/4, c_2 = -\pi\lambda_2\lambda_3^3/24$. By Corollary 3.2, $\mathcal{N}_f(\boldsymbol{\lambda}) = 2$.

In the case that $\lambda_1\lambda_2 \neq 0$, we can prove $\mathcal{N}_f(\boldsymbol{\lambda}) = 1$ similarly. \square

This proposition shows that, for $\boldsymbol{\lambda} \notin \Gamma_1 \cup \Gamma_2$, the origin O is a weak focus of multiplicity at most 2, and there are small perturbations such that exactly j limit cycles bifurcate from the weak focus of multiplicity j for $j = 1, 2$.

On the other hand, the origin O is a center of (4.1) if and only if $\alpha = 0$ and $\boldsymbol{\lambda} \in \Gamma_1 \cup \Gamma_2$. Clearly, every point in Γ_1 and Γ_2 can be written as $\boldsymbol{\lambda}^{(1)} := (0, \lambda_2^{(1)}, 0, \lambda_4^{(1)})$ and $\boldsymbol{\lambda}^{(2)} := (\lambda_1^{(2)}, 0, \lambda_3^{(2)}, \lambda_4^{(2)})$ respectively. In order to avoid a double discussion at the intersection of $\Gamma_1 \cap \Gamma_2$, we assume either $\lambda_1^{(2)} \neq 0$ or $\lambda_3^{(2)} \neq 0$.

Proposition 4.2. For $\boldsymbol{\lambda}$ equal to $\boldsymbol{\lambda}^{(1)}$ or $\boldsymbol{\lambda}^{(2)}$ in \mathcal{C} , the cyclicity $\mathcal{N}_c(\boldsymbol{\lambda})$ and the multiplicity $\iota(\boldsymbol{\lambda})$ of the origin O in the family (4.1) have the results given in Table 1:

For	if	then
$\boldsymbol{\lambda}^{(1)} := (0, \lambda_2^{(1)}, 0, \lambda_4^{(1)})$	$\lambda_4^{(1)} \neq 0, \vartheta > 0$	$\iota(\boldsymbol{\lambda}^{(1)}) = 4, \mathcal{N}_c(\boldsymbol{\lambda}^{(1)}) = 4$
	$\lambda_4^{(1)} \neq 0, \vartheta \leq 0$	$\iota(\boldsymbol{\lambda}^{(1)}) = 4, \mathcal{N}_c(\boldsymbol{\lambda}^{(1)}) \geq 2$
	$\lambda_4^{(1)} = 0$	$\iota(\boldsymbol{\lambda}^{(1)}) = 4, \mathcal{N}_c(\boldsymbol{\lambda}^{(1)}) \geq 2$
$\boldsymbol{\lambda}^{(2)} := (\lambda_1^{(2)}, 0, \lambda_3^{(2)}, \lambda_4^{(2)})$ where either $\lambda_1^{(2)} \neq 0$ or $\lambda_3^{(2)} \neq 0$	$\lambda_1^{(2)} = 0, \lambda_3^{(2)} \neq 0$	$\iota(\boldsymbol{\lambda}^{(2)}) = 2, \mathcal{N}_c(\boldsymbol{\lambda}^{(2)}) = 2$
	$\lambda_1^{(2)} \neq 0$	$\iota(\boldsymbol{\lambda}^{(2)}) = 1, \mathcal{N}_c(\boldsymbol{\lambda}^{(2)}) = 1$
Remark: $\vartheta := 5(\lambda_4^{(1)})^2 - 8\Omega(\lambda_2^{(1)})$ and Ω is given in (4.2).		

Table 1: The number of limit cycles bifurcating from the center O .

Proof. By the definition of Γ_1 and Γ_2 every focal value of family (4.1) is of the form

$$\lambda_2(f_1(\boldsymbol{\lambda})\lambda_1 + f_2(\boldsymbol{\lambda})\lambda_3), \quad (4.4)$$

where $f_1, f_2 \in \mathbb{R}[\boldsymbol{\lambda}]$. From (4.4) we see for $\boldsymbol{\lambda} = \boldsymbol{\lambda}^{(1)}$ that the least integer j such that

$$\langle g_3, \dots, g_{2j+1} \rangle_{\boldsymbol{\lambda}^{(1)}} = \langle g_3, \dots, g_{2i+1}, \dots \rangle_{\boldsymbol{\lambda}^{(1)}}$$

in $\mathbb{R}\{\boldsymbol{\lambda}\}_{\boldsymbol{\lambda}^{(1)}}$ is 4, because every focal value lies in $\langle \lambda_1\lambda_2, \lambda_2\lambda_3 \rangle_{\boldsymbol{\lambda}^{(1)}}$, $\lambda_1\lambda_2, \lambda_2\lambda_3 \in \langle g_3, g_5, g_7, g_9 \rangle_{\boldsymbol{\lambda}^{(1)}}$ and $\lambda_2\lambda_3 \notin \langle g_3, g_5, g_7 \rangle_{\boldsymbol{\lambda}^{(1)}}$. This implies $\iota(\boldsymbol{\lambda}^{(1)}) = 4$. For $\boldsymbol{\lambda} = \boldsymbol{\lambda}^{(2)}$ we consider the case that $\lambda_1^{(2)} = 0$ but $\lambda_3^{(2)} \neq 0$, and the case that $\lambda_1^{(2)} \neq 0$ separately. In the first case $\langle g_3 \rangle_{\boldsymbol{\lambda}^{(2)}} \neq \langle g_3, \dots, g_{2i+1}, \dots \rangle_{\boldsymbol{\lambda}^{(2)}}$ and $\langle g_3, g_5 \rangle_{\boldsymbol{\lambda}^{(2)}} = \langle g_3, \dots, g_{2i+1}, \dots \rangle_{\boldsymbol{\lambda}^{(2)}}$ in $\mathbb{R}\{\boldsymbol{\lambda}\}_{\boldsymbol{\lambda}^{(2)}}$, which implies that $\iota(\boldsymbol{\lambda}^{(2)}) = 2$. In the second case $\lambda_2^{(2)} \in \langle g_3 \rangle_{\boldsymbol{\lambda}^{(2)}}$ in $\mathbb{R}\{\boldsymbol{\lambda}\}_{\boldsymbol{\lambda}^{(2)}}$ by the expression of g_3 . Thus, $\langle g_3 \rangle_{\boldsymbol{\lambda}^{(2)}} = \langle g_3, \dots, g_{2i+1}, \dots \rangle_{\boldsymbol{\lambda}^{(2)}}$ in $\mathbb{R}\{\boldsymbol{\lambda}\}_{\boldsymbol{\lambda}^{(2)}}$, which implies $\iota(\boldsymbol{\lambda}^{(2)}) = 1$.

First, consider $\boldsymbol{\lambda}^{(1)}$ in the case that $\lambda_4^{(1)} \neq 0$. Let

$$\begin{aligned} \alpha(\eta) &:= 0 + d_0\eta^{\alpha_0}, & \lambda_1(\eta) &:= 0 + d_1\eta^{\alpha_1}, & \lambda_2(\eta) &:= \lambda_2^{(1)} + d_2\eta^{\alpha_2}, \\ & & \lambda_3(\eta) &:= 0 + d_3\eta^{\alpha_3}, & \lambda_4(\eta) &:= \lambda_4^{(1)} + d_4\eta^{\alpha_4}, \end{aligned} \quad (4.5)$$

where the α_j 's and the d_j 's are undetermined. Then, we obtain

$$\begin{aligned} g_1(\alpha(\eta)) &= 2\pi d_0\eta^{\alpha_0}, \\ g_3(\boldsymbol{\lambda}(\eta)) &= \left(\pi d_1\lambda_2^{(1)} + \pi d_1d_2\eta^{\alpha_2} \right) \eta^{\alpha_1}/4, \\ g_5(\boldsymbol{\lambda}(\eta)) &= - \left(\pi\lambda_2^{(1)}d_3^3 + \pi d_2d_3^3\eta^{\alpha_2} \right) \eta^{3\alpha_3}/24, \\ g_7(\boldsymbol{\lambda}(\eta)) &= \left(\pi\lambda_2^{(1)}\lambda_4^{(1)}d_3^2 + \pi\lambda_2^{(1)}d_3^2d_4\eta^{\alpha_4} + \pi d_2d_3^2\lambda_4^{(1)}\eta^{\alpha_2} + \pi d_2d_3^2d_4\eta^{\alpha_2+\alpha_4} \right) \eta^{2\alpha_3}/96, \\ g_9(\boldsymbol{\lambda}(\eta)) &= - \left(\Omega(\lambda_2^{(1)} + d_2\eta^{\alpha_2})\pi\lambda_2^{(1)}d_3 + \Omega(\lambda_2^{(1)} + d_2\eta^{\alpha_2})\pi d_2d_3\eta^{\alpha_2} \right) \eta^{\alpha_3}/960, \end{aligned}$$

which give the power sequence $\{w_i\}$ and the leading coefficients c_i 's as follows:

$$\begin{aligned} w_0 &= \alpha_0, & w_1 &= \alpha_1, & w_2 &= 3\alpha_3, & w_3 &= 2\alpha_3, & w_4 &= \alpha_3, \\ c_0 &= 2\pi d_0, & c_1 &= \frac{\pi}{4}d_1\lambda_2^{(1)}, & c_2 &= -\frac{\pi}{24}d_3^3\lambda_2^{(1)}, & c_3 &= \frac{\pi}{96}d_3^2\lambda_2^{(1)}\lambda_4^{(1)}, & c_4 &= -\Omega(\lambda_2^{(1)})\frac{\pi}{960}d_3\lambda_2^{(1)}, \end{aligned}$$

when $\lambda_2^{(1)} \neq 0$, or

$$\begin{aligned} w_0 &= \alpha_0, & w_1 &= \alpha_1 + \alpha_2, & w_2 &= 3\alpha_3 + \alpha_2, & w_3 &= 2\alpha_3 + \alpha_2, & w_4 &= \alpha_3 + \alpha_2, \\ c_0 &= 2\pi d_0, & c_1 &= \frac{\pi}{4}d_1d_2, & c_2 &= -\frac{\pi}{24}d_2d_3^3, & c_3 &= \frac{\pi}{96}d_2d_3^2\lambda_4^{(1)}, & c_4 &= -\Omega(0)\frac{\pi}{960}d_2d_3, \end{aligned}$$

when $\lambda_2^{(1)} = 0$. We claim that there exist positive numbers α_i 's in (4.5) such that

$$w_0 - w_1 > w_1 - w_2 > w_2 - w_3 = w_3 - w_4 > 0. \quad (4.6)$$

In fact, (4.6) is equivalent to either $\alpha_0 - \alpha_1 > \alpha_1 - 3\alpha_3 > \alpha_3$ when $\lambda_2^{(1)} \neq 0$ or $\alpha_0 - \alpha_1 - \alpha_2 > \alpha_1 - 3\alpha_3 > \alpha_3$ when $\lambda_2^{(1)} = 0$, from which we can choose

$$\alpha_0 = 100, \quad \alpha_1 = 50, \quad \alpha_2 = 10, \quad \alpha_3 = 10,$$

for both cases. Thus, by (4.6) our system falls into the case **(D1)**. From the definition (2.6) we compute

$$\begin{aligned} \hat{c}_0 &= c_0 = 2\pi d_0, & \hat{c}_1 &= c_1 = \pi d_1\lambda_2^{(1)}/4, & \hat{c}_2 &= c_2 = -\pi d_3^3\lambda_2^{(1)}/24, \\ \hat{c}_3 &= c_2 + c_3 + c_4 = -\frac{\pi}{960}d_3\lambda_2^{(1)} \left(40d_3^2 - 10\lambda_4^{(1)}d_3 + \Omega(\lambda_2^{(1)}) \right), & \hat{c}_4 &= c_4 = -\Omega(\lambda_2^{(1)})\frac{\pi}{960}d_3\lambda_2^{(1)}, \end{aligned}$$

when $\lambda_2^{(1)} \neq 0$, or

$$\begin{aligned} \hat{c}_0 &= c_0 = 2\pi d_0, & \hat{c}_1 &= c_1 = \pi d_1d_2/4, & \hat{c}_2 &= c_2 = -\pi d_2d_3^3/24, \\ \hat{c}_3 &= c_2 + c_3 + c_4 = -\pi d_2d_3 \left(40d_3^2 - 10\lambda_4^{(1)}d_3 + \Omega(0) \right) / 960, & \hat{c}_4 &= c_4 = -\Omega(0)\pi d_2d_3/960, \end{aligned}$$

when $\lambda_2^{(1)} = 0$. Then, if $\vartheta > 0$ we can check that

$$\hat{c}_0\hat{c}_1 < 0, \quad \hat{c}_1\hat{c}_2 < 0, \quad \hat{c}_2\hat{c}_3 < 0, \quad \hat{c}_3\hat{c}_4 < 0,$$

where we choose in (4.5) either

$$d_0 = -\text{sign}(\lambda_2^{(1)}\lambda_4^{(1)}), \quad d_1 = \text{sign}(\lambda_4^{(1)}), \quad d_2 = 0, \quad d_3 = \lambda_4^{(1)}/8, \quad d_4 = 0$$

when $\lambda_2^{(1)} \neq 0$, or

$$d_0 = -\text{sign}(\lambda_4^{(1)}), \quad d_1 = \text{sign}(\lambda_4^{(1)}), \quad d_2 = 1, \quad d_3 = \lambda_4^{(1)}/8, \quad d_4 = 0$$

when $\lambda_2^{(1)} = 0$. Therefore $\mathcal{N}_c(\boldsymbol{\lambda}^{(1)}) = 4$ by Corollary 3.2. Similarly, if $\vartheta \leq 0$ we can check that

$$\hat{c}_0\hat{c}_1 < 0, \quad \hat{c}_1\hat{c}_2 < 0,$$

where we choose in (4.5) either

$$d_0 = -\text{sign}(\lambda_2^{(1)}), \quad d_1 = 1, \quad d_2 = 0, \quad d_3 = 1, \quad d_4 = 0$$

when $\lambda_2^{(1)} \neq 0$, or

$$d_0 = -1, \quad d_1 = 1, \quad d_2 = 1, \quad d_3 = 1, \quad d_4 = 0$$

when $\lambda_2^{(1)} = 0$. Therefore $\mathcal{N}_c(\boldsymbol{\lambda}^{(1)}) \geq 2$ by Corollary 3.1.

Next, consider $\boldsymbol{\lambda}^{(1)}$ in the case that $\lambda_4^{(1)} = 0$. Let

$$\begin{aligned} \alpha(\eta) &:= 0 + d_0\eta^{\alpha_0}, & \lambda_1(\eta) &:= 0 + d_1\eta^{\alpha_1}, & \lambda_2(\eta) &:= \lambda_2^{(1)} + d_2\eta^{\alpha_2}, \\ \lambda_3(\eta) &:= 0 + d_3\eta^{\alpha_3}, & \lambda_4(\eta) &:= 0 + d_4\eta^{\alpha_4}. \end{aligned} \quad (4.7)$$

Then

$$\begin{aligned} g_1(\alpha(\eta)) &= 2\pi d_0\eta^{\alpha_0}, \\ g_3(\boldsymbol{\lambda}(\eta)) &= \left(\pi d_1 \lambda_2^{(1)} \eta^{\alpha_1} + \pi d_1 d_2 \eta^{\alpha_1 + \alpha_2} \right) / 4, \\ g_5(\boldsymbol{\lambda}(\eta)) &= - \left(\pi \lambda_2^{(1)} d_3^3 \eta^{3\alpha_3} + \pi d_2 d_3^3 \eta^{3\alpha_3 + \alpha_2} \right) / 24, \\ g_7(\boldsymbol{\lambda}(\eta)) &= \left(\pi \lambda_2^{(1)} d_3^2 d_4 \eta^{2\alpha_3 + \alpha_4} + \pi d_2 d_3^2 d_4 \eta^{2\alpha_3 + \alpha_2 + \alpha_4} \right) / 96, \\ g_9(\boldsymbol{\lambda}(\eta)) &= - \left(\Omega(\lambda_2^{(1)} + d_2\eta^{\alpha_2}) \pi \lambda_2^{(1)} d_3 \eta^{\alpha_3} + \Omega(\lambda_2^{(1)} + d_2\eta^{\alpha_2}) \pi d_2 d_3 \eta^{\alpha_2 + \alpha_3} \right) / 960, \end{aligned}$$

which gives the power sequence $\{w_i\}$ and the leading coefficients c_i 's as follows:

$$\begin{aligned} w_0 &= \alpha_0, & w_1 &= \alpha_1, & w_2 &= 3\alpha_3, & w_3 &= 2\alpha_3 + \alpha_4, & w_4 &= \alpha_3, \\ c_0 &= 2\pi d_0, & c_1 &= \frac{\pi}{4} d_1 \lambda_2^{(1)}, & c_2 &= -\frac{\pi}{24} d_3^3 \lambda_2^{(1)}, & c_3 &= \frac{\pi}{96} d_3^2 d_4 \lambda_2^{(1)}, & c_4 &= -\Omega(\lambda_2^{(1)}) \frac{\pi}{960} d_3 \lambda_2^{(1)}, \end{aligned}$$

when $\lambda_2^{(1)} \neq 0$ and

$$\begin{aligned} w_0 &= \alpha_0, & w_1 &= \alpha_1 + \alpha_2, & w_2 &= 3\alpha_3 + \alpha_2, & w_3 &= 2\alpha_3 + \alpha_2 + \alpha_4, & w_4 &= \alpha_3 + \alpha_2, \\ c_0 &= 2\pi d_0, & c_1 &= \frac{\pi}{4} d_1 d_2, & c_2 &= -\frac{\pi}{24} d_2 d_3^3, & c_3 &= \frac{\pi}{96} d_2 d_3^2 d_4, & c_4 &= -\Omega(0) \frac{\pi}{960} d_2 d_3, \end{aligned}$$

when $\lambda_2^{(1)} = 0$. We claim that there exist positive numbers α_i 's in (4.7) such that

$$w_0 - w_1 > w_1 - w_2 > w_2 - w_3 < w_3 - w_4. \quad (4.8)$$

In fact, (4.8) is equivalent to either $\alpha_0 - \alpha_1 > \alpha_1 - 3\alpha_3 > \alpha_3 - \alpha_4 < \alpha_3 + \alpha_4$ when $\lambda_2^{(1)} \neq 0$, or $\alpha_0 - \alpha_1 - \alpha_2 > \alpha_1 - 3\alpha_3 > \alpha_3 - \alpha_4 < \alpha_3 + \alpha_4$ when $\lambda_2^{(1)} = 0$, from which we can choose

$$\alpha_0 = 13, \quad \alpha_1 = 8, \quad \alpha_2 = 3, \quad \alpha_3 = 2, \quad \alpha_4 = 1,$$

when $\lambda_2^{(1)} \neq 0$, and

$$\alpha_0 = 18, \quad \alpha_1 = 10, \quad \alpha_2 = 3, \quad \alpha_3 = 2, \quad \alpha_4 = 1,$$

when $\lambda_2^{(1)} = 0$. Thus **(LD)** holds on the scale $(i_0, i_1, i_2) := (0, 1, 2)$, where

$$h_{i_0} := 10, \quad h_{i_1} := 3.5, \quad h_{i_2} := 1.5,$$

when $\lambda_2^{(1)} \neq 0$, and

$$h_{i_0} := 10, \quad h_{i_1} := 4.5, \quad h_{i_2} := 2.5,$$

when $\lambda_2^{(1)} = 0$. Moreover “=” in **(LD)** holds only for $j = i_\nu$ ($\nu = 0, 1, 2$). From the definition of \hat{c}_{i_ν} we compute

$$\hat{c}_{i_0} = c_0 = 2\pi d_0, \quad \hat{c}_{i_1} = c_1 = \pi d_1 \lambda_2^{(1)}/4, \quad \hat{c}_{i_2} = c_2 = -\pi d_3^3 \lambda_2^{(1)}/24,$$

when $\lambda_2^{(1)} \neq 0$, or

$$\hat{c}_{i_0} = c_0 = 2\pi d_0, \quad \hat{c}_{i_1} = c_1 = \pi d_1 d_2/4, \quad \hat{c}_{i_2} = c_2 = -\pi d_2 d_3^3/24,$$

when $\lambda_2^{(1)} = 0$. We can check that

$$\hat{c}_{i_0} \hat{c}_{i_1} < 0, \quad \hat{c}_{i_1} \hat{c}_{i_2} < 0,$$

where we choose in (4.7) either

$$d_0 = -\text{sign}(\lambda_2^{(1)}), \quad d_1 = d_2 = d_3 = d_4 = 1,$$

when $\lambda_2^{(1)} \neq 0$, or

$$d_0 = -1, \quad d_1 = d_2 = d_3 = d_4 = 1,$$

when $\lambda_2^{(1)} = 0$. Therefore $\mathcal{N}_c(\boldsymbol{\lambda}^{(1)}) \geq 2$ by Theorem 2.1.

We similarly consider $\boldsymbol{\lambda}^{(2)}$ and obtain $\mathcal{N}_c(\boldsymbol{\lambda}^{(2)}) = 2$ in the case that $\lambda_1^{(2)} = 0, \lambda_3^{(2)} \neq 0$ and $\mathcal{N}_c(\boldsymbol{\lambda}^{(2)}) = 1$ in the case that $\lambda_1^{(2)} \neq 0$. \square

Proposition 4.2 implies that for some $\boldsymbol{\lambda}' \in \mathcal{C}$ there are 4 limit cycles bifurcating from the center O of system (4.1) $_{(\alpha, \boldsymbol{\lambda})=(0, \boldsymbol{\lambda}')}$. However, it is impossible to obtain 4 limit cycles bifurcating from the center by either the well-known method of independent focal values, or the method given in [11, Theorem 6.6]. In fact, the focal values g_3, g_5, g_7, g_9 are not independent because $g_7 = 0$ when $g_5 = 0$. We can compute

$$\frac{\partial(g_3, g_5, g_7, g_9)}{\partial(\lambda_1, \dots, \lambda_4)} = \begin{pmatrix} \frac{\pi}{4}\lambda_2 & \frac{\pi}{4}\lambda_1 & 0 & 0 \\ 0 & -\frac{\pi}{24}\lambda_3^3 & -\frac{\pi}{8}\lambda_2\lambda_3^2 & 0 \\ 0 & \frac{\pi}{96}\lambda_3^2\lambda_4 & \frac{\pi}{48}\lambda_2\lambda_3\lambda_4 & \frac{\pi}{96}\lambda_2\lambda_3^2 \\ 0 & -\frac{F(\lambda_2)\pi}{960}\lambda_3 & -\frac{G(\lambda_2)\pi}{960}\lambda_2 & 0 \end{pmatrix},$$

where $F(\lambda) := 6720 + 3795\lambda^2 + 305\lambda^4$ and $G(\lambda) := 6720 + 1265\lambda^2 + 61\lambda^4$, and obtain

$$\text{rank} \left(\frac{\partial(g_3, g_5, g_7, g_9)}{\partial(\lambda_1, \dots, \lambda_4)} \Big|_{\boldsymbol{\lambda}=\boldsymbol{\lambda}^{(1)}} \right) \leq 2, \quad \text{rank} \left(\frac{\partial(g_3, g_5, g_7, g_9)}{\partial(\lambda_1, \dots, \lambda_4)} \Big|_{\boldsymbol{\lambda}=\boldsymbol{\lambda}^{(2)}} \right) \leq 1,$$

which implies by [3, Theorem 1.3] that 2 limit cycles can be bifurcated from the center O for the case that $\boldsymbol{\lambda} = \boldsymbol{\lambda}^{(1)}$, and 1 limit cycle for the case that $\boldsymbol{\lambda} = \boldsymbol{\lambda}^{(2)}$. Moreover, from the expressions of the g_{2i+1} 's given in (4.2) we find that “=” always appears in **(D)**, which implies that the result of [11, Theorem 6.6] cannot be used to find 4 limit cycles bifurcating from the center O .

References

- [1] A. A. Andronov, E. A. Leontovich, I. I. Gordon, A. G. Maier, *Qualitative Theory of Second-Order Dynamic Systems*, Wiley, New York, 1973.
- [2] C. Chicone, M. Jacobs, Bifurcation of critical periods for plane vector fields, *Trans. Amer. Math. Soc.* **312**(1989) 433-486.
- [3] C. Christopher, C. Li, *Limit Cycles of Differential Equations*, Birkhäuser Verlag, Basel, 2007.
- [4] C. Christopher, N. G. Lloyd, Small-amplitude limit cycles in polynomial Liénard systems, *Nonlin. Diff. Eqns. Appl.* **3**(1996) 183-190.
- [5] F. Dumortier, J. Llibre, J. C. Artés, *Qualitative Theory of Planar Differential Systems*, Springer, Berlin, 2006.
- [6] A. Gasull, J. Torregrosa, Small-amplitude limit cycles in Liénard systems via multiplicity, *J. Diff. Eqns.* **159**(1999) 186-211.
- [7] R. Hartshorne, *Algebraic Geometry*, Springer, New York, 1977.
- [8] I. D. Iliev, Perturbations of quadratic centers, *Bull. Sci. Math.* **122**(1998) 107-161.
- [9] S. Lang, *Algebra*, Addison-Wesley Publishing Comp. Inc., New York, 1965.
- [10] J. Li, Hilbert's 16th problem and bifurcations of planar polynomial vector fields, *Int. J. Bifurc. Chaos* **13**(2003) 47-106.
- [11] Y. Liu, Theory of center-focus for a class of higher-degree critical points and infinite points, *Science in China: Series A* **44**(2001) 365-377.
- [12] J. Llibre, C. Valls, Classification of centers, their cyclicity and isochronicity for a class of polynomial differential systems generalizing the linear systems with cubic homogeneous nonlinearities, *J. Diff. Eqns.* **246**(2009) 2192-2204.
- [13] J. Llibre, C. Valls, On the number of limit cycles of a class of polynomial differential systems, *Proc. Roy. Soc. London Ser. A* **468**(2012) 2347-2360.
- [14] N. G. Lloyd, S. Lynch, Small amplitude limit cycles of certain Lienard systems, *Proc. Roy. Soc. London Ser. A* **418**(1988) 199-208.
- [15] A. M. Lyapunov, Problème général de la stabilité du mouvement, *Ann. Math. Stud.* **17**, Princeton University Press, Princeton, 1947.
- [16] H. Poincaré, Mémoire sur les courbes définies par les équations différentielles, *J. Math. Pure. Appl.* **4**(1885) 167-244.
- [17] R. Roussarie, *Bifurcation of Planar Vector Fields and Hilbert's 16th Problem*, IMPA, 1995.
- [18] S. Shi, A method of constructing cycles without contact around a weak focus, *J. Diff. Eqns.* **41**(1981) 313-319.
- [19] Z. Zhang, T. Ding, W. Huang, Z. Dong, *Qualitative Theory of Differential Equations*, Transl. Math. Monogr., Amer. Math. Soc., Providence, RI, 1992.
- [20] H. Zoladek, Eleven small limit cycles in a cubic vector field, *Nonlinearity* **8**(1995) 843-860.