## GIAO-PCM Calculations on Alanine Diamide Models Aimed at Predicting Protein Secondary Structures

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# GIAO-PCM Calculations on Alanine Diamide Models Aimed at Predicting Protein Secondary Structures 

Ilona Hudáky, ${ }^{[a]}$ Imre Jákii, ${ }^{[b]}$ Gábor Náray-Szabó ${ }^{[a]}$ and András Perczel ${ }^{[a, b]]^{*}}$


#### Abstract

In this paper we extend our theoretical studies dealing with the dependence of relative proton and carbon chemical shifts (CSs) of protein backbone atoms on their conformational position. In an earlier paper (A. Czajlik, I. Hudáky, A. Perczel, J Comp Chem 2011, 32, 3362) we reported on a fair agreement between calculated and observed backbone CSs as a function of backbone conformation. Applying the polarizable continuum model (PCM) in this work, we compare relative CSs of fully optimized alanine diamide conformers with gas phase calculations and experimental results. Along a path on the Ramachandran surface, we collated calculated relative CSs obtained with and without explicit water molecules, as well as with and without considering the PCM reaction field. Furthermore, we traced the energetically relevant reaction paths along the torsional angle $\psi$ connecting the lowest energy minima (helical, extended, polyproline II and inverse $\gamma$-turn) on the Ramachandran plot, with the prospect to facilitate identifying them by their relative CSs. We found that consideration of the solvent effect of the environment around a diamide model improves the agreement with experimental findings on abundant conformers. This agreement is of the level achieved previously by a thorough gas phase investigation on considerably larger oligoalanine models. By relating $\Delta \delta C^{\alpha}, \Delta \delta H^{\alpha}$ and $\Delta \delta C^{\beta}$ values of polyproline II and inverse $\gamma$-turn to the experimentally well characterized helical and extended data, our calculations contribute to protein secondary structure prediction based on nuclear magnetic CSs.


Keywords: protein secondary structure, relative chemical shift, PCM, ab initio, DFT

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## Introduction

There are intensive attempts to use nuclear magnetic resonance (NMR) chemical shifts (CS) in the determination of protein 3D structures either exclusively, or additionally to other parameters (such as nuclear Overhauser effects, NOE, or residual dipolar couplings, RDC). ${ }^{[1]}$ The first step in protein structure refinement by NMR is the assignment of each CS to a certain nucleus of an amino acid residue along the primary sequence. Beyond the identity of the nucleus and the molecular fold, CSs are influenced by several factors such as the proximity of charged sites, ring current effects, solvent, pH , temperature, etc.. Nevertheless those nuclei that either constitute the polypeptide backbone or are very close to it (Figure 1/A) show such differences from nucleus and amino acid specific reference values that depend on the backbone fold of the residues. Accordingly, these data can be used for the estimation of backbone conformations. ${ }^{[2-6]}$ This is why we attempted earlier the identification of the nine typical backbone folds ${ }^{[7]}$ of chiral amino acid residues on the basis of CS calculated for homoconformer polyalanine models in the gas phase. ${ }^{[8]}$ That was a sequel of the description of the CS properties of a simple diamide model. ${ }^{[9,10]}$

The benefit of the inclusion of solvent effects through the PCM reaction field ${ }^{[11]}$ into the quantum mechanical calculations on small model peptides aimed at predicting the geometry of amino acid residues in proteins was demonstrated in one of our previous studies. ${ }^{[12]}$ As published there, the conformational search of the Ramachandran map using the HCO-L-Ala- $\mathrm{NH}_{2}$ model peptide at either the RHF/PCM/6-31+G(d) or the B3LYP/PCM/6$31+G(d)$ level of theory readily delivered two low energy regions, one around conformer $\alpha_{\mathrm{L}}$ (characteristic for the right-handed $\alpha$-helix) and another around conformers $\beta_{\mathrm{L}} / \varepsilon_{\mathrm{L}}$ (elongated backbone fold). Without restriction on $\phi$ or $\psi$ torsional angles, conformers $\alpha_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$ are not genuine energy minima in the gas phase, ${ }^{[7]}$ however, these regions are well populated in experimentally determined protein structures. Good correlation was found between the relative energies, $\Delta E$, of the representative minima at $\mathrm{RHF} / \mathrm{PCM} / 6-31+\mathrm{G}(\mathrm{d})$, $\mathrm{B} 3 \mathrm{LYP} / \mathrm{PCM} / 6-$ $31+\mathrm{G}(\mathrm{d})$, as well as B3LYP/PCM/6-311++G(d,p) levels of theory and the logarithm of experimental abundances of alanine residues retrieved from the Protein Data Bank. At the RHF/6-31+G(d) level more alanine residues were assigned around the calculated minima than at other levels. The PCM results on the model peptide HCO-L-Ala- $\mathrm{NH}_{2}$ referring to an aqueous solution proved to match experimental data of internal alanine residues of proteins just as well as the data of surface alanines, even though the former ones are not surrounded by
solvent molecules. It was concluded that the most important effect of considering solvation within the model is that polarizability of the environment is respected. ${ }^{[12]}$

Without constrain on the backbone torsional angles $\phi$ and $\psi$ (Figure 1/B), $\alpha_{\mathrm{L}}$ or $\varepsilon_{\mathrm{L}}$ conformations of amino acid residues can be optimized in the gas phase only as parts of larger model peptides. In these, several amide groups are engaged in hydrogen bond formation. A single right-handed helix ${ }^{[13]}$ and a collagen triple helix ${ }^{[14]}$ consist of successive amino acid residues folded into $\alpha_{\mathrm{L}}$ conformation in the former and into $\varepsilon_{\mathrm{L}}$ one in the latter case, respectively, and both of these homoconformers can be optimized in the gas phase. Intra- and intermolecular H-bonds within the single right-handed helix as well as the three polyproline II helices of the collagen model stabilize these structures. By reducing the model to a single diamide, however, the H -bonded partners disappear, thus conformers $\alpha_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$ converge to other minima. Gas phase optimizations deliver conformer $\gamma_{\mathrm{L}}$ as the global minimum. This structure is an inverse $\gamma$-turn and has an H -bond within itself. In proteins, however, it is occurs seldom compared to helical and polyproline II structures.

As the application of the PCM reaction field ${ }^{[11]}$ improves the performance of a diamide model for an amino acid residue as far as geometry and conformational abundance in proteins is concerned, ${ }^{[12]}$ it seems to be plausible that PCM may help to reflect the dependence of certain CSs on the backbone fold, too. Applying the PCM reaction field, full and partial optimizations of diamide models and subsequent calculations of CSs were performed for the following specific aims:
i) to characterize all fully optimized alanine conformers by relative CSs at various $a b$ initio and DFT levels;
ii) to compare relative CSs along a path on the Ramachandran surface calculated with different approaches such as with and without explicit water molecules as well as with and without the PCM reaction field;
iii) to trace the energetically relevant reaction paths along the torsional angle $\psi$ connecting the lowest energy minima, i.e. $\alpha_{\mathrm{L}}, \beta_{\mathrm{L}}, \gamma_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$, and to characterize them by relative CSs.

## Methods

Calculations were carried out using the Gaussian $09^{[15]}$ (and occasionally the Gaussian $03^{[16]}$ ) program package (Table 1). The PCM reaction field ${ }^{[11]}$ was applied with water as solvent. The model peptide was HCO-L-Ala- $\mathrm{NH}_{2}$ (Figure 1/A) in most cases, and $\mathrm{CH}_{3} \mathrm{CO}$-L-$\mathrm{Ala}^{-} \mathrm{NHCH}_{3}$ (Figure 1/B) in the few other instances. The HCO-L-Ala- $\mathrm{NH}_{2}$ model peptide was surrounded by four water molecules (Figure 1/C) in the case of a path on the Ramachandran surface. During optimizations, the GDIIS algorithm ${ }^{[17]}$ was opted for. The PCM reaction field may not be applied with the basis set 3-21G according to an earlier observation indicating that the RHF/PCM/3-21G level reflects gas phase calculations rather than models with explicit solvent molecules. ${ }^{[12]}$ Chemical shielding tensors were calculated with the GIAO method. Isotropic chemical shielding values ( $\sigma$ ) were collected for the selected nuclei (amide nitrogen and hydrogen, $\alpha$ - and $\beta$ - carbon, $\alpha$-hydrogen, as well as carbonyl carbon, see Figure 1/A). CSs ( $\delta$ ) were not referenced to molecules such as TMS or $\mathrm{NH}_{3}$. Instead, relative CSs $(\Delta \delta)$ were deduced directly from chemical shielding values:

$$
\begin{equation*}
\Delta \delta=\sigma_{r e f}-\sigma \tag{1}
\end{equation*}
$$

The preference of relative values over absolute ones for correlating experimental chemical shifts in peptide analogues to computed ab initio or DFT results was suggested by Pulay and coworkers. ${ }^{[18]}$ In the present study, the reference $\left(\sigma_{\text {ref }}\right)$ is the right-handed helical $\alpha_{L}$ conformer ${ }^{[8]}$ except when otherwise stated.

## Detection and characterization of fully optimized alanine conformers

A thorough conformational search ${ }^{[12]}$ of the $\mathrm{HCO}-\mathrm{L}-\mathrm{Ala}-\mathrm{NH}_{2}$ model peptide using the PCM reaction field as implemented in Gaussian 03 revealed the existence of minima in seven out of the nine conformational regions of the Ramachandran map. Multiplication of minima was typical. In the present study, optimization of all nine ideal amino acid conformers ( $\alpha_{\mathrm{L}}, \beta_{\mathrm{L}}$, $\gamma_{\mathrm{L}}, \delta_{\mathrm{L}}, \varepsilon_{\mathrm{L}}, \alpha_{\mathrm{D}}, \gamma_{\mathrm{D}}, \delta_{\mathrm{D}}$ and $\varepsilon_{\mathrm{D}}$, Figure 2$)^{[7]}$ was attempted by four different levels of theory using the polarizable continuum model (RB3LYP/PCM/6-311+G(d,p), RB3LYP/PCM/6$311++G(2 d, 2 p), \quad$ RHF/PCM/6-311++G(d,p) and RHF/PCM/6-311++G(2d,2p)). The performances of program packages Gaussian 03 and Gaussian 09 were compared at the level RB3LYP/PCM/6-311++G(2d,2p). Difference caused by taking $\mathrm{CH}_{3} \mathrm{COAlaNHCH}_{3}$ instead of
$\mathrm{HCOAlaNH}_{2}$ as the model peptide was tested at levels RB3LYP/PCM/6-311+G(d,p) and RHF/PCM/6-311++G(d,p).

Chemical shielding values $(\sigma)$ were calculated usually at the theoretical level identical to that of the optimization. At levels RB3LYP/6-311++G(2d,2p) and RHF/6-311++G(2d,2p), however, the effect of omission of PCM from CS calculations was investigated. Relative CSs $(\Delta \delta)$ at each inspected level of theory were referenced to the respective CS in the right-handed helical $\alpha_{\mathrm{L}}$ conformer.

## Comparison of reaction paths calculated with different approaches

Greatly differing conformers accommodate surrounding water molecules dissimilarly. ${ }^{[19]}$ In a thorough conformational search, the solvent molecules should be placed around the solute according to a preceding molecular mechanical investigation. However, the H-bond pattern of $\mathrm{CH}_{3} \mathrm{CO}-\mathrm{L}-\mathrm{Ala}-\mathrm{NHCH}_{3}$ and surrounding four water molecules is very similar in the cases of conformers $\beta_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}\left(\beta_{2}^{\prime}\right.$ and $P_{\mathrm{II}}$ by the original notation). ${ }^{[19]}$ Namely, water molecules form two pairs, and each pair connects an amide H to the carbonyl O of the other peptide group (Figure 1/C). It was therefore expected and later positively confirmed that four water molecules placed in such a way around HCO-L-Ala- $\mathrm{NH}_{2}$ would keep the H-bond pattern along the whole reaction path between conformational regions $\beta_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$ (Figure S-1). Hence a $\beta_{\mathrm{L}}-\varepsilon_{\mathrm{L}}$ path allows the same polarized moieties to remain engaged in H -bonds with explicit water molecules during the whole scan. It is thus ideal for the comparison of approaches with and without explicit water molecules.

The $\beta_{\mathrm{L}}-\varepsilon_{\mathrm{L}}$ path was traced at five values of torsional angle $\phi: 180^{\circ}, 210^{\circ}, 240^{\circ}, 270^{\circ}$ and $300^{\circ}$ (i.e. $-180^{\circ},-150^{\circ},-120^{\circ},-90^{\circ}$ and $-60^{\circ}$ ). Five different RHF approaches (methods I to M in Table 1) were used. Method J is the only one where the level of optimization (RHF/321 G ) differs from that of the CS calculation (RHF/6-311++g(d,p)). Methods J, K and L involve 4 explicit water molecules in the model. In methods I and L the PCM reaction field is applied. Method M refers to a pure gas phase calculation. This is the only one where more restriction had to be applied than that of the torsional angle $\phi$. Because the conformation at $\phi=270^{\circ}$ without any constrain on $\psi$ converged to the $\gamma_{\mathrm{L}}$ region, $\psi$ was set equal to $144^{\circ}$. Relative CSs $(\Delta \delta)$ at each inspected level of theory were referenced to the respective chemical shielding values of the conformation with $\phi=270^{\circ}$ (i.e. $-90^{\circ}$ ).

## The reaction path from conformer $\alpha_{L}$ through $\gamma_{L}$ reaching either $\beta_{L}$ or $\varepsilon_{L}$

Reaction path along the torsional angle $\psi$ was traced by method I (RHF/PCM/6$311+\mathrm{G}(\mathrm{d}, \mathrm{p})$ level of theory) because it performed among the best methods ( $\mathrm{F}, \mathrm{H}$ and I ) in reflecting experimental protein data (see section bellow). As far as initial conformations are concerned, the tracing of two paths was attempted: $\delta_{\mathrm{D}}-\delta_{\mathrm{L}}-\beta_{\mathrm{L}}$ and $\alpha_{\mathrm{L}}-\gamma_{\mathrm{L}}-\beta_{\mathrm{L}}$. At the grid points, torsional angle $\psi$ was fixed at $-90^{\circ},-60^{\circ},-30^{\circ}, 0^{\circ}, 30^{\circ}, 60^{\circ}, 90^{\circ}, 120^{\circ}, 150^{\circ}$ and $180^{\circ}$, and all other parameters were let relax. Relative CSs $(\Delta \delta)$ were referenced to the respective chemical shielding values of the right-handed helical $\alpha_{\mathrm{L}}$ conformer.

## Results and Discussion

## Detection and characterization of fully optimized alanine conformers

Optimizations of all the nine ideal conformers were attempted by five different methods using the polarizable continuum model (Table 2, Figures 2 and 3/A). Only seven of them were found to be genuine minima, while conformations $\delta_{\mathrm{D}}$ and $\delta_{\mathrm{L}}$ converged to other ones. Additionally, conformer $\gamma_{\mathrm{L}}$ disappeared on optimization with Gaussian 03 (method A: at PCM/RB3LYP/6-311++G(2d,2p) level of theory). This result is in accordance with the previously published ${ }^{[12]}$ sets of conformers obtained with Gaussian 03 at levels RHF/PCM/6$31+G(d)$, B3LYP/PCM/6-31+G(d) and B3LYP/PCM/6-311++G(d,p). On the contrary, conformer $\gamma_{\mathrm{L}}$ proved to be a genuine minimum with Gaussian 09 and has low relative energy (between 1 and $1.5 \mathrm{kJ.mol}^{-1}$ by RB3LYP methods and between 4 and $4.4 \mathrm{kJ.mol}^{-1}$ by RHF methods, see Table 2).

The position of a conformer relative to $\alpha_{L}$ on the 2D relative chemical shielding plots $\left(\Delta \delta H^{\alpha}-\Delta \delta \mathrm{C}^{\alpha}, \Delta \delta \mathrm{C}^{\alpha}-\Delta \delta C^{\beta}\right.$ and $\left.\Delta \delta \mathrm{H}^{\alpha}-\Delta \delta C^{\beta}\right)$ is qualitatively independent of the applied method (Figures 3/B-C). Numerical differences, however, occur (Table S-I). When PCM is not included in the CS calculation (methods C and G) $\Delta \delta H^{\alpha}$ values shift upfield for all conformers without exception (Figures 3/B and C). This is because the absolute chemical shielding ( $\sigma$ ) decreases for $\alpha_{L}$ while increases for all other conformers. When compared to $a b$ initio (methods F to G), DFT results (methods A to E) give downfield shifted $\Delta \delta \mathrm{C}^{\alpha}$ for conformer $\gamma_{D}$, and upfield shifted $\Delta \delta C^{\beta}$ for conformers $\alpha_{D}, \varepsilon_{D}$ and $\gamma_{L}$. With all methods, the great shifts of the rare conformers ( $\alpha_{D}$ and $\gamma_{D}$ and $\varepsilon_{D}$ ) are striking (downfield $\delta C^{\alpha}$, upfield $\delta C^{\beta}$ and $\Delta \delta H^{\alpha}$ ), however, not observed in proteins.

Nuclei $\mathrm{C}^{\alpha}, \mathrm{H}^{\alpha}$ and $\mathrm{C}^{\beta}$ were previously selected as those that best reflect the secondary structure of amino acid residues by their relative CSs. ${ }^{[8]}$ Relevant experimental results published there are now quoted and compared to $\Delta \delta \mathrm{C}^{\alpha}, \Delta \delta \mathrm{H}^{\alpha}$ and $\Delta \delta \mathrm{C}^{\beta}$ of the seven optimized conformers (methods B to I, Table S-II). The best performance is obtained by methods F, H and I, where $R^{2}$ is $0.42-0.43$ for $\Delta \delta C^{\alpha}, 0.69-0.71$ for $\Delta \delta H^{\alpha}$ and 0.69-0.75 for $\Delta \delta C^{\beta}$. Correlation coefficients are highly similar to relevant data $(0.42,0.80$, and 0.73 , respectively) obtained by a very extensive gas phase investigation carried out on oligoalanine models which involved even double-stranded $\beta$-sheets. ${ }^{[8]}$ Without differentiating between
inward and outward looking $\mathrm{H}^{\alpha}$ nuclei, correlation coefficients for the nine conformers of single stranded oligoalanines would lower to $0.37,0.65$, and 0.42 , respectively.

There are analogous gas phase data on the alanine diamide model available in the literature. ${ }^{[9]}$ Collating them to the above quoted experimental results, one theoretical level proves to be adequate $\left(R^{2}\right.$ is 0.41 for $\Delta \delta C^{\alpha}, 0.69$ for $\Delta \delta \mathrm{H}^{\alpha}$ and 0.56 for $\Delta \delta \mathrm{C}^{\beta}$ at GIAO-RHF/TZ2P//B3LYP/6-311++G** when conformer $\alpha_{\mathrm{L}}$ is represented by a model with $\alpha$-helixlike backbone constrains). This level predicts, however, that the $\delta \mathrm{H}^{\alpha}$ difference between $\beta_{\mathrm{L}}$ and $\alpha_{\mathrm{L}}$, is as small as $0.04 \mathrm{ppm} .{ }^{[9]}$ This is in poor accordance with the experimental data of $0.57 \mathrm{ppm} .{ }^{[8]}$ The same value ranges between 0.15 and 0.49 when the CSs are calculated with PCM (methods A-B, D-F and H-I, Table S-II) and becomes negative when PCM is not involved in the CS calculation (methods C and G). All in all, the very important advantage of the application of the PCM model lies in its simplicity coupled with fair accuracy of the obtained results even on a model as small as a diamide.

Correlations between the applied theoretical levels are tabulated in the Supporting Information (Table S-III). As mentioned above, conformer $\gamma_{\mathrm{L}}$ disappears when treated with G03 instead of G09. Correlation between the two program packages (methods A and B), $R^{2}$ is the worst (i.e. 0.29 ) for $\Delta \delta \mathrm{H}^{\mathrm{NH}}, 0.93$ for $\Delta E$ and $\Delta \delta \mathrm{N}^{\mathrm{NH}}$, and 0.98 -1.00 for all other relative CSs and the torsional angles. No significant difference is detected when the model is changed from HCO-L-Ala- $\mathrm{NH}_{2}$ to $\mathrm{CH}_{3} \mathrm{CO}-\mathrm{L}-\mathrm{Ala}-\mathrm{NHCH}_{3}$ (methods E vs. D, and I vs. H) as $R^{2}$ is 1.000 for $\phi$ and $\psi, 1.00$ for $\Delta E, 0.99$ for $\Delta \delta \mathrm{C}^{\alpha}, \Delta \delta \mathrm{H}^{\alpha}$ and $\Delta \delta \mathrm{C}^{\beta}$, and not less than 0.95 for other CSs. Ab initio and DFT results can be scaled excellently to each other (methods B vs. F, D vs. H) with $R^{2}$ as large as 0.96 , except for $\Delta E, \Delta \delta C^{\beta}$ and $\Delta \delta C^{\prime}\left(R^{2}=0.8-0.9\right)$. When PCM is neglected in the single point calculation of energies and CSs (methods B vs. D, F vs. G), correlation is poor for $\Delta E$ and $\Delta \delta C^{\prime}\left(R^{2}=0.4-0.6\right)$, but excellent for aliphatic carbons ( $R^{2}=0.99$ ).

## Comparison of a reaction path calculated with different approaches

Along the $\beta_{\mathrm{L}}-\varepsilon_{\mathrm{L}}$ path (Table S-IV), the unrestricted torsional angle $\psi$ remained within the interval $130^{\circ}-165^{\circ}$ (Figure 3/A). The only exception is the grid point at $\phi=270^{\circ}$ (i.e. $-90^{\circ}$ ) by pure gas phase calculation (method M), where $\psi$ had to be constrained, and thus set to $144^{\circ}$, in order to avoid convergence to the $\gamma_{\mathrm{L}}$ region.

The local minimum with four explicit water molecules in the gas phase along the path is $\varepsilon_{\mathrm{L}}$ as obtained with the basis sets $3-21 \mathrm{G}$ and $6-311++\mathrm{G}(\mathrm{d}, \mathrm{p})\left(\phi=-92^{\circ}\right.$ and $-89^{\circ}$,
respectively (see methods J and K in Table S-IV and Figure 4/A). It is interesting to note that the RHF/6-311++g(d,p)/RHF/3-21G single point energy minimum is closer to $\phi=-120^{\circ}$ than to $-90^{\circ}$. The reason for this may be the difference in optimal values of parameters (bond lengths and angles) obtained by these two basis sets. Though the $\Delta E$ curves by methods J and K show only light increase up to $\phi=-150^{\circ}\left(\Delta E\right.$ is not greater than $\left.3.5 \mathrm{~kJ} \cdot \mathrm{~mol}^{-1}\right)$, the $\beta_{\mathrm{L}}$ minimum could not be optimized. A greater effort on optimally placing solvent molecules around the solute might have rendered a $\beta_{\mathrm{L}}$ minimum, as reported for $\mathrm{CH}_{3} \mathrm{CO}$-L-Ala$\mathrm{NHCH}_{3} \cdot{ }^{[19]}$ In that case, conformer $\beta_{\mathrm{L}}\left(\phi=-151^{\circ}\right)$ has an energy of $7.89 \mathrm{~kJ} . \mathrm{mol}^{-1}$ higher than that of the global minimum $\varepsilon_{\mathrm{L}}\left(\phi=-94^{\circ}\right)$. This energy difference is the double of the "single point" $\Delta E$ obtained by methods J and K .

Retaining the explicit solvent molecules and applying the PCM reaction field (method L) gives very similar results ( $\varepsilon_{\mathrm{L}}$ minimum at $\phi=-82^{\circ}$ ). On the contrary, the local minimum without explicit water molecules is $\beta_{\mathrm{L}}$ (methods I and M ). In the gas phase (method M ), neither conformer $\varepsilon_{\mathrm{L}}$, nor the grid point at $\phi=270^{\circ}$ (i.e. $-90^{\circ}$ ) without constraining $\psi$ could be optimized, instead both converge to the $\gamma_{\mathrm{L}}$ region. Applying the PCM (method I), however, $\varepsilon_{\mathrm{L}}$ is a low energy minimum ( $0.48 \mathrm{~kJ} . \mathrm{mol}^{-1}$ above conformer $\beta_{\mathrm{L}}$, method I).

Clearly, the PCM $\beta_{\mathrm{L}}-\varepsilon_{\mathrm{L}}$ path does not reflect exactly the $\Delta E$ curve obtained with four explicit water molecules. This should not be considered as a deficiency, however, when the goal is to interpret experimental protein Ramachandran plots, because $\beta_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$ appear there as two "conformational attractors" ${ }^{\text {" }}$ [20] of fusing spots. ${ }^{[12]}$

Contrary to energetic differences, relative CS, $\Delta \delta \mathrm{C}^{\alpha}, \Delta \delta \mathrm{C}^{\beta}$ and $\Delta \delta \mathrm{H}^{\alpha}$ appear very similar as obtained by the five inspected methods (Figures 4/B-D). $\Delta \delta \mathrm{C}^{\alpha}$ gives $\Lambda$-shaped, while $\Delta \delta \mathrm{H}^{\alpha}$ shows V -shaped curves with their respective maximum and minimum in the vicinity of $\phi=240^{\circ}$ (i.e. $-120^{\circ}$ ). $\Delta \delta C^{\beta}$ curves seem to have a minimum at a somewhat different value of $\phi$, between $210^{\circ}$ and $240^{\circ}$, where the $\mathrm{C}^{\prime}$ formy1 $-\mathrm{N}^{\mathrm{NH}}-\mathrm{C}^{\alpha}-\mathrm{C}^{\beta}$ torsion is almost $90^{\circ}$. The various methods give very similar $\Delta \delta C^{\beta}$ data, except for method J. After fitting a line on points (either chemical shielding, or absolute or relative chemical shielding, $\sigma, \delta, \Delta \delta$ ) of either $\mathrm{C}^{\alpha}$ or $\mathrm{H}^{\alpha}$ with $\phi$ greater than $240^{\circ}$ and another line on those with $\phi$ not less than $240^{\circ}$, these lines intercept at $241^{\circ}<\phi<250^{\circ}$. By method I, the intercept for $\Delta \delta \mathrm{C}^{\alpha}$ and $\Delta \delta \mathrm{H}^{\alpha}$ is equally at $246^{\circ}\left(-114^{\circ}\right)$. When $\phi \approx-114^{\circ}, \mathrm{H}^{\alpha}$ lies in the plane of the first amide group ( $\mathrm{C}^{\prime}$ formy1- $\mathrm{N}^{\mathrm{NH}}-\mathrm{C}^{\alpha}-$ $\mathrm{H}^{\alpha}$ torsion is about zero). At lower values, $\mathrm{H}^{\alpha}$ is on the same side of the first amide plane as $\mathrm{C}^{\prime}$, while at greater values, $\mathrm{H}^{\alpha}$ moves to the side of $\mathrm{C}^{\beta}$. For values $-180^{\circ}<\phi<-114^{\circ}$ (given in
degree), $\sigma \mathrm{C}^{\alpha}$ is approximately $(0.05698 \phi+159.4) \mathrm{ppm}$, while $\sigma \mathrm{H}^{\alpha}$ is $(-0.008155 \phi+27.14) \mathrm{ppm}$ by method I. For values $-114^{\circ}<\phi<-60^{\circ}, \sigma \mathrm{C}^{\alpha}$ is about $(-0.07446 \phi+144.4) \mathrm{ppm}$, while $\sigma \mathrm{H}^{\alpha}$ is $(-0.01520 \phi+29.80) \mathrm{ppm}$. As a consequence, a conformation in the $\beta_{\mathrm{L}}$ region and another in $\varepsilon_{\mathrm{L}}$ may be equally far away from $\phi=114^{\circ}$, and therefore both their $\delta \mathrm{C}^{\alpha}$ and $\delta \mathrm{H}^{\alpha}$ values are identical. In such case, their assignation to either of these regions has to rely on $\delta \mathrm{C}^{\beta}$.

## Reaction path from conformer $\alpha_{L}$ through $\gamma_{L}$ reaching either $\beta_{L}$ or $\varepsilon_{L}$

According to the previous section, a path of low energy barrier exists between conformers $\beta_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$. In contrast, the search along the torsional angle $\psi$ (Figures 5 and S-3) does not indicate whether both conformers $\beta_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$ are directly connected to $\gamma_{\mathrm{L}}$, or only one of them. Torsional angles $\phi$ and $\psi$ may differ from the ideal values obtained by diamide optimizations. It is therefore very important to know, what overlaps are to be expected among the most populated conformers. According to Figures 5/C-E, characteristic $\mathrm{C}^{\alpha}, \mathrm{C}^{\beta}$ and $\mathrm{H}^{\alpha} \mathrm{CSs}$ of conformers $\varepsilon_{\mathrm{L}}$ and $\gamma_{\mathrm{L}}$ can be positioned between typical helical and extended values:

$$
\begin{align*}
& \Delta \delta C^{\alpha}\left(\alpha_{\mathrm{L}}\right) \gg \Delta \delta \mathrm{C}^{\alpha}\left(\varepsilon_{\mathrm{L}}\right)>\Delta \delta \mathrm{C}^{\alpha}\left(\beta_{\mathrm{L}}\right)>\Delta \delta C^{\alpha}\left(\gamma_{\mathrm{L}}\right)  \tag{2}\\
& \Delta \delta C^{\beta}\left(\beta_{\mathrm{L}}\right)>\Delta \delta \mathrm{C}^{\beta}\left(\alpha_{\mathrm{L}}\right) \approx \Delta \delta C^{\beta}\left(\varepsilon_{\mathrm{L}}\right)>\Delta \delta C^{\beta}\left(\gamma_{\mathrm{L}}\right)  \tag{3}\\
& \Delta \delta \mathrm{H}^{\alpha}\left(\beta_{\mathrm{L}}\right) \approx \Delta \delta \mathrm{H}^{\alpha}\left(\gamma_{\mathrm{L}}\right)>\Delta \delta \mathrm{H}^{\alpha}\left(\varepsilon_{\mathrm{L}}\right) \approx \Delta \delta \mathrm{H}^{\alpha}\left(\alpha_{\mathrm{L}}\right) \tag{4}
\end{align*}
$$

Thus both $\varepsilon_{\mathrm{L}}$ and $\gamma_{\mathrm{L}}$ give rather extended-like $\delta C^{\alpha}$ (upfield), and helical-like $\delta C^{\beta}$ (upfield), while $\delta \mathrm{H}^{\alpha}$ is extended-like for $\gamma_{\mathrm{L}}$ (downfield) and helical-like for $\varepsilon_{\mathrm{L}}$ (upfield). These tendencies may explain several cases of secondary structure prediction where it is contradicting between the nuclei $\mathrm{C}^{\alpha}, \mathrm{C}^{\beta}$ or $\mathrm{H}^{\alpha}$. When confronting the above inequalities with experimental results of reference [8], only the upfield shift of $\Delta \delta C^{\beta}\left(\gamma_{\mathrm{L}}\right)$ in (3) cannot be traced on the protein alanine 2D CS plots.

## Conclusion

We present a simple computational scheme for the estimation of the secondary structure dependence of backbone chemical shifts (CS) in proteins. The application of the polarizable continuum model (PCM) to CS calculations seems plausible as this reaction field proved to improve the performance of a diamide model for an amino acid residue as far as geometry and conformational abundance in proteins is concerned. ${ }^{[12]}$ Relative energies calculated on diamides in gas phase do not succeed in reflecting experimental propensities of protein amino acid backbone folds as conformations $\alpha_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$ ( $\alpha$-helical and polyproline II) converge to other minima, and the global energy minimum is $\gamma_{\mathrm{L}}$ (inverse $\gamma$-turn). This problem arises from unsatisfied polarized groups, and can be settled either by longer models, or adding explicit solvent molecules, but the computationally simplest solution is treating the surroundings as a polarizable continuum by the PCM reaction field.

Prediction of relative chemical shifts also improves on applying PCM, as the experimentally well characterized $\delta H^{\alpha}$ difference between $\beta_{\mathrm{L}}$ and $\alpha_{\mathrm{L}}$ conformers (extended and $\alpha$-helical) is now approximated, while it is negligible in gas phase results of diamides. Protein relative CSs are best reflected by GIAO-RHF/PCM methods at $6-311++\mathrm{G}(2 \mathrm{~d}, 2 \mathrm{p})$ or $6-311++\mathrm{G}(\mathrm{d}, \mathrm{p})$ level of theory calculated on the seven conformers optimized at identical levels, where $R^{2}$ is $0.42-0.43$ for $\Delta \delta \mathrm{C}^{\alpha}, 0.69-0.71$ for $\Delta \delta H^{\alpha}$ and $0.69-0.75$ for $\Delta \delta \mathrm{C}^{\beta}$. Correlation coefficients are highly similar to relevant previously published data ( $0.42,0.80$, and 0.73 , respectively ${ }^{[8]}$ obtained by a very extensive gas phase investigation carried out on several residue long oligoalanine models involving even double-stranded $\beta$-sheets. Inclusion of explicit solvent molecules appears unnecessary when PCM is applied.

Through path calculation connecting the lowest energy minima, i.e., $\beta_{\mathrm{L}}$, $\gamma_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$, the well known and widely used distinction between helical $\left(\alpha_{\mathrm{L}}\right)$ and extended $\left(\beta_{\mathrm{L}}\right)$ residues on the ground of relative CSs is extended to the less frequent $\gamma_{\mathrm{L}}$ and $\varepsilon_{\mathrm{L}}$ conformers: Thus both $\varepsilon_{\mathrm{L}}$ and $\gamma_{\mathrm{L}}$ give rather extended-like $\delta \mathrm{C}^{\alpha}$ (upfield), and helical-like $\delta \mathrm{C}^{\beta}$ (upfield), while $\delta \mathrm{H}^{\alpha}$ is extended-like for $\gamma_{\mathrm{L}}$ (downfield) and helical-like for $\varepsilon_{\mathrm{L}}$ (upfield). This explains several cases of secondary structure prediction where it is contradicting between the nuclei $\mathrm{C}^{\alpha}, \mathrm{C}^{\beta}$ or $\mathrm{H}^{\alpha}$.

## Figure captions

Figure 1. The investigated model peptides. A) Model peptide $\mathrm{HCOAlaNH}_{2}$ with the labels of the most important nuclei. B) Model peptide $\mathrm{CH}_{3} \mathrm{COAlaNHCH}_{3}$ with the definition of torsional angles $\phi$ and $\psi . \mathbf{C}$ ) Model composed of the peptide $\mathrm{HCOAlaNH}_{2}$ and four explicit water molecules.

Figure 2. The Ramachandran map defined by torsional angles $\phi$ and $\psi$. The nine regions are named according to Perczel et al. ${ }^{[7]}$ The seven minima of the $\mathrm{HCOAlaNH}_{2}$ model optimized with PCM is presented to characterize the relevant conformational regions.

Figure 3. Geometry and relative chemical shielding of optimized alnine diamide conformers: $\alpha_{\mathrm{D}}\left(\right.$ black + ), $\alpha_{\mathrm{L}}($ cyan x$), \beta_{\mathrm{L}}$ (black triangle), $\gamma_{\mathrm{D}}$ (red hollow square), $\gamma_{\mathrm{L}}$ (red square), $\varepsilon_{\mathrm{D}}$ (green hollow diamond), $\varepsilon_{\mathrm{L}}$ (green diamond). For the applied levels of theory see Table I. Several data points are signed by the letter of the applied method. A) Ramachandran plot, B) $\Delta \delta H^{\alpha}-\Delta \delta C^{\alpha}$ plot, C) $\Delta \delta C^{\alpha}-\Delta \delta C^{\beta}$ plot, D) $\Delta \delta H^{\alpha}-\Delta \delta C^{\beta}$ plot. The origin of every relative chemical shift scale is set at the right-handed helical $\alpha_{L}$ conformer (see Table S-I).

Figure 4. Relative energy and chemical shifts of alanine diamide conformations optimized along the $\beta_{\mathrm{L}}-\varepsilon_{\mathrm{L}}$ path. Data points at $\phi=180^{\circ}, 210^{\circ}, 240^{\circ}, 270^{\circ}$ and $300^{\circ}$ (i.e. at $-180^{\circ},-150^{\circ}$, $-120^{\circ}$, $-90^{\circ}$ and $-60^{\circ}$ refer to structures of constrained $\phi$. All other structures are fully optimized $\beta_{\mathrm{L}}$ or $\varepsilon_{\mathrm{L}}$ conformers. A) Relative energy over the (lower) minimum of the path. B) $\Delta \delta \mathrm{C}^{\alpha}$, C) $\Delta \delta C^{\beta}$, D) $\Delta \delta \mathrm{H}^{\alpha}$. Relative chemical shifts are referenced to the data point at $270^{\circ}$ (i.e. $-90^{\circ}$ ).

Figure 5. Geometry, energy and relative chemical shifts of alanine diamide conformations optimized along the $\alpha_{L}-\gamma_{\mathrm{L}}-\varepsilon_{\mathrm{L}}-\beta_{\mathrm{L}}$ paths. Data points at $\psi=-90^{\circ},-60^{\circ},-30^{\circ}, 0^{\circ}, 30^{\circ}, 60^{\circ}, 90^{\circ}$, $120^{\circ}, 150^{\circ}$ and $180^{\circ}$ refer to structures of constrained $\psi$. All other structures are fully optimized $\alpha_{\mathrm{L}}, \gamma_{\mathrm{L}}, \varepsilon_{\mathrm{L}}$ or $\beta_{\mathrm{L}}$ conformers. A) Ramachandran plot. B) Relative energy over that of the $\alpha_{\mathrm{L}}$ conformer. C) $\left.\left.\Delta \delta \mathrm{C}^{\alpha}, \mathbf{D}\right) \Delta \delta \mathrm{C}^{\beta}, \mathbf{E}\right) \Delta \delta \mathrm{H}^{\alpha}$. Relative chemical shifts are referenced to the right-handed helical $\alpha_{\mathrm{L}}$ conformer.

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Table 1. Methods for chemical shielding calculations applying different levels of theory.

| Method | Gaussian program package | Model peptide | Level of theory |
| :---: | :---: | :---: | :---: |
| A | G03 | $\mathrm{HCOAlaNH}_{2}$ | GIAO-RB3LYP/PCM/6-311++G(2d,2p)//RB3LYP/PCM/6-311++G(2d,2p) |
| B | G09 | HCOAlaNH2 | GIAO-RB3LYP/PCM/6-311++G(2d,2p)//RB3LYP/PCM/6-311++G(2d,2p) |
| C | G09 | $\mathrm{HCOAlaNH}_{2}$ | GIAO-RB3LYP/6-311++G(2d,2p)//RB3LYP/PCM/6-311++G(2d,2p) |
| D | G09 | $\mathrm{CH}_{3} \mathrm{COAlaNHCH}_{3}$ | GIAO-RB3LYP/PCM/6-311+G(d,p)//RB3LYP/PCM/6-311+G(d,p) |
| E | G09 | $\mathrm{HCOAlaNH}_{2}$ | GIAO-RB3LYP/PCM/6-311+G(d,p)//RB3LYP/PCM/6-311+G(d,p) |
| F | G09 | HCOAlaNH2 | GIAO-RHF/PCM/6-311++G(2d,2p)//RHF/PCM/6-311++G(2d,2p) |
| G | G09 | HCOAlaNH2 | GIAO-RHF/6-311++G(2d,2p)//RHF/PCM/6-311++G(2d,2p) |
| H | G09 | $\mathrm{CH}_{3} \mathrm{COAlaNHCH}_{3}$ | GIAO-RHF/PCM/6-311++G(d,p)//RHF/PCM/6-311++G(d,p) |
| I | G09 | $\mathrm{HCOAlaNH}_{2}$ | GIAO-RHF/PCM/6-311++G(d,p)//RHF/PCM/6-311++G(d,p) |
| J | G09 | $\mathrm{HCOAlaNH}_{2}+4 \mathrm{H}_{2} \mathrm{O}$ | GIAO-RHF/6-311++G(d,p)//RHF/3-21G |
| K | G09 | $\mathrm{HCOAlaNH}_{2}+4 \mathrm{H}_{2} \mathrm{O}$ | GIAO-RHF/6-311++G(d,p)//RHF/6-311++g(d,p) |
| L | G09 | $\mathrm{HCOAlaNH}_{2}+4 \mathrm{H}_{2} \mathrm{O}$ | GIAO-RHF/PCM/6-311++G(d,p)//RHF/PCM/6-311++g(d,p) |
| M | G09 | $\mathrm{HCOAlaNH}_{2}$ | GIAO-RHF/6-311++G(d,p)//RHF/6-311++g(d,p) |

[^1]Table 2. Chemical shielding and other data of optimized alnine diamide conformers.

| Conf | $\Delta E$ | $\phi$ | $\psi$ | $\sigma \mathbf{N}^{\text {NH }}$ | $\sigma \mathbf{H}^{\mathrm{NH}}$ | $\sigma C^{\alpha}$ | $\sigma \mathbf{H}^{\alpha}$ | $\sigma \mathrm{C}^{\beta}$ | $\sigma \mathrm{C}^{\prime}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Method: | A |  |  |  |  |  |  |  |  |
| $\alpha_{\text {D }}$ | 8.71 | 61.8 | 37.1 | 107.2 | 24.26 | 127.4 | 27.61 | 168.2 | 1.7 |
| $\alpha_{L}$ | 0.00 | -86.4 | -17.8 | 102.4 | 24.72 | 129.1 | 27.10 | 164.1 | -0.2 |
| $\beta_{\mathrm{L}} \mathrm{I}$ | 0.34 | -131.8 | 143.1 | 104.7 | 24.53 | 133.3 | 26.60 | 161.2 | 0.8 |
| $\varepsilon_{\text {D }}$ | 15.12 | 59.5 | -144.1 | 107.4 | 24.38 | 127.3 | 27.73 | 167.1 | 2.1 |
| $\varepsilon_{L}$ | 0.64 | -89.7 | 143.7 | 102.3 | 24.77 | 132.2 | 26.89 | 163.2 | -0.5 |
| $\gamma_{\text {D }}$ | 11.49 | 72.5 | -49.3 | 106.8 | 24.42 | 119.8 | 27.84 | 165.2 | -1.9 |
| Method: | B |  |  |  |  |  |  |  |  |
| $\alpha_{\text {D }}$ | 10.13 | 63.7 | 34.4 | 110.7 | 25.48 | 126.9 | 27.92 | 168.4 | 3.3 |
| $\alpha_{L}$ | 0.80 | -85.6 | -15.5 | 104.7 | 25.95 | 128.7 | 27.41 | 164.4 | 1.7 |
| $\beta_{L}$ | 0.00 | -153.6 | 158.6 | 109.0 | 25.00 | 131.3 | 27.12 | 161.5 | 2.6 |
| $\varepsilon_{\text {D }}$ | 14.40 | 58.8 | -141.6 | 110.1 | 25.56 | 127.0 | 28.08 | 167.1 | 3.6 |
| $\varepsilon_{L}$ | 3.02 | -75.2 | 146.9 | 104.6 | 26.00 | 130.6 | 27.43 | 164.3 | 1.2 |
| $\gamma_{\mathrm{D}}$ | 8.73 | 72.1 | -49.0 | 109.6 | 25.65 | 119.3 | 28.13 | 165.2 | -0.2 |
| $\gamma_{L}$ | 1.51 | -84.1 | 69.5 | 100.7 | 26.06 | 132.0 | 27.09 | 168.4 | 2.1 |
| Method: | C |  |  |  |  |  |  |  |  |
| $\alpha_{\text {D }}$ | 24.41 | 63.7 | 34.4 | 116.1 | 26.45 | 126.4 | 28.31 | 167.5 | 8.0 |
| $\alpha_{L}$ | 14.85 | -85.6 | -15.5 | 108.2 | 26.96 | 129.5 | 27.19 | 163.7 | 5.9 |
| $\beta_{\mathrm{L}}$ | 3.62 | -153.6 | 158.6 | 114.7 | 25.45 | 131.4 | 27.25 | 161.3 | 4.4 |
| $\varepsilon_{\text {D }}$ | 28.00 | 58.8 | -141.6 | 116.7 | 26.42 | 127.4 | 28.41 | 166.3 | 8.3 |
| $\varepsilon_{L}$ | 15.27 | -75.2 | 146.9 | 109.8 | 26.80 | 131.4 | 27.60 | 164.0 | 4.9 |
| $\gamma_{\text {D }}$ | 9.08 | 72.1 | -49.0 | 112.2 | 26.40 | 119.0 | 28.27 | 165.0 | 4.0 |
| $\gamma_{\mathrm{L}}$ | 0.00 | -84.1 | 69.5 | 102.7 | 26.82 | 132.4 | 27.18 | 167.7 | 5.6 |
| Method: | D |  |  |  |  |  |  |  |  |
| $\alpha_{D}$ | 10.28 | 64.5 | 33.3 | 113.8 | 25.85 | 126.0 | 28.32 | 168.6 | 3.8 |
| $\alpha_{L}$ | 0.00 | -89.6 | -12.2 | 107.6 | 26.40 | 127.4 | 27.80 | 164.1 | 2.6 |
| $\beta_{\mathrm{L}}$ | 0.88 | -151.9 | 154.8 | 111.0 | 25.44 | 129.7 | 27.64 | 161.6 | 2.9 |
| $\varepsilon_{\text {D }}$ | 14.96 | 58.5 | -139.3 | 111.6 | 25.78 | 126.3 | 28.47 | 167.7 | 4.1 |
| $\varepsilon_{\text {L }}$ | 3.08 | -74.5 | 142.5 | 105.8 | 26.24 | 128.3 | 27.94 | 164.3 | 1.2 |
| $\gamma_{\text {D }}$ | 9.18 | 73.4 | -53.2 | 111.9 | 25.89 | 118.5 | 28.47 | 165.6 | 0.7 |
| $\gamma_{L}$ | 1.04 | -85.1 | 71.0 | 101.7 | 26.36 | 130.5 | 27.56 | 168.1 | 2.8 |
| Method: | E |  |  |  |  |  |  |  |  |
| $\alpha_{\text {D }}$ | 9.29 | 64.2 | 33.7 | 110.5 | 25.94 | 126.7 | 28.17 | 168.1 | 3.7 |
| $\alpha_{L}$ | 0.31 | -91.4 | -10.6 | 104.5 | 26.42 | 129.4 | 27.58 | 163.9 | 2.2 |
| $\beta_{L}$ | 0.00 | -152.9 | 158.4 | 108.8 | 25.50 | 131.5 | 27.41 | 161.1 | 3.2 |
| $\varepsilon_{\text {D }}$ | 13.86 | 59.3 | -142.4 | 109.7 | 25.98 | 127.0 | 28.32 | 166.9 | 4.2 |
| $\varepsilon_{\text {L }}$ | 2.70 | -75.5 | 147.2 | 104.5 | 26.39 | 130.5 | 27.70 | 164.1 | 1.6 |
| $\gamma_{\text {D }}$ | 8.46 | 72.6 | -52.3 | 109.5 | 26.04 | 119.0 | 28.37 | 165.0 | 0.3 |
| $\gamma_{\mathrm{L}}$ | 1.22 | -84.4 | 70.1 | 100.5 | 26.42 | 132.0 | 27.41 | 168.1 | 2.8 |

(Table 2. continued)

| Method: | $\mathbf{F}$ |  |  |  |  |  |  |  |  |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\alpha_{\mathrm{D}}$ | 9.72 | 62.5 | 37.9 | 147.8 | 26.70 | 147.3 | 28.90 | 179.7 | 13.7 |
| $\boldsymbol{\alpha}_{\mathrm{L}}$ | $\mathbf{0 . 0 0}$ | $\mathbf{- 8 0 . 8}$ | $\mathbf{- 2 1 . 0}$ | 143.5 | 27.08 | 147.6 | 28.46 | 177.7 | 11.9 |
| $\beta_{\mathrm{L}}$ | 0.54 | -150.7 | 153.8 | 146.5 | 26.32 | 150.5 | 28.13 | 174.8 | 13.6 |
| $\varepsilon_{\mathrm{D}}$ | 15.20 | 58.6 | -142.7 | 146.7 | 26.62 | 147.3 | 29.01 | 179.0 | 13.9 |
| $\varepsilon_{\mathrm{L}}$ | 1.24 | -71.8 | 148.6 | 143.1 | 27.07 | 149.4 | 28.49 | 177.8 | 12.5 |
| $\gamma_{\mathrm{D}}$ | 13.36 | 75.7 | -45.1 | 146.8 | 26.88 | 141.4 | 28.99 | 178.3 | 10.6 |
| $\gamma_{\mathrm{L}}$ | 3.96 | -88.7 | 76.2 | 140.8 | 27.20 | 152.0 | 28.11 | 179.7 | 13.7 |
| Method: | $\mathbf{G}$ |  |  |  |  |  |  |  |  |
| $\alpha_{\mathrm{D}}$ | 22.89 | 62.5 | 37.9 | 152.4 | 27.74 | 146.8 | 29.36 | 179.0 | 18.4 |
| $\alpha_{\mathrm{L}}$ | 13.91 | -80.8 | -21.0 | 147.0 | 28.15 | 148.2 | 28.21 | 177.0 | 16.1 |
| $\beta_{\mathrm{L}}$ | 0.88 | -150.7 | 153.8 | 151.0 | 26.74 | 150.5 | 28.30 | 174.6 | 15.6 |
| $\varepsilon_{\mathrm{D}}$ | 25.91 | 58.6 | -142.7 | 152.0 | 27.50 | 147.4 | 29.36 | 178.5 | 18.2 |
| $\varepsilon_{\mathrm{L}}$ | 11.12 | -71.8 | 148.6 | 147.4 | 27.89 | 149.7 | 28.73 | 177.4 | 16.2 |
| $\gamma_{\mathrm{D}}$ | 11.20 | 75.7 | -45.1 | 148.9 | 27.69 | 141.2 | 29.14 | 178.0 | 14.5 |
| $\gamma_{\mathrm{L}}$ | $\mathbf{0 . 0 0}$ | $\mathbf{- 8 8 . 7}$ | $\mathbf{7 6 . 2}$ | 142.2 | 27.97 | 152.2 | 28.21 | 179.1 | 16.9 |
|  |  |  |  |  |  |  |  |  |  |
| Method: | $\mathbf{H}$ |  |  |  |  |  |  |  |  |
| $\alpha_{\mathrm{D}}$ | 10.19 | 62.9 | 37.6 | 148.7 | 27.17 | 147.1 | 29.25 | 180.0 | 14.3 |
| $\boldsymbol{\alpha}_{\mathrm{L}}$ | $\mathbf{0 . 0 0}$ | -80.6 | -21.6 | 143.9 | 27.55 | 146.6 | 28.81 | 177.9 | 12.4 |
| $\beta_{\mathrm{L}}$ | 1.87 | -151.7 | 153.0 | 146.9 | 26.70 | 149.4 | 28.56 | 175.1 | 13.9 |
| $\varepsilon_{\mathrm{D}}$ | 16.14 | 58.4 | -140.7 | 146.3 | 26.96 | 146.8 | 29.32 | 179.5 | 14.5 |
| $\varepsilon_{\mathrm{L}}$ | 1.75 | -70.2 | 145.9 | 142.5 | 27.39 | 148.1 | 28.93 | 178.0 | 12.5 |
| $\gamma_{\mathrm{D}}$ | 14.89 | 76.8 | -51.7 | 147.4 | 27.19 | 140.8 | 29.25 | 178.8 | 11.7 |
| $\gamma_{\mathrm{L}}$ | 4.40 | $\mathbf{- 8 9 . 5}$ | $\mathbf{7 7 . 9}$ | 140.1 | 27.56 | 150.9 | 28.52 | 179.8 | 14.4 |
|  |  |  |  |  |  |  |  |  |  |
| Method: | $\mathbf{I}$ |  |  |  |  |  |  |  |  |
| $\alpha_{\mathrm{D}}$ | 9.32 | 62.7 | 37.4 | 147.8 | 27.10 | 147.6 | 29.11 | 179.9 | 14.4 |
| $\boldsymbol{\alpha}_{\mathrm{L}}$ | $\mathbf{0 . 0 0}$ | $\mathbf{- 8 1 . 3}$ | $\mathbf{- 2 1 . 2}$ | 143.4 | 27.45 | 147.9 | 28.68 | 177.9 | 12.2 |
| $\beta_{\mathrm{L}}$ | 1.11 | -149.9 | 153.9 | 146.6 | 26.72 | 150.9 | 28.37 | 174.9 | 14.2 |
| $\varepsilon_{\mathrm{D}}$ | 15.18 | 59.2 | -144.0 | 146.6 | 26.99 | 147.6 | 29.18 | 179.1 | 14.7 |
| $\varepsilon_{\mathrm{L}}$ | 1.59 | -71.7 | 149.2 | 143.1 | 27.39 | 149.7 | 28.73 | 178.1 | 13.0 |
| $\gamma_{\mathrm{D}}$ | 13.45 | 76.0 | -50.4 | 147.1 | 27.19 | 141.3 | 29.16 | 178.6 | 11.6 |
| $\gamma_{\mathrm{L}}$ | 4.15 | -89.2 | 76.0 | 140.7 | 27.52 | 152.2 | 28.39 | 179.9 | 14.5 |

[^2]

Figure 1. The investigated model peptides. A) Model peptide $\mathrm{HCO}-\mathrm{Ala}-\mathrm{NH}_{2}$ with the labels of the most important nuclei. B) Model peptide $\mathrm{CH}_{3} \mathrm{CO}-\mathrm{Ala}-\mathrm{NHCH}_{3}$ with the definition of torsional angles $\varphi$ and $\psi$. C) Model composed of the peptide HCO-Ala- $\mathrm{NH}_{2}$ and four explicit water molecules.
$185 \times 148 \mathrm{~mm}$ ( $300 \times 300$ DPI)

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Figure 2. The Ramachandran map defined by torsional angles $\varphi$ and $\psi$. The nine regions are named according to Perczel et al. ${ }^{[7]}$ The seven minima of the HCO-Ala- $\mathrm{NH}_{2}$ model optimized with PCM is presented to characterize the relevant conformational regions. $153 \times 144 \mathrm{~mm}(300 \times 300$ DPI)


Figure 3. Geometry and relative chemical shielding of optimized alnine diamide conformers: $a_{D}$ (black + ), $a_{L}$ (cyan $x$ ), $\beta_{L}$ (black triangle), $Y_{D}$ (red hollow square), $Y_{L}$ (red square), $\varepsilon_{D}$ (green hollow diamond), $\varepsilon_{L}$ (green diamond). For the applied levels of theory see Table I. Several data points are signed by the letter of the applied method. A) Ramachandran plot, B) $\Delta \delta H^{a}-\Delta \delta C^{a}$ plot, C) $\Delta \delta C^{a}-\Delta \delta C^{\beta}$ plot, D) $\Delta \delta H^{a}-\Delta \delta C^{\beta}$ plot. The origin of every relative chemical shift scale is set at the right-handed helical $a_{\mathrm{L}}$ conformer (see Table S-I).

## $202 \times 209 \mathrm{~mm}$ ( $300 \times 300$ DPI)



Figure 4. Relative energy and chemical shifts of alanine diamide conformations optimized along the $\beta_{\mathrm{L}}-\varepsilon_{\llcorner }$ path. Data points at $\varphi=180^{\circ}, 210^{\circ}, 240^{\circ}, 270^{\circ}$ and $300^{\circ}$ (i.e. at $180^{\circ}, 150^{\circ}, 120^{\circ}, 90^{\circ}$ and $60^{\circ}$ refer to structures of constrained $\varphi$. All other structures are fully optimized $\beta_{\llcorner }$or $\varepsilon_{\llcorner }$conformers. A) Relative energy over the (lower) minimum of the path. B) $\Delta \delta \mathrm{C}^{a}$, C) $\Delta \delta \mathrm{C}^{\beta}$, D) $\Delta \delta \mathrm{H}^{a}$. Relative chemical shifts are referenced to the data point at $270^{\circ}$ (i.e. $90^{\circ}$ ).
$209 \times 148 \mathrm{~mm}(300 \times 300$ DPI)


E


Figure 5. Geometry, energy and relative chemical shifts of alanine diamide conformations optimized along the $a_{L}-\gamma_{L}-\varepsilon_{L}-\beta_{L}$ paths. Data points at $\psi=-90^{\circ},-60^{\circ},-30^{\circ}, 0^{\circ}, 30^{\circ}, 60^{\circ}, 90^{\circ}, 120^{\circ}, 150^{\circ}$ and $180^{\circ}$ refer to structures of constrained $\psi$. All other structures are fully optimized $a_{L}, \gamma_{L}, \varepsilon_{L}$ or $\beta_{L}$ conformers. A) Ramachandran plot. B) Relative energy over that of the $a_{L}$ conformer. C) $\Delta \delta C^{a}$, D) $\left.\Delta \delta C^{\beta}, ~ E\right) ~ \Delta \delta H^{a}$. Relative chemical shifts are referenced to the right-handed helical $a_{L}$ conformer.
$297 \times 209 \mathrm{~mm}(300 \times 300$ DPI)


Graphical abstract
$115 \times 79 \mathrm{~mm}$ ( $300 \times 300$ DPI)


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[^1]:    John Wiley \& Sons, Inc.

[^2]:    Method: A-I: as defined in Table I. Conf: conformer defined according to Perczel et al. ${ }^{[7]}$
    Global minima are set in bold (A: $\boldsymbol{\alpha}_{\mathbf{L}}, E_{\mathrm{h}}=-417.399952, \mathbf{B}: \boldsymbol{\beta}_{\mathbf{L}}, E_{\mathrm{h}}=-417.387894, \mathbf{C}: \boldsymbol{\gamma}_{\mathbf{L}}, E_{\mathrm{h}}=-417.372090, \mathbf{D}:$
    $\alpha_{\mathrm{L}}, E_{\mathrm{h}}=-496.026422, \mathbf{E}: \boldsymbol{\beta}_{\mathrm{L}}, E_{\mathrm{h}}=-417.375569, \mathbf{F}: \boldsymbol{\alpha}_{\mathrm{L}}, E_{\mathrm{h}}=-414.943855, \mathbf{G}: \boldsymbol{\gamma}_{\mathrm{L}}, E_{\mathrm{h}}=-414.924922, \mathbf{H}: \boldsymbol{\alpha}_{\mathrm{L}}$,
    $E_{\mathrm{h}}=-493.018634$, I: $\boldsymbol{\alpha}_{\mathbf{L}}, E_{\mathrm{h}}=-414.928398$, in hartree) $\Delta \boldsymbol{E}$ : relative energy above the global minimum in kJ.mol
    ${ }^{1}$. At methods $\mathbf{C}$ and $\mathbf{G}$, relative energies deduced from the relevant single point calculations are given.
    Backbone torsional angles $\phi$ and $\psi$ are given in degree, chemical shielding in ppm.

